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Search for tWZ production in the Full Run 2 ATLAS
dataset using events with four leptons

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Abstract

The search for tWZ production using 139 fb^{-1} of pp collision data at a centre-of-mass energy of $\sqrt{s} = 13 \text{ TeV}$, recorded by the ATLAS experiment at CERN, is presented. Events containing exactly four electrons or muons (tetralepton) are selected with additional criteria based on the number of jets, the number of b -tagged jets, and the number of Z boson candidates are used to define signal and control regions. The large $t\bar{t}Z$ and ZZ backgrounds are distinguished from signal by a BDT-based algorithm. Inputs to the BDT-based algorithm include, a kinematic reconstruction algorithm which reconstructs leptonically decaying top quarks and a BDT-based algorithm which aims to classify ℓb systems originating from top quarks. The expected signal strength is extracted via a blinded maximum-likelihood fit to multiple signal and control regions. The measured signal strength is $\mu(tWZ) = 1.91^{+0.95}_{-0.82}$, leading to an expected significance of 1.44σ . An expected upper limit on the signal strength is set and is given by $\mu_{up}^{exp} = 1.61^{+2.35}_{-1.16}$. Furthermore, a combined blinded maximum-likelihood fit was performed across the tetralepton and trilepton (studied in an independent analysis) channels, to further increase the sensitivity of $\sigma(tWZ)$. The measured signal strength is $\mu(tWZ) = 1.80^{+0.70}_{-0.65}$, leading to an expected significance of 1.61σ . An expected upper limit on the signal strength is set and is given by $\mu_{up}^{exp} = 1.43^{+2.04}_{-1.03}$.

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Acknowledgements

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Chapter 1

76

Introduction

77 The production of a single top quark in association with a W^\pm and Z boson (tWZ) at the CERN LHC is sensitive
 78 to both the neutral and charged electroweak couplings of the top quark as the process involves the simultaneous
 79 production of a W boson and a Z boson in association with the top quark. Due to the very large coupling of the top
 80 quark to the Higgs boson, the electroweak couplings of the top quark are a theoretically well-motivated area in which
 81 to search for the first signs of new physics [**TopQuarkCouplings-searchNewPhysics, top-a-tool-for-disc**] that
 82 could offer a resolution to the Hierarchy Problem [**Burdman·2015**, 54]. The recent lack of signs of new physics
 83 from LHC data tells us that new physics is either very heavy, or is very weakly coupled to Standard Model particles,
 84 therefore signs of new physics might only be observed in anomalous rates of well-chosen processes. A prime example
 85 of such a process is tWZ . This has an extremely low production cross section (≈ 160 fb for $\sqrt{s} = 13$ TeV [25]),
 86 meaning that it is an extremely rare process to observe and subsequently, it has never been observed by any particle
 87 physics experiment.

88

89 The latest datasets recorded by the ATLAS experiment at the CERN LHC are sufficiently large to allow a potential
 90 observation of tWZ production. In this analysis, the Full Run 2 dataset recorded by ATLAS is used to search for tWZ production in the tetralepton channel (tWZ with exactly four final state leptons). A kinematic
 91 reconstruction technique is used which aims to discriminate between tWZ and our most prominent background
 92 process, $t\bar{t}Z$. In addition to this, Machine Learning techniques are implemented to further isolate our tWZ signal.
 93 Backgrounds from Standard Model (SM) processes including those in which one or more leptons originate from the
 94 semileptonic decay of a heavy hadron or a photon conversion are estimated by fitting predictions from simulation to
 95 data in dedicated control regions. As this work forms the basis of an official ATLAS analysis, only blinded results
 96 are shown. A maximum likelihood fit is performed over our two tWZ signal regions and three control regions, to
 97 measure the cross section of tWZ in the tetralepton channel. In this thesis only the tetralepton channel is explored,
 98 however, the results of a combined fit that includes an entirely independent analysis of the trilepton channel is
 99 presented with no overlap of events between the regions defined for these channels.

101

Chapter 2

102

Theory

2.1 Standard Model of Particle Physics

104 The SM is a model based on Quantum Field Theory (QFT) which classifies all known elementary particles and
 105 describes their interactions. It is a well-tested model and has shown to be hugely successful in describing experi-
 106 mental data to great precision [**ALTARELLI 1998**, 8]. For example, in the top quark sector, the $t\bar{t}$ cross section
 107 predictions have been confirmed to 3.9% accuracy [**cms-ttbar**, **ATLAS-CONF-2019-041**]. It incorporates three
 108 of the four fundamental forces of nature: the electromagnetic, the weak and the strong forces. In Figure 1, all
 109 known elementary particles described by the SM, are shown.

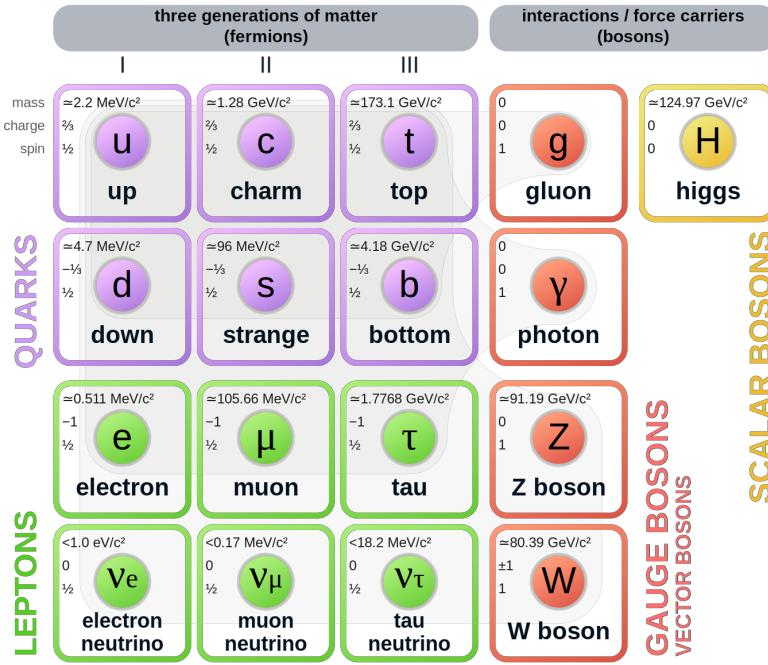


Figure 1: A summary of the elementary particles described by the SM [64] is shown. Fermions are shown on the left, with quarks shown in purple and leptons shown in green. Bosons are shown on the right, with gauge bosons shown in red and the Higgs boson shown in yellow. The mass, electric charge and spin of each particle is shown on the top left of each particle's block.

110 Particles in the SM are uniquely described by their quantum numbers: electric charge and spin. The SM particles
 111 are split into main two classes, based off their spin quantum numbers. Particles which have half-integer spin are
 112 called fermions, and those which have integer spin are called bosons. Fermions are further divided into three genera-
 113 tions, each comprising of two quarks, one charged lepton and one neutrino. In a generation, the more massive quark
 114 has an electric charge of $+2/3$ (up-type) and the less massive quark has an electric charge of $-1/3$ (down-type). All

115 charged leptons have an electric charge of -1 and all neutrinos are electrically neutral. The masses of the particles
 116 in a generation increase with increasing generation number, with generation 1 particles being the least massive and
 117 generation 3 particles being the most massive. Quarks carry electric and colour charge, and can therefore interact
 118 via the electromagnetic, weak and strong forces. Colour charge can take on three values: red, green and blue.
 119 It is important to note that colour charge is completely unrelated to the everyday meaning of colour, and it just
 120 represents the quantum state of the particle. Due to colour confinement [82], quarks cannot be isolated from one
 121 another. They exist in colourless bound states, called hadrons, consisting of two or more quarks. Hadrons consisting
 122 of an even number of quarks are known as mesons and those consisting of an odd number of quarks are known as
 123 baryons. On the other hand, charged leptons (electron (e), muon (μ) and tau (τ)) only carry electric charge and
 124 can therefore interact electromagnetically and weakly, but not through the strong interaction. The electric and
 125 colour neutral fermions, neutrinos, can only interact via the weak force.

126
 127 Particles are able to interact with one-another via the exchange of a gauge boson (boson with spin-1). Photons
 128 are massless, spin-1 gauge bosons which mediate electromagnetic interactions between particles which carry electric
 129 charge, such as quarks and charged leptons (e , μ and τ). The weak interaction is mediated by three massive gauge
 130 bosons, the electrically charged W^+ and W^- bosons and the electrically neutral Z boson. Gluons are massless,
 131 spin-1 gauge bosons which mediate strong interactions between particles which carry colour charge, such as quarks.
 132 Since gluons carry colour charge, they interact with themselves. The massive, spin-0, electrically neutral Higgs
 133 boson mediates the Higgs field which gives mass to the W^\pm and Z bosons via the so-called Brout-Englert-Higgs
 134 mechanism [42, 50, 49]. The Brout-Englert-Higgs mechanism induces spontaneous electroweak symmetry breaking
 135 to provide mass terms for the W^\pm and Z bosons in the electroweak Lagrangian of the SM. All particles described in
 136 the SM have their own antiparticle, with the same mass, but opposite charges. Some particles, such as the photon,
 137 are their own antiparticle.

138
 139 Although the SM has shown to be hugely successful, it is incomplete and fails to describe certain observed phenomena.
 140 The most notable example being the absence of gravity from the SM. The gravitational force is $\approx 10^{29}$ [75]
 141 weaker than the weak force, therefore quantum gravitational effects are expected to only become significant at
 142 energies much larger than that currently accessible by the LHC (known as the Planck scale $\approx 10^9$ GeV) [54]. This
 143 large difference in strength between the weak force and gravity is known as the Hierarchy Problem. Cosmological
 144 observations infer that around 84% of the matter in the universe consists of gravitationally interacting matter known
 145 as dark matter [53]. None of the particles described in the SM are good dark matter candidates, therefore the SM
 146 only accounts for a small fraction of the total matter of the universe. The large discrepancy between the observed
 147 amount of matter and antimatter in the universe, sometimes referred to as the matter-antimatter asymmetry, is not
 148 fully explained by the SM. Neutrinos in the SM are assumed to be massless, however observations of neutrino oscillations
 149 (neutrinos undergoing flavour change as they travel through space) imply that neutrinos do have mass [45].
 150 Beyond the Standard Model (BSM) theories attempt to explain the phenomena which the SM cannot. For example,
 151 a popular extension to the SM, Supersymmetry (SUSY) introduces new particles to the SM which are counterparts
 152 to the existing SM particles with the same quantum numbers, except for their spins [56]. SUSY provides elegant
 153 explanations to many shortcomings of the SM, however none of the supersymmetric particles described by SUSY
 154 have been observed experimentally [30].

155 2.1.1 The Top Quark

156 The top quark is the heaviest particle in the SM, with a mass of 172.76 ± 0.30 GeV [48]. According to the SM, since
 157 the coupling to the Higgs boson is proportional to the the mass of the interacting particle, the top quark is strongly
 158 coupled to the Higgs boson. Physics processes involving top quarks is therefore a theoretically well-motivated area
 159 to search for new physics, since it is the most likely particle to couple to new physics theories at the TeV scale. Its
 160 large mass also makes it highly unstable, with a mean lifetime of $\approx 0.5 \times 10^{-24}$ s [48]. The top quark's lifetime
 161 is shorter than that of the hadronisation process, and it therefore decays before hadronising. The top quark can
 162 therefore be measured indirectly via its decay products. Top quarks almost always decay to a W boson and a
 163 b -quark ($\frac{\Gamma(Wb)}{\Gamma(Wq(q=b,s,d))} = 0.957 \pm 0.034$ [48]). The b -quark is the second heaviest quark in the SM, however its
 164 lifetime is still longer than the hadronisation time scale [48]. In hadron collider experiments, b -quarks travel a short
 165 distance in the detector before hadronising to form jets. In Table 1, the dominant final state branching fractions of
 166 the top quark are shown.

167 Hadronic final states are more than twice as likely than leptonic final states. Final state decays to different lepton
 168 flavours are roughly equally probable.

Decay Mode	Branching Fraction ($\frac{\Gamma_i}{\Gamma}$)
$t \rightarrow Wb \rightarrow e\nu_e b$	(11.10 \pm 0.30)%
$t \rightarrow Wb \rightarrow \mu\nu_\mu b$	(11.40 \pm 0.20)%
$t \rightarrow Wb \rightarrow \tau\nu_\tau b$	(10.70 \pm 0.50)%
$t \rightarrow Wb \rightarrow q\bar{q}b$	(66.50 \pm 1.40)%

Table 1: The dominant final state branching fractions of the top quark [48] are shown.

169

Top quark production can be placed into two main categories: pair production ($t\bar{t}$) and single-top production (t) [73]. In the LHC, top quarks are mainly produced in pairs via strong interactions in gluon-gluon fusion ($gg \rightarrow t\bar{t}$) or quark annihilation ($q\bar{q} \rightarrow t\bar{t}$). Top quark production via gluon-gluon fusion is the dominating process [24]. The production cross section for $t\bar{t}$ (leptonic final state) in pp collisions with $\sqrt{s} = 13$ TeV was measured by ATLAS with a value of $830 \pm 0.4(\text{stat}) \pm 36(\text{syst}) \pm 14(\text{lumi})$ pb [4], with good agreement between measurement and theoretical prediction.

176

Single top production occurs via the weak interaction. The most abundant production mechanisms leading to single top production are the s -, t - and Wt - channels [48]. In the s -channel, an initial quark annihilates with an anti-quark of different flavour, producing a virtual W boson which decays to a top quark and anti-bottom quark. In the t -channel, an initial b quark interacts with a different flavour quark via the exchange of a W boson. This interaction produces a top quark and another quark. In the Wt -channel, an initial gluon interacts with a b quark to produce a top quark and a W boson, either via the absorption of the gluon by the b quark or via the exchange of a top quark. In Table 2, single top production cross sections in pp collisions at $\sqrt{s} = 13$ TeV for the s -, t - and Wt -channels, are shown.

Channel	Process	Total Cross Section [pb]
s	$q\bar{q}' \rightarrow W \rightarrow \bar{b}t$	$10.32^{+0.40}_{-0.36}$
t	$bq' \rightarrow W \rightarrow tq$	$216.99^{+9.04}_{-7.71}$
Wt	$bg \rightarrow b/t \rightarrow Wt$	71.7 ± 3.85

Table 2: Single top production cross sections in pp collisions at $\sqrt{s} = 13$ TeV for the s -, t - and Wt -channels [31] are shown. The prime superscript on q' indicates that the quark has a different flavour to q .

Single top production is suppressed compared to pair produced top production, with $t\bar{t}$ production (leptonic final state) being around three times as likely to occur than single top production across all decay channels.

2.1.1.1 Motivation for the search for tWZ production in the tetralepton channel

The recent lack of signs of new physics from LHC data [74] tells us that new physics is either very heavy, or is very weakly coupled to SM particles. We therefore might only observe signs of new physics in anomalous rates of well-chosen processes. tWZ is a prime example of such a process. It has an extremely low production cross section (0.7 fb for $\sqrt{s} = 13$ TeV [25]), and has subsequently never been observed by any particle physics experiment. Since tWZ involves a charged W boson and neutral Z boson, its cross section is sensitive to the charged and neutral couplings to the top quark. In turn, the top quark is strongly coupled to the Higgs boson, due to its large mass. Due to the top quark's large coupling to the Higgs boson, corrections to the Higgs boson mass diverge in the SM. The top quark's couplings are modified, in order to remove this divergence, in many scenarios of new physics that aim to resolve the Hierarchy Problem. Since the Z boson may be radiated from the initial-state b -quark, the final-state top quark, or the final-state Z boson, the tWZ process embeds the $b - Z$, $t - Z$ and $W - Z$ electroweak couplings which are often modified in BSM physics. Therefore tWZ is an important process in the search for signs of new physics and BSM physics.

200

One such BSM theory which is sensitive to tWZ production [57, 63] is the Standard Model Effective Field Theory (SMEFT) [26]. The SMEFT attempts to describe physics at large energy scales which we have not yet been able to probe experimentally. The SMEFT inherits the same QFT framework as the SM, and adds Lagrangian terms to the SM Lagrangian which describe the interactions of SM particles at higher energy scales. Analogous to the coupling constants found in the SM Lagrangian, which indicate the interaction strengths between different particles, SMEFT contains scalar coefficients which operate in the same way. These scalar coefficients are known as Wilson coefficients. It has been shown that the cross section of tWZ is sensitive to many Wilson coefficients. An experimental constraint on the cross section of tWZ is therefore expected to be impactful on a global fit on all the Wilson coefficients in SMEFT.

Prior to this analysis, only three experimental studies of tWZ in ATLAS have been performed. The first and third studies utilised the trilepton channel to search for tWZ production, whereas the second study utilised both the tri- and tetralepton channels. The first search utilised 36 fb^{-1} of ATLAS data and an upper limit on the cross section of tWZ was set at a value of ≈ 6 times the SM cross section [67]. The second study investigated the feasibility of a cross section measurement of tWZ production with CMS Run 3 data (300 fb^{-1}) [Tschida:2020ftz]. The study showed that it is possible to exclude $\mu(tWZ)$ at the 7σ significance level using 300 fb^{-1} of data. This study needs to be further investigated, since its findings seem improbable given the results obtained in this thesis. The third search utilised 139 fb^{-1} of ATLAS data and an expected upper limit on the cross section of tWZ was set at a value of ≈ 2.6 times the SM cross section [81]. In Section 4.10.2, the latter analysis will be used in combination with this analysis, in order to further increase the sensitivity of the cross section of tWZ .

2.2 tWZ

2.2.1 Tetralepton Channel

In Figure 2, the Leading Order (LO) Feynman diagram for tWZ in the tetralepton channel, is shown.

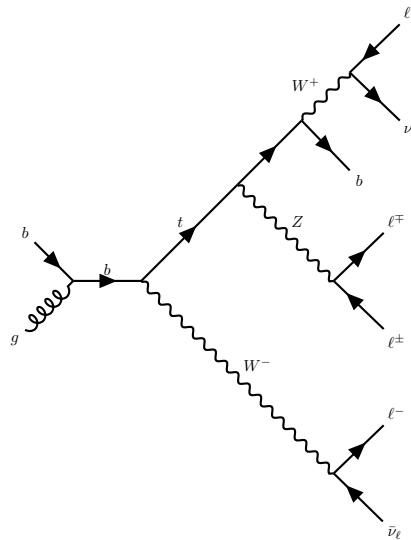


Figure 2: The LO Feynman diagram of tWZ production in the tetralepton channel is shown.

2.2.1.1 Backgrounds

The main backgrounds for tWZ (tetralepton channel) are the production of a two tops, both in the $\ell\nu b^1$ final state channel, together with a Z boson ($t\bar{t}Z$) and diboson production with fully leptonic final states (ZZ). In Figure 2.2.1.1, LO Feynman diagrams for $t\bar{t}Z$ and ZZ in the tetralepton channel, are shown.

¹In this thesis, ℓ refers to an electron or muon, ν refers to a neutrino or anti-neutrino and b refers to a bottom quark or anti-bottom quark

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229

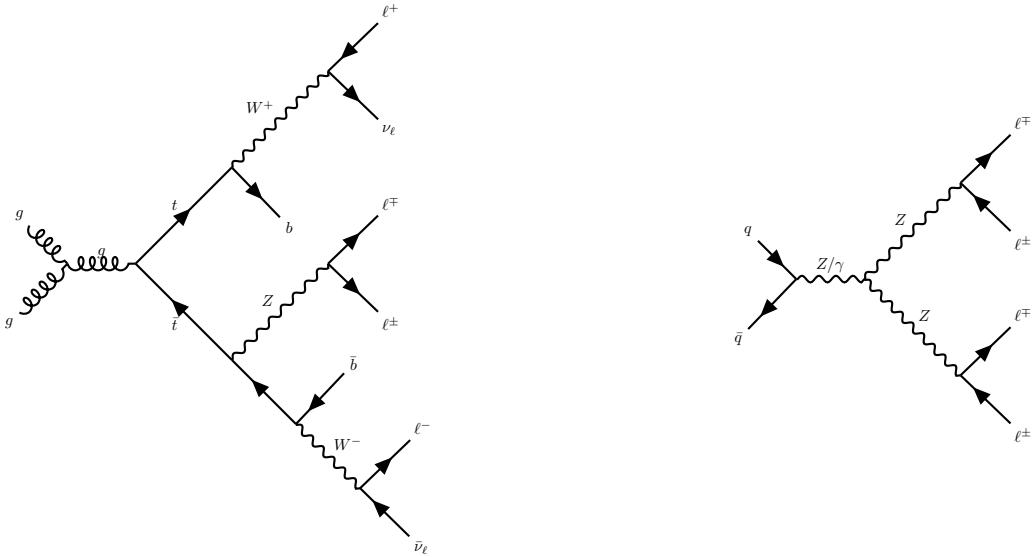


Figure 3: LO Feynman diagrams for $t\bar{t}Z$ (left) and ZZ (right) in the tetralepton channel are shown.

230 The $t\bar{t}Z$ process contains four leptons and two b -quarks in its final state (inclusive $\sigma(t\bar{t}Z) = 0.95 \pm 0.08_{\text{stat}} \pm 0.10_{\text{syst}}$
231 pb at $\sqrt{s} = 13$ TeV [2]) and can easily mimic the tWZ signal process, for instance, by one of its b -jets getting
232 missed during detection. The ZZ process contains four leptons and zero b -quarks in its final state (inclusive
233 $\sigma(ZZ) = 14.6^{+1.9}_{-1.8}(\text{stat})^{+0.5}_{-0.3}(\text{syst}) \pm 0.2(\text{theo}) \pm 0.4(\text{lumi})$ pb at $\sqrt{s} = 13$ TeV [55]). One way in which ZZ can mimic
234 the tWZ signal process is by reconstruction of a non-prompt b -jet.

235 2.2.2 Comparison to Trilepton Channel

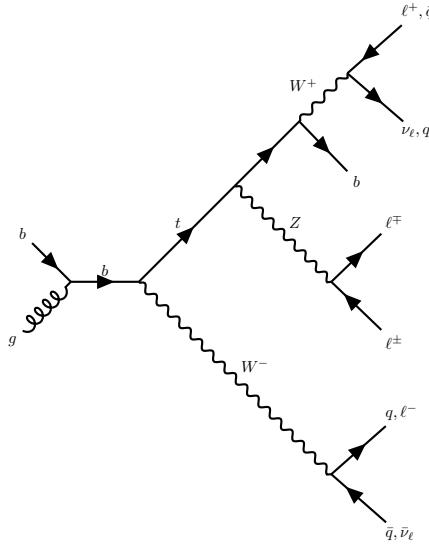


Figure 4: Example Feynman diagram of tWZ production in the tri-lepton channel.

