# PHY 300 HW 2, Carl Liu

### 3-1

Since F = ma = -kx. That means -ma/x = k. Thus

$$k = \frac{-(1*10^{-3}kg(-1.7514*10^{-2})\frac{m}{s^2})}{43.785*10^{-2}m} = 4*10^{-5}\frac{kg}{s^2} = 4*10^{-5}\frac{N}{m}$$

### 3-2

A) We have the period of oscillation is  $T = 2\pi/\omega$ . But  $\omega = \sqrt{k/m}$ . So the period of oscillation is  $T = 2\pi/\sqrt{k/m} = 2\pi\sqrt{m/k}$ 

B)

- 1) The force applied by the springs to the mass would be -kx + (-kx). This is essentially doubling the spring constant and so  $\omega = \sqrt{2k/m}$  making  $T = 2\pi\sqrt{m/2k}$
- 2) The spring constant stays the same for the whole spring because it is a longer version of the same spring. But because x describes only the displacement of one of the springs, we have F = -kx/2. So  $\omega = \sqrt{k/2m}$ . Meaning  $T = 2\pi\sqrt{2m/k}$ .

### 3-3

- A) It will leave the platform when x''<-g=-9.8m/s. This is because the force on the block is that of the platform and gravity where the platform will enact a force of g+x'' as long as x''>-g. So we have  $x''=-g=-5*10^{-2}\mathrm{m}\omega^2\cos(\omega t+\pi)$ , where  $\omega=2\pi(10/\pi)=20rad/s$ . So  $\cos(\omega t+\pi)=9.8/(5*10^{-2}*20^2)=9.8/(5*10^{-2}*20^2)=0.49$  Then  $x=5*10^{-2}\mathrm{m}\cos(\omega t+\pi)=5*10^{-2}\mathrm{m}*0.49=0.0245\mathrm{m}$ . This is the point where the acceleration of the platform is g and when the block will stop touching the platform.
- B) The height of the block will be dependent on the energy imparted onto it by the spring. Since  $K = mv^2/2$ , and  $v = x' = -A\omega\sin(\omega t + \pi)$  we have  $K = mA^2\omega^2\sin^2(\omega t + \pi)/2 = mA^2\omega^2(1-\cos(\omega t + \pi))/2$ . But K = mgh thus  $h = A^2\omega^2(1-\cos(\omega t + \pi))/g = 5*10^{-2}*20^2\mathrm{m}(1-0.49)/9.8 = 0.0388m$ . Since the initial position is 0.0245m, we have a final height of 0.0245m + 0.0388m = 0.0633m above displacement. Since the platform has an amplitude of 0.05m, we thus conclude that the block goes 0.0633m 0.05m = 0.0133m above the

max height of the platform.

### 3-13

We consider first a solution of the form  $Ce^{irt} = x$ . Then

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = -Cr^2 e^{rt} + iC\gamma r e^{rt} + C\omega_0^2 e^{irt} = 0$$

So for  $Ce^{rt}$  to be a solution we must have  $r^2 - i\gamma r - \omega_0^2 = 0$ . Meaning

$$r = \frac{i\gamma}{2} \pm \frac{\sqrt{-\gamma^2 + 4\omega_0^2}}{2}$$

We therefore have a solution of the form

$$C_1 e^{(\gamma t/2) + (i\sqrt{-\gamma^2 + 4\omega_0^2}t/2)} + C_2 e^{(\gamma t/2) - (i\sqrt{-\gamma^2 + 4\omega_0^2}t/2)} =$$

$$(C_1 + C_2)e^{\gamma t/2}(e^{i\sqrt{-\gamma^2 + 4\omega_0^2}t/2} + e^{-i\sqrt{-\gamma^2 + 4\omega_0^2}t/2}) = x$$

But we know that

$$(e^{i\sqrt{-\gamma^2 + 4\omega_0^2}t/2} + e^{-i\sqrt{-\gamma^2 + 4\omega_0^2}t/2}) = 2\cos(\frac{\sqrt{-\gamma^2 + 4\omega_0^2}}{2}t)$$

So we have

$$(C_1 + C_2)2e^{\gamma t/2}\cos(\frac{\sqrt{-\gamma^2 + 4\omega_0^2}}{2}t) = x$$

But we have  $(C_1 + C_2)2$  is an arbitrary constant so lets have  $A = (C_1 + C_2)2$ . Lets define  $\alpha = \gamma/2$  and  $\omega = \frac{\sqrt{-\gamma^2 + 4\omega_0^2}}{2} = \sqrt{\omega_0^2 - \gamma^2/4}$ . Thus

$$x = (C_1 + C_2)2e^{\gamma t/2}\cos(\frac{\sqrt{-\gamma^2 + 4\omega_0^2}}{2}t) = Ae^{\alpha t}\cos(\omega t)$$

is a solution where  $\alpha$  and  $\omega$  has been defined as above.

