# Blendstrings: An environment for computing with smooth functions

Robert M. Corless 27 July 2023, ISSAC Tromsø

Maple Transactions

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#### What this talk is about

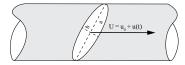


Figure 1: Pulsatile flow of blood in a vessel deformed to elliptic cross section

[Well, it's actually about my Maple code to evaluate Mathieu functions and the theory behind the code, but that hemodynamic problem was the motivating example.]

# Some of my papers related to this topic

- 0. Piers Lawrence & RMC, "Numerical stability of barycentric Hermite rootfinding" [Link to ACM Digital Library version] SNC 2011
- RMC & Erik Postma, "Blends in Maple" [arXiv Link] Maple Conference 2020
- Chris Brimacombe, RMC, & Mair Zamir, "Computation and applications of Mathieu functions: A historical perspective" SIAM Review 2021 [Open Access Link, courtesy MZ]
- RMC, "Sobre la iteración cúbica inversa", Gaceta de la Real Sociedad Matematica Española, 2021 [arXiV link] English version
- RMC, "Blends have decent numerical properties" [Link] Maple Transactions 2023
- Chris Brimacombe, RMC, & Mair Zamir, "Elliptic cross sections in blood flow regulation" Mathematics, Science, & Industry July 2023 [Open Access Link, courtesy MZ]

#### I wrote new code in Maple for this

I wrote an Hermite–Obreshkov ODE solver in Maple in order to approximate both Mathieu and modified Mathieu functions and provide a spectrally convergent expansion for the solution:

$$v(\xi,\eta) = \sum_{m\geq 0} b_{2m} \operatorname{ce}_{2m}(\eta) \operatorname{Ce}_{2m}(\xi)$$
 (1)

and thus solve the original model equation. The code I wrote uses blendstrings to approximate the Mathieu functions  $ce_{2m}(\eta)$  and the modified Mathieu functions  $Ce_{2m}(\xi)$  to arbitrary precision.

#### But why did I write my own special-purpose IVP solver?

- The Mathieu equations have double eigenvalues. All the codes that I
  knew of could not handle the double eigenvalue problem explicitly.
- Mathieu functions are useful: spectral convergence so only six terms needed for double precision accuracy
- I wanted an independent method whose solutions could be verified a posteriori
- Purely imaginary  $q = i\rho\omega d^2/4$  is needed [2.]
- The case of "near circularity" is actually very hard and one needs high precision.

## Could there be other uses for blendstrings?

- Efficient high-accuracy solution of IVP (or BVP) for D-finite ODE
- Solving delay differential equations by the method of steps! You can
  use Chebfun for this, so you could use blendstrings, instead!
  Quadrature is so easy with blendstrings<sup>1</sup>.
- Ned Nedialkov and John Pryce have already implemented a similar method in a quite general way, for solving DAE. Their code DAETS works well. Perhaps some experimental features of this Maple code could influence future development of DAETS.

http://www.cas.mcmaster.ca/~nedialk/daets/

 $<sup>^{\,\,\</sup>mathrm{I}}$  l haven't tried this yet but I will soon. I need to implement discontinuity handling first.

## **Blends and Blendstrings**

To implement this code, I first wrote (with Erik Postma) an **efficient** and **numerically stable** evaluator for what we call "blends" (arbitrary degree two-point Hermite interpolational polynomials), together with routines for manipulating them: integration, differentiation, rootfinding, addition, multiplication, etc. [1]

A "string of blends" is a particular kind of piecewise interpolational polynomial, that has some interesting properties, especially in the context of ODE solving.

Blendstrings are analogous to cubic splines<sup>2</sup>, but much smoother. Approximations on each subinterval are grade m+n+1 (usually m=n) and are m-times differentiable at the knots.

<sup>&</sup>lt;sup>2</sup>Splines are not blends because the derivatives at the knots are approximations, but pure Hermite cubic interpolation with exact derivatives  $f'(x_i)$  at the knots is a blend.

#### **Blends**

Suppose that we know some Taylor coefficients of a function at two distinct points, say z=a and z=b. Then put z=a+s(b-a) and the interval  $0 \le s \le 1$  determines a line segment in the z-plane.

Then (Hermite, Cours d'Analyse 1873)

$$H(s) = \sum_{j=0}^{m} p_j \sum_{k=0}^{m-j} {n+k \choose k} s^{k+j} (1-s)^{n+1}$$

$$+ \sum_{j=0}^{n} (-1)^j q_j \sum_{k=0}^{n-j} {m+k \choose k} s^{m+1} (1-s)^{k+j}$$
(2)

has  $H^{(j)}(0)/j!=p_j$  for  $0\leq j\leq m$  and  $H^{(j)}(1)/j!=q_j$  for  $0\leq j\leq n$ . Here differentiation is wrt s so one has to be careful about bookkeeping.