236 The most apparent difference between the tri and tetralepton channels is the number of events present, with the
237 tetralepton channel having far less events in its phase space than that of the tri-lepton channel. The lack of statistics
238 in the tetralepton channel can be attributed to its low production cross section, $\sigma_{(tW^\pm Z).Br(4\ell)}^{\text{NLO}} = 0.7$ fb [25].
239 The tri-lepton channel has a production cross section ($\sigma_{(tW^\pm Z).Br(3\ell)}^{\text{NLO}} = 3.9$ fb [25]) around a factor of 4 larger
240 than that of the tetralepton channel. This difference between the production cross section of the two decay
241 channels can be largely attributed to the difference in branching ratios ($\frac{\Gamma_i}{\Gamma}$) between a hadronically decaying W

²⁴² boson, $\frac{\Gamma_{W \rightarrow had}}{\Gamma_W} = (67.41 \pm 0.27)\%$ [48], present in the tri-lepton channel and a leptonically decaying W boson,
²⁴³ $\frac{\Gamma_{W \rightarrow \ell\nu}}{\Gamma_W} = (10.86 \pm 0.09)\%$ [48], present in the tetralepton channel.

²⁴⁴
²⁴⁵ Despite the tetralepton channel's low statistics, it is not subject to the large WZ background present in the trilepton
²⁴⁶ channel [81]. The tetralepton channel has a substantial amount of ZZ background (not present in the trilepton
²⁴⁷ channel), fortunately this can be easily suppressed due to the full reconstructability of the two leptonically decaying
²⁴⁸ Z -bosons.

249

Chapter 3

The ATLAS Experiment and Detector

250

251 3.1 The ATLAS Experiment

252 The ATLAS (A Toroidal LHC ApparatuS) detector is one of two general purpose detectors at CERN (the European
 253 Organization for Nuclear Research) near Geneva in Switzerland. These detectors collect data from the collisions
 254 provided by the worlds highest energy particle accelerator [27], the Large Hadron Collider (LHC) situated at CERN.

255 In this section, information about the LHC and the ATLAS detector are given. This includes technical aspects
 256 of the ATLAS detector and the processing of data into meaningful physics objects¹ to be used in analyses. The
 257 following chapter consists of information from "The LHC Design Report" [27], "LHC Machine" [44] and "The
 258 ATLAS Experiment at the CERN Large Hadron Collider" [37] unless otherwise stated.

260 3.1.1 Large Hadron Collider (LHC)

261 The LHC is a circular 27 km particle accelerator located in an underground tunnel on the border between France
 262 and Switzerland. The accelerator consists of supercooled, superconducting magnets which accelerate and collide
 263 beams of protons at centre-of-mass energies up to $\sqrt{s} = 13$ TeV at instantaneous luminosities of $\mathcal{L} \sim 10^{34}$ cm $^{-2}$ s $^{-1}$.
 264 In the LHC, pp beams consist of bunches of protons which collide every 25 ns, corresponding to a frequency of 40
 265 MHz. Several accelerator systems are used to accelerate protons and heavy ions to such high energies. Protons are
 266 extracted from a tank of ionised hydrogen gas and are injected into the Linear Accelerator 2 (LINAC), where they
 267 are linearly accelerated to momenta of 50 MeV. The proton bunches are then sequentially accelerated by a chain
 268 of circular accelerators. The chain starts with the Booster which accelerates the protons to momenta of up to 1.4
 269 GeV. The proton bunches are then fed through to the Proton Synchrotron (PS) and the Super Proton Synchrotron
 270 (SPS) which accelerate the protons to momenta of up to 25 GeV and 450 GeV respectively. The protons are
 271 then transferred to two beam pipes of the LHC where they travel in opposite directions. Both proton beams are
 272 accelerated to their final momenta of 6.5 TeV, resulting in a centre-of-mass energy of 13 TeV. These proton beams
 273 then collide at one of the four main interaction points (positions along the beam pipe where collisions occur)
 274 situated along the LHC.

275 The four main experiments located at the interaction points are ATLAS, the Compact Muon Solenoid (CMS),
 276 Large Hadron Collider Beauty (LHCb) Experiment and A Large Ion Collider Experiment (ALICE). The ATLAS
 277 and CMS detectors are general-purpose detectors which investigate a wide range of physics processes. Since both
 278 ATLAS and CMS can measure the same processes, they are able to cross-check and validate measurements taken
 279 by one another. The LHCb detector is specifically designed to study decays of particles containing b -quarks. The
 280 ALICE detector is designed to study the strongly interacting quark-gluon plasma which is formed at extremely
 281 high energy densities. At the interaction points, the two proton beams which consist of protons in closely packed
 282 bunches, travel in opposite directions to one another and collide. Many hard pp collisions (events) can occur per
 283 bunch crossing, however it is the most energetic collision in the bunch crossing that is interesting for discovery
 284 potential. The most energetic collision is therefore chosen to be studied and any additional collisions are aimed
 285 to be rejected. These additional collisions are referred to as *pile-up*. Pileup complicates the reconstruction of the
 286 particles originating from the hard collision of interest.

¹meaningful physical systems which can be reconstructed from detector information (e.g leptons and jets)

288 **3.1.1.1 Luminosity**

289 This section consists of information from "Modern Particle Physics" [75], unless otherwise stated.

290
291 The event production rate at the LHC, $R(t)$, for a certain process of interest is given by,

$$R(t) = \mathcal{L}(t)\sigma \quad (3.1)$$

292 where $\mathcal{L}(t)$ is the instantaneous luminosity and σ and is the cross section of the process of interest. The instantaneous
293 luminosity, $\mathcal{L}(t)$, is independent on the process of interest, and depends on various collider and beam parameters.
294 $\mathcal{L}(t)$ can be written in terms of these parameters as,

$$\mathcal{L}(t) = f \frac{N n_1 n_2}{4\pi \sigma_x \sigma_y} \quad (3.2)$$

295 where f is the beam revolution frequency, N is the number of proton bunches colliding per second, n_1 and n_2 are the
296 number of protons in the colliding bunches, σ_x and σ_y are the beam spread in the x and y directions respectively.
297 The total integrated luminosity, L , across some time interval, is given by,

$$L = \int \mathcal{L} dt. \quad (3.3)$$

298 The units of L are inverse area, and are given by fb^{-1} at the LHC and the ATLAS detector. In Figure 5, the
299 total integrated luminosity delivered to ATLAS, recorded by ATLAS, and certified to be good enough for physics
300 analyses (the data passes certain quality control criteria) for $\sqrt{s} = 13$ TeV pp collisions at the LHC is shown [77].

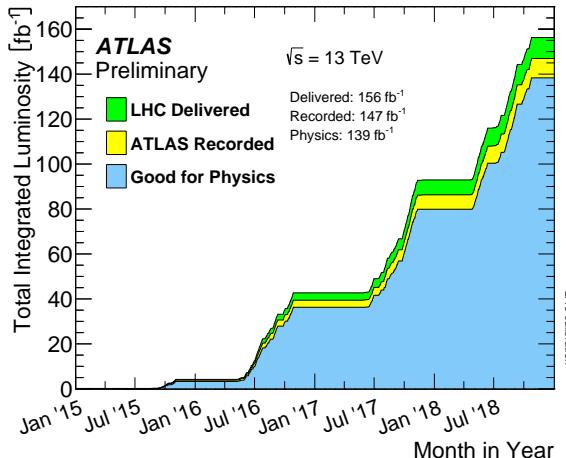


Figure 5: The total integrated luminosity delivered to ATLAS, recorded by ATLAS, and certified to be good enough for physics analyses (the data passes certain quality control criteria) for $\sqrt{s} = 13$ TeV pp collisions at the LHC is shown [77]. The total integrated luminosity delivered by the LHC, recorded by ATLAS and certified to be good quality data are shown by the green, yellow and blue histograms respectively. The month and year of data taking is shown on the x-axis and the total integrated luminosity (in fb^{-1}) is shown on the y-axis.

301 A total integrated luminosity of 139 fb^{-1} of data certified as good for physics was recorded by ATLAS between
302 2015 and 2018. This data taking period is referred to as Run 2, since it proceeds the Run 1 data taking period
303 (2011 and 2012) and the Long Shutdown 1 LHC upgrade period (2013 and 2014). In this analysis, we use the Full
304 Run 2 dataset.

305 **3.2 The ATLAS Detector**

306 In Figure 6, the schematic of the ATLAS detector, is shown.

307 The detector is cylindrically shaped which covers close to 4π in solid angle. It has a length of 44 m, a diameter
308 of 25 m and a mass of 7000 tons. The ATLAS detector consists of four main sub-detectors arranged in concentric
309 cylindrical layers around the beam pipe. These include the inner detector, the electromagnetic calorimeter, the

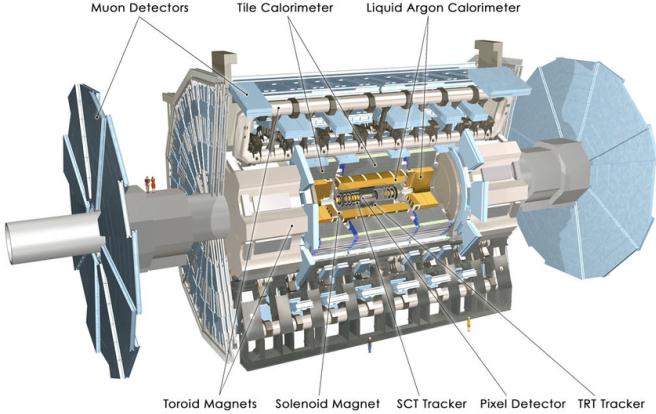


Figure 6: Schematic of the ATLAS detector [37]

hadronic calorimeters and the muon spectrometer. The sub-detectors record the momenta, energies and trajectories of different particles produced in the collider, allowing for the reconstruction and identification of these particles to be used in physics analyses.

3.2.1 Coordinate System and Kinematics

The ATLAS detector adopts a right-handed coordinate system. The origin is at the nominal interaction point with the z -axis defined to be counter-clockwise along the beam line. The $x - y$ plane (or transverse plane) is perpendicular to the beam line, with the x -axis pointing towards the centre of the LHC ring and the y -axis pointing upwards towards the Earth's surface. The azimuthal angle, $\phi \in [-\pi, \pi]$, is measured in the transverse plane with respect to the positive x -axis. The polar angle, $\theta \in [0, \pi]$, is measured in the $z - y$ plane with respect to the positive y -axis. A quantity called the pseudorapidity, $\eta \in [0, \infty]$ is defined as,

$$\eta = -\ln \tan\left(\frac{\theta}{2}\right) \quad (3.4)$$

η is often used as a measure of the polar angle, instead of θ , since the difference in η between two particles, $\Delta\eta$, is invariant under a Lorentz boost in the z -direction [78]. The angular distance between two physics objects, ΔR , can be written as,

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} \quad (3.5)$$

where $\Delta\phi$ is the difference in ϕ between the two physics objects of interest. Quantities defined in the transverse plane are often used to describe the kinematics of physics objects in hadron collider experiments. The transverse momentum, p_T , is defined as,

$$p_T = \sqrt{(p_x)^2 + (p_y)^2} \quad (3.6)$$

where p_x and p_y are the x and y components of the physics object's momenta, respectively. The transverse energy, E_T , is defined as,

$$E_T = \sqrt{m^2 + p_T^2} \quad (3.7)$$

where m is the invariant mass of the physics object.

3.2.2 Inner Detector

The inner detector is the first layer of concentric cylindrical sub-detector layers in the ATLAS detector. It is used to identify charged particles and reconstruct the trajectories of charged particles produced in the collisions via energy deposition in semiconductor material (hits) and the ionisation of gas. It consists of three complementary sub-detectors (in order from nearest to farthest from the beam pipe): the Pixel Detector, the Semiconductor Tracker (SCT) and the Transition Radiation Detector (TRT). The Pixel Detector and SCT are based on semiconductor technology and have the highest granularity of any sub-detector in ATLAS, in order to cope with the high frequency

336 of collisions near the interaction point. The TRT consists of drift tubes (straws) containing a gas mixture, which
 337 allows measurement of the energy deposited by charged particles through the ionisation of the gas. Solenoid magnets
 338 surround the inner detector and bend the trajectories of charged particles. The charges and transverse momenta
 339 of particles can be inferred from their bent trajectories, which are reconstructed by the hits produced via energy
 340 deposition in the Inner Detector.

341 3.2.3 Electromagnetic and Hadronic Calorimeters

342 The Electromagnetic Calorimeter (ECAL) and Hadronic Calorimeter (HCAL) surround the Inner Detector, with
 343 the ECAL nearer to the beam line. The ECAL and HCAL provide accurate measurements of the energy of
 344 particles which interact electromagnetically (e.g. photons and electrons) and hadronically (e.g. jets), respectively.
 345 Particles entering the calorimeters interact with the detector material and create either a electromagnetic shower
 346 (in the ECAL) or a hadronic shower (in the HCAL), depositing all their energy in the calorimeter cells. The
 347 primary mechanism of energy deposition in the ECAL is through bremsstrahlung (for electrons) and pair production
 348 (photons). Hadrons usually deposit a small amount of their energy in the ECAL, and interact via inelastic scattering
 349 with the nuclei of the detector material. The hadronic showers (jets) produced in these nuclear interactions travel
 350 much further than an electromagnetic shower, and for that reason, the volume of the HCAL is designed to occupy
 351 a much larger space than that of the ECAL.

352 3.2.4 Muon Spectrometer

353 The Muon Spectrometer (MS) is the outermost sub-detector of ATLAS and surrounds the HCAL. Muons traverse
 354 through the inner detector and calorimeters, with minimal energy loss, before reaching the MS. The MS consists
 355 of trigger and high-precision tracking systems. Large superconducting toroid shaped magnets deflect the incoming
 356 muons to measure their trajectories and subsequently their momenta via the curvature of the trajectories. The MS
 357 measures muon trajectories as they ionize gas (filled with Ar and CO₂ gas) in the MS drift chambers.

358 3.2.5 Trigger and Data Acquisition System

359 The Trigger and Data Acquisition System (TDAQ) manages and handles the large amount of data produced within
 360 the ATLAS detector. In Run 2, pp bunch crossings occur every 25 ns, corresponding to an event rate of 40 MHz.
 361 The TDAQ system performs a fast preliminary reconstruction to select events with signatures which are interesting
 362 for physics analyses. The information collected from these events are permanently stored for offline reconstruction
 363 and analysis, and the rest (the vast majority of events) are discarded. The trigger system reduces the 40 MHz data
 364 rate to around 1 kHz.

365 3.2.6 Particle Identification and Object Reconstruction

366 Particles originating from pp collisions, or from their subsequent decays, traverse through the ATLAS detector and
 367 interact with its different sub-detectors, producing characteristic electronic signals. These signals are then processed
 368 by various algorithms to reconstruct and identify the physics objects (e.g. electrons, muons, jets) in the event. This
 369 section outlines the procedures used to define these physics objects.

370 3.2.6.1 Tracks and primary vertices

371 The trajectories of charged particles, or tracks, are reconstructed in the ID. First, energy is deposited by charged
 372 particles (hits) in pixels or strips, in the Pixel and SCT detectors respectively. Adjacent pixels or strips are grouped
 373 together in *energy clusters*. Energy clusters define 3D space-points indicating the location where the charged
 374 particle traversed. Track seeds are then defined as sets of three space-points, in either the Pixel or SCT detectors.
 375 A Kalman filter [10] is then used to build track candidates from the track seeds. Often, multiple track candidates
 376 are built per track seed, therefore an ambiguity solver [35] is needed for finding the track which best represents the
 377 traversal of the charged particle. The ambiguity solver ranks each track from a given seed based on, the number of
 378 associated hits, the number of holes (expected hits which are absent), track momenta and the χ^2 of the track fit.
 379 Low ranked tracks are then discarded. High ranked tracks are refitted, introducing information from the TRT.

380 The primary vertex is the location of the pp collision of interest (i.e. from the hard scatter). The primary vertex
 382 from the hard scatter needs to be identified, to isolate the event of interest from unwanted pile-up events. In the

383 event reconstruction procedure [62], the primary vertex is defined as the vertex of the event with the largest sum
 384 of $(p_T)^2$ (corresponding to the measured $(p_T)^2$ of the particle from its reconstructed track) of its associated tracks.
 385 Furthermore, the primary vertex is required to have at least two associated tracks. To reduce contamination from
 386 fake tracks used in primary vertex reconstruction, only tracks which pass certain tight selection criteria are used
 387 in the reconstruction procedure. An iterative fitting procedure is then used to reconstruct the primary vertex by
 388 finding a set of reconstructed tracks which have a common vertex.

389 3.2.6.2 Electrons

390 Since electrons are charged particles, they give rise to tracks in the Inner Detector. They also deposit energy in the
 391 ECAL via electromagnetic showering. Electrons are therefore reconstructed and identified from signals in the Inner
 392 Detector and ECAL. Electrons are reconstructed using a dynamic clustering algorithm [1] which matches electron
 393 candidate tracks in the Inner Detector to energy clusters in the ECAL. The dynamic clustering algorithm matches
 394 tracks to energy clusters which have local maxima, to form electron candidates.

395 A likelihood discriminant is used to identify electrons. Quantities measured in the Inner Detector and ECAL
 396 are used as input, such that they discriminate well between prompt isolated electrons and other physics objects
 397 (e.g. jets, electron from a photon conversion, electron from a semi-leptonically decaying hadron). Important input
 398 variables include the shape of the electromagnetic shower, track quality in the Inner Detector and information from
 399 the TRT.

400 3.2.6.3 Muons

401 Muons leave tracks in the Inner Detector and the MS. They traverse the ECAL and HCAL with no significant
 402 energy loss. Muons are therefore reconstructed and identified from information in the Inner Detector and MS.
 403 Tracks are reconstructed [5] in the Inner Detector and MS independently. Both tracks are combined, using a global
 404 χ^2 fit, resulting in reconstructed muon candidates.

405 Similar to electron identification, muons use a likelihood discriminant to identify prompt muons and suppress
 406 background contamination (mainly from pion and kaon decays).

408 3.2.6.4 Jets and b -tagging

409 Coloured particles emerging from the interaction point result in collimated streams of colourless particles, known
 410 as jets. Jets can deposit energy in the Inner Detector and in the HCAL. Jets in ATLAS are reconstructed from
 411 topological clusters using the anti- k_t algorithm [28]. Topological clusters are groups of adjacent calorimeter cells
 412 which contain energy deposition above the average amount of noise expected in the cell. Adjacent cells are grouped
 413 together under certain criteria to form topological clusters which form jets.

414 Different tagging algorithms are used to identify the quark flavour which initiated a jet. b -quark tagging is used
 415 extensively in top physics, due to the b -quark present in the top quark's dominant decay channel (See Table
 416 1). Hadrons arising from b -quark hadronisation have mean lifetimes ~ 1.5 ps and travel (on average) a few
 417 millimetres before decaying. This creates a secondary vertex within the jet (See Figure 7). This characteristic decay
 418 signature, along with several other unique features of b -jets, are exploited in b -tagging algorithms to distinguish
 419 b -jets from c - or light flavour jets. In Figure 7, an illustration of the production of a b -jet, is shown.

420 In this analysis, we use the recommended DL1r (Deep-Learning Flavour Tagger) tagging algorithm [18]. The DL1r
 421 algorithm combines outputs from several low-level tagging algorithms using a Deep Neural Network and outputs
 422 the probability that a given input jet is identified as a b , c or light flavoured jet.

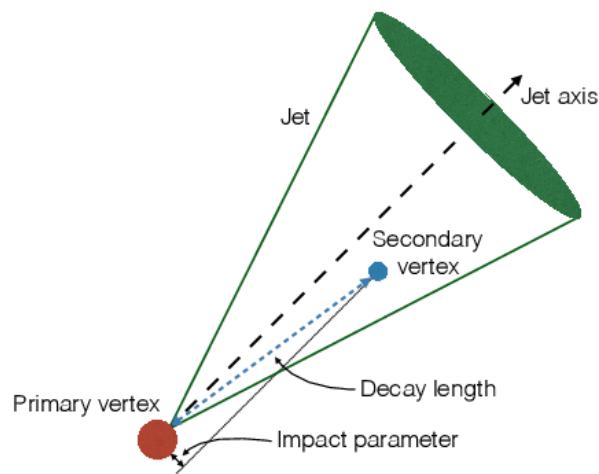


Figure 7: An illustration [38] of the production of a b -jet is shown. This illustrates the presence of a secondary vertex within a b -jet.

424

Chapter 4

425

Analysis Setup and Strategy

426 In this chapter, the setup of the analysis and the analysis strategy is presented for tWZ production in the tetralepton
 427 channel based on an integrated luminosity of 139 fb^{-1} of data recorded by ATLAS. This includes a description of the
 428 data and simulated Monte Carlo (MC) samples, definitions of physics objects, event selection and the subsequent
 429 definition of signal and control regions used in this analysis. Furthermore, estimation of the fake lepton component
 430 using the MC template method is described. Various Machine Learning techniques and a kinematic reconstruction
 431 algorithm, used to discriminate between signal and background events, are described. An outline of the systematic
 432 uncertainties affecting the measurement are presented. Finally, the results of the analysis are presented.

