

# 1 Maxwell Equation and Electromagnetic Induction

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**Motivation:** Ten most important equations that changed the world. Among those:

$$a^2 + b^2 = c^2 \quad (\text{Pythagoras})$$

$$e^{i\pi} + 1 = 0 \quad (\text{Euler})$$

$$\vec{F} = m\vec{a}; \vec{F} = G \frac{mM}{r^2} \quad (\text{Newton})$$

(...Einstein, Schrödinger)

and Maxwell:

$$\nabla \cdot \vec{D} = \rho \quad (\text{Gauss})$$

$$\nabla \cdot \vec{B} = 0 \quad (\text{Gauss})$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (\text{Faraday})$$

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \quad (\text{Ampère-Maxwell})$$

$$\vec{D} = \varepsilon \varepsilon_0 \vec{E} \quad (\text{materials: electric field, dielectric field})$$

$$\vec{H} = \mu \mu_0 \vec{B} \quad (\text{materials: magnetizing field, magnetic induction})$$

$$\vec{j} = \sigma \vec{E} \quad (\text{Ohm's law})$$

and Integral:

$$\int \int_{\partial\Omega} \vec{D} \cdot d\vec{S} = \int \int \int_{\Omega} \rho \quad (\text{Gauss})$$

$$\int \int_{\partial\Omega} \vec{B} \cdot d\vec{S} = 0 \quad (\text{Gauss})$$

$$\int_{\partial\Sigma} \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int \int_{\Sigma} \vec{B} \cdot d\vec{S} \quad (\text{Faraday})$$

$$\int_{\partial\Sigma} \vec{H} \cdot d\vec{l} = \int \int_{\Sigma} \vec{J} \cdot d\vec{S} + \frac{\partial}{\partial t} \int \int_{\Sigma} \vec{D} \cdot d\vec{S} \quad (\text{Ampère-Maxwell})$$

$\Omega$  : Volume

$\Sigma$  : Surface

$\partial\Omega$  : Surface of volume

$\partial\Sigma$  : Edge of surface

Have a picture in mind for each one of them: (1) Source of a static E-field. (2) Dipole b-Field. (3) Induction with Magnet. (4) Displacement Currents and Bio-Savart law

Integral vs. Differential form: EMF, loops, LENZ law

Principle of self-induction

Table resistance frequency dependency

$$Z_L = j\omega L \quad V \text{ lags } I \quad Z_C = \frac{1}{j\omega C} \quad I \text{ lags } V \quad Z_R = R \text{ in phase}$$

Classification of electrical methods

Principles of the Slingram Method

$$I_{l1} = I_1 e^{i\omega t}$$

$$V_{l3} = L_{13} \frac{dI_{l1}}{dt} = i\omega L_{13} I_1 e^{i\omega t} \quad V_{l2} = L_{12} \frac{dI_{l1}}{dt} = i\omega L_{12} I_1 e^{i\omega t}$$

What is  $I_2$  with a R-L Subsurface model? Ohms law:

$$Z_{l2} = R + i\omega L_{l2} \quad V_{l2} = (R + i\omega L_{l2}) I_2$$

Induced voltage is balanced by inductance and resistance:

$$i\omega L_{12} I_1 e^{i\omega t} + (R + i\omega L_{l2}) I_2 = 0$$

$$I_{l2} = I_2 e^{i\omega t} = \frac{-i\omega L_{12}}{R_{l2} + i\omega L_{l2}} I_1 e^{i\omega t} \quad I_{l2} = I_2 e^{i\omega t} = \frac{-i\omega \frac{L_{12}}{R_{l2}}}{1 + i\omega \frac{L_{l2}}{R_{l2}}} \frac{L_{12}}{L_{l2}} I_1 e^{i\omega t}$$

$I_{l2}$  will produce a secondary B. How does that appear in Loop 3:

$$V_{l3} = L_{13} \frac{dI_{l1}}{dt} \text{ primary} \quad V_{l3} = L_{23} \frac{dI_{l2}}{dt} \text{ secondary}$$

$$\frac{U_s}{U_p} = -\frac{L_{12} L_{23}}{L_{13} L_{l2}} \left( \frac{i\omega \frac{L_{12}}{R_{l2}}}{1 + i\omega \frac{L_{l2}}{R_{l2}}} \right)$$

induction number:  $\alpha = \omega \frac{L_{l2}}{R_{l2}}$  helps to write the complex number in standard form:  $\frac{U_s}{U_p} = -\frac{L_{12} L_{23}}{L_{13} L_{l2}} \left( \frac{1}{1 + \alpha^2} (\alpha^2 + i\alpha) \right)$

## 2 Seismics

$$\frac{\partial^2 \Psi}{\partial t^2} = v^2 \frac{\partial^2 \Psi}{\partial x^2}$$

- This is an (undamped) wave equation.
- This is a second order (hyperbolic) partial differential equation.
- It is a function of space  $x$  and time  $t$ .

### 2.1 General solution

$$\Psi = f(x + vt) + f(x - vt)$$

Exampel:  $\Psi(x, t) = e^{k(x-vt)}, \sin(x - vt), (x - vt)^3$ .

However, only one solution complies with the given initial and boundary conditions. What is  $v$ ?