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## **Numerical stability**

Blends are *ridiculously* good numerically. The presence of those potentially large binomial coefficients suggests the opposite, but the doubly-recursive Horner implementation has been used successfully for degrees up to about 1000, even for very hard-to-approximate functions.

There's a couple of reasons for that success.

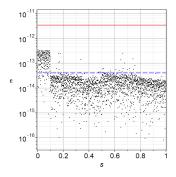
#### Backward stability result

- double Horner gives the *exact* value of a blend with Taylor coefficients  $p_j(1+\theta_j)$  and  $q_j(1+\psi_j)$ , where each  $|\theta_j| \leq \gamma_{3M+N}$  and  $|\psi_j| \leq \gamma_{3M+N}$ , with  $M = \max(m,n)$  and  $N = \min(m,n)$ .
- ullet proof uses the positivity of  $s^a(1-s)^b$  on  $0\leq s\leq 1$

**Zero coefficients**  $p_j = 0$  and  $q_j = 0$  are not disturbed. This could work for *lacunary* series. This is a *very strong* backward stability result (cf. those of Alicja Smoktunowicz for the Clenshaw–Curtis algorithm).

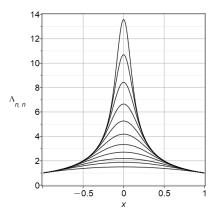
Here  $\gamma_n = nu/(1-nu) = nu + O(u^2)$  where u is the unit roundoff. For double precision,  $u \approx 10^{-16}$ . See e.g. Higham's Accuracy and Stability of Numerical Algorithms. This proof and all its details are in the Maple Transactions paper [4].

# A grade 1000 example



**Figure 2:** The backward error for a random example, computed in double precision. The Taylor coefficients  $p_j$  and  $q_j$  were pseudorandomly generated; m=368 and n=999-m=631 so the grade is m+n+1=1000. The red line is the theoretical bound  $\gamma_{3M+N}$ , the blue dashed line is a crude random walk model of the error  $\gamma_{\sqrt{3M+N}}$ , and the black dots are the actual backward error computed by the Oettli–Prager theorem.

# Lebesgue constants on $-1 \le x \le 1$



**Figure 3:** The Lebesgue function of balanced blends on [-1,1] grows like  $2\sqrt{n/\pi}$  where the grade is 2n+1.

Lebesgue constants *must* be unbounded on [-1,1]; optimal is  $O(\ln n)$ , like Lagrange interpolation on Chebyshev nodes. But  $\sqrt{n}$  is pretty decent.

#### On $0 \le s \le 1$ the Lebesgue constant is 2

The Lebesgue function, which gives a bound for the condition number, is (for blends) the polynomial you get with all series coefficients 1 at the left and all coefficients  $(-1)^j$  on the right:

$$L_{m,n}(s) = \sum_{j=0}^{m} \sum_{k=0}^{m-j} {n+k \choose k} s^{k+j} (1-s)^{n+1}$$

$$+ \sum_{j=0}^{n} \sum_{k=0}^{n-j} {m+k \choose k} s^{m+1} (1-s)^{k+j}$$
(3)

We can show that on  $0 \le s \le 1$ ,  $L_{m,m}(s) < 2$  and indeed  $L_{m,m}(s) \le L_{m,m}(1/2) = 2 - 2^{-2m+1} {2m+2 \choose m+1} \sim 2 - 2/\sqrt{\pi m} + O(1/m^{3/2})$ .

Bernstein polynomial bases  $\phi_j^m = {m \choose j} s^j (1-s)^{m-j}$  are better, with Lebesgue constant 1. But 2 is good.

#### The proof is very pretty.

One first proves by induction (and a contour integral) that

$$L_{m,m}(s) - L_{m-1,m-1}(s) = \frac{1}{m+1} {2m \choose m} s^m (1-s)^m.$$
 (4)

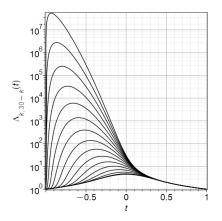
Then the Lebesgue function can be written

$$L_{m,m}(s) = \sum_{j=0}^{m} \frac{1}{j+1} {2j \choose j} s^{j} (1-s)^{j}$$
 (5)

and this is, with x=s(1-s), a truncation of the ordinary generating function for Catalan numbers. It is maximal when s=1/2 and bounded above by 2.

Then we translate this result to the interval  $-1 \le t \le 1$  for "fair" comparison to other bases. Details in the Maple Transactions paper [4].

# Unbalanced blends are bad, though



**Figure 4:** The Lebesgue function of *unbalanced* blends on [-1,1] grows exponentially. Just as bad as equally-spaced interpolation nodes, really.