4.1 Data and Monte Carlo Simulation

4.1.1 Data Samples

433 The pp collision data used in this analysis was collected with the ATLAS detector at the LHC from 2015 to 2018. This
 434 data period of data taking is referred to as Run 2. During this period, pp collisions at $\sqrt{s} = 13 \text{ TeV}$, corresponding
 435 to an integrated luminosity (\mathcal{L}) of 156 fb^{-1} , were delivered by the LHC. The ATLAS detector managed to record
 436 147 fb^{-1} of this total delivered data. 139 fb^{-1} of the data recorded by ATLAS is considered to be good enough
 437 for physics analyses (the data passes certain quality control criteria) and placed into the *Good Runs List* [47].
 438 Specific time-intervals where the luminosity is assumed to be constant are known as *lumi-blocks*. Only events from
 439 lumi-blocks in which LHC beams were stable and all ATLAS detectors were operational are selected. The list of
 440 suitable run and lumi-blocks is summarised in the official Good Runs Lists for 2015-2018 data, as specified in
 441 Ref. [goodrunlist]. The integrated luminosities corresponding to the individual datasets for years 2015, 2016,
 442 2017 and 2018 are 3.2 fb^{-1} , 33.0 fb^{-1} , 44.3 fb^{-1} and 58.5 fb^{-1} respectively.
 443

4.1.2 Monte Carlo Samples

444 Simulated MC samples were generated and used to model the SM tWZ signal and its backgrounds.
 445

446 The following background processes are considered:

- 447 • **$t\bar{t}Z$** : $t\bar{t}$ with an associated Z -boson, in the tetralepton final state. Therefore, both top-quarks decay leptonically
 448 (e.g. $t \rightarrow W^+ b \rightarrow \ell^+ \nu b$) and of these top-quarks emits a Z -boson which decays leptonically ($Z \rightarrow \ell^\pm \ell^\mp$
 (OSSF lepton pair)). This results in a final state with 4 leptons and 2 b-quarks.
- 449 • **ZZ** : Diboson production with a tetralepton final state, therefore both Z -bosons decay leptonically ($Z \rightarrow \ell^\pm \ell^\mp$
 450 (OSSF lepton pair)).
- 451 • **other**: Processes with a relatively minimal, but non-negligible background contribution
 - 452 - VVV ($V = W/Z$)
 - 453 - $t\bar{t}$
 - 454 - $t\bar{t}W$
 - 455 - $t\bar{t}WW$

- 459 - $t\bar{t}H$
 460 - WZ
 461 - $t\bar{t}t$
 462 - $t\bar{t}\bar{t}\bar{t}$
 463 - tZq

464 The MC simulations are achieved via the use of event generators and parton shower generators. Event generators
 465 simulate the pp collisions (hard events) by sampling the proton's Parton Distribution Functions (PDFs) at the
 466 desired energy scale. The parton shower generators simulate any incoming or outgoing particles from the hard
 467 process, which carry QCD color charge and can therefore lead to parton showers.

468
 469 The production of tWZ events is simulated with the **MADGRAPH5_AMC@NLO 2.3.3** generator providing matrix element
 470 (ME) calculations at NLO (next-to-leading order) in QCD. The events are interfaced with **PYTHIA 8.235** for the
 471 parton shower. The production of $t\bar{t}Z$ and $t\bar{t}W$ events are simulated with the **MADGRAPH5_AMC@NLO 2.3.3** generator
 472 providing ME calculations at NLO in QCD. The events are interfaced with **PYTHIA 8.210** for the parton shower.
 473 Event generation of tWZ and $t\bar{t}Z$ results in diagrams which overlap with one another, that is, these diagrams contain
 474 the same initial and final state particles. Several methods exist in order to separate between the two processes,
 475 by removing the overlap, therefore avoiding double counting. There are two different diagram removal procedures,
 476 diagram removal procedure 1 (DR1) [40] and diagram removal procedure 2 (DR2). The DR1 scheme is used to
 477 remove the overlap (interference) between tWZ and $t\bar{t}Z$. A comparison to the DR2 scheme is used to estimate
 478 part of the theoretical systematic on the modelling of the tWZ signal (See Section 4.8.2). Diboson processes which
 479 feature the three charged leptons and one neutrino or four charged lepton in their final states, such as WZ and
 480 ZZ , are simulated using **SHERPA 2.2.2** at NLO in QCD precision. The events are interfaced with **SHERPA** for the
 481 parton shower. Triboson processes such as WWW , WWZ , WZZ , and ZZZ containing up to six leptons in their
 482 final states are simulated using **SHERPA 2.2.2** at NLO in QCD precision. The events are interfaced with **SHERPA** for
 483 the parton shower. The production of $t\bar{t}$ events are simulated with the **POWHEG** generator providing ME calculations
 484 at NLO in QCD. The events are interfaced with **PYTHIA 8.210** for the parton shower. The production of $t\bar{t}t$, $t\bar{t}\bar{t}$
 485 and $t\bar{t}WW$ are simulated using the **MADGRAPH5_AMC@NLO 2.2.2** generator at LO in QCD precision. The events are
 486 interfaced with **PYTHIA 8.186** for the parton shower. The production of $t\bar{t}$ with an associated Higgs boson, $t\bar{t}H$, are
 487 generated using the **MADGRAPH5_AMC@NLO 2.6.0** generator at NLO in QCD precision. The events are showered using
 488 **PYTHIA 8.230**. The production of $t\bar{t}$ events are simulated with the **POWHEG** generator providing ME calculations at
 489 NLO in QCD. The events are showered using **PYTHIA 8.230**. The production of a single top quark in association
 490 with a Z -boson and an extra parton, tZq , is simulated using **MADGRAPH5_AMC@NLO 2.3.3** at NLO in QCD precision.
 491 The events are interfaced with **PYTHIA 8.230** for the parton shower. In Table 3, the event generator and parton
 492 shower used for each process's sample are shown.

Process	Event Generator	Cross section calculation	Parton Shower
tWZ	MADGRAPH5_AMC@NLO 2.3.3	NLO	PYTHIA 8.235
$t\bar{t}Z$	MADGRAPH5_AMC@NLO 2.3.3	NLO	PYTHIA 8.210
ZZ, WZ	SHERPA 2.2.2	NLO	SHERPA
$VVV (V = W/Z)$	SHERPA 2.2.2	NLO	SHERPA
$t\bar{t}$	POWHEG	NLO	PYTHIA 8.230
$t\bar{t}W$	MADGRAPH5_AMC@NLO 2.3.3	NLO	PYTHIA 8.210
$t\bar{t}WW$	MADGRAPH5_AMC@NLO 2.2.2	LO	PYTHIA 8.186
$t\bar{t}H$	MADGRAPH5_AMC@NLO 2.6.0	NLO	PYTHIA 8.230
$t\bar{t}t, t\bar{t}\bar{t}$	MADGRAPH5_AMC@NLO 2.2.2	LO	PYTHIA 8.186
tZq	MADGRAPH5_AMC@NLO 2.3.3	NLO	PYTHIA 8.230

Table 3: The event generator and parton shower used for the signal and background process's MC samples is shown.

4.2 Objects

494 In this section the physics objects (leptons, jets and b -tagged jets) used in this analysis are outlined.

4.2.1 Leptons

In this analysis only e and μ leptons are considered, since τ leptons are difficult to detect in the ATLAS detector. τ leptons are challenging to detect since they have an extremely short lifetime (290.3 ± 0.5 fs [48]) which causes them to decay before reaching any detector components and therefore can only be reconstructed via their decay products. In addition to our selection criteria of exactly four leptons, it is required that the Leading (L), Next-to-Leading (NL), Next-to-Next-to-Leading (NNL) and Next-to-Next-to-Next-to-Leading (NNNL) leptons have p_T greater than 28, 18, 10 and 10 GeV respectively. Relatively loose object-level cuts are chosen in an attempt to maximize our signal statistics, since the analysis is heavily statistically limited. Reconstructed electrons are required to be within $|\eta| < 2.47$ and excluding the transition region between the barrel and end-cap calorimeters at $1.37 < |\eta| < 1.52$. Reconstructed muons are required to be within $|\eta| < 2.5$. The transverse impact parameter, d_0 , is defined as the minimal spatial distance between the object's (referring to leptons) trajectory and the primary vertex (the vertex associated with the p - p hard scatter). The longitudinal impact parameter, z_0 , is defined as the value of z of the point on the object's trajectory which determines d_0 . To ensure consistency between the lepton and the primary vertex, it is required that $|\frac{d_0}{\sigma(d_0)}| < 5$, $|z_0 \sin \theta| < 0.5$ mm for electrons and $|\frac{d_0}{\sigma(d_0)}| < 3$, $|z_0 \sin \theta| < 0.5$ mm for muons, following the current recommendations [76]. To avoid instances where one detector signal can result in multiple different reconstructed objects, an overlap removal is applied which ignores all but one of these objects (See Section 4.2.4). Electrons are selected using a likelihood based discriminant [1] which takes measurements from the tracking system, calorimeter system and quantities derived from both the tracking and calorimeter system as input. Muons are selected using the Muon Selection Tool [66]. Loose electrons are defined with the criteria above, using the `LooseAndBLayerLH` ($\sim 91\%$ selection efficiency for electrons with $E_T > 30$ GeV [41]) identification algorithm (which has a certain cut applied). Similarly, tight electrons are defined with the criteria above, using the `TightLH` ($\sim 80\%$ selection efficiency for electrons with $E_T > 30$ GeV [41]) algorithm (which has a certain cut applied). Both loose and tight muons use the `Medium` ($\sim 95\%$ selection efficiency [5]) algorithm (which has a certain cut applied). Tight leptons additionally require that they are sufficiently isolated from other particles produced in the collision. This is done by defining a cone of radius $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$ around the particle of interest and summing the p_T of all the reconstructed particles surrounding the particle of interest, situated within the cone. A quantity, I_{rel} , is then defined as $I_{rel} = \frac{\sum p_T(\text{surrounding candidate})}{p_T(\text{candidate})}$, the ratio of this sum to the p_T of the lepton candidate. If this value is large, it is likely that the particle of interest originated from a jet (together with many other particles), whereas a prompt decay product resulting from the hard scatter will have little to no energy surrounding it ($I_{rel} \ll 1$). The `IsolationSelectionTool` with the `PLVTight` ($\sim 70\%$ efficiency [7]) and `PLVTight` ($\sim 70\%$ efficiency at $p_T = 30$ GeV [34]) algorithm are used for tight electrons and tight muons respectively (following the current recommendations [70]). In Table 4, a summary of the selection criteria for leptons is shown.

	Electrons		Muons	
	Tight	Loose	Tight	Loose
p_T cuts		$p_T(\ell_1, \ell_2, \ell_3, \ell_4) > (28, 18, 10, 10)$ GeV		
Overlap Removal		Described in Section 4.2.4		
η cuts	$ \eta(\ell_e) < 2.47$ excluding $1.37 < \eta(\ell_e) < 1.52$		$ \eta(\ell_\mu) < 2.5$	
Impact Parameters	$ \frac{d_0}{\sigma(d_0)} < 5$, $ z_0 \sin \theta < 0.5$ mm		$ \frac{d_0}{\sigma(d_0)} < 3$, $ z_0 \sin \theta < 0.5$ mm	
Identification WP	TightLH	LooseAndBLayerLH	Medium	Medium
Isolation WP	PLVTight	Not Used	PLVTight	Not Used

Table 4: A summary of the requirements applied for selecting tight and loose leptons (e, μ) is shown.

4.2.2 Jets

Jets are reconstructed using the anti- k_t algorithm (See Section 3.2.6.4). The `AntiKt4EMPFflowjets` ($\sim 97\%$ average efficiency with JVT (outlined in the subsequent paragraph) > 0.2 [76]) algorithm (which has a certain cut applied) is used, following the current recommendations [76]. The jet-vertex-tagger (JVT) and the forward jet-vertex-tagger (fJVT) are likelihood discriminant which aim to suppress pile-up jets. The `Medium` algorithm (which has a certain cut applied) is used for the JVT and the fJVT (following the current recommendations [69]). Additionally, a requirement that jets have a JVT value greater than 0.5 is applied. In the same way as with leptons, ambiguities are removed where one detector signal can result in multiple different reconstructed objects, via overlap removal

(See Section 4.2.4). Jets are required to be within $p_T(\text{jet}) > 20 \text{ GeV}$. Looser p_T cuts are applied in an attempt to increase our limited signal statistics. A forward jet is a signature of single top production, jets are therefore required to have $|\eta| < 4.5$ in order to include these forward jets.

4.2.3 b -tagging

The DL1r b -tagger [65] was used to identify jets as b -jets (See Section 3.2.6.4). Different DL1r working points are used to identify b -jets in our event selection (See Section 4.4). The working points are defined based off a cut on the DL1r score corresponding to a b -jet tagging efficiency of 60%, 70%, 77% and 85%.

Since this analysis is heavily statistically limited, the amount of statistics in our regions are aimed to be maximized. In an attempt to achieve this goal in the $t\bar{t}Z$ CR, b -tagged jets were placed under *tight* and *loose* definitions. A tight b -tagged jet is defined as a jet which passes the 77%, 70%, 65% or 60% DL1r b -tagger working point. A loose b -tagged jet is defined as a jet which passes 85% DL1r b -tagger working point, but not the 77%, 70%, 65% or 60% DL1r b -tagger working points. Different numbers (and definitions) of tight and loose b -tagged jets were tried in each region, with the final selection criteria being chosen which maximised the expected significance of $\sigma(tWZ)$ (See Section 4.4.1).

4.2.4 Overlap Removal Procedure

The overlap removal procedure is used on pre-selected leptons and jets. It is performed sequentially, in the following steps:

1. If the separation between a pre-selected electron and pre-selected muon is within $\Delta R < 0.01$, or they share a track, the pre-selected electron is discarded.
2. If the separation between a jet and a pre-selected electron is within $\Delta R < 0.2$, the jet is discarded.
3. Any remaining electron or muon closer than $\Delta R = 0.4$ to a jet, is discarded.
4. If the distance between a jet and a pre-selected muon is $\Delta R < 0.4$ and the jet has more than two associated tracks, then the muon is discarded, otherwise the jet is discarded.

4.3 Kinematic cuts

In order to suppress potential fakes and quarkonia (low mass resonances such as J/ψ and upsilon) a requirement that all OSSF lepton pairs have an invariant mass, m_{OSSF} , greater than 10 GeV is applied. The final state lepton charges must sum to zero. Therefore a requirement of $\sum_{i=1}^4 \text{charge}(\ell_i) = 0$ is applied. The invariant mass of the OSSF lepton pair coming from the Z boson must equal the invariant mass of the Z boson, and noting that e,μ reconstruction and identification in the ATLAS detector has a high efficiency [58], these OSSF leptons are used to reconstruct Z bosons with relatively high confidence. A Z candidate is defined in this analysis as an OSSF lepton pair with an invariant mass, m_{OSSF} , satisfying the condition, $|m_{\text{OSSF}} - m_Z| < 30 \text{ GeV}$, where $m(Z)$ is the nominal Z boson mass (91.1876 GeV [48]). This wider mass window is used in order to cover the full range of the $m(Z)$ distribution, in an attempt to increase the number of events which pass our baseline selections. Multiple Z candidates can be present in certain decay channels (e.g. $eeee$, $\mu\mu ee$, $\mu\mu\mu\mu$). In these cases, the Z candidate which has an invariant mass closest to the nominal Z boson mass is chosen.

4.4 Regions and Event Selection

Two tWZ SRs are defined in an attempt to suppress and constrain the ZZ background. Both tWZ SRs are required to have exactly four tight leptons, exactly one Z -boson candidate, exactly one tight b -tagged jet (from the decay of the top quark) and greater than or equal to one jet. The two tWZ SR's differ by the flavours of their leptons which don't originate from the decay of a Z -boson (non- Z leptons). The ZZ background has two Z -bosons which decay into a pair of OSSF lepton pairs, in order to mimic the tWZ signal. This is taken advantage of, to define a tWZ region enrich in ZZ background and one with a minimal ZZ background component. This is done by requiring that one of the tWZ SRs has its two non- Z leptons to have opposite flavour and the other tWZ SR

Baseline selections				
$N_\ell = 4$ $p_T(\ell_1, \ell_2, \ell_3, \ell_4) > (28, 10, 10, 10)$ GeV $p_T(\text{jet}) > 20$ GeV, $ \eta(\text{jet}) < 4.5$, $\text{jvt} > 0.5$ $ \eta(\ell_e) < 2.47$ excluding $1.37 < \eta(\ell_e) < 1.52$ $ \eta(\ell_\mu) < 2.5$ $\sum_{i=1}^4 \text{charge}(\ell_i) = 0$ All OSSF lepton pairs require $m_{\text{OSSF}} > 10$ GeV				
Regions				
tWZ OF SR	tWZ SF SR	$t\bar{t}Z$ CR	ZZb CR	$(tWZ)_{\text{fake}}$ CR
$N_\ell(\text{tight}) = 4$	$N_\ell(\text{tight}) = 4$	$N_\ell(\text{tight}) = 4$	$N_\ell(\text{tight}) = 4$	$N_\ell(\text{tight}) = 3$ $N_\ell(\text{loose and NOT tight}) = 1$
N_Z candidate = 1	N_Z candidate = 1	N_Z candidate = 1	N_Z candidate = 2	N_Z candidate = 1
$N_{\text{jet}} \geq 1$	$N_{\text{jet}} \geq 1$	$N_{\text{jet}} \geq 2$	$N_{\text{jet}} \geq 1$	$N_{\text{jet}} \geq 1$
$N_{\text{b-jet}}(\text{tight}) = 1$	$N_{\text{b-jet}}(\text{tight}) = 1$	$N_{\text{b-jet}}(\text{tight}) \geq 1$ $N_{\text{b-jet}}(\text{loose}) \geq 0$ $N_{\text{b-jet}}(\text{tight}) + N_{\text{b-jet}}(\text{loose}) = 2$	$N_{\text{b-jet}}(\text{tight}) = 1$	$N_{\text{b-jet}}(\text{tight}) = 1$
Opp. Flavour non-Z leptons	Same Flavour non-Z leptons	-	-	-

Table 5: A summary of the requirements applied for selecting events in the signal and control regions is shown.

is required to have its non- Z leptons to have the same flavour. These two disjoint tWZ SRs are named tWZ OF SR and tWZ SF SR respectively. It is therefore expected that the tWZ SF SR contains the majority of the ZZ background events across both tWZ SRs.

In order to check the modelling of the most dominant background components in our signal region, $t\bar{t}Z$ and ZZb control regions are defined. The $t\bar{t}Z$ control region has the same requirement on the number of reconstructed Z boson candidates in the signal region (due to a commonality on the number of Z bosons present in both processes), however it is required that there are at least two jets and that exactly two of these jets are b -tagged (corresponding to the b -quark jets originating from the two top-quark decays). A ZZb region is defined, as opposed to a ZZ region, since the ZZ background present in the tWZ signal region contains exactly one b -tagged jet. Therefore defining a region with ZZ plus exactly one b -jet more closely resembles the ZZ background present in the signal region. In addition to this, mis-modelling of ZZ has been seen in other analyses [3, 36], further motivating the use of a ZZb control region over a ZZ CR. The ZZb CR requires exactly two Z boson candidates and exactly one b -tagged jet, resulting in an implicit requirement on the number of jets ($N_{\text{jet}} \geq 1$).

Fake leptons are objects reconstructed as leptons, but do not correspond to the leptons that are of interest in our analysis. Fake leptons can be split up into two main categories, irreducible (prompt) fakes and reducible (non-prompt) fakes. Irreducible fakes are true leptons which do not come from the process of interest. Reducible fakes are objects which are mis-identified or incorrectly reconstructed as leptons. In the ATLAS detector, the probability for a fake to occur is very low. In order to constrain the fake lepton component contained within the $t\bar{t}Z$ sample, a $(tWZ)_{\text{fake}}$ CR is defined which is as similar as possible to the tWZ CR but is enhanced in fakes. This is achieved by defining the $(tWZ)_{\text{fake}}$ CR to inherit the same selection criteria as the tWZ SRs however, in this case, a requirement of exactly 3 tight leptons and exactly 1 loose (and NOT tight) lepton (since looser leptons are more likely to be fakes, compared to tighter leptons) is applied. A $(t\bar{t}Z)_{\text{fake}}$ CR (requiring exactly 3 tight leptons and exactly 1 loose (and NOT tight) lepton) was tried as an alternative to the $(tWZ)_{\text{fake}}$ CR, however a much larger suppression of fakes were observed in this region (compared to the $(tWZ)_{\text{fake}}$ CR). This suppression of fakes can be explained by the extra b -tagged jet requirement (exactly two b -tagged jets are required in the $t\bar{t}Z$ region, compared to exactly one in the tWZ SRs (See Table 5)) which causes suppression of fakes via the overlap removal procedure (See Section 4.2.4).

In Table 5, a summary of the final selection criteria and region definitions is shown.

4.4.1 Optimization studies for event selection

In order to find the selection criteria for jets and leptons which maximized sensitivity to the tWZ signal, studies were performed by plotting the expected significance (Z_μ^{exp}) and expected upper limit (μ_{up}^{exp}) for different selection criteria. The fitting procedure as described in Section 4.9.1 was used to calculate the expected upper limits and expected significances in this study (these metrics are fully described later in Section 4.9.1). The same selection criteria and regions defined in Table 5 was used (unless otherwise specified), except for the selection(s) which were being optimised in each case.

In Figure 8 the expected significance (Z_μ^{exp}) and expected upper limits (μ_{up}^{exp}) for different $\eta(jet)$ cuts are shown.

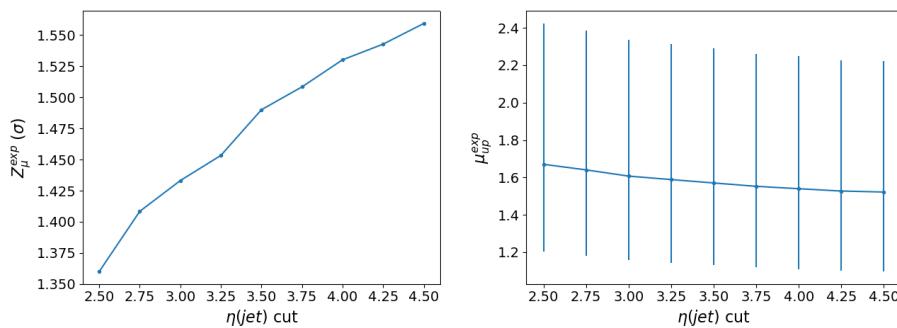


Figure 8: **Left:** Expected significance (Z_μ^{exp}) for different $\eta(jet)$ cuts is shown. The cuts applied on the $\eta(jet)$ are shown on the x-axis and corresponding expected significance from the likelihood fit is shown on the y-axis. **Right:** Expected upper limit (μ_{up}^{exp}) for different $\eta(jet)$ cuts is shown. The cuts applied on the $\eta(jet)$ are shown on the x-axis and corresponding expected upper limits are shown on the y-axis. Error bars representing the total uncertainty on the expected upper limits are shown as vertical lines.

From Figure 8, it can be seen that the $\eta(jet)$ cut which maximises the sensitivity of tWZ in the tetralepton channel is requiring that $\eta(jet) < 4.5$. This selection criteria was set for the $\eta(jet)$ across all regions.

In Figure 9 the expected significance (Z_μ^{exp}) and expected upper limits (μ_{up}^{exp}) for different $p_T(jet)$ cuts are shown.

