

**MAT 341 HW3, Carl Liu**

**1.**

A)

$$\begin{aligned} f(x) &= \int f'(x) dx = \int \sum_{n=1}^{\infty} \frac{n}{3^n} \sin(nx) dx = \sum_{n=1}^{\infty} \int \frac{n}{3^n} \sin(nx) dx = \\ &= -\sum_{n=1}^{\infty} \frac{1}{3^n} \cos(nx) + C \end{aligned}$$

So we have

$$2 = f(0) = -\sum_{n=1}^{\infty} \frac{1}{3^n} \cos(n \cdot 0) + C = -\sum_{n=1}^{\infty} \frac{1}{3^n} + C$$

But we have by geometric series

$$\sum_{n=1}^{\infty} \frac{1}{3^n} = \sum_{n=0}^{\infty} \frac{1}{3^n} - 1 = \frac{1}{1 - 3^{-1}} - 1 = \frac{3}{2} - 1 = \frac{1}{2}$$

Thus  $2 = -(1/2) + C$ . Resulting in  $C = 5/2$  and

$$f(x) = -\sum_{n=1}^{\infty} \frac{1}{3^n} \cos(nx) + 5/2$$

B) Using Theorem 7 of section 1-5, we have continuous derivatives  $f^{(k)}$  for all  $k$  when

$$\sum_{n=1}^{\infty} |n^k a_n| + |n^k b_n|$$

converges for all  $k$ . Since  $b_n = 0$ , we only need convergences of

$$\sum_{n=1}^{\infty} \left| \frac{n^k}{3^n} \right|$$

for all  $k$ . So let  $k \geq 1$ . Using the ratio test we have

$$\frac{x_{n+1}}{x_n} = \frac{(n+1)^k / 3^{n+1}}{n^k / 3^n} = \frac{(n+1)^k}{3n^k} = \left( \frac{n+1}{3n} \right)^k$$

But for  $n \geq 1$ , we have  $2n \geq n+1$ . So  $2/3 \geq (n+1)/3n$  for all  $n \geq 1$ . Meaning  $(2/3)^k \geq ((n+1)/3n)^k$  for all  $n \geq 1$ . Since  $2/3 < 1$ , we have  $(2/3)^k < 1$ . So by the squeeze test  $1 > \lim_{n \rightarrow \infty} 2/3^k = 2/3^k \geq \lim_{n \rightarrow \infty} ((n+1)/3n)^k \geq \lim_{n \rightarrow \infty} 0 = 0$ . Thus

$$1 > \lim_{n \rightarrow \infty} \frac{(n+1)^k}{3n^k} = \lim_{n \rightarrow \infty} \frac{x_{n+1}}{x_n}$$

So by the ratio test we have convergence of the series

$$\sum_{n=1}^{\infty} \left| \frac{n^k}{3^n} \right|$$

and since  $k \geq 1$  was arbitrary, this is true for all  $k \geq 1$  and thus we can conclude that  $f$  has derivative  $f^{(k)}$  for all  $k \geq 1$  as required.

**2.**

A) We have

$$\frac{\partial u}{\partial t} = -\lambda^2 k e^{-\lambda^2 k t} \cos(\lambda x)$$

and

$$\frac{\partial u}{\partial x} = -\lambda e^{-\lambda^2 k t} \sin(\lambda x) \quad \frac{\partial^2 u}{\partial x^2} = -\lambda^2 e^{-\lambda^2 k t} \cos(\lambda x)$$

So

$$\frac{\partial^2 u}{\partial x^2} = -\lambda^2 e^{-\lambda^2 k t} \cos(\lambda x) = \frac{-\lambda^2 k e^{-\lambda^2 k t} \cos(\lambda x)}{k} = \frac{1}{k} \frac{\partial u}{\partial t}$$

as required.