## Making it jump

If we put the series at the left to be  $1,0,0,\ldots$  and the series at the right to be  $-1,0,0,\ldots$ , no analytic function can do this. Truncating, say with m=368 and n=631 (so the grade m+n+1=1000) we get a polynomial, which plots below.

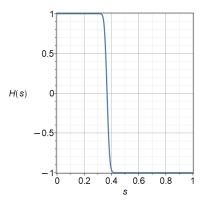


Figure 5: Look how smooth the plot is. No "Gibbs phenomenon" at all!

# Blendstrings

A string of blends, or blendstring, is a set of the form

$$\mathcal{B} := \left\{ L_k \right\}_{k=0}^M \tag{6}$$

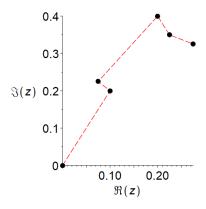
where each  $L_k$  is a list of the form

$$L_k := [\alpha_k, C_{k,0}, C_{k,1}, \dots, C_{k,m_k}]$$
 (7)

intended to represent the known Taylor coefficients  $C_{k,j}$  at the point  $z=\alpha_k$ .

Two blendstrings are *compatible* if they have the same knots in the same order and with the same degrees  $m_k$  at each knot. Then they can be added together, etc. If A and B are compatible blendstrings, then A op B can be a blendstring, where "op" is any of +, -, \*, / (if B is not zero), or even  $\hat{}$ .

# A possible blendstring



**Figure 6:** The knots  $\alpha_k$  are plotted with solid circles. Taylor coefficients are known at those knots. On the segment between any two knots, Hermite's formula gives a "blend" approximating the underlying function.

#### Quadrature

There is a very interesting (exact!) quadrature formula for a blend, which allows *indefinite* integration of blendstrings. Obreshkov (1940) had a version of it written with ratios of binomial coefficients<sup>3</sup>.

$$\int_{s=0}^{1} H(s) ds = \frac{(m+1)!}{(m+n+2)!} \sum_{j=0}^{m} \frac{(n+m-j+1)!}{(m-j)!(j+1)} p_{j} + \frac{(n+1)!}{(m+n+2)!} \sum_{j=0}^{n} \frac{(-1)^{j}(n+m-j+1)!}{(n-j)!(j+1)} q_{j}$$
(8)

There is a five-line Maple proof of this formula.

 $[{\sf MapleProofHermiteQuadrature.mw}] \ {\sf Or} \ {\sf you} \ {\sf can} \ {\sf use} \ {\sf contour} \ {\sf integration}.$ 

 $<sup>^3</sup>$ lt  $\emph{might}$  have been a reinvention. Hermite must have known it, surely? Or Darboux or Gauss.

#### Indefinite integration?

The definite integral (8) from the last slide allows one to (trivially) compute the Taylor coefficients for  $I(s) = \int_0^s H(\sigma) \, d\sigma$  at s=1: the zeroth order coefficient is now known (it's just the definite integral  $\int_0^1 H(\sigma) \, d\sigma$ ) and all the derivatives are simply related to the (known) derivatives of H(s) at s=1:

$$I'(s) = H(s) = \sum_{j=0}^{m} q_j (s-1)^j$$

$$I(s) = I(1) + \sum_{j=0}^{m} \frac{q_j}{j+1} (s-1)^{j+1}.$$
(9)

Likewise, I(0) = 0 is known, and all its derivatives at s = 0 are related in the same way.

Therefore, we can immediately find a blend for I(s) on  $0 \le s \le 1$ .

By propagating information from the previous subinterval, this gives us a blendstring for the integral of the function being approximated by the blendstring.

# Computational Version (new for this talk)

Put  $c_0 = (m+1)/(m+n+2)$  and  $d_0 = (n+1)/(m+n+2)$  and define

$$c_{j} = \frac{j(m-j+1)}{(j+1)(m+n+2-j)}c_{j-1} \quad \text{for} \quad 1 \le j \le m$$
 (10)

$$d_j = -\frac{j(n-j+1)}{(j+1)(m+n+2-j)}d_{j-1}$$
 for  $1 \le j \le n$ . (11)

(Notice the sign alternation in the  $d_j$ s.) Then

$$I(1) = \sum_{j=0}^{m} c_j p_j + \sum_{j=0}^{n} d_j q_j.$$
 (12)

#### Continued

If m = n then  $|d_j| = c_j$  and the formula becomes

$$I(1) = \sum_{j=0}^{m} c_j \left( p_j + (-1)^j q_j \right)$$

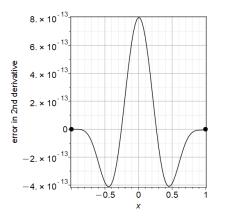
$$= \frac{1}{2} \left( p_0 + q_0 \right) + \frac{m}{4(2m+1)} \left( p_1 - q_1 \right) + \dots + c_m \left( p_m + (-1)^m q_m \right) .$$
(13)

The final coefficient  $c_m = 1/((m+1)\binom{2m+2}{m+1})$  is asymptotic to  $2^{-2m-2}\sqrt{\pi/m}$ . This method is *beautifully* stable numerically: componentwise backward error using IEEE 854 standard floats is bounded by  $\gamma_{m+n+2}$ .