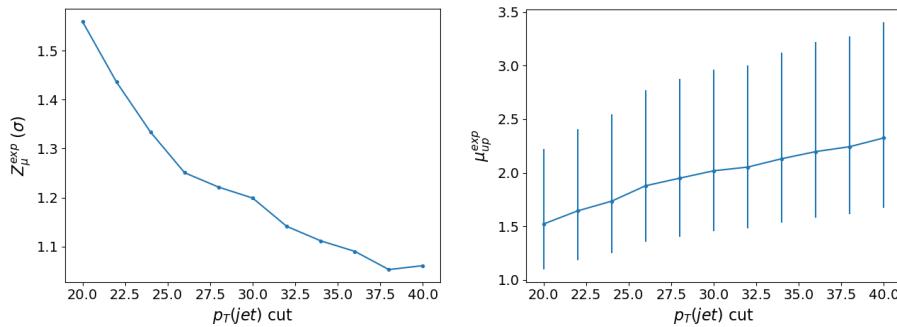


Figure 9: **Left:** Expected significance (Z_μ^{exp}) for different $p_T(jet)$ cuts is shown. The cuts applied on the $p_T(jet)$ are shown on the x-axis and corresponding expected significance from the likelihood fit is shown on the y-axis. **Right:** Expected upper limit (μ_{up}^{exp}) for different $p_T(jet)$ cuts is shown. The cuts applied on the $p_T(jet)$ are shown on the x-axis and corresponding expected upper limits are shown on the y-axis. Error bars representing the total uncertainty on the expected upper limits are shown as vertical lines.

From Figure 9, it can be seen that the $p_T(jet)$ cut which maximises the sensitivity of tWZ is requiring that $p_T(jet) > 20$ GeV. This selection criteria was set for the $p_T(jet)$ across all regions.

In Figure 10 the expected significance (Z_μ^{exp}) and expected upper limits (μ_{up}^{exp}) for a range of different configurations of DL1r b -tagged jet working points across different regions. From Figure 10, it can be seen that requiring that b -tagged jets pass the 77% DL1r WP in the tWZ SR, $(tWZ)_{fake}$ CR and the ZZb CR and that at least one b -tagged jet in the $t\bar{t}Z$ SR passes the 77% DL1r WP (the other jet is

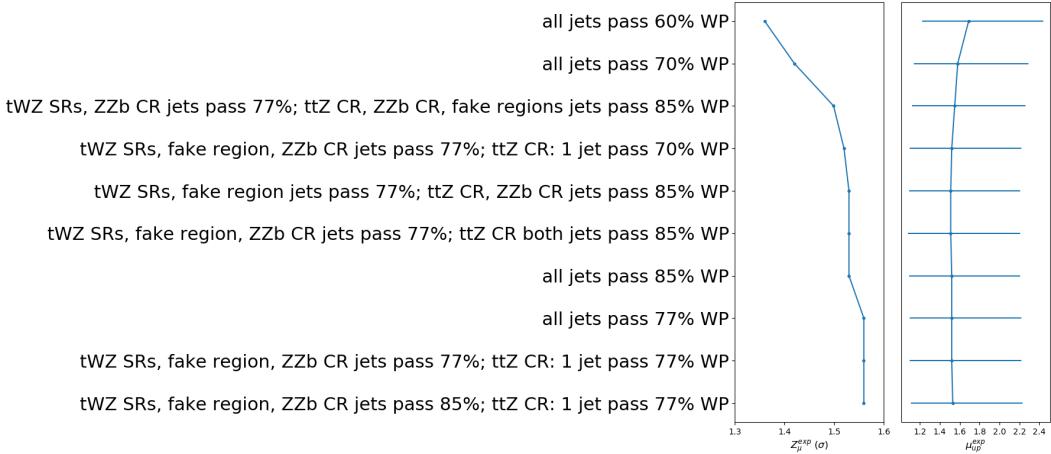


Figure 10: Expected significance (Z_μ^{exp}) and expected upper limit ($\mu_{\text{up}}^{\text{exp}}$) for different configurations of DL1r b -tagged jet working points is shown. The common y-axis shows the different configurations of DL1r b -tagged jet working points. On the left panel, the expected significance from the likelihood fit is shown on the x-axis. On the right panel, the expected upper limit from the likelihood fit is shown on the x-axis (with the corresponding total uncertainty represented by horizontal lines).

629 just required to pass the 85% DL1r WP) maximises the sensitivity overall (compared to the other investigated
 630 configurations). This configuration was chosen b -tagged jets.

631
 632 The p_T (L Lepton) is constrained by the single lepton triggers (Table ??). A cut was chosen to be applied on the
 633 p_T (NL Lepton) slightly tighter than the tightest single lepton p_T cut in the trigger. The p_T (NL Lepton) cut can
 634 be optimized by comparing the expected significance and limit for a range of p_T (NL Lepton) cuts to determine the
 635 cut which maximizes sensitivity.

636
 637 In Figure 11 the expected significance (Z_μ^{exp}) and expected upper limits ($\mu_{\text{up}}^{\text{exp}}$) for different p_T (NL Lepton) cuts is
 shown.

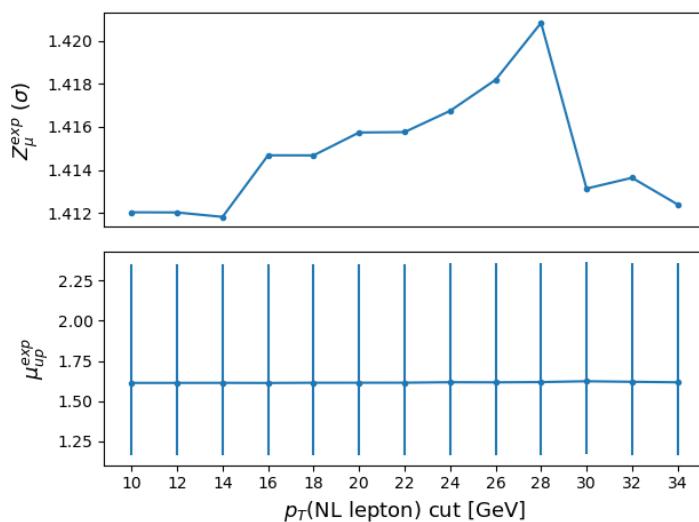


Figure 11: Expected significance (Z_μ^{exp}) and expected upper limit ($\mu_{\text{up}}^{\text{exp}}$) for different p_T (NL Lepton) cuts is shown. The common x-axis shows cut applied to the p_T of the next-to-leading lepton. On the top panel, the expected significance from the likelihood fit is shown on the y-axis. On the bottom panel, the expected upper limit from the likelihood fit is shown on the y-axis (with the corresponding total uncertainty represented by vertical lines).

639 Since there is a very small change between the different p_T (NL Lepton) cuts on the sensitivity of tWZ , a
 640 p_T (NL Lepton) cut is applied at 18 GeV (avoiding a p_T cut near the sharp drop in expected significance after
 641 28 GeV), therefore applying a cut above the tightest, looser dilepton trigger p_T cut (17 GeV) to suppress any
 642 systematic from the modelling of the trigger efficiency.

643 4.5 Signal and Control Regions

644 In this section, expected number of events of variables in each region are shown. For each figure in this
 645 section, the data is given by the black points and the MC predictions for each process are given by the filled
 646 histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the
 647 diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The
 648 lower panel in each plot shows the ratios of the data to the theoretical predictions. In order to suppress a bias
 649 towards large signal observations in the development of the analysis, data has not been analysed in bins where the
 650 expected $\frac{\text{signal}}{\text{background}}$ exceeds 0.1. Blinded bins are shaded with black diagonal lines and their data points are omitted.
 651

In Table 6, the expected number of events for each sample in each region are shown. The finite number of events

	tWZ OF SR	tWZ SF SR	$t\bar{t}Z$ CR	ZZb CR	$(tWZ)_{\text{fake}}$ CR
$t\bar{t}Z$	13.9325 ± 1.84643	10.1343 ± 1.36039	31.7149 ± 4.46776	5.26303 ± 0.696828	19.1224 ± 2.50011
$t\bar{t}Z$ fakes	0.0687541 ± 0.0482172	0.032827 ± 0.026286	0.0709734 ± 0.043509	0.0474576 ± 0.0301512	4.94775 ± 2.48939
tWZ	3.81359 ± 0.392241	2.57584 ± 0.326401	2.61991 ± 0.861557	1.4023 ± 0.156686	4.93485 ± 0.692143
ZZ	0.546045 ± 0.18975	8.76232 ± 2.66871	1.22357 ± 0.376889	46.0616 ± 13.9203	7.76724 ± 2.36894
other	$t\bar{t}$	$6e-06 \pm 3.04506e-06$	0.250783 ± 0.44226	0.269883 ± 0.223373	$6e-06 \pm 3.04506e-06$
	tZq	0.0827265 ± 0.0399222	0.0757694 ± 0.0355101	0.0637132 ± 0.0293762	0.0590199 ± 0.0244576
	$t\bar{t}W$	$0.00674747 \pm 0.00793546$	$0.00279491 \pm 0.00287747$	$6e-06 \pm 3.04506e-06$	$0.00221727 \pm 0.00562041$
	WZ	0.0439316 ± 0.0241635	0.0397876 ± 0.0154764	0.0134837 ± 0.0128327	0.0474188 ± 0.0330635
	$t\bar{t}t$	$0.000987429 \pm 0.000768187$	$0.00249801 \pm 0.00138007$	0.0141085 ± 0.00486102	$6e-06 \pm 3.04506e-06$
	$t\bar{t}\bar{t}$	0.00934516 ± 0.0080725	0.0107503 ± 0.00852049	0.0570846 ± 0.0206271	$6e-06 \pm 3.04506e-06$
	$t\bar{t}WW$	0.0293456 ± 0.0263573	0.0296011 ± 0.0196075	0.26412 ± 0.0936908	0.013096 ± 0.0323943
	$VVV (V = W/Z)$	0.280384 ± 0.0866421	0.191257 ± 0.0595588	0.0696624 ± 0.0228108	0.171171 ± 0.0526519
	$t\bar{t}H$	0.854064 ± 0.177974	0.674566 ± 0.141771	1.98187 ± 0.406211	0.151447 ± 0.0357703
Total	19.6684 ± 1.95158	22.7832 ± 3.10338	38.3633 ± 4.6342	53.2187 ± 13.9618	49.5163 ± 4.77745
data	-	-	36	49	57

Table 6: The expected number of events for each sample in each region is shown.

652 expected to be observed in data (MC simulation) carries an associated statistical uncertainty. To first order, this
 653 uncertainty can be written as the square root of the expected number of events to be observed in data. In contrast
 654 to this, predictions based on MC simulation carry uncertainties due to the finite number of simulated events utilised.
 655 This uncertainty can be quantified by the Number of Equivalent Events [43], N_{equiv} , which relates the sample of
 656 N events (weighted by MC event weights) to N_{equiv} events with all MC event weights equal to 1, that would have
 657 the same relative statistical fluctuation. The Number of Equivalent Events, N_{equiv} , can be written as,

$$N_{\text{equiv}} = \frac{(\sum_i^N w_i)^2}{\sum_i^N w_i^2} \quad (4.1)$$

658 where w_i is the MC event weight for event i . The standard uncertainty of N_{equiv} is given by $u(N_{\text{equiv}}) = \sqrt{N_{\text{equiv}}}$.

659 In Table 7, the number of equivalent events, N_{equiv} and its percentage uncertainty ($\frac{u(N_{\text{equiv}})}{N_{\text{equiv}}} \times 100$), is shown for
 660 each sample in each region.

661 N_{equiv} is relatively large, compared to the background processes, for the tWZ signal in all regions. This tells us
 662 that we have a large number of raw MC events for the tWZ sample. An extended tWZ sample was generated to
 663 increase the number of signal events used to train the event-level BDT (See Section 4.7.3).

664 4.5.1 tWZ OF SR

665 In this section, expected number of events of variables in the tWZ OF SR are shown.

666 In Figure 12 MC predictions for p_T , η and ϕ for leading (L) leptons and next-to-leading (NL) leptons in the tWZ
 667 OF SR region is shown.

	tWZ OF SR	tWZ SF SR	$t\bar{t}Z$ CR	ZZb CR	(tWZ) fake CR			
	N_{equiv}	Uncertainty [%]	N_{equiv}	Uncertainty [%]	N_{equiv}	Uncertainty [%]	N_{equiv}	Uncertainty [%]
tWZ	6463.29 ± 80.39	1.24	4153.0 ± 64.44	1.55	4800.67 ± 69.29	1.44	2497.07 ± 49.97	2.0
$t\bar{t}Z$	1363.87 ± 36.93	2.71	1031.04 ± 32.11	3.11	3237.01 ± 56.89	1.76	561.41 ± 23.69	4.22
ZZ	50.89 ± 7.13	14.02	975.61 ± 31.23	3.2	267.85 ± 16.37	6.11	7023.35 ± 83.81	1.19
other	748.0 ± 27.35	3.66	2.47 ± 1.57	63.58	4.23 ± 2.06	48.6	255.32 ± 15.98	6.26
$t\bar{t}Z$ fakes	6.72 ± 2.59	38.56	1.31 ± 1.14	87.34	16.05 ± 4.01	24.96	7.15 ± 2.67	37.41
Total	8632.77 ± 92.91	1.08	6163.44 ± 78.51	1.27	8325.81 ± 91.25	1.1	10344.3 ± 101.71	0.98
							12044.27 ± 109.75	0.91

Table 7: The number of equivalent events, N_{equiv} and its percentage uncertainty ($\frac{u(N_{equiv})}{N_{equiv}} \times 100$), is shown for each sample in each region.

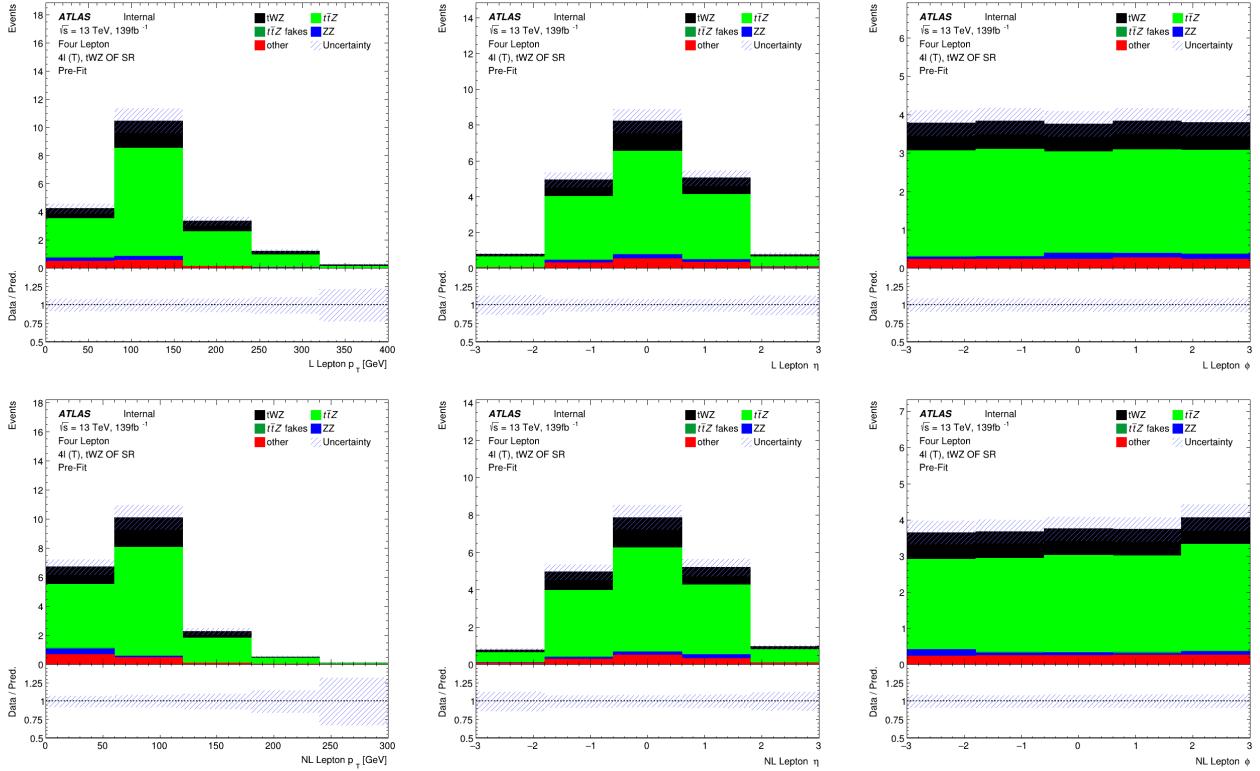


Figure 12: MC predictions for p_T , η and ϕ for leading (L) leptons (top row) and next-to-leading (NL) leptons (bottom row) in the tWZ OF SR region is shown.

- 671 In Figure 13 MC predictions for p_T , η and ϕ for leading (L) jets and next-to-leading (NL) jets in the tWZ OF SR
672 region is shown.
673 In Figure 14 MC predictions for p_T , η and ϕ of the next-to-next-to-leading (NNL) jets, H_T (scalar sum of Jet p_T)
674 and the Number of jets in the tWZ OF SR region is shown.
675 In Figure 15 MC predictions for p_T , η and ϕ of the leading b-tagged jets, the scalar sum of b-tagged jet p_T and the
676 Number of b-tagged jets in the tWZ OF SR region is shown.

677 4.5.2 tWZ SF SR

- 678 In this section, expected number of events of variables in the tWZ SF SR are shown.
679
680 In Figure 16 MC predictions for p_T , η and ϕ for leading (L) leptons and next-to-leading (NL) leptons in the tWZ
681 SF SR region is shown.
682 In Figure 17 MC predictions for p_T , η and ϕ for leading (L) jets and next-to-leading (NL) jets in the tWZ SF SR
683 region is shown.
684 In Figure 18 MC predictions for p_T , η and ϕ of the next-to-next-to-leading (NNL) jets, H_T (scalar sum of Jet p_T)
685 and the Number of jets in the tWZ SF SR region is shown.
686 In Figure 19 MC predictions for p_T , η and ϕ of the leading b-tagged jets, the scalar sum of b-tagged jet p_T and the

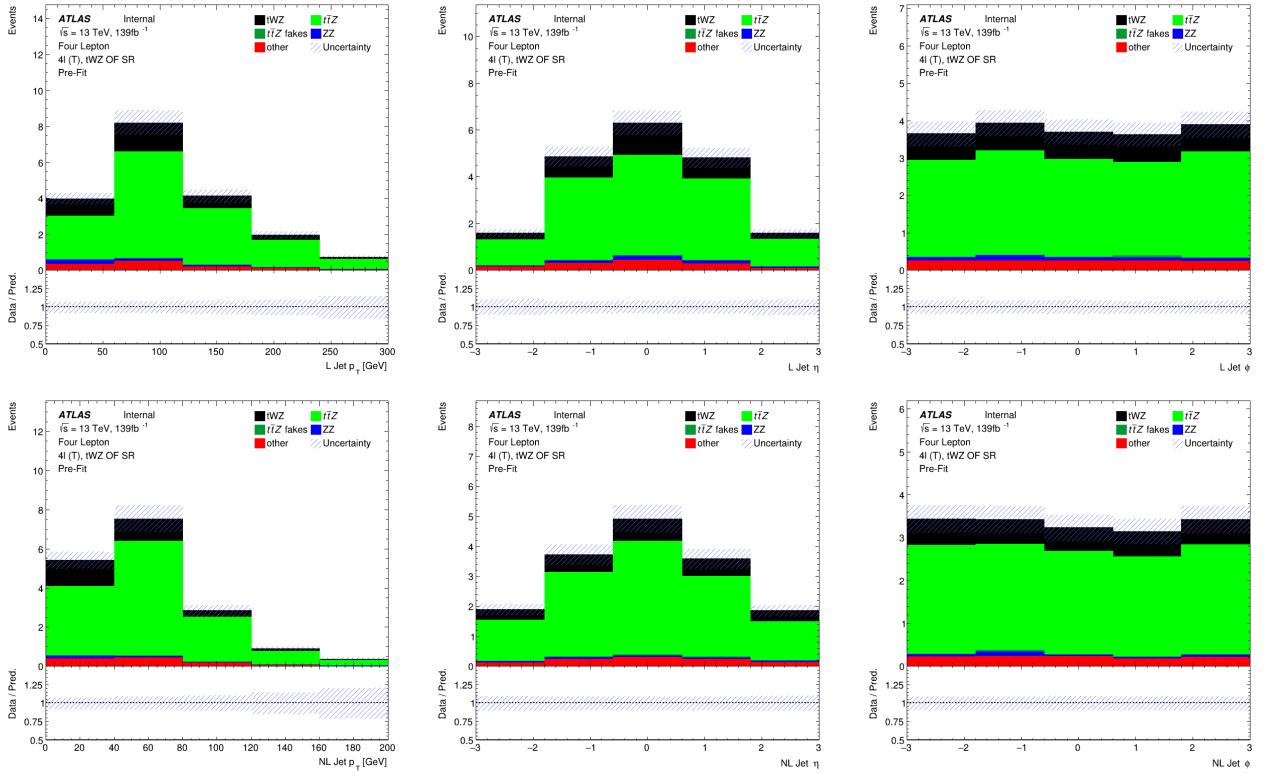


Figure 13: MC predictions for p_T , η and ϕ for leading (L) jets (top row) and next-to-leading (NL) jets (bottom row) in the tWZ OF SR region is shown.

687 Number of b-tagged jets in the tWZ SF SR region is shown.

688 4.5.3 $t\bar{t}Z$ CR

689 In this section, expected number of events of variables in the $t\bar{t}Z$ CR are shown.

690
691 In Figure 20 MC predictions for p_T , η and ϕ for leading (L) leptons and next-to-leading (NL) leptons in the $t\bar{t}Z$
692 CR region is shown.

693 In Figure 21 MC predictions for p_T , η and ϕ for leading (L) jets and next-to-leading (NL) jets in the $t\bar{t}Z$ CR region
694 is shown.

695 In Figure 22 MC predictions for p_T , η and ϕ of the next-to-next-to-leading (NNL) jets, H_T (scalar sum of Jet p_T)
696 and the Number of jets in the $t\bar{t}Z$ CR region is shown.

697 In Figure 23 MC predictions for p_T , η and ϕ of the leading b-tagged jets, the scalar sum of b-tagged jet p_T and the
698 Number of b-tagged jets in the $t\bar{t}Z$ CR region is shown.

699 4.5.4 ZZb CR

700 In this section, expected number of events of variables in the ZZb CR are shown.

701
702 In Figure 24 MC predictions for p_T , η and ϕ for leading (L) leptons and next-to-leading (NL) leptons in the ZZb
703 CR region is shown.

704 In Figure 25 MC predictions for p_T , η and ϕ for leading (L) jets and next-to-leading (NL) jets in the ZZb CR region
705 is shown.