B) We have

$$\frac{\partial u}{\partial t} = -\lambda^2 k e^{-\lambda^2 k t} \sin(\lambda x) \quad \frac{\partial^2 u}{\partial x^2} = -\lambda^2 e^{-\lambda^2 k t} \sin(\lambda x)$$

So

$$\frac{\partial^2 u}{\partial x^2} = -\lambda^2 e^{-\lambda^2 k t} \sin(\lambda x) = \frac{-\lambda^2 k e^{-\lambda^2 k t} \sin(\lambda x)}{k} = \frac{1}{k} \frac{\partial u}{\partial t}$$

as required.

**3.**

A) We have at the steady state

$$\frac{d^2v}{dx^2} = 0 \quad v(a) = T_0 \quad \frac{dv}{dx}(0) = 0$$

Integrating we then get

$$\frac{dv}{dx}(x) = A \quad v(x) = Ax + B$$

So

$$\frac{dv}{dx}(0) = 0 = A$$

making  $v(x) = B$ . Since  $v(a) = T_0$ , we then have  $T_0 = B$  making  $v(x) = T_0$  the steady state solution

B) We have at the steady state

$$\frac{d^2v}{dx^2}(x) = 0 \quad v(0) = T_0 + \frac{dv}{dx}(0) \quad \frac{dv}{dx}(a) = 0$$

Integrating results us in

$$\frac{dv}{dx}(x) = A \quad v(x) = Ax + B$$

Since  $v'(a) = 0$ , we must have  $A = 0$  meaning  $v'(x) = 0$ . Thus  $v'(0) = 0$  and  $v(x) = B$ . Since  $v(0) = T_0 + v'(0) = T_0 + 0 = T_0$ , we have  $B = T_0$  and so the steady state solution is  $v(x) = T_0$

C) We have at the steady state

$$\frac{d^2v}{dx^2}(x) = 0 \quad v(0) = T_0 + \frac{dv}{dx}(0) \quad v(a) = T_1 - \frac{dv}{dx}(a)$$

Integrating results us in

$$\frac{dv}{dx}(x) = A \quad v(x) = Ax + B$$

We then have  $v'(0) = A = v'(a)$ . Thus  $v(0) = T_0 + A = A * 0 + B$ . Meaning  $T_0 + A = B$ . Then  $v(a) = T_1 - A = Aa + B$ . So  $T_1 - A = Aa + T_0 + A$ . Then

$T_1 = A(2+a) + T_0$ . Thus  $A = (T_1 - T_0)/(2+a)$  and  $B = T_0 + (T_1 - T_0)/(2+a)$ . So

$$v(x) = \frac{T_1 - T_0}{2+a}x + T_0 + \frac{T_1 - T_0}{2+a} = \frac{T_1 - T_0}{2+a}(x+1) + T_0$$

4.

The amount of heat going into the rod over a period of time is  $q(0, t)A\Delta t = -C(u(0, t + \Delta t) - u(0, t))$  where  $q$  is the heat flux. This is because the water and rod has the same temperature at 0 meaning that the change in temperature of the rod is equivalent to the change in temperature of the water. But the water holds  $C$  heat per unit of temperature and this must go into the rod when the temperature decreases. So a decreasing temperature means an increasing flux. Then by dividing through with  $\Delta t$  we obtain

$$Aq(0, t) = \frac{-C(u(0, t + \Delta t) - u(0, t))}{\Delta t}$$

Taking the limit as  $\Delta t \rightarrow 0$ , we obtain

$$Aq(0, t) = -C \frac{\partial u}{\partial t}(0, t)$$

But by Fourier's law, we have  $q(0, t) = -\kappa \frac{\partial u}{\partial x}$  and so we can conclude

$$\kappa \frac{\partial u}{\partial x} = C \frac{\partial u}{\partial t}(0, t)$$

There is no negative because the flux into the rod increases at  $x = 0$  when the temperature decreases at that point, the reason was mentioned above. This gets carried throughout the derivation resulting in the above equation.