It is not very accurate if m and n are greatly different, but if  $m \approx n$  it's wonderfully accurate when the underlying function has no nearby singularities.

#### **Differentiation**

The numerical routine Erik Postma and I wrote to evaluate blends also has the ability to evaluate arbitrary derivatives, by what I call "semi-automatic differentiation".



**Figure 7:** Error in 2nd derivative with m=7 at either end

#### Rootfinding

There is a **companion pencil** for a blend [0]: its eigenvalues are the roots of the polynomial expressed by the blend. Conversion to the monomial basis is avoided. It works well if the degrees are not too large (binomial coefficients *do* bother this method). There may be better methods.

One other method is to use Inverse Cubic Iteration [3] (an English version is on the arXiv). This method has convergence order  $1+\sqrt{3}=2.732$  with the same effort as Newton's method. It was inspired by thinking about low-order blends.

Since blends are most accurate on  $0 \le s \le 1$ , methods which concentrated on that interval might be the best of all. Contour integrals, perhaps.

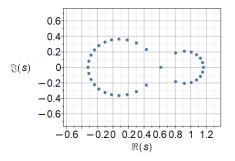
#### **Example**

Suppose m=2 and we know the Taylor coefficients up to the second derivative at s=0 and at s=1. Then the following matrix has H(s) as determinant:

$$\begin{bmatrix} 0 & \frac{p''(1)}{2} & p'(1) & p(1) & \frac{p''(0)}{2} & p'(0) & p(0) \\ -1 & \lambda - 1 & 0 & 0 & 0 & 0 & 0 \\ 3 & -1 & \lambda - 1 & 0 & 0 & 0 & 0 \\ -6 & 0 & -1 & \lambda - 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & \lambda & 0 & 0 \\ 3 & 0 & 0 & 0 & -1 & \lambda & 0 \\ 6 & 0 & 0 & 0 & 0 & -1 & \lambda \end{bmatrix}$$
 (14)

# **Hermite Jump Zeros**

If m=21 and n=13, with series  $1,0,0,\ldots$  at the left and  $-1,0,0,\ldots$  at the right, we get the following zeros using the companion pencil method in double precision. (NB: the very high degree example earlier needs 450 digit multiple precision to find eigenvalues of the companion pencil).



**Figure 8:** Zeros of the grade 21 + 13 + 1 = 55 polynomial fitting the flat series

#### In summary

#### With blendstrings we can

- Approximate smooth functions on an interval in a perfectly stable, decently conditioned, and a nearly spectrally efficient manner.
- Combine functions very efficiently:  $f \pm g$ ,  $f \cdot g$ ,  $f^g$ , and if  $g \neq 0$  we can find f/g.
- Differentiate and integrate the functions, because these operations are closed on sets of compatible blendstrings.
- Find zeros.
- Solve differential equations.

Therefore, as one can with Chebfun and Approxfun, one can do a kind of symbolic computation with smooth functions, sometimes *much* more rapidly than one can do exact computation.

#### Where next?

I would like to add Laurent series and Puiseux series, in order to explore the singularity detection and location facilities afforded by these interpolational polynomials.

I want to add *two-sided* blendstrings, so as to allow for the possibility of jump discontinuities at the knots.

I would like to add approximate *composition* (and *decomposition*) to the things one can do with blendstrings. Seems complicated, but interesting. [Link] Some recent work of Erik Postma

I promised to try to solve delay DE with blendstrings. That should be a simple job<sup>4</sup>. (Once we can handle discontinuities.)

Rootfinding is important. What's the best method to find the roots of a blendstring? (J.P. Boyd & L.N. Trefethen subdivision?)

<sup>&</sup>lt;sup>4</sup> "There's no such thing as a simple job." —Tim Daly

#### Thank you for listening.

This work supported by NSERC, by the Spanish MICINN, and by the Isaac Newton Institute in Cambridge.

I thank Michael Monagan for comments and Ned Nedialkov for the same and for sending me a copy of Obreshkov's 1940 paper. I also thank Erik Postma and my other co-authors, and especially I thank John C. Butcher for teaching me the contour integral technique for interpolation, by which I (re)derived all these formulae.