706 In Figure 26 MC predictions for p_T , η and ϕ of the next-to-next-to-leading (NNL) jets, H_T (scalar sum of Jet p_T)
707 and the Number of jets in the ZZb CR region is shown.

708 In Figure 27 MC predictions for p_T , η and ϕ of the leading b-tagged jets, the scalar sum of b-tagged jet p_T and the
709 Number of b-tagged jets in the ZZb CR region is shown.

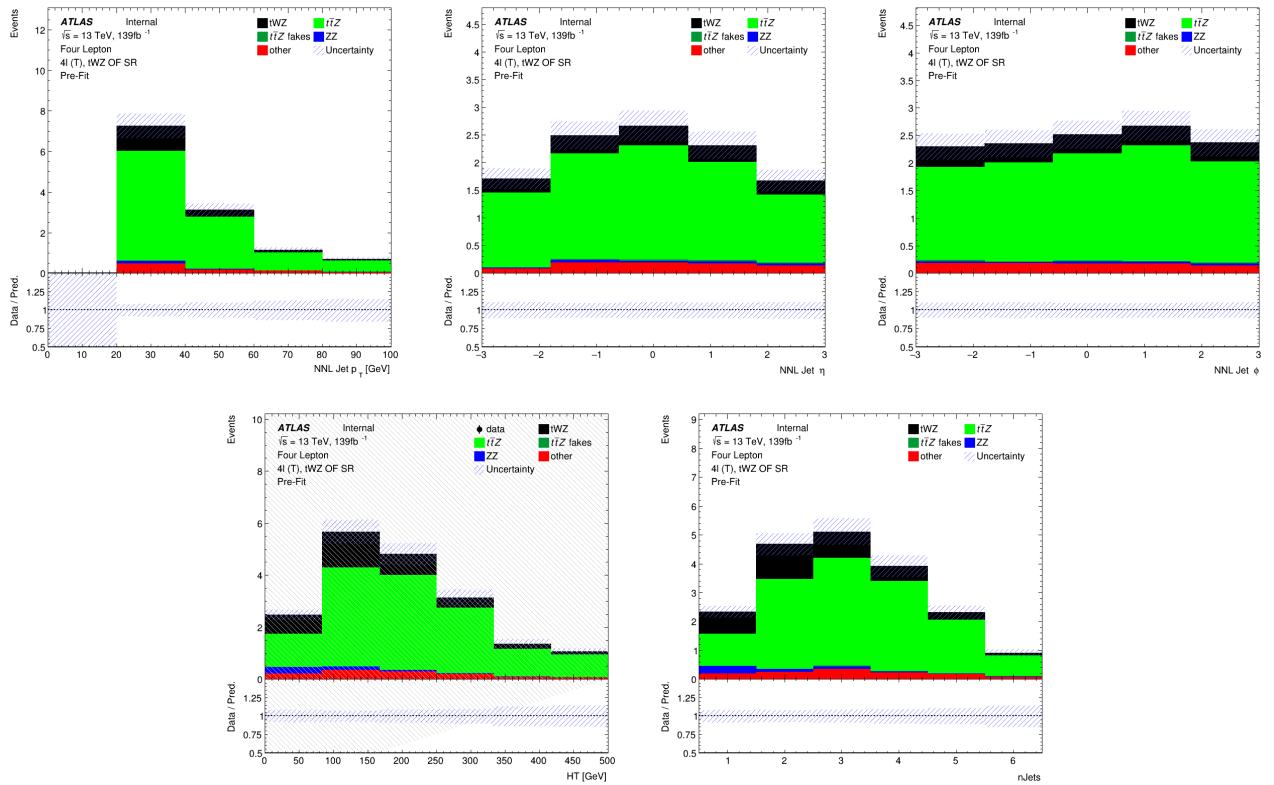


Figure 14: **Top row:** MC predictions for p_T , η and ϕ for next-to-next-to-leading (NNL) jets in the tWZ OF SR region is shown. **Bottom row:** MC predictions for H_T (scalar sum of Jet p_T) (left) and the Number of jets (right) in the tWZ OF SR region is shown.

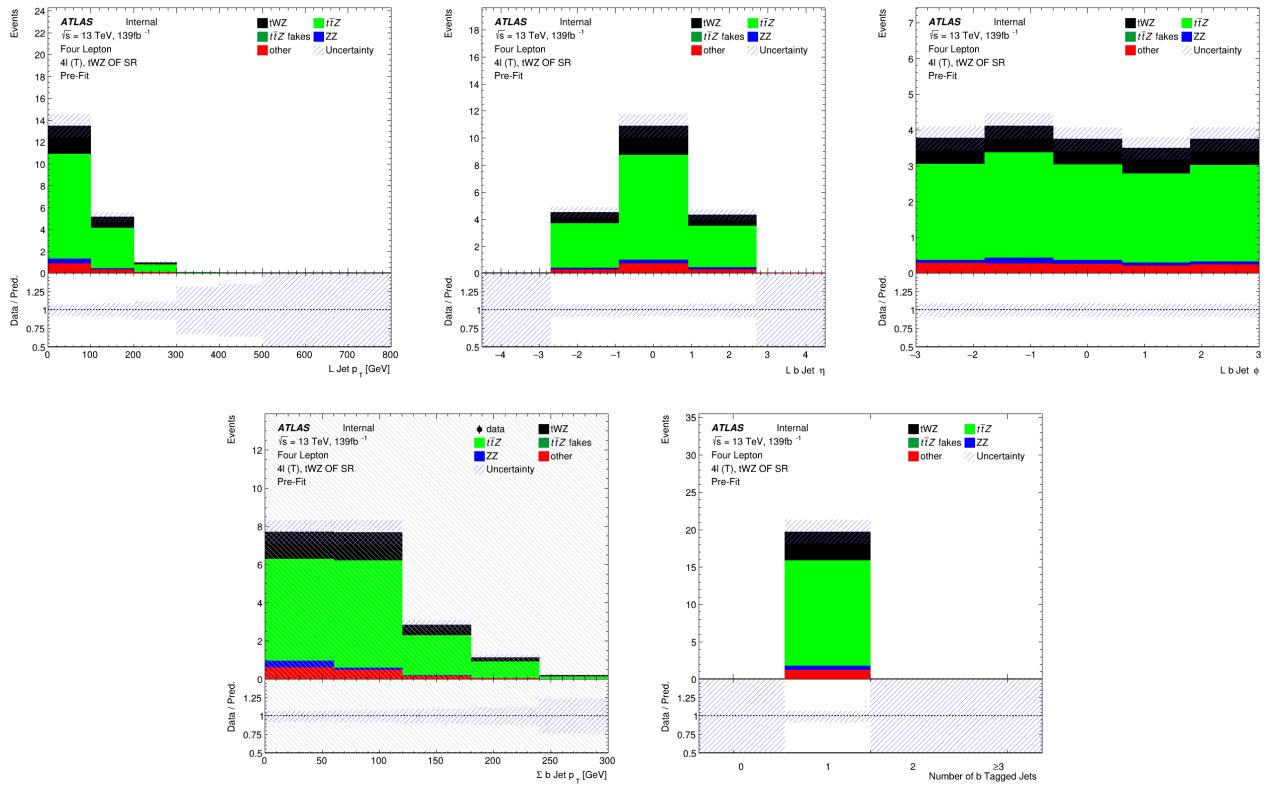


Figure 15: **Top row:** MC predictions for p_T , η and ϕ for leading b-tagged jets in the tWZ OF SR region is shown. **Bottom row:** MC predictions for the scalar sum of b-tagged jet p_T (left) and the Number of b-tagged jets (right) in the tWZ OF SR region is shown.

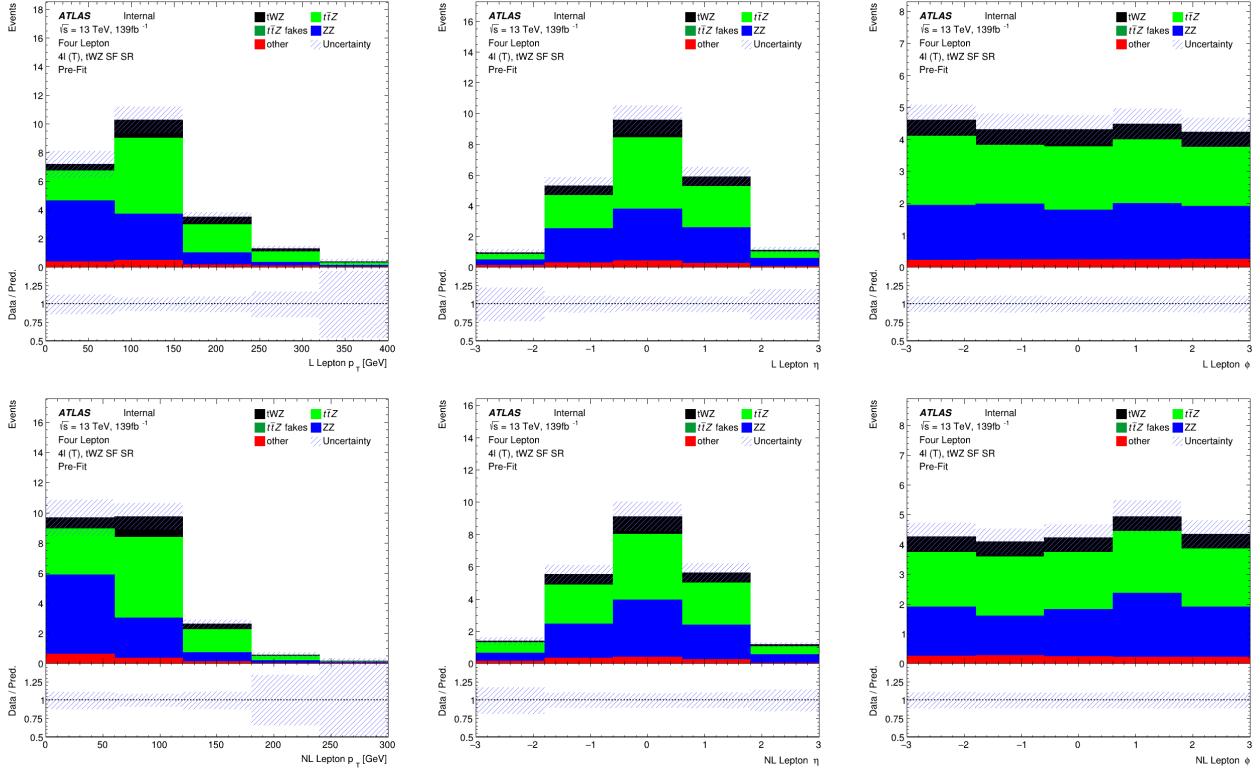


Figure 16: MC predictions for p_T , η and ϕ for leading (L) leptons (top row) and next-to-leading (NL) leptons (bottom row) in the tWZ SF SR region is shown.

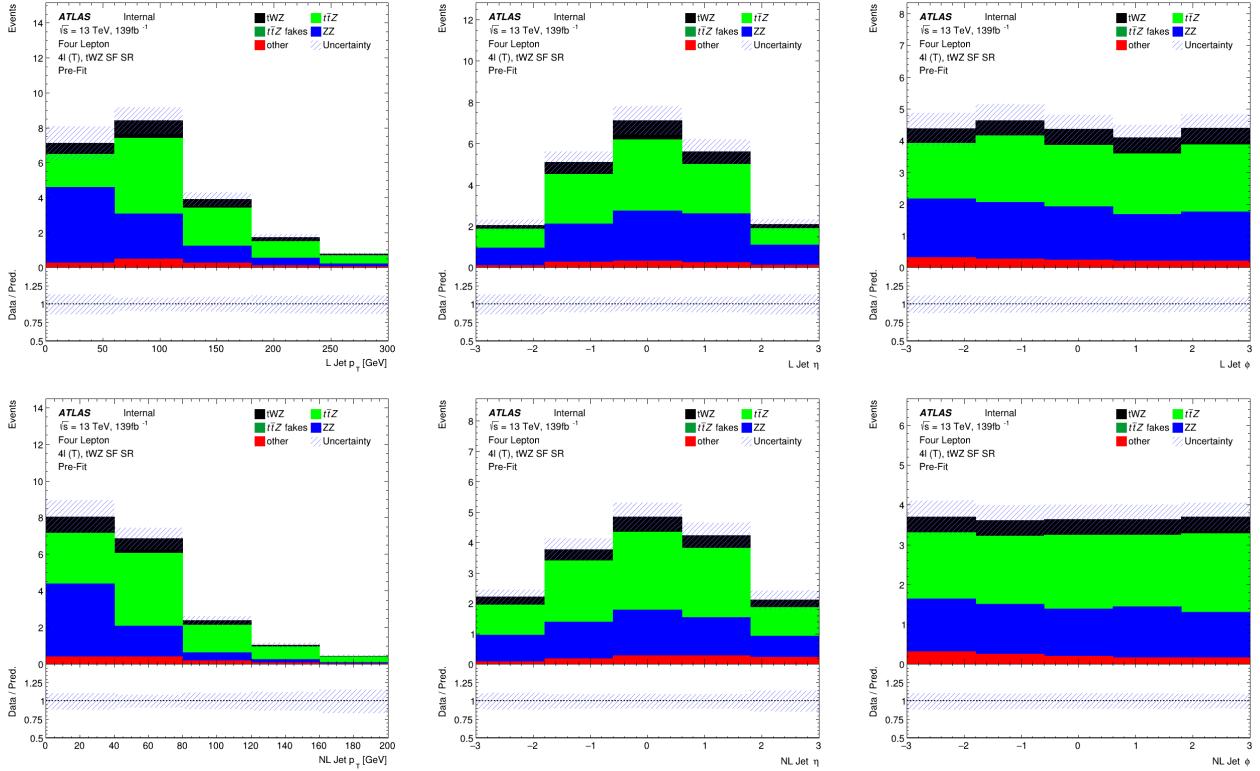


Figure 17: MC predictions for p_T , η and ϕ for leading (L) jets (top row) and next-to-leading (NL) jets (bottom row) in the tWZ SF SR region is shown.

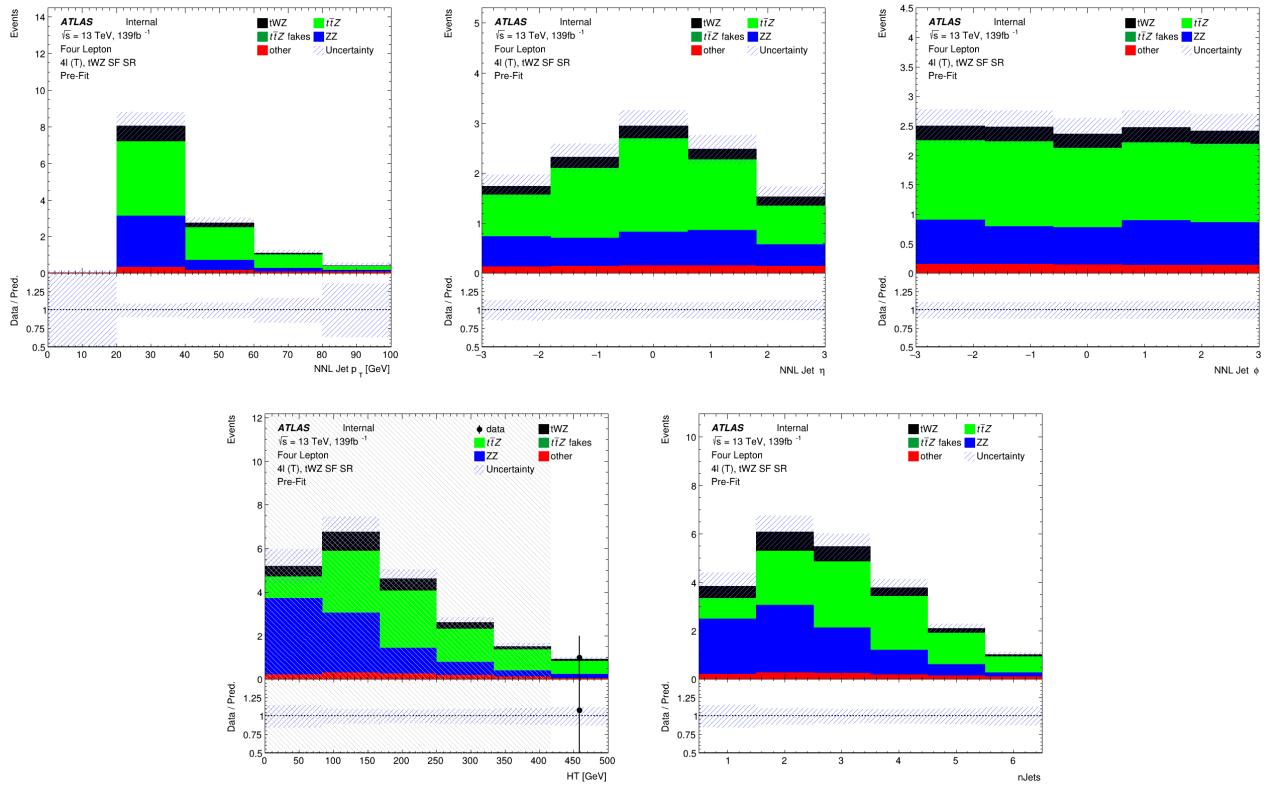


Figure 18: **Top row:** MC predictions for p_T , η and ϕ for next-to-next-to-leading (NNL) jets in the tWZ SF SR region is shown. **Bottom row:** MC predictions for H_T (scalar sum of Jet p_T) (left) and the Number of jets (right) in the tWZ SF SR region is shown.

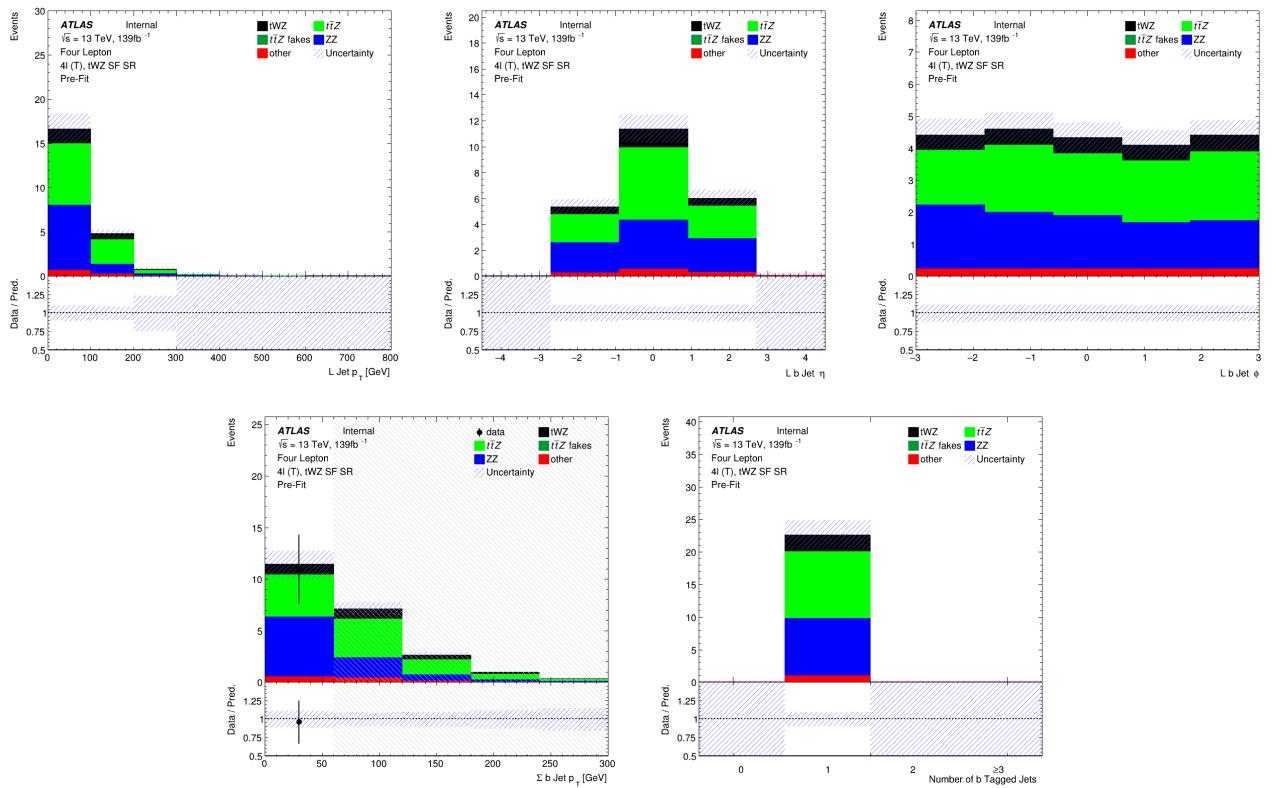


Figure 19: **Top row:** MC predictions for p_T , η and ϕ for leading b-tagged jets in the tWZ SF SR region is shown. **Bottom row:** MC predictions for the scalar sum of b-tagged jet p_T (left) and the Number of b-tagged jets (right) in the tWZ SF SR region is shown.

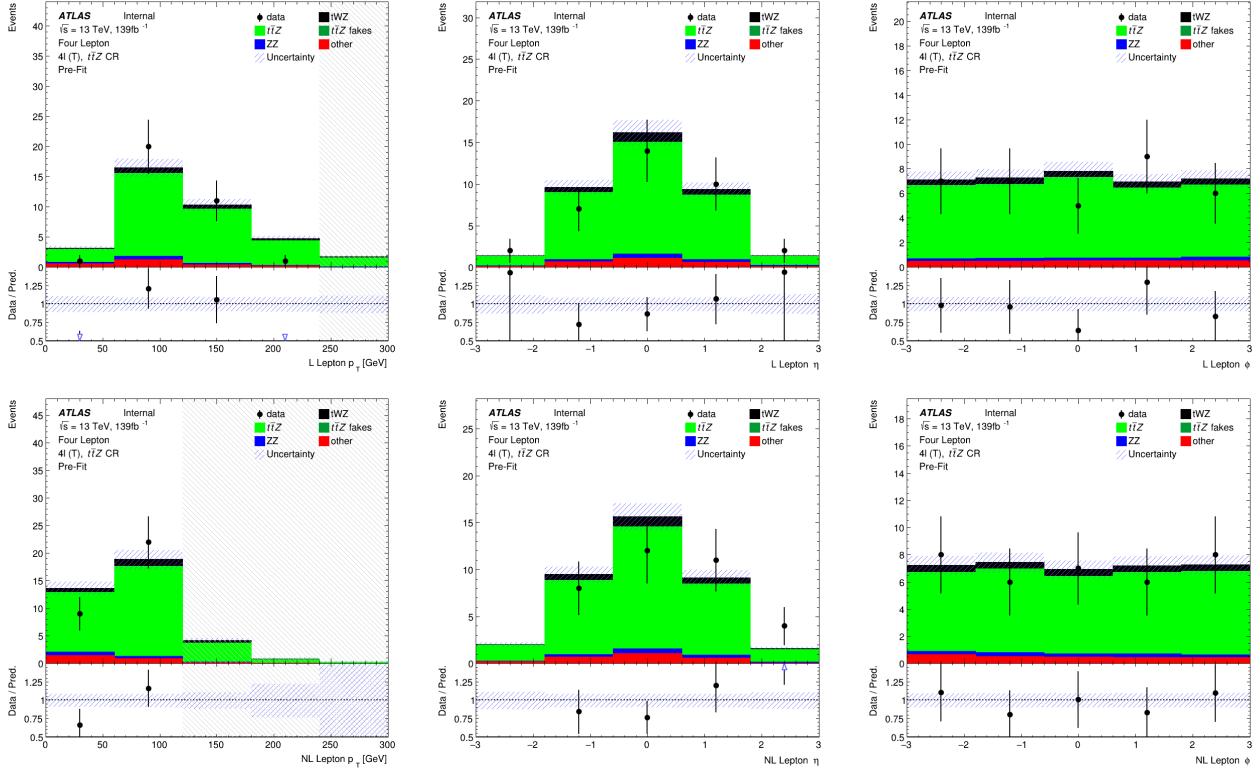


Figure 20: MC predictions for p_T , η and ϕ for leading (L) leptons (top row) and next-to-leading (NL) leptons (bottom row) in the $t\bar{t}Z$ CR region is shown.

