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PHY 981

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Exercise 5

a) S.p. states ϕ_λ ($\lambda=1, \dots, A$), new basis ψ_a ($a=1, \dots, A$)

$$\psi_a = \sum_\lambda C_{a\lambda} \phi_\lambda \quad \text{when } C_{a\lambda} \text{ are elements of a unitary matrix}$$

New basis is orthogonal

$$\langle \psi_b | \psi_a \rangle = \sum_\lambda \sum_{\lambda'} C_{b\lambda'}^* C_{a\lambda} \underbrace{\langle \phi_{\lambda'} | \phi_\lambda \rangle}_{\delta_{\lambda'\lambda}}$$

$$= \sum_\lambda C_{b\lambda}^* C_{a\lambda} = \delta_{ab}$$

using property of unitary matrices

$$b) |\Psi\rangle^{As} = \frac{1}{\sqrt{A!}} \sum_P (-1)^P P \prod_{i=1}^A \psi_i$$

$$= \frac{1}{\sqrt{A!}} \sum_P (-1)^P P \prod_{i=1}^A \left(\sum_\lambda C_{i\lambda} \phi_\lambda \right)$$

$$= \frac{1}{\sqrt{A!}} \sum_{P_i} (-1)^{P_i} P_i \prod_{i=1}^A \sum_\lambda (-1)^{P_\lambda} P_\lambda \prod_{\lambda=1}^A C_{i\lambda} \phi_\lambda$$

since more than one p.t. cannot occupy a single state

e.g. for $A=2$:

$$\psi_a^{As} = \frac{1}{\sqrt{2!}} (|a\rangle|b\rangle - |b\rangle|a\rangle) = \frac{1}{\sqrt{2}} \left((c_1^a|1\rangle + c_2^a|2\rangle)(c_1^b|1\rangle + c_2^b|2\rangle) - (c_1^b|1\rangle + c_2^b|2\rangle)(c_1^a|1\rangle + c_2^a|2\rangle) \right)$$

$$= \left(\cancel{c_1^a c_1^b |1\rangle|1\rangle} + c_1^a c_2^b |1\rangle|2\rangle + c_2^a c_1^b |2\rangle|1\rangle + \cancel{c_2^a c_2^b |2\rangle|2\rangle} - \cancel{c_1^b c_1^a |1\rangle|1\rangle} - c_1^b c_2^a |1\rangle|2\rangle - c_2^b c_1^a |2\rangle|1\rangle - \cancel{c_2^b c_2^a |2\rangle|2\rangle} \right)$$

$$(c_1^a c_2^b - c_1^b c_2^a) |12\rangle + (c_2^a c_1^b - c_2^b c_1^a) |21\rangle$$

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$$\underbrace{(c_1^a c_2^b - c_1^b c_2^a)}_{\det C} |12\rangle - \underbrace{(c_2^b c_1^a - c_2^a c_1^b)}_{\det C} |21\rangle$$

$$\det(C) \equiv \begin{vmatrix} c_1^a & c_2^a \\ c_1^b & c_2^b \end{vmatrix} = c_1^a c_2^b - c_1^b c_2^a$$

$$\text{So, } \psi_a^{As} = \det C \phi_a^{ns}$$

c) Since C is unitary, $CC^\dagger = C^\dagger C = \mathbb{1}$

$$|\det(C)|^2 = \det(C) \det(C^\dagger) = \det(CC^\dagger) = \det(\mathbb{1}) = 1$$

$$\therefore |\det(C)| = 1 \quad \text{and so} \quad \det(C) = e^{i\theta}$$

Exercise 6

6-1

A-particle Slater determinants $|\phi_0\rangle$ $|\phi_i^a\rangle$ $|\phi_{ij}^{ab}\rangle$

