#### Electrostatics Basics

Line Charge:	$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\vec{r'})}{  r_{ m sep}  ^2} \hat{r}_{ m sep} dt'$
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Surface Charge: 
$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(\vec{r'})}{||\vec{r_{\rm sep}}||^2} \hat{r_{\rm sep}} \, da'$$
Volume Charge: 
$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(r')}{||\vec{r_{\rm sep}}||^2} \hat{r_{\rm sep}} \, d\tau'$$

Volume Charge: 
$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{||\vec{r}_{\rm sep}||^2} \hat{r}_{\rm sep} d\vec{r}$$

Potential Difference: 
$$V(\vec{b}) - V(\vec{a}) = -\int_{\vec{a}}^{\vec{b}} \vec{E} \cdot d\vec{l}$$

Potential of a Volume 
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r'})}{||\vec{r_{\rm sp}}||} d\tau' t$$
 Charge:

Potential of a Collection of Point Charges: 
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \sum_{i=0}^{n} \frac{q_i}{||\vec{r}_{\text{sep}}||}$$

Work to Move a Charge: 
$$W = \int_{2}^{\mathbf{b}} \vec{F} \cdot d\vec{l} = -Q \int_{2}^{\mathbf{b}} \vec{E} \cdot d\vec{l}$$

Work to Move a Charge (Potential): 
$$W = Q[V(\vec{b}) - Q(\vec{a})]$$

Energy of a Collection of Point Charges: 
$$W = \frac{1}{2} \sum_{i=1}^{n} q_i V(\vec{r_i})$$

Energy of a Continuous Charge Distribution: 
$$W = \frac{\epsilon_0}{2} \int\limits_{\rm univ} E^2 d\tau$$

Field at a Charged 
$$\vec{E}_{surf} = \frac{1}{2} \left( \vec{E}_{above} + \vec{E}_{below} \right)$$

Parallel Plate Voltage: 
$$V = \frac{Q}{A\epsilon_0} d$$

Parallel Plate Voltage: 
$$V = \frac{Q}{A\epsilon_0} d$$
 Capacitance: 
$$C \equiv \frac{Q}{V}$$
 Energy of a Capacitor: 
$$W = \int_0^Q \left(\frac{q}{C}\right) dq = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C V^2$$

# Electrostatic Multipoles

Generalized Multipole 
$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int_V ||\vec{r'}||^n P_n(\cos\alpha) \rho(\vec{r'}) d\tau'$$
 Expansion:

$$V(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{\frac{r}{r}}$$
 Dipole Voltage: 
$$V(r) = \frac{1}{4\pi\epsilon_0} \frac{p}{\frac{r}{r^2}} \cdot \hat{r}$$
 Dipole Moment (Continuous): 
$$\vec{p} \equiv \int_{V} \vec{r'} \rho(\vec{r'}) d\tau'$$

Dipole Moment (Continuous): 
$$\vec{p} \equiv \int_{V}^{T} \vec{r'} \rho(\vec{r'}) d\tau$$
Dipole Moment (Discrete):  $\vec{r} = \sum_{n=0}^{\infty} \sigma_{i} \vec{r'}$ 

Dipole Moment (Discrete): 
$$\vec{p} = \sum_{i=1}^{n} q_i \vec{r'}_i$$
 Dipole Moment (Change of 
$$\vec{r'} = \vec{n} - O\vec{q}$$

Dipole Moment (Change of Origin): 
$$\vec{p'} = \vec{p} - Q\vec{a}$$

Torsus of a Dipole  $\vec{N} = \vec{r} \times \vec{E}$ 

Torque of a Dipole: 
$$\vec{N} = \vec{p} \times \vec{E}$$
  
Energy of a Dipole:  $U = -\vec{p} \cdot \vec{E}$ 

Field of a Dipole: 
$$V = \vec{p} \times \vec{E}$$
  
Energy of a Dipole:  $U = -\vec{p} \cdot \vec{E}$   
Force on a Dipole:  $\vec{F} = (\vec{p} \cdot \nabla) \vec{E}$   
Field of a Dipole:  $\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{p}{r^2} \left(2\cos\theta \hat{r} + \sin\theta \hat{\theta}\right)$ 

# Electrostatics in Matter

Electric Displacement: 
$$\vec{D} \equiv \epsilon_0 \vec{E} + \vec{P}$$

Bound Surface Charge: 
$$\sigma_{\text{bound}} = \vec{P} \cdot \hat{n}$$

Bound Volume Charge: 
$$\rho_{\text{bound}} \equiv -\nabla \cdot \vec{P}$$

Linear Dielectrics – 
$$\vec{P} = \epsilon_0 \chi_e \vec{E}$$

Polarization: 
$$P = \epsilon_0 \chi_e E$$

Linear Dielectrics – 
$$\vec{D} = \epsilon \vec{E}$$

Electric Permittivity: 
$$\epsilon = \epsilon_0 (1 + \chi_e) = \epsilon_0 \epsilon_r$$