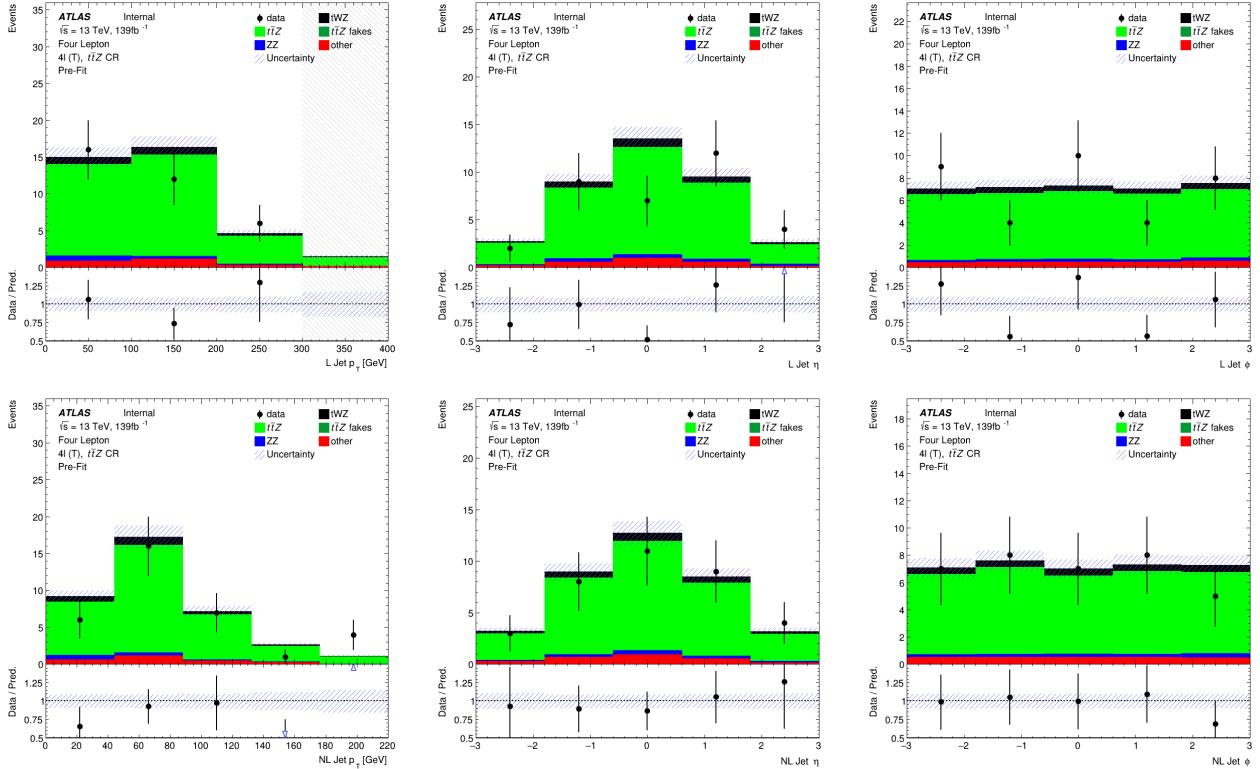


Figure 21: MC predictions for p_T , η and ϕ for leading (L) jets (top row) and next-to-leading (NL) jets (bottom row) in the $t\bar{t}Z$ CR region is shown.

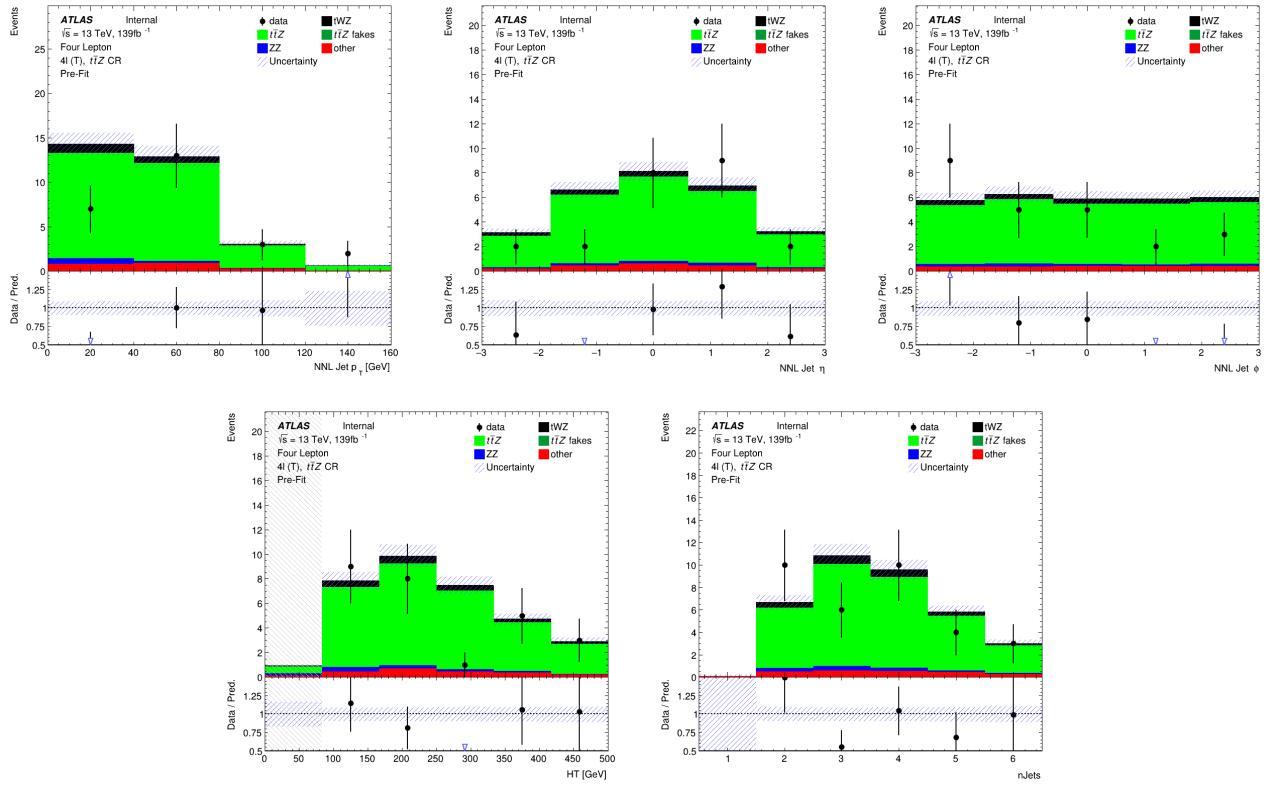


Figure 22: **Top row:** MC predictions for p_T , η and ϕ for next-to-next-to-leading (NNL) jets in the $t\bar{t}Z$ CR region is shown. **Bottom row:** MC predictions for H_T (scalar sum of Jet p_T) (left) and the Number of jets (right) in the $t\bar{t}Z$ CR region is shown.

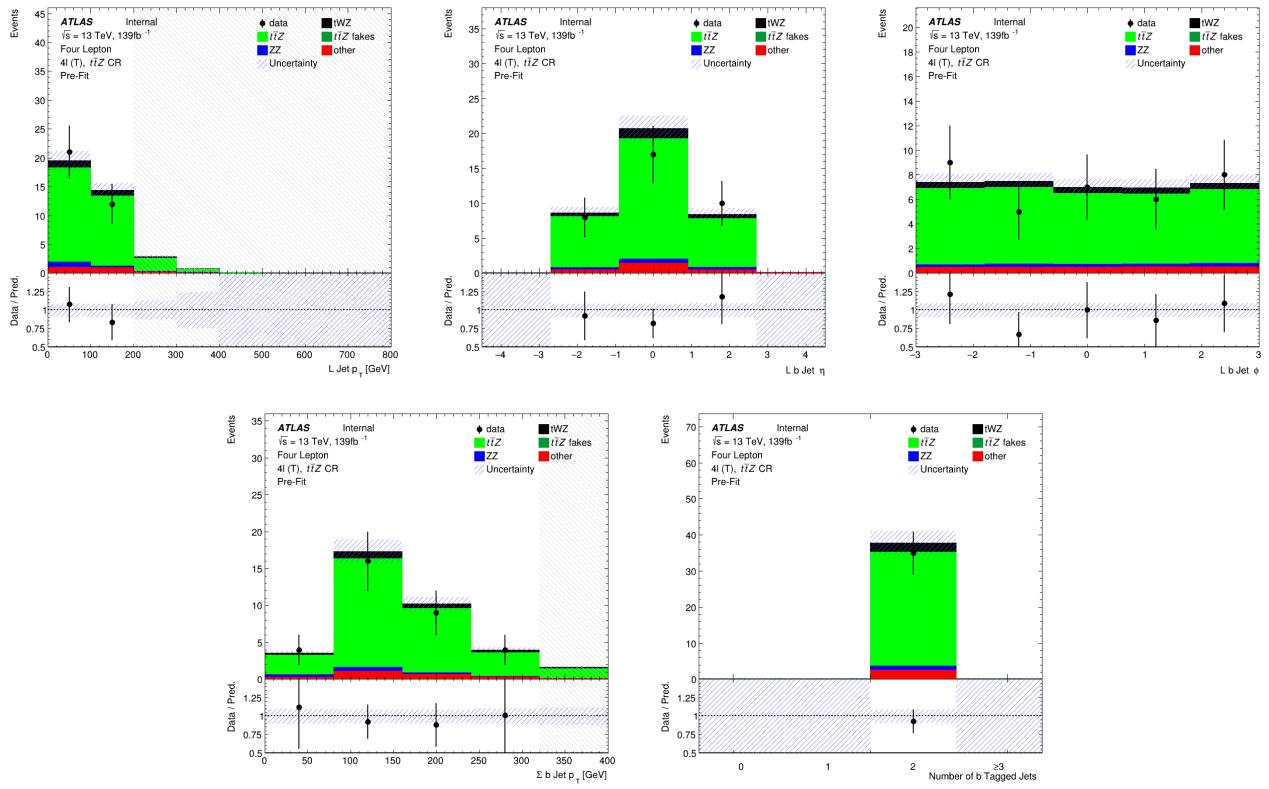


Figure 23: Top row: MC predictions for p_T , η and ϕ for leading b-tagged jets in the $t\bar{t}Z$ CR region is shown.
Bottom row: MC predictions for the scalar sum of b-tagged jet p_T (left) and the Number of b-tagged jets (right) in the $t\bar{t}Z$ CR region is shown.

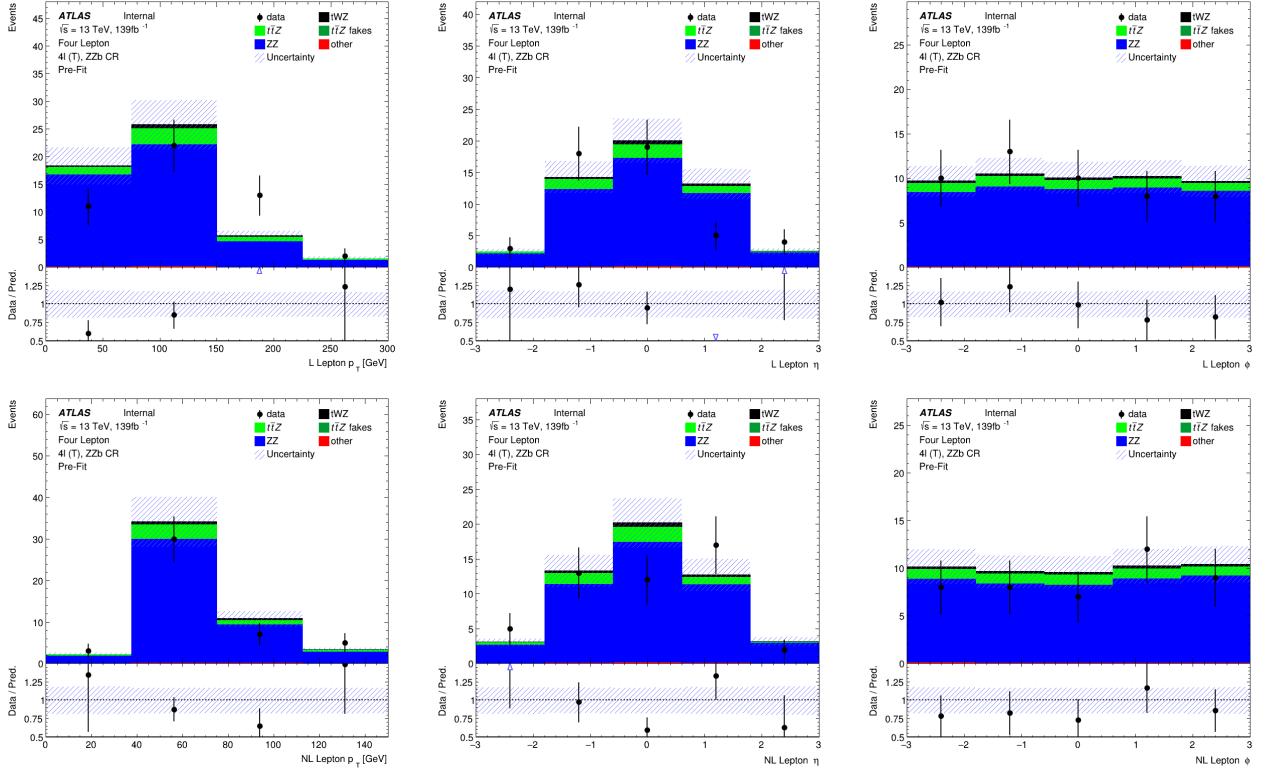


Figure 24: MC predictions for p_T , η and ϕ for leading (L) leptons (top row) and next-to-leading (NL) leptons (bottom row) in the ZZb CR region is shown.

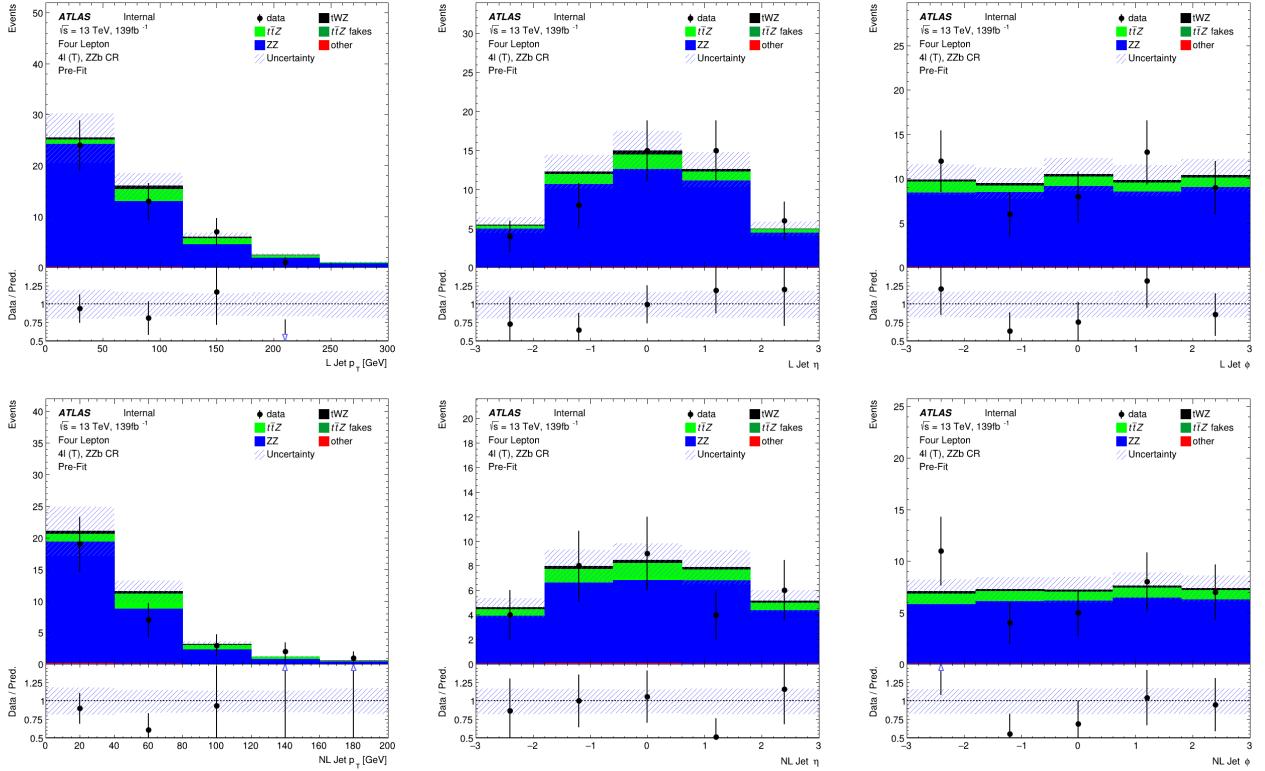


Figure 25: MC predictions for p_T , η and ϕ for leading (L) jets (top row) and next-to-leading (NL) jets (bottom row) in the ZZb CR region is shown.

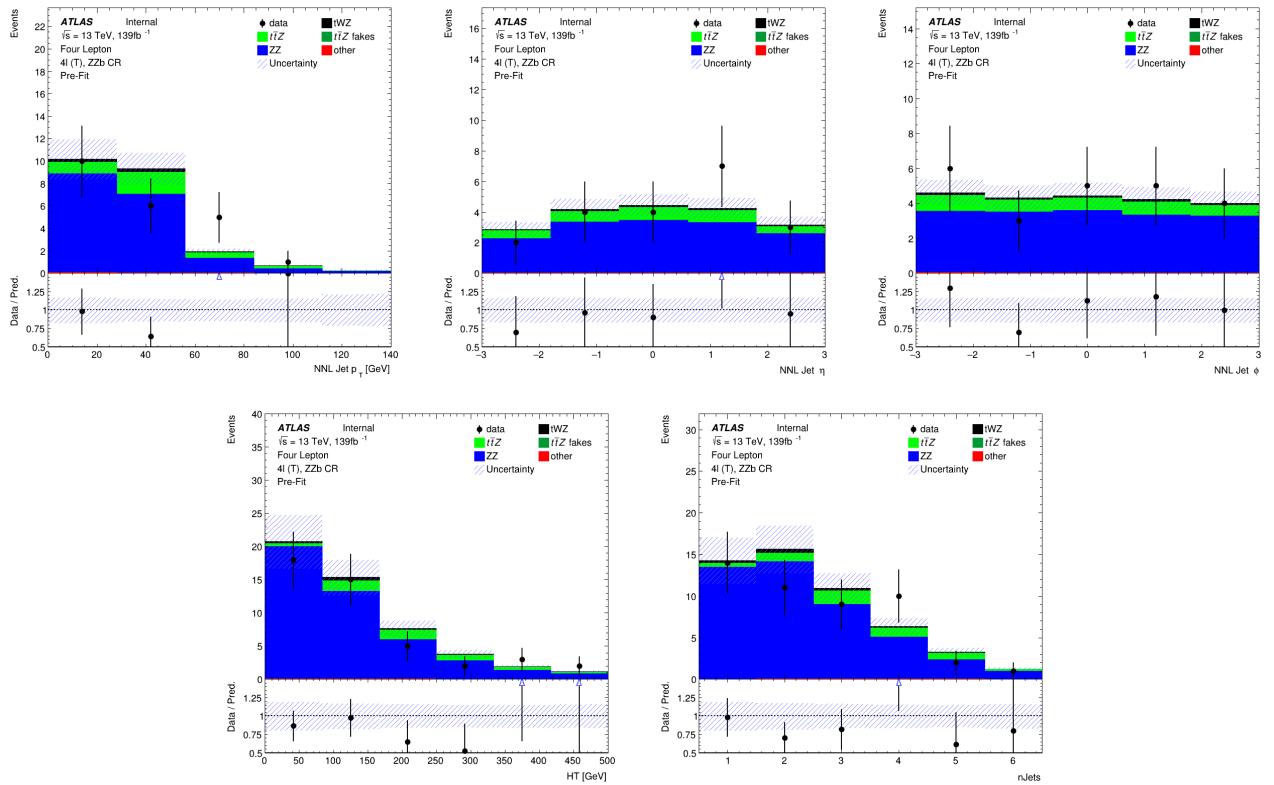


Figure 26: **Top row:** MC predictions for p_T , η and ϕ for next-to-next-to-leading (NNL) jets in the ZZb CR region is shown. **Bottom row:** MC predictions for H_T (scalar sum of Jet p_T) (left) and the Number of jets (right) in the ZZb CR region is shown.