## 3-18

A) The increase in potential energy can be seen as moving water up a distance of y. Since the top plane is tilted but flat, we have  $y = \frac{y_0}{L/2}x = 2y_0x/L$ . Now consider a slice of water in the pond of that height that was also moved up

that height. It would have a potential energy increase of (dm)gy. But we have dm = p(dv) = pby(dx). So

$$(dm)gy = pby^2g(dx) = \frac{4pgby_0^2}{L^2}x^2(dx)$$

Integrating this from 0 to L/2 results in

$$U = \int_{0}^{L/2} \frac{4pgby_0^2}{L^2} x^2(dx) = \frac{4pgby_0^2}{L^2} \int_{0}^{L/2} x^2(dx) = \frac{4pgby_0^2 L^3}{24L^2} = \frac{pgby_0^2 L}{6}$$

as required.

B) Since vhb - (v + dv)hb = b(dx)(dy/dt), we have -h(dv) = (dx)(dy/dt). But from earlier, we have  $y = 2y_0x/L$  So  $dy = 2dy_0x/L$ . Thus  $(dv) = -(dx)2dy_0(dx)/((dt)hL)$ . Then integrating both sides we obtain

$$v(x) = \int \frac{-2dy_0x}{((dt)hL)} dx = \frac{-x^2}{hL} \frac{dy_0}{dt} + C$$

Since v(L/2) = 0, we have

$$0 = \frac{-L^2}{4hL}\frac{dy_0}{dt} + C$$

So

$$\frac{L}{4h}\frac{dy_0}{dt} = C$$

Thus we have

$$v(x) = \frac{L}{4h} \frac{dy_0}{dt} - \frac{x^2}{hL} \frac{dy_0}{dt}$$

But

$$v(0) = \frac{L}{4h} \frac{dy_0}{dt}$$

Thus

$$v(x) = v(0) - \frac{x^2}{hL} \frac{dy_0}{dt}$$

C)

The kinetic energy of a slice of water lying between x and x + dx is  $mv^2/2$ . But m = p(dv) = pbh(dx) and  $v(x) = v(0) - \frac{x^2}{hL} \frac{dy_0}{dt}$ . So

$$dK = pbh(dx)(v(0) - \frac{x^2}{hL}\frac{dy_0}{dt})^2/2$$

Integrating between  $\pm L/2$  results in

$$\begin{split} \int_{-L/2}^{L/2} pbh(v(0) - \frac{x^2}{hL} \frac{dy_0}{dt})^2 / 2(dx) &= \frac{pbh}{2} \int_{-L/2}^{L/2} (v(0) - \frac{x^2}{hL} \frac{dy_0}{dt})^2 (dx) = \\ \frac{pbh}{2} (v(0)^2 L + \int_{-L/2}^{L/2} (-x^2 2v(0) \frac{1}{hL} \frac{dy_0}{dt} + x^4 (\frac{1}{hL} \frac{dy_0}{dt})^2) dx) &= \\ \frac{pbh}{2} (v(0)^2 L - \frac{2}{3} (\frac{L}{2})^3 2v(0) \frac{1}{hL} \frac{dy_0}{dt} + \frac{2}{5} (\frac{L}{2})^5 (\frac{1}{hL} \frac{dy_0}{dt})^2) &= \\ \frac{pbhL^3}{2} ((\frac{1}{4h} \frac{dy_0}{dt})^2 - \frac{2}{3} (\frac{1}{2})^3 2\frac{1}{4h} \frac{dy_0}{dt} \frac{1}{h} \frac{dy_0}{dt} + \frac{2}{5} (\frac{1}{2})^5 (\frac{1}{h} \frac{dy_0}{dt})^2) &= \\ \frac{pbL^3}{2h} \left(\frac{dy_0}{dt}\right)^2 (\frac{1}{16} - \frac{1}{24} + \frac{1}{80}) &= \frac{pbL^3}{2h} \left(\frac{dy_0}{dt}\right)^2 (\frac{15 - 10 + 3}{240}) &= \\ \frac{pbL^3}{60h} \left(\frac{dy_0}{dt}\right)^2 \end{split}$$

as required

D) So  $\frac{pbL^3}{60h} \left(\frac{dy_0}{dt}\right)^2 + \frac{pgby_0^2L}{6} = E$ . Where  $\frac{pbL^3}{60h} \left(\frac{dy_0}{dt}\right)^2 = A$  and  $\frac{pgby_0^2L}{6} = B$ . Since E is a constant, we can conclude  $\omega = \sqrt{B/A}$ . Since  $T = 2\pi/\omega$ , we get

$$T = 2\pi \sqrt{A/B} = 2\pi \sqrt{\frac{pbL^36}{60hpgbL}} = 2\pi \sqrt{\frac{L^2}{10hg}}$$

E

$$T = 2\pi \sqrt{\frac{L^2}{10hg}} = 2\pi * 70*10^3 m \sqrt{\frac{1}{10*150m*9.8m/s}} = 3627.6s = 60.46mins$$

This is about an 18% error in calculations.