One body operator $\hat{F} = \sum_{i=1}^A \hat{f}(x_i)$

Two body operator $\hat{G} = \sum_{i,j}^A \hat{g}(x_i, x_j)$, $\hat{g}(x_i, x_j) = \hat{g}(x_j, x_i)$

$$a) \langle \phi_0 | \hat{F} | \phi_0 \rangle = \frac{1}{A!} \sum_{i=1}^A \langle i | \hat{f} | i \rangle \cdot A!$$

$$\langle \phi_0 | \hat{G} | \phi_0 \rangle = \sum_{i,j}^A \langle i j | g | i j \rangle - \langle i j | g | j i \rangle$$

since $\hat{g}(x_i, x_j) = \hat{g}(x_j, x_i)$

$$b) \langle \phi_0 | \hat{F} | \phi_i^a \rangle = \sum_j^A \langle i | f_j | a \rangle \left(\frac{1}{A!} \right)$$

since when $\langle i |$ is not "linked" to $|a\rangle$ (in other words) when i and a are not in the same slot, it is guaranteed that one of the other slots will introduce an inner product between orthogonal s.p. states

e.g. for $A=2$

$$\begin{aligned} \langle \phi_0^{A=2} | \hat{F} | \phi_1^{A=2, a} \rangle &= \langle 1 \cancel{2} | f_n | 1 a \rangle - \langle 1 \cancel{2} | f_n | a 1 \rangle - \langle \cancel{2} 1 | f_n | 1 a \rangle + \langle \cancel{2} 1 | f_n | a 1 \rangle \\ &+ \langle 1 \cancel{2} | f_n | 1 a \rangle - \langle 1 \cancel{2} | f_n | a 1 \rangle - \langle \cancel{2} 1 | f_n | 1 a \rangle + \langle \cancel{2} 1 | f_n | a 1 \rangle \\ &= \frac{1}{2} (\langle 2 | f | a \rangle + \langle 2 | f | a \rangle) \end{aligned}$$

since the ptebs are identical, we set $f_n = f_n = f$

$$\langle \Phi_0 | \hat{G} | \Phi_i^a \rangle$$

(6-2)

- as with the one-body case, whenever $\langle i |$ is not involved in a matrix element with $|a\rangle$, the term will be 0 since there will be an inner product between orthogonal states due to the fact that $|a\rangle$ has replaced $|i\rangle$ only on the right of the operator

$$\hookrightarrow 2\langle i j | g | a m \rangle - 2\langle i j | g | m a \rangle \delta_{jm}$$

- factor of 2 since $g(x_i x_j) = g(x_j x_i)$
- for the same reason as point "1", we must have $j = m$

$$2\langle i j | g | a j \rangle - 2\langle i j | g | j a \rangle$$

- Sum over states i, j

$$\langle \Phi_0 | \hat{G} | \Phi_i^a \rangle = \frac{2}{A!} \sum_{i > j} \langle i j | g | a j \rangle - \langle i j | g | j a \rangle$$

- c) $\langle \Phi_0 | \hat{F} | \Phi_{ij}^{ab} \rangle = 0$ since f_0 connects only $\langle i |$ and $|a\rangle$ via $\langle i | f_0 | a \rangle$ or $\langle j |$ and $|b\rangle$ via $\langle j | f_0 | b \rangle$ the remaining "mis match" between the left and right will make each term 0

$$\langle \Phi_0 | \hat{G} | \Phi_{ij}^{ab} \rangle = 2\langle i j | g | a b \rangle - 2\langle i j | g | b a \rangle$$

the only surviving term is when the replacement terms $|a b\rangle$ meet the terms they replace $\langle i j |$

$$\langle \Phi_0 | \hat{G} | \Phi_{ijk}^{abc} \rangle = 0$$

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since \hat{G} can connect, at most, 2 s.p. states there will be an $|a\rangle$ $|b\rangle$ or $|c\rangle$ matched with a bra from the original set of s.p. states in each term causing each term to be 0.

d) For a two body Hamiltonian, the expectation value

$$\langle \Psi_i | \hat{H} | \Psi_i \rangle \text{ where } \Psi_i = \sum_{\lambda=0}^{\infty} C_{i\lambda} \Phi_{\lambda}^{AS}$$

will only involve terms where Φ_{λ} on the right differs from $\Phi_{\lambda'}$ on the left by no more than 2 s.p. states