Draw Picture 4.4 Telford of  $\Psi$  of wave form  $\Psi$  on y-axis with an  $x_0$  and  $t_0 + \Delta t$  marked.

$$t_0 + \Delta t : \Psi_1(x_0 + \Delta x, t + t_0) = f(x_0 + \Delta x - v(t_0 + \Delta t))$$

$$t_0 : \Psi_1(x_0, t) = f(x_0 + \Delta x - vt_0)$$

$$x_0 - vt_0 = x_0 + \Delta x - v(t_0 + \Delta t)$$

- $x \pm vt$  is the phase
- $x \pm vt = \text{const.}$  wavefronts
- normal to the wavefront is the raypath

## 2.2 Specific solutions

### 2.2.1 Spherical waves

$$\frac{\partial^2 \Psi}{\partial t^2} = \frac{v^2}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \Psi}{\partial r} \right)$$

Solution:

$$\psi(x, t) = \frac{1}{r} f(r - vt)$$

Spherical waves the circles are wavefronts (i.e. lines of constant phase). Raypath are thus along  $r$ . In the farfield spherical waves can be approximated with plane waves.

*draw Figure 4.6 plane waves and spherical waves*

### 2.2.2 Harmonic waves

$$\psi(x, t) = A \cos(k(x - vt)) \quad (1)$$

$$\psi(x, t) = A/r \cos(k(x - vt)) \quad (2)$$

At fixed  $t$  if  $x$  increases by  $2\pi/k$  then everything repeats so

$$\lambda = \frac{2\pi}{k} \text{wavelength}$$

At fixed  $x$  it varies harmonically with time from  $-A$  to  $A$  (or  $-A/r$  to  $A/r$ ).  $A$  is the amplitude.

At fixed  $x$  repetition by  $T$  if  $kVT = 2\pi = 2\pi(vT/\lambda)$

$$T = \lambda/v \quad (3)$$

$$f = 1/T \quad (4)$$

$$v = \lambda f \quad (5)$$

With  $T$  period,  $f$  frequency,  $v$  phase velocity. Commonly used is the angular frequency  $\omega = 2\pi f$  so that

$$\Psi(x, t) = A \cos(k(x - vt)) = A \cos(kx - \omega t) \quad (6)$$

$$(7)$$

Sometimes an additional phase offset is used:

$$\Psi(x, t) = A \cos(k(x - vt)) = A \cos(kx - \omega t + \varepsilon) \quad (8)$$

$$(9)$$

### 2.3 Some critical steps in dipped layer refraction seismics

Goal: write  $t^-$  as a function of velocities and dip. Isolate slope and simplify y-offset.  $t^- = \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} + \frac{x \cos(\theta) - (d^- + d^+) \tan(i_c)}{v_2}$

$$\begin{aligned}
 t^- &= \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} + \frac{x \cos(\theta) - (d^- + d^+) \tan(i_c)}{v_2} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} - \frac{(d^- + d^+) \tan(i_c)}{v_2} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} - \frac{(d^- + d^+) \sin(i_c)}{\cos(i_c) v_2} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} - \frac{(d^- + d^+) v_1}{\cos(i_c) v_2^2} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{d^-}{v_1 \cos(i_c)} + \frac{d^+}{v_1 \cos(i_c)} - \frac{(d^- + d^+) v_1^2}{v_1 \cos(i_c) v_2^2} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{1}{v_1 \cos(i_c)} (d^- + d^+ - (d^- + d^+) \frac{v_1^2}{v_2^2}) \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{1}{v_1 \cos(i_c)} ((d^- + d^+) (1 - \frac{v_1^2}{v_2^2})) \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{(d^- + d^+) \cos(i_c)}{v_1}
 \end{aligned}$$

eliminate  $d^+$  using use:  $d^+ = d^- + x \sin(\theta)$  which is clear from geometry when you put  $\theta$  at the surface.

$$\begin{aligned}
 t^- &= \frac{x \cos(\theta)}{v_2} + \frac{(d^- + d^+) \cos(i_c)}{v_1} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{(2d^- + x \sin(\theta)) \cos(i_c)}{v_1} \\
 &= \frac{x \cos(\theta)}{v_2} + \frac{2d^- \cos(i_c)}{v_1} + \frac{x \sin(\theta) \cos(i_c)}{v_1} \\
 &= \frac{x \cos(\theta) \sin(i_c)}{v_2 \sin(i_c)} + \frac{2d^- \cos(i_c)}{v_1} + \frac{x \sin(\theta) \cos(i_c)}{v_1} \\
 &= \frac{x \cos(\theta) \sin(i_c)}{v_1} + \frac{2d^- \cos(i_c)}{v_1} + \frac{x \sin(\theta) \cos(i_c)}{v_1} \\
 &= \frac{x \sin(\theta + i_c)}{v_1} + \frac{2d^- \cos(i_c)}{v_1}
 \end{aligned}$$

$\theta$  and  $i_c$  are yet unknown but can be derived from the arcsin relation straightforward.

$$\frac{1}{v_2} + \frac{1}{v_2^+} = \frac{1}{v_1} \sin(i_c) \cos(\theta) = \frac{2}{v_2} \cos(\theta)$$