Gauss's Law for 
$$\vec{D}$$
 (Derivative):  $\nabla \cdot \vec{D} = \rho_f$ 

Gauss's Law for 
$$\vec{D}$$
 
$$\oint \vec{D} \cdot d\vec{A} = Q_{f_{\rm encl}}$$

(Integral): 
$$\mathcal{Y}^{D \cdot uA} = \mathcal{Y}_{f_{en}}$$

Force on a Dielectric: 
$$F = -\nabla U$$

Energy of a Dielectric: 
$$W = \frac{1}{2} \int \vec{D} \cdot \vec{E} d\tau$$

#### Names of Stuff

- $\epsilon_0$ : Permittivity of Free Space
- $\epsilon_r$ : Dielectric Constant or Relative Permittivity
- $\epsilon$ : Permittivity of a Material

### Techniques for Solving Problems

Cylindrical Laplacian Solution:

$$V(s,\phi) = A \ln s + B + \sum_{n=1}^{\infty} \left( A_n s^n + \frac{B_n}{s^n} \right) \left( C_n \cos(\phi n) + D_n \sin(\phi n) \right)$$

Solution for Spherical 
$$f(r,\theta) = \sum_{l=0}^{\infty} \left(A_l r^l + \frac{B_l}{r^{l+1}}\right) P_l(\cos\theta)$$
 Laplacians:

Method of Images General 
$$\sum \frac{q_i}{||\vec{r_{sep,i}}||} = 0$$
 Principle:

$$\begin{array}{ll} \text{Method of Images for Two} \\ \text{Points:} & -\frac{q_1}{q_2} = \frac{||\vec{r_{\text{sep}_1}}||}{||\vec{r_{\text{sep}_2}}||} \end{array}$$

Image Charge Surface 
$$\sigma = -\epsilon_0 \frac{\partial V}{\partial n}$$
 Charge:

Binomial Expansion: 
$$(1+x)^k = \sum_{n=0}^{\infty} \left(\frac{k!}{n!(k-n)!}\right) x^n$$
 Taylor Expansion: 
$$f(x) \approx \sum_{n=0}^{\infty} \frac{f^{(n)}(0)}{n!}$$

Taylor Expansion: 
$$f(x) \approx \sum_{n=0}^{\infty} \frac{f^{-(n)}(0)}{n!}$$

Because I can't remember: 
$$f(x) \approx f(0) + f'(0) \cdot x + f''(0) \cdot \frac{x^2}{2} + \cdots$$

Because I can't remember: 
$$f(x) \approx f(0) + f'(0) \cdot x + f'(0) \cdot \frac{1}{2} + \cdots$$
  
 $P_0(\cos \theta) = 1$ 

Legendre Polynomials: 
$$P_1(\cos \theta) = \cos \theta$$

$$P_2(\cos \theta) = \frac{3}{2}\cos^2 \theta - \frac{1}{2}$$

$$P_3(\cos \theta) = \frac{5}{2}\cos^3 \theta - \frac{3}{2}\cos \theta$$

#### Magnetostatics

Lorentz Force Law: 
$$\vec{F}_{\text{mag}} = Q(\vec{v} \times \vec{B})$$

Surface Current Density: 
$$\vec{K} \equiv \frac{d\vec{l}}{dl_{\perp}}$$

Volume Current Density: 
$$J \equiv \frac{dI}{da}$$
$$\vec{I} = a\vec{v}$$

Magnetic Force (General): 
$$\vec{F}_{\text{mag}} = \int I(d\vec{I} \times \vec{B}) = I \int (d\vec{I} \times \vec{B})$$

Magnetic Force (Surface): 
$$\vec{F}_{\text{mag}} = \int (\vec{K} \times \vec{B}) da$$

Magnetic Force (Volume): 
$$\vec{F}_{\text{mag}} = \int (\vec{J} \times \vec{B}) d\tau$$

Biot-Savart Law: 
$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{I} \times \hat{r}_{\text{sep}}}{||r_{\text{sep}}||^2} dt'$$
 
$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} I \int \frac{d\vec{I}' \times \hat{r}_{\text{sep}}}{||r_{\text{sep}}||^2}$$

Ampere's Law: 
$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} \\ \phi \ \vec{B} \cdot d\vec{l} = \mu_0 I_{\rm encl}$$