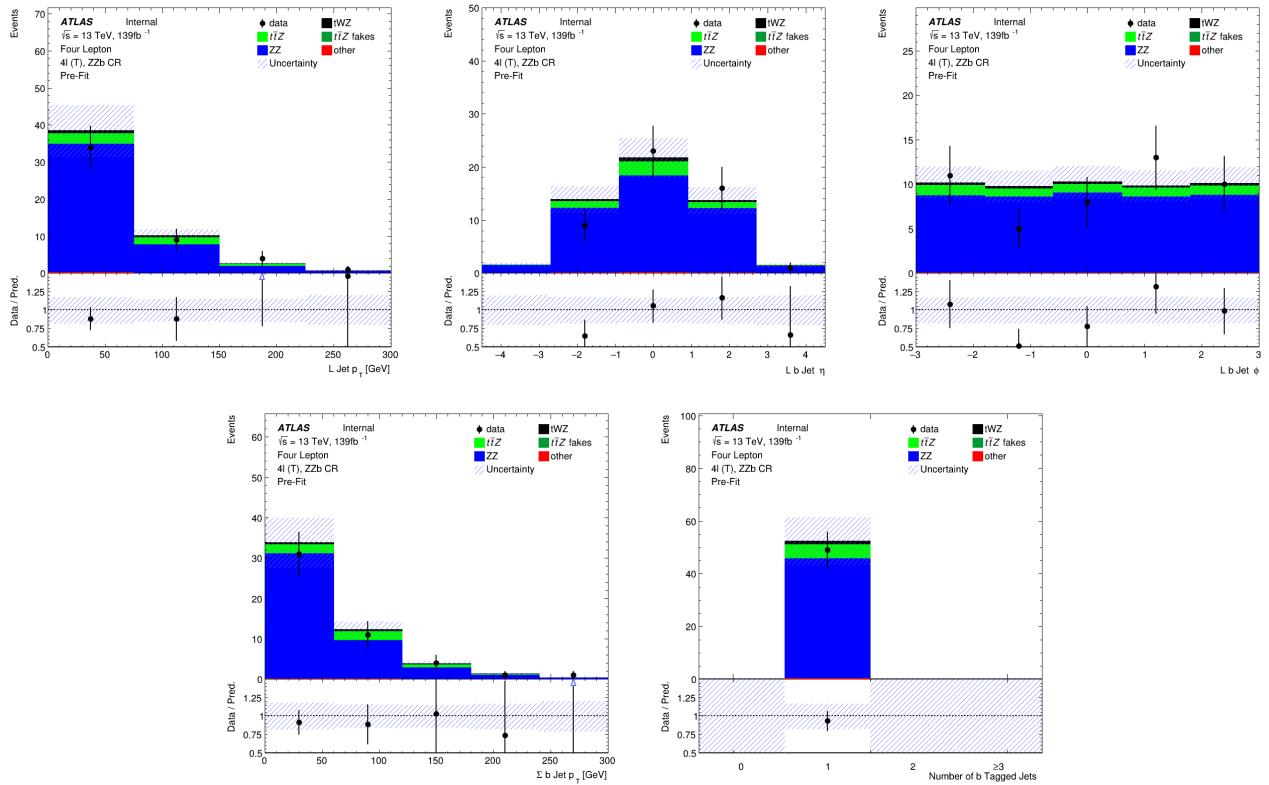


Figure 27: **Top row:** MC predictions for p_T , η and ϕ for leading b-tagged jets in the ZZb CR region is shown. **Bottom row:** MC predictions for the scalar sum of b-tagged jet p_T (left) and the Number of b-tagged jets (right) in the ZZb CR region is shown.

4.5.5 $(tWZ)_{\text{fake}}$ CR

In this section, expected number of events of variables in the $(tWZ)_{\text{fake}}$ CR are shown.

In Figure 28 MC predictions for p_T , η and ϕ for leading (L) leptons and next-to-leading (NL) leptons in the $(tWZ)_{\text{fake}}$ CR region is shown.

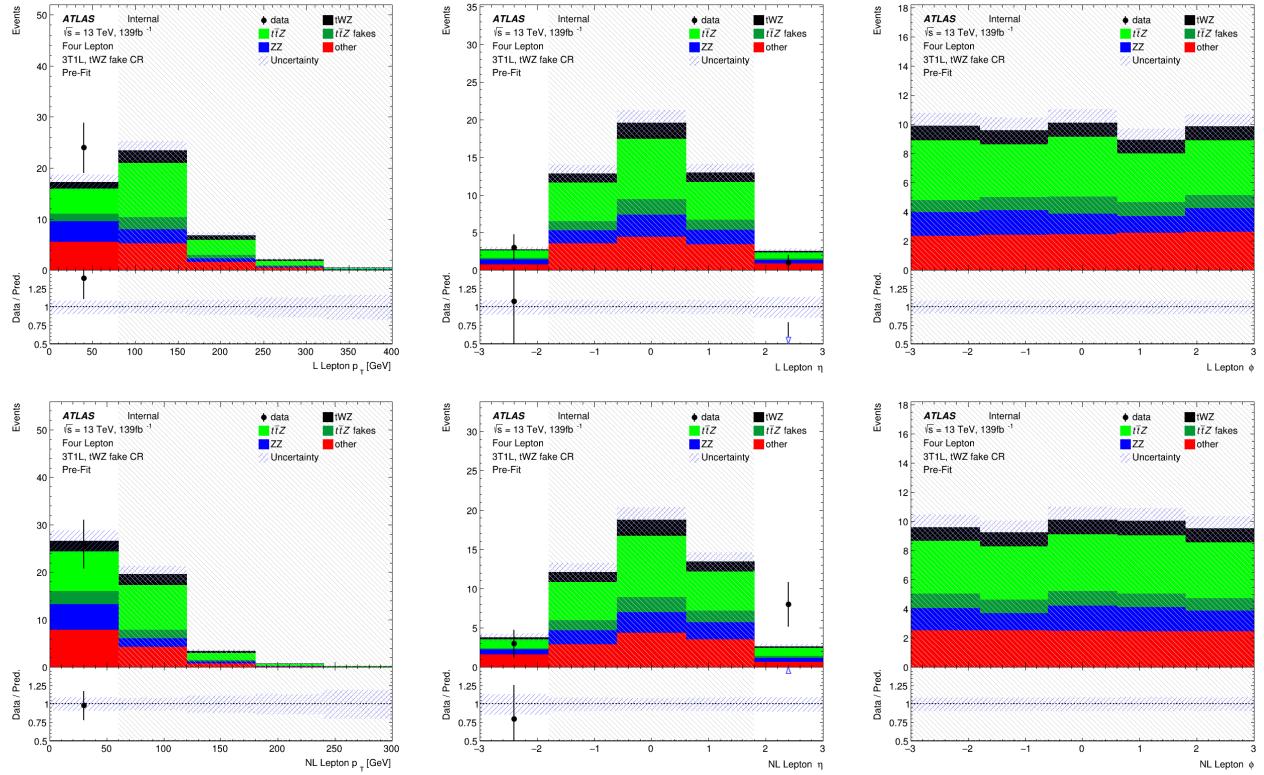


Figure 28: MC predictions for p_T , η and ϕ for leading (L) leptons (top row) and next-to-leading (NL) leptons (bottom row) in the $(tWZ)_{\text{fake}}$ CR region is shown.

- In Figure 29 MC predictions for p_T , η and ϕ for leading (L) jets and next-to-leading (NL) jets in the $(tWZ)_{\text{fake}}$ CR region is shown.
 In Figure 30 MC predictions for p_T , η and ϕ of the next-to-next-to-leading (NNL) jets, H_T (scalar sum of Jet p_T) and the Number of jets in the $(tWZ)_{\text{fake}}$ CR region is shown.
 In Figure 31 MC predictions for p_T , η and ϕ of the leading b-tagged jets, the scalar sum of b-tagged jet p_T and the Number of b-tagged jets in the $(tWZ)_{\text{fake}}$ CR region is shown.
 Despite the limited statistics present in the tetralepton channel, overall relatively good agreement is observed between data and simulation.

4.6 Fake Lepton Estimation

Fake leptons are objects reconstructed as leptons, but do not correspond to the leptons which originate from the process of interest. The sources of fake leptons include those originating from heavy hadron decays, light hadron decays or via the conversion of a photon to a lepton. In the ATLAS detector, the probability for a fake to occur is very low. In this section, the method used to estimate the fake lepton contribution is described.

As $t\bar{Z}$ is the dominant background process ($\sim 75\%$ of the total background contribution), it is assumed that $t\bar{Z}$ will also dominate the events containing fake leptons. The fake lepton efficiency, ϵ , can be written as $\epsilon = \frac{N_{\text{fake}}^{\text{tight}}}{N_{\text{fake}}^{\text{loose}}}$, where $N_{\text{fake}}^{\text{tight}}$ is the number of fake leptons which pass the tight lepton selection (See Section 4.2.1) and $N_{\text{fake}}^{\text{loose}}$ is the number of fake leptons which pass the loose lepton selection (See Section 4.2.1). The probability of one

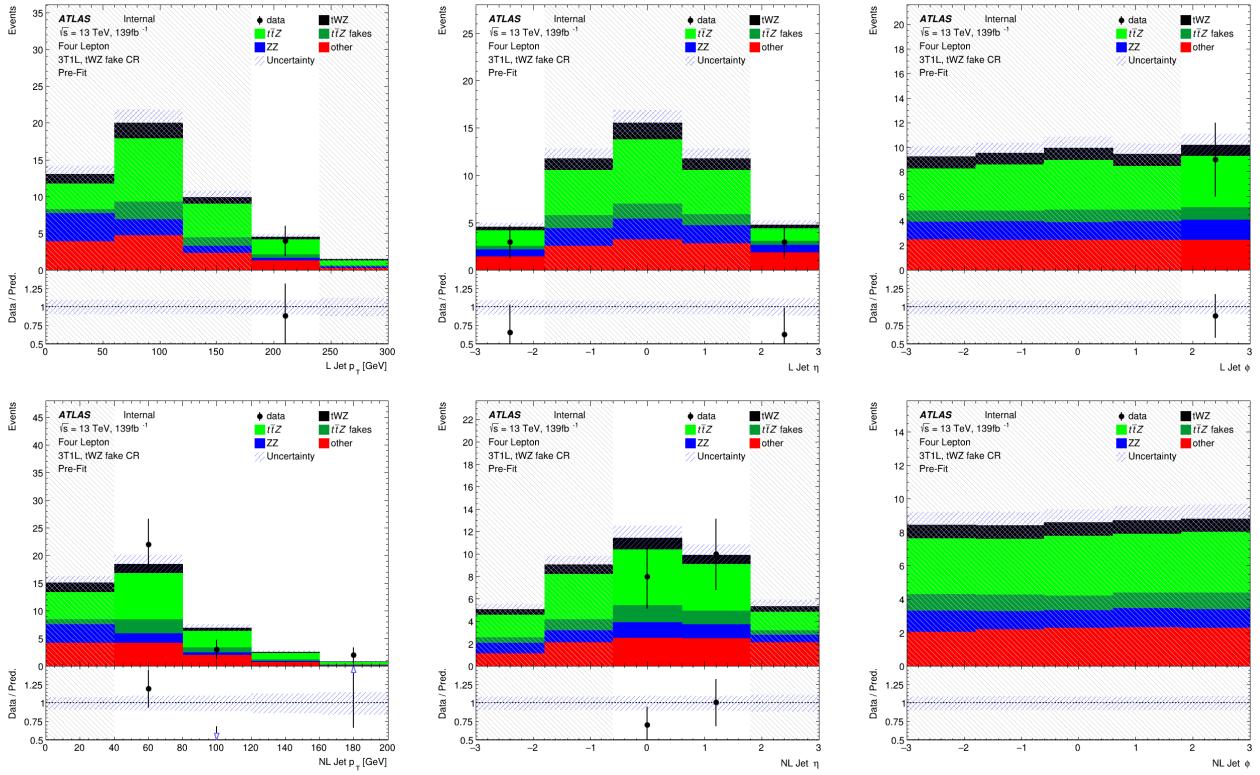


Figure 29: MC predictions for p_T , η and ϕ for leading (L) jets (top row) and next-to-leading (NL) jets (bottom row) in the (tWZ)_{fake} CR region is shown.

733 fake lepton to occur, $P(\text{one fake } \ell)$, is proportional to $\epsilon_1 \ll 1$ [lesage2017lepton, ATLAS-CONF-2012-048]
 734 and the probability for two fakes to occur is, $P(\text{two fakes } \ell)$, is proportional to $\epsilon_2 < \epsilon_1 \ll 1$. In this analysis, an
 735 estimation of the fake lepton component to the highest order is investigated and therefore the case where at least
 736 one fake lepton occurs in a $t\bar{t}Z$ event is considered.

737 Firstly, the dominant $t\bar{t}Z$ background is split up into $t\bar{t}Z$ and ($t\bar{t}Z$)_{fake} components. Secondly, a (tWZ)_{fake} CR
 738 (See Section 4.4) is defined which is enhanced in fakes and aims to constrain the ($t\bar{t}Z$)_{fake} background in the SR.
 739 All events which contribute to the ($t\bar{t}Z$)_{fake} background are determined by the IFF Truth Classifier [51]. The
 740 IFF Truth Classifier is a tool which aims to classify leptons based off their truth information. It uses the more
 741 general MCTruthClassifier [61] tool's output as input and returns one of the following lepton categories: Unknown,
 742 KnownUnknown (leptons which can (in principle) be classified, but the MCTruthClassifier fails to classify the
 743 lepton's truth type or origin), IsoElectron, ChargeFlipIsoElectron, PromptMuon, PromptPhotonConversion,
 744 ElectronFromMuon, TauDecay, BHadronDecay, CHadronDecay or LightFlavorDecay (More details [52]). Given
 745 these categories, leptons are considered as fake if they are classified as PromptPhotonConversion, BHadronDecay,
 746 CHadronDecay or LightFlavorDecay (i.e. a lepton originating from the decay of a b -Hadron, c -Hadron or
 747 light-flavour jet). Events which contribute to the ($t\bar{t}Z$)_{fake} background are those where at least one lepton from
 748 the $t\bar{t}Z$ sample are classified by the IFF Truth Classifier with one of the four aforementioned categories.

750 The (tWZ)_{fake} CR aims to be as similar as possible to the tWZ SRs, but enhanced in fakes. This CR can
 751 then be used to constrain the normalisation of the ($t\bar{t}Z$)_{fake} template. To ensure that this region is enhanced
 752 in fakes, it is required that it contains 3 tight leptons and 1 loose lepton, since loose leptons are more likely
 753 to be fakes. Leptons from heavy decays are produced in jets and are typically surrounded by other ener-
 754 getic particles. Since the loose lepton definition relaxes the isolation requirement, leptons satisfying the loose
 755 criteria are more enhanced in these fake leptons. By using the p_T of the loose lepton ($p_T(\text{Loose Lepton})$) in
 756 this region as the variable used in the fit, the shape (and normalisation) of the ($t\bar{t}Z$)_{fake} template can be constrained.

758
 759 In Figure 32, the number of leptons classified as fake, split up by their IFF Truth classification, in each region is

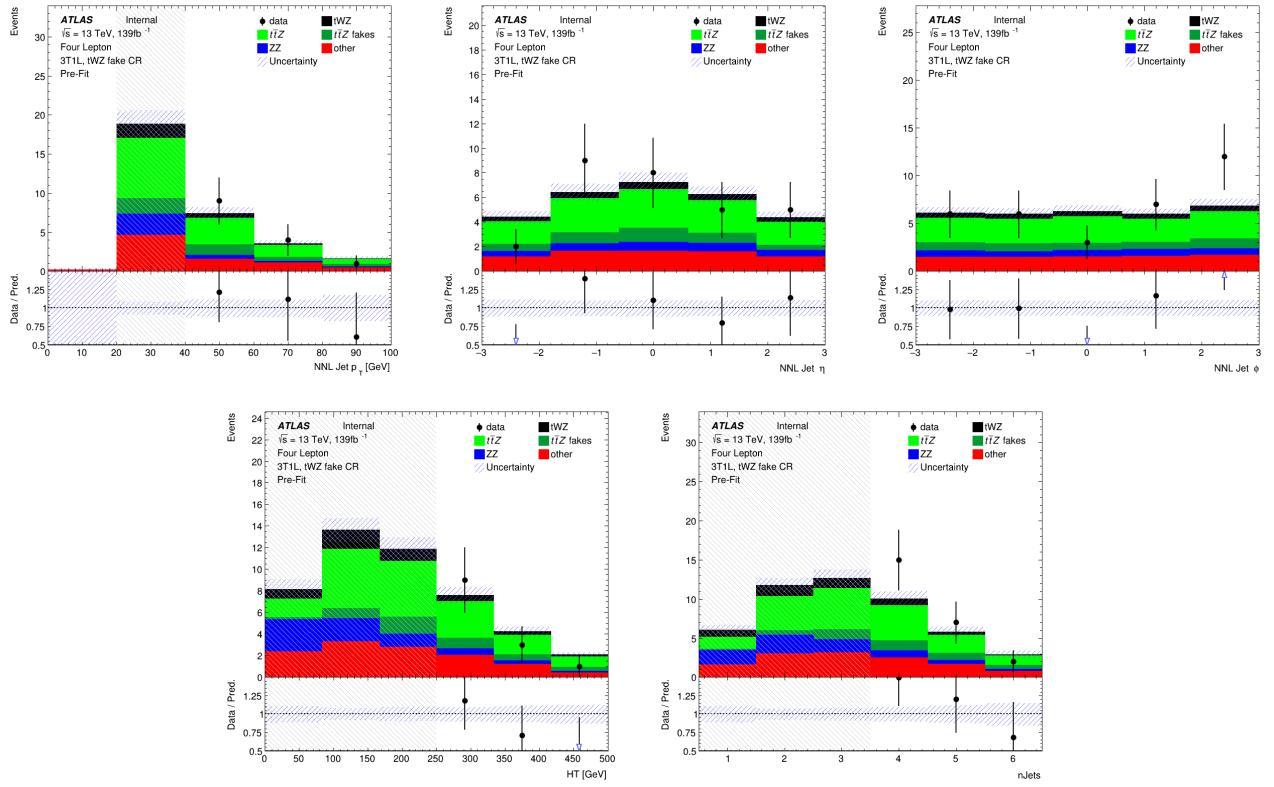


Figure 30: **Top row:** MC predictions for p_T , η and ϕ for next-to-next-to-leading (NNL) jets in the $(tWZ)_{\text{fake}}$ CR region is shown. **Bottom row:** MC predictions for H_T (scalar sum of Jet p_T) (left) and the Number of jets (right) in the $(tWZ)_{\text{fake}}$ CR region is shown.

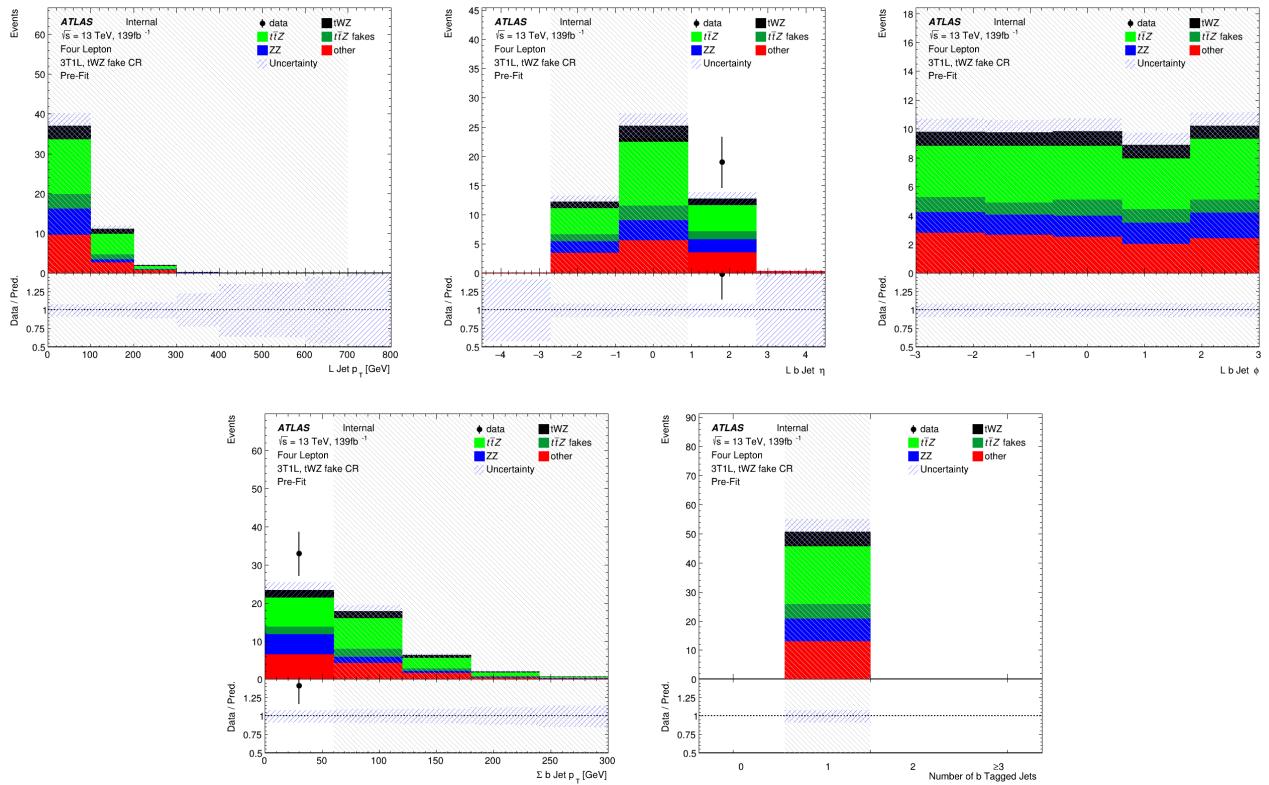


Figure 31: **Top row:** MC predictions for p_T , η and ϕ for leading b-tagged jets in the $(tWZ)_{\text{fake}}$ CR region is shown. **Bottom row:** MC predictions for the scalar sum of b-tagged jet p_T (left) and the Number of b-tagged jets (right) in the $(tWZ)_{\text{fake}}$ CR region is shown.

760 shown.

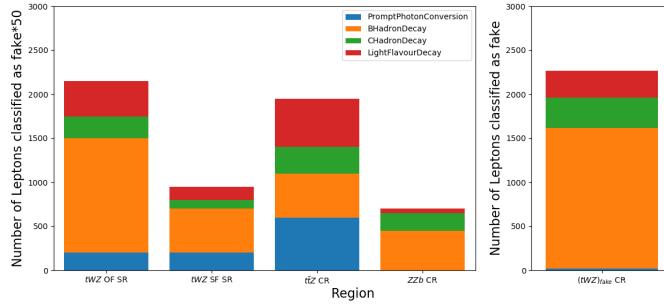


Figure 32: The number of leptons classified as fake, split up by their IFF Truth classification, in each region is shown. The left panel shows the number of leptons classified as fakes, scaled by a factor of 50, on the y-axis. The right panel shows the number of leptons classified as fakes (unscaled), on the y-axis. The different signal and control regions are shown on the x-axes of the left and right panels. The IFF truth classification of the leptons are shown in the legend and correspond to the different coloured stacked histograms.

761 Around 50 times more fake leptons pass our selection criteria for the $(tWZ)_{\text{fake}}$ CR, compared to remaining four
 762 regions. This relative abundance of fake leptons present in the $(tWZ)_{\text{fake}}$ CR further justifies our use of this region
 763 to constrain the fake lepton component. In Figure 33, the relative dominance of the different classifications for fake
 764 leptons (classified by the IFF truth classified) in each region, is shown.

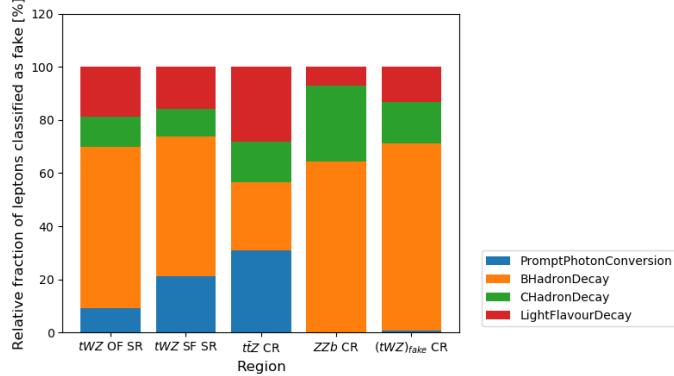


Figure 33: The relative dominance of the different classifications for fake leptons (classified by the IFF truth classified) in each region, is shown. The relative dominance of leptons classified as fakes, as a fraction of the total number of fake leptons (in each region), is shown on the y-axis. The different signal and control regions are shown on the x-axis. The IFF truth classification of the leptons are shown in the legend and correspond to the different coloured stacked histograms.

765 The majority of fake leptons which pass our selection criteria originate from the decay of b -hadrons, in all regions
 766 but the ttZ CR. The smaller proportion of fake leptons originating from b -hadron decays in the ttZ CR could
 767 possibly be due to statistical fluctuations resulting from the low number of fake leptons which pass our selection
 768 criteria in this region (~ 40 fake leptons). In Figure 34, the amount of fake and real $t\bar{t}Z$ events which pass our
 769 selection criteria, in each region, is shown.

770 Around 20% of all $t\bar{t}Z$ events are classified as fake events (having one or more of its leptons being classified as fake)
 771 in the $(tWZ)_{\text{fake}}$ CR. The tWZ OF SR, tWZ SF SR, ttZ CR and ZZb CR have less than 1% of their total $t\bar{t}Z$
 772 events being fake. The non-negligible amount of fake $t\bar{t}Z$ events present in the $(tWZ)_{\text{fake}}$ CR, allows the $t\bar{t}Z$ fake
 773 background to be constrained by the $(tWZ)_{\text{fake}}$ CR.

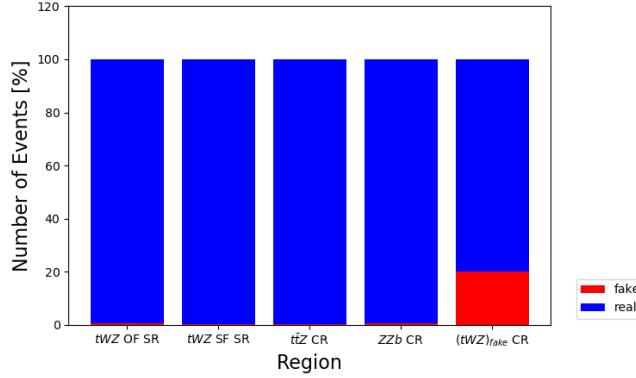


Figure 34: The percentage of fake and real $t\bar{t}Z$ events which pass our selection criteria, in each region, is shown. The relative number of fake and real events (in % of the total number of events in the nominal and fake $t\bar{t}Z$ background samples) is shown on the y-axis. The different signal and control regions are shown on the x-axis. The blue and red histograms represent the percentage of real and fake events (out of the total number of events in the nominal and fake $t\bar{t}Z$ background samples), respectively.

774 4.7 Machine Learning Techniques

775 The presence of different numbers of top quarks is a key discriminator between signal and the dominant background
 776 process, $t\bar{t}Z$. This information is aimed to be exploited by reconstructing ℓb systems as a proxy for top quarks (since,
 777 $t \rightarrow W(\rightarrow \ell\nu)b$). This is done in two ways, firstly by implementing a Boosted Decision Tree (BDT) which is used
 778 to classify ℓb systems originating from top quarks, and secondly, by implementation of a kinematic reconstruction
 779 algorithm (Two Neutrino Scanning Method) which aims to determine the likelihood of an event containing two top
 780 quarks. In this thesis, this BDT is referred to as an *object-level* classifier.

781 Certain variables constructed from event information show discrimination between signal and background events.
 782 This information can be exploited to discriminate between signal and background events by constructing an
 783 algorithm which takes advantage of these discriminating variables. A BDT is implemented and is used to
 784 discriminate between tWZ events and its major backgrounds, $t\bar{t}Z$ and ZZ . In this thesis, this BDT is referred to
 785 as an *event-level* classifier. The discriminator output from the object-level BDT can be converted to a variable
 786 which can then be used as an input to the event-level BDT.

787 BDTs are chosen, as opposed to another ML algorithm, since they are not very sensitive to overtraining
 788 and perform well with minimal/no optimisation or tweaking of the hyper parameters. A multi-layered se-
 789 quential neutral network was tried, however, it was out-performed by a BDT. More specifically, Scikit-Learn's
 790 `GradientBoostingClassifier` [71] was used.

792 4.7.1 Object-level BDT

793 The object-level BDT was trained on a $t\bar{t}$ sample (simulated using the same generator, parton shower and to the
 794 same order of QCD as the $t\bar{t}$ sample described in Section 4.1.2) with a baseline selection of exactly 1 tight lepton
 795 with $p_T > 28$ GeV. Additionally, jets in this sample are required to have $p_T > 20$ GeV. Jets are identified as
 796 b -tagged jets by the 77% DL1r working point. These baseline selections were chosen to mimic those used in the
 797 event selection of the analysis (outlined in Table 5). The leptons and b -jets used in training the object-level BDT
 798 are required to pass the aforementioned baseline selections. This $t\bar{t}$ sample was utilised in training the BDT to
 799 avoid using a subset of events from the MC samples used in the rest of the analysis, therefore maximizing the
 800 amount of generated events available to use in other parts of the analysis.

801
 802 The signal class is defined to consist of reconstructed ℓb systems (defined as the sum of the 4-vectors of a lepton
 803 and a b -tagged jet) originating from top quarks which are well matched to their truth counterparts. All possible
 804 combinations of ℓ and b -tagged jets are selected from the events. In particular, we require that ΔR between the
 805 reconstructed and truth ℓb system is less than 0.05. We additionally require that the reconstructed lepton and the
 806 truth top have charges with the same sign (since $t \rightarrow b\ell^+\bar{\nu}_\ell$ and $\bar{t} \rightarrow b\ell^-\nu_\ell$). The background class is defined to

807 consist of all reconstructed ℓb systems which fail to pass the criteria for ℓb systems which are labelled as signal.
 808 These definitions for the signal and background classes ensure that the signal class consists of mostly ℓb systems
 809 originating from top quarks and the background class consists of mostly ℓb systems which do not originate from
 810 top quarks.

811 Different observables corresponding to an ℓb system were used as input to training. The optimum values
 812 for the hyper-parameters used were determined by training the BDT with a range of different values for the
 813 hyper-parameters and choosing the set of values which maximized the mean accuracy (based off 5 fold kfold
 814 cross-validation). This method is more commonly referred to as hyper-parameter optimisation or tuning. After
 815 hyper-parameter optimisation, the mean accuracy of each fold increased from 0.76 to 0.77 ($\sim 1\%$ increase).
 816 Variables can be assigned a score called *variable importance*, based on their usefulness on predicting a target
 817 variable (in this case, a signal or background event). The variable importance for any given variable was obtained
 818 by computing the mean accuracy of the model, removing the variable from training, retraining the model and
 819 computing the mean accuracy of this new model. The difference between mean accuracies of the unaltered model
 820 and the retrained model (after removal of the variable) gives us the variable importance of the variable of interest.
 821 This method returns positive values for variables which increase the mean accuracy of the model and negative
 822 values for variables which decrease the mean accuracy of the model. Variables with negative variable importances
 823 were completely removed from training.

825 In Table 8, the variables used in training the object-level BDT are shown.

Observable	Description
$m(\ell b)$	Invariant mass of the ℓb system
$p_T(\ell b)$	p_T of the ℓb system
$\Delta\eta(\ell, b)$	$\Delta\eta$ between the ℓ and b -tagged jet
$\Delta\phi(\ell, b)$	$\Delta\phi$ between the ℓ and b -tagged jet
$\Delta R(\ell, b)$	ΔR between the ℓ and b -tagged jet

826 Table 8: A list of the observables used in the object-level BDT, ordered by variable importance (descending, top to bottom) is shown.

827 In Figure 35, normalised distributions of the variables used in the object-level BDT, for the signal and background
 828 classes are shown.

829 Overall the BDT input variables show a large amount of discrimination.

830 We can check the modelling of the input variables to the object-level BDT by studying the agreement between data
 831 and simulation in the $t\bar{t}Z$ CR. In Figure 36, MC predictions for the input variables to the object-level BDT in the
 832 $t\bar{t}Z$ CR are shown.

833 Overall, there is good agreement between data and simulation for the variables used in the object-level BDT, in
 834 the $t\bar{t}Z$ CR. This suggests that the variables used in the object-level BDT are well-modelled and are reasonable to
 835 include as inputs to the object-level BDT.

836 A final check can be done to study the similarity of the ℓb systems present in the $t\bar{t}$ sample which we use for training
 837 the object-level BDT, and the ℓb systems which we aim to identify using the object-level BDT. More specifically, we
 838 want to ensure that the modelling of the ℓb systems in the $t\bar{t}$ sample are sufficiently similar to those in the tWZ and
 839 $t\bar{t}Z$ samples (see Table 3). This is done to understand how well the BDT (trained on ℓb systems in the $t\bar{t}$ sample)
 840 generalises to classifying ℓb systems in the analysis (tWZ and $t\bar{t}Z$ samples). In Figure 37, normalised distributions
 841 of the variables used in the object-level BDT for the $t\bar{t}$, tWZ and $t\bar{t}Z$ samples, are shown.

842 The $m(\ell b)$, $p_T(\ell b)$ and $\Delta\eta(\ell, b)$ distributions show little discrimination between the samples. However, the $\Delta\phi(\ell, b)$
 843 and $\Delta R(\ell, b)$ distributions show a large amount of discrimination between the three samples. This suggests that
 844 the use of the $t\bar{t}$ sample in training the object-level BDT may be sub-optimal. However, it still represents the best
 845 option available, since our other options involve utilising of a subset of generated events used in the other parts of
 846 the analysis. This would result in a smaller number of generated events used in the background prediction, leading
 847 to larger statistical uncertainties, in an already statistically limited analysis.

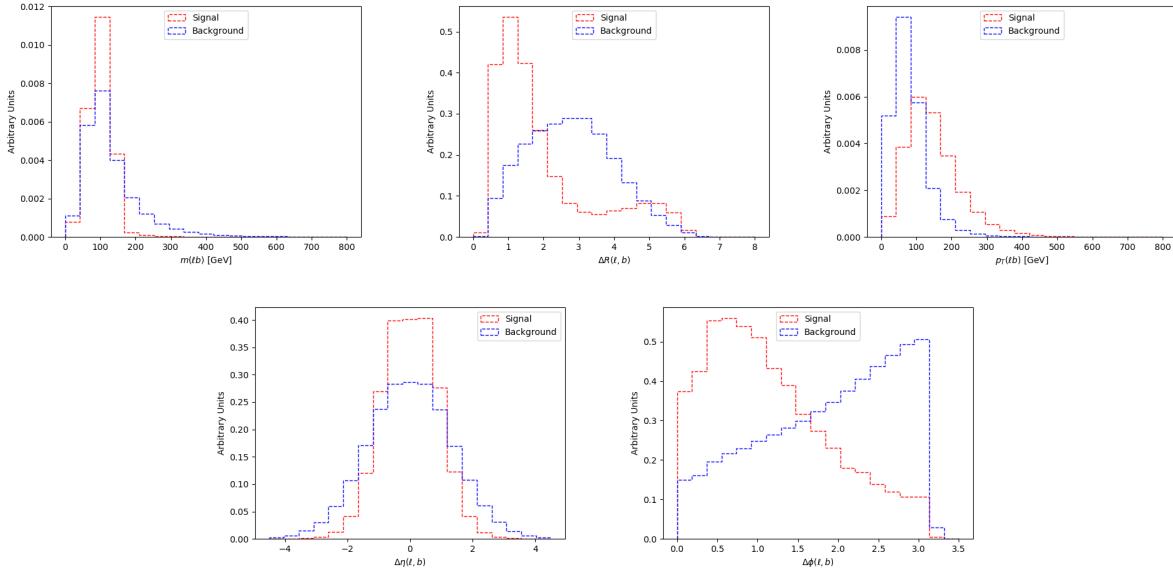


Figure 35: Normalised distributions of the variables used in the object-level BDT (ordered from top left to bottom right via decreasing variable importance), for the signal and background classes are shown. **From top left to bottom right:** Invariant mass of the ℓb system. ΔR between the ℓ and b -tagged jet. The p_T of the lb system. $\Delta\eta$ between the ℓ and b -tagged jet. $\Delta\phi$ between the ℓ and b -tagged jet. The red and blue dotted lined histograms represent the signal and background classes events (from the training set), respectively. These histograms are normalised to an area of 1. The variable used in training is shown on the x-axis. The y-axis shows the relative number of events for the signal and background classes (in arbitrary units).

850

In Table 9, the hyper-parameters used in the object-level BDT is shown.

Hyper-parameter	Value	Description
loss	deviance	The loss function to be optimised
criterion	friedman_mse	The function used to measure the quality of a split
n_estimators	200	The number of boosting stages to perform
learning_rate	0.1	The step size at each iteration during optimisation
max_depth	6	The maximum depth of the individual regression estimators
min_samples_split	2	The minimum number of samples (events) required to split an internal node
min_samples_leaf	1	The minimum number of samples (events) required to be at a leaf node
validation_fraction	0.1	The proportion of training data to set aside as validation set for early stopping
n_iter_no_change	20	Training terminates when the validation score (determined by the validation set) does not improve in all of the previous

Table 9: A list of the hyper-parameters used in the object-level BDT is shown. Hyperparameters not listed in this table use the default values as stated in the Scikit-learn Documentation[72].

851

852 The number of events used in training for the signal and background classes were 49871 and 384152 respectively.
 853 Imbalanced datasets can cause ML classifiers to ignore small classes while concentrating on classifying large classes
 854 more accurately, which may result in the trained classifier performing sub-optimally. In order to correct this dataset
 855 imbalance, we ensure that the relative weighting of each event is such that the sum of the signal weights is equal to
 856 the sum of the background weights. In Figure 38 the normalised histograms of the training and test sets (extracted
 857 from fold 5 from a 5 fold kfold cross validation) for signal and background is shown.

858

We can see that the shapes of the training and test sets for both signal and background are very similar. This
 859 is a good indicator that no over-training occurred. Another over-training check is performed using 5 fold kfold
 860 cross validation. We ensure that the variance of the mean accuracy of each folds' test set in cross validation is
 861 substantially small. This indicates that fluctuations in features from different training sets are not learnt by the
 862 classifier. For the object-level classifier, a variance of 3.24×10^{-7} was calculated for the mean accuracies of each
 863 folds' test set in cross validation. This tells us that the BDT does not perform better on one subset of a dataset
 864 over another and it is therefore not prone to learning statistical fluctuations of a subset, which would result in a
 865 BDT which does not generalise well to unseen datasets. This small variance therefore provides further evidence

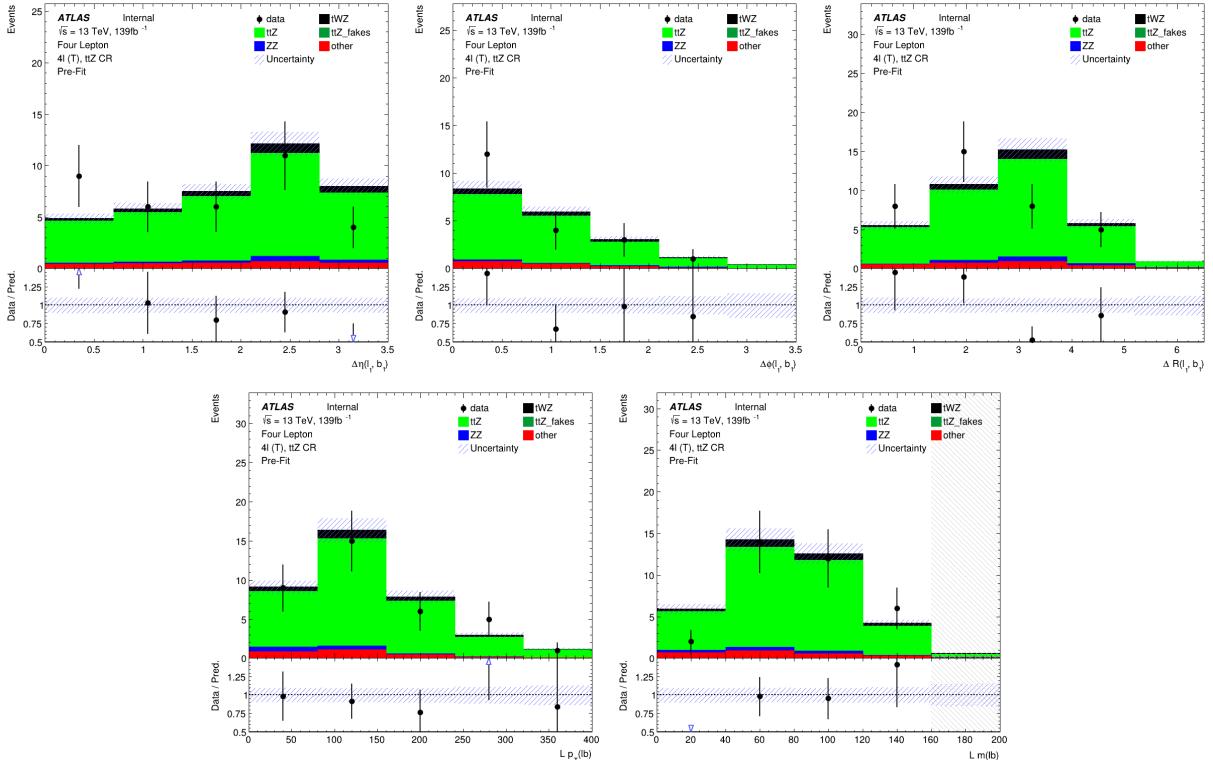


Figure 36: Pre-fit distributions of variables used as input to the object-level BDT (ordered from top left to bottom right via decreasing variable importance), in the $t\bar{t}Z$ CR, are shown. **From top left to bottom right:** $\Delta\eta$ between the lepton and b -jet of the leading ℓb system. $\Delta\phi$ between the lepton and b -jet of the leading ℓb system. ΔR between the lepton and b -jet of the leading ℓb system. p_T of the leading ℓb system. Mass of the leading ℓb system. The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the blue diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. Bins with $\frac{\text{signal}}{\text{background}}$ greater than 0.1 are kept blinded. Blinded bins are shaded with black diagonal lines and their data points are omitted.

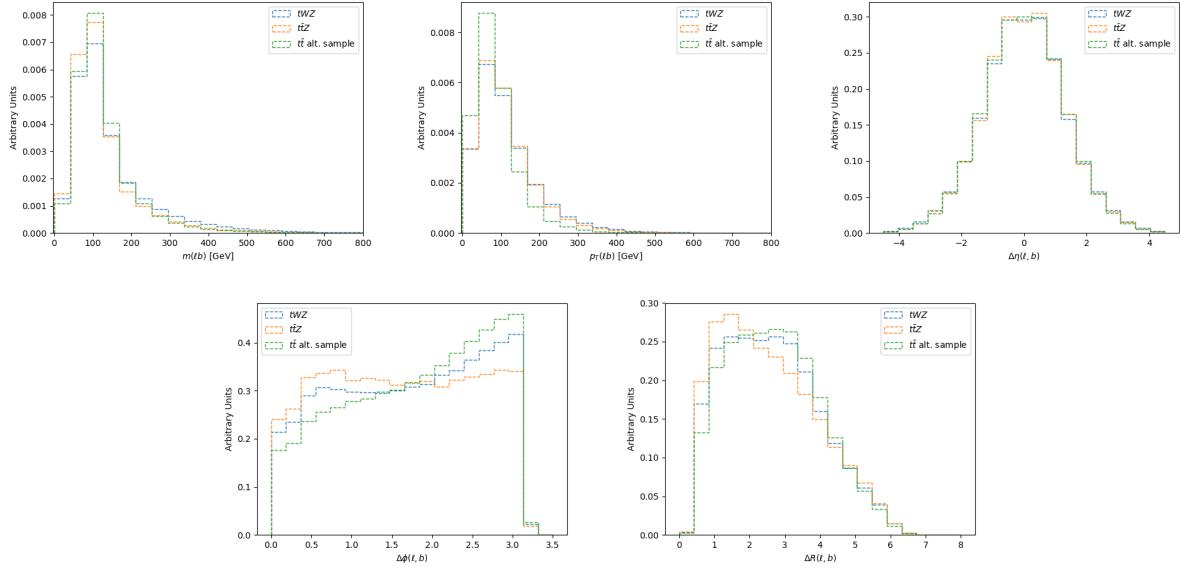


Figure 37: Normalised distributions of the variables (ordered from top left to bottom right via decreasing variable importance) used in the object-level BDT for the $t\bar{t}$, tWZ and ttZ samples, are shown. **From top left to bottom right:** Invariant mass of the ℓb system. ΔR between the ℓ and b -tagged jet. The p_T of the ℓb system. $\Delta\eta$ between the ℓ and b -tagged jet. $\Delta\phi$ between the ℓ and b -tagged jet. The green, blue and orange dotted lined histograms represent events from the $t\bar{t}$, tWZ and ttZ samples, respectively. These histograms are normalised to an area of 1. The variable used in training is shown on the x-axis. The y-axis shows the relative number of events (in arbitrary units).