#### 4-3

- A) We have  $m\frac{d^2x}{dt^2} + b\frac{dx}{dt} + kx = 0$  and  $\omega = \sqrt{\frac{k}{m} \frac{b^2}{4m^2}} = 17.321 rad/s$ . So the period of oscillation is  $T = 2\pi/\omega = 0.363s$
- B) We have the amplitude of the steady state being

$$A = \frac{F_0/m}{[(\omega_0^2 - \omega^2)^2 + (\gamma\omega)^2]^{1/2}}$$

But

$$\omega_0 = \sqrt{k/m} = 20s^{-1}$$

and

$$\gamma = b/m = 20s^{-1}$$

So

$$A = \frac{2N/0.2kg}{[((20s^{-1})^2 - (30s^{-1})^2)^2 + (20s^{-1}30s^{-1})^2]^{1/2}} = 0.0128m$$

### 4-5

- A) The forces on a pendulum driven as so would depend on the angle between a vertical line going through the pivot of the pendulum and the bob. Specifically we have a force on the bob of  $mg\sin(\theta)=-mg(x-\xi)/l$  for small angles. But  $m(d^2x)/(dt^2)=-mg(x-\xi)/l-b(dx)/(dt)$ . So  $(d^2x)/(dt^2)+\gamma(dx)/(dt)=g(\xi-x)/l$  meaning  $(d^2x)/(dt^2)+\gamma(dx)/(dt)+gx/l=g\xi/l$  as required. The steady state solution of such an equation is of the form  $A\cos(\omega t+\alpha)$ . Where  $A=\xi_0/[(\omega_0^2-\omega^2)^2+(\gamma\omega)^2]^{1/2}$  and  $\tan(\alpha)=\gamma\omega/(\omega_0^2-\omega^2)$
- B) We have  $A = \xi_0/[(\omega_0^2 \omega_0^2)^2 + (\gamma\omega_0)^2]^{1/2} = \xi_0/(\gamma\omega_0)$  at resonance. We know that  $\omega_0^2 = g/l = 9.8s^{-2}$  and we have  $e^{-1} = e^{-\gamma 50T/2}$  So  $2/50T = \gamma$ . So  $Q = \omega_0/\gamma = 50T\omega_0/2$ . But  $T = 2\pi/\omega_0$  so  $Q = 50\pi$ . Thus at resonance we have  $A = \xi_0/(\gamma\omega_0) = \xi_0Q/\omega_0^2 = 1*10^{-3}m*50\pi/9.8 = 0.016m$

### 4-10

A) Since  $W = \int F dx$ , we have dW = F dx. Because P = dW/dt we then have P = F dx/dt, but we have F = bv and dx/dt = v. Thus  $P = bv * v = bv^2$  as required,

- B) Since  $x = A\cos(\omega t \delta)$ , we have  $v = -A\omega\sin(\omega t \delta)$ . So  $P = bA^2\omega^2\sin^2(\omega t \delta)$ . The average of  $\sin^2(\omega t \delta) = 1/2$  and thus we have  $\overline{P} = bA^2\omega^2/2$  as required.
- C) We have

$$A = \frac{F_0}{k} \frac{\omega_0/\omega}{\left[\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}\right]^{1/2}}$$

Substituting it in results in

$$\overline{P} = b \frac{F_0^2}{2k^2} \frac{\omega_0^2}{\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}} =$$

$$\frac{F_0^2}{2kQ} \frac{\omega_0}{\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}}$$

as required.

## 4-13

- A)  $\omega_0$  is when P is at max power so we have  $\omega_0 = 40s^{-1}$ . Since  $Q = \omega_0/\gamma$  and we have  $\gamma = 41 39 = 2s^{-1}$ , Q = 20
- B) Since  $E = E_0 e^{-\gamma t}$ , we set  $E = E_0/e^5$ . Then  $E_0/e^5 = E_0 e^{-\gamma t}$  and thus  $e^{-5} = e^{-\gamma t}$ . So  $-5 = -\gamma t$  and we must therefore have  $t = 5/\gamma = 2.5s$ . Since a period is  $2\pi/\omega_0 = 0.157s$ , we have 2.5s/0.157s = 15.9 cycles.