Magnetic Vector Potential: 
$$\vec{B} = \vec{\nabla} \times \vec{A}$$

Magnetic Vector Potential: 
$$\vec{B} = \vec{\nabla} \times \vec{A}$$
  
 $\vec{\nabla} \cdot \vec{A} = 0$   
 $\nabla^2 \vec{A} = -\mu_0 \vec{J}$   
Vector Potential – Line:  $\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{I}}{||\vec{r_{\rm sep}}||} dl' = \frac{\mu_0 I}{4\pi} \int \frac{1}{||\vec{r_{\rm sep}}||} d\vec{l}'$   
Vector Potential – Surface:  $\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{K}}{||\vec{r_{\rm sep}}||} da'$   
Vector Potential – Volume:  $\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{||\vec{r_{\rm sep}}||} d\tau'$   
Magnetic Dipole Moment:  $\vec{r} = I \int d\vec{q} - I \vec{q}$ 

Vector Potential – Line: 
$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{1}{||r_{\text{sep}}^-||} dl' = \frac{\mu_0 I}{4\pi} \int \frac{1}{||r_{\text{sep}}^-||} dl'$$

Vector Potential – Surface: 
$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{K}{||r_{\text{sep}}||} da'$$

Vector Potential – Volume: 
$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{J(r)}{||r_{\text{sep}}||} d\tau'$$

Magnetic Dipole Moment: 
$$\vec{m} \equiv I \int d\vec{a} = I\vec{a}$$

Magnetic Dipole Potential 
$$\vec{A}_{\rm dip}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \hat{r}}{r^2}$$
 Expansion:

Magnetic Torque: 
$$\vec{N} = \vec{m} \times \vec{B}$$

Force on a Magnetic Dipole: 
$$\vec{F} = \nabla(\vec{m} \cdot \vec{B})$$

Force Between Two Current-Carrying Wires: 
$$\frac{F}{l} = \frac{\mu_0 I_1 I_2}{2\pi r}$$

# Maxwell's Equations

Gauss's Law: 
$$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho$$

(unnamed) 
$$\vec{\nabla} \cdot \vec{B} = 0$$

Faraday's Law: 
$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Ampere's Law with Maxwell's Correction: 
$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \,\, \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

#### Magnetostatics in Matter

magnetization is **parallel** to  $\vec{B}$ Paramagnets:

Diamagnets: magnetization is **opposite** to  $\vec{B}$ 

magnetization holds outside the presence of Ferromagnets:

an external magnetic field.

 $\vec{M} \equiv$  magnetic dipole per unit volume Magnetization:

 $\vec{J}_{\rm bound} = \vec{\nabla} \times \vec{M}$ Bound Volume Current:

Bound Surface Current:  $\vec{K}_{\text{bound}} = \vec{M} \times \hat{n}$ 

 $\vec{J} = \vec{J}_b + \vec{J}_f$ Free Current:

 $\vec{H} \equiv \frac{1}{\mu_0} \vec{B} - \vec{M}$ Auxillary Field:

Ampere's Law for  $\vec{\nabla} \times \vec{H} = \vec{J}_f$ Auxillary Fields:  $\oint \vec{H} \cdot d\vec{I} = I_{f_{\text{encl}}}$ 

Linear Magnetics:

Permeability:  $\mu \equiv \mu_0 (1 + \chi_m)$ 

Linear Magnetics with E  $\vec{B} = \mu \vec{H}$ . and H Fields:

### Electrodynamics

V = IROhm's Law:

Power in a Circuit:

$$\begin{split} P &= VI = I^2 R = \frac{V^2}{R} \\ \mathcal{E} &= \frac{F_{\text{mag, tot}}}{Q} = \int (\vec{v} \times \vec{B}) d\vec{l} \end{split}$$
Electromotive Force:

 $\mathcal{E} = \oint \vec{f} \cdot d\vec{I}$   $\mathcal{E} = -\frac{d\Phi}{dt}$   $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ 

Faraday's Law:

Inductance:  $\Phi = LI$ 

 $\mathcal{E} = -L\frac{dI}{dt}$ 

Work of a Magnetic Field:

 $W = \frac{1}{2\mu_0} \int_{\text{all space}} B^2 d\tau$ 

Displacement Current:  $\vec{J}_d \equiv \epsilon_0 \frac{\partial \vec{E}}{\partial t}$ 

### Maxwell's Equations in Matter

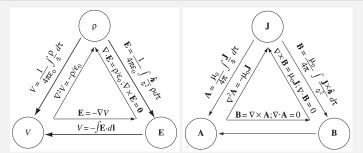
 $\vec{\nabla} \cdot \vec{D} = \rho_f$ Gauss's Law:

 $\vec{\nabla} \cdot \vec{B} = 0$ (unnamed)

 $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ Faraday's Law:

 $\vec{\nabla} \times \vec{H} = \vec{J}_f + \frac{\partial \vec{D}}{\partial t}$ Ampere's Law with Maxwell's Correction:

# Triangles



# Takeaways from Practice

- Gauss's law inside a dielectric always includes both  $\rho_b$  and  $\sigma_b$ .
- Use the definition for displacement where ever possible:  $\vec{D} \equiv$  $\epsilon_0 \vec{E} + \vec{P}$ .
- The potential is constant below an image plane and the field is

### Boundary Conditions in Electrostatics

$$\vec{D}_{\text{above}}^{\perp} - \vec{D}_{\text{below}}^{\perp} = \sigma_f$$

$$\vec{D}_{\mathrm{above}}^{\parallel} - \vec{D}_{\mathrm{below}}^{\parallel} = \vec{P}_{\mathrm{above}}^{\parallel} - \vec{P}_{\mathrm{below}}^{\parallel}$$

$$\vec{E}_{\text{above}}^{\perp} - \vec{E}_{\text{below}}^{\perp} = \sigma_{\text{tot}}/\epsilon_0$$

$$\vec{E}_{\text{above}}^{\parallel} - \vec{E}_{\text{below}}^{\parallel} = \vec{0}$$

$$V_{\rm above} = V_{\rm below}$$

$$\epsilon_{\text{above}} \vec{E}_{\text{above}}^{\perp} - \epsilon_{\text{below}} \vec{E}_{\text{below}}^{\perp} = \sigma_f$$

$$\begin{split} & \epsilon_{\text{above}} \vec{E}_{\text{above}}^{\perp} - \epsilon_{\text{below}} \vec{E}_{\text{below}}^{\perp} = \sigma_f \\ & \epsilon_{\text{above}} \frac{\partial V_{\text{above}}}{\partial n} - \epsilon_{\text{below}} \frac{\partial V_{\text{below}}}{\partial n} = -\sigma_f \end{split}$$

### Boundary Conditions in Magnetostatics

$$\vec{B}_{\text{above}}^{\parallel} - \vec{B}_{\text{below}}^{\parallel} = \mu_0 \vec{K}$$

$$\vec{B}_{\text{above}} - \vec{B}_{\text{below}} = \mu_0(\vec{K} \times \hat{\mathbf{n}})$$

$$\vec{\Lambda}$$
,  $-\vec{\Lambda}$ ,

$$\begin{split} \vec{A}_{\rm above} &= \vec{A}_{\rm below} \\ \frac{\partial A_{\rm above}}{\partial n} &- \frac{\partial A_{\rm below}}{\partial n} = -\mu_0 \vec{K} \end{split}$$

$$\vec{H}_{\mathrm{above}}^{\perp} - \vec{H}_{\mathrm{below}}^{\perp} = - (\vec{M}_{\mathrm{above}}^{\perp} - \vec{M}_{\mathrm{below}}^{\perp})$$

$$\vec{H}_{\mathrm{above}}^{\parallel} - \vec{H}_{\mathrm{below}}^{\parallel} = \vec{K}_f \times \hat{\mathbf{n}}$$

## More Boundary Conditions

Equations 7.60-63 are the general boundary conditions for electrodynamics. In the case of *linear* media, they can be expressed in terms of E and B alone:

(i) 
$$\epsilon_1 E_1^{\perp} - \epsilon_2 E_2^{\perp} = \sigma_f$$
,

(iii) 
$$\mathbf{E}_1^{\parallel} - \mathbf{E}_2^{\parallel} = \mathbf{0},$$

(ii) 
$$B_1^{\perp} - B_2^{\perp} = 0$$
,

(ii) 
$$B_1^{\perp} - B_2^{\perp} = 0$$
, (iv)  $\frac{1}{\mu_1} \mathbf{B}_1^{\parallel} - \frac{1}{\mu_2} \mathbf{B}_2^{\parallel} = \mathbf{K}_f \times \hat{\mathbf{n}}$ .  $\right\}$  (7.64)

(7.65)

In particular, if there is no free charge or free current at the interface, then

(i) 
$$\epsilon_1 E_1^{\perp} - \epsilon_2 E_2^{\perp} = 0$$
,

(iii) 
$$\mathbf{E}_{1}^{\parallel} - \mathbf{E}_{2}^{\parallel} = \mathbf{0}$$
,

(ii) 
$$B_1^{\perp} - B_2^{\perp} = 0$$

(ii) 
$$B_1^{\perp} - B_2^{\perp} = 0$$
, (iv)  $\frac{1}{\mu_1} \mathbf{B}_1^{\parallel} - \frac{1}{\mu_2} \mathbf{B}_2^{\parallel} = \mathbf{0}$ .

As we shall see in Chapter 9, these equations are the basis for the theory of reflection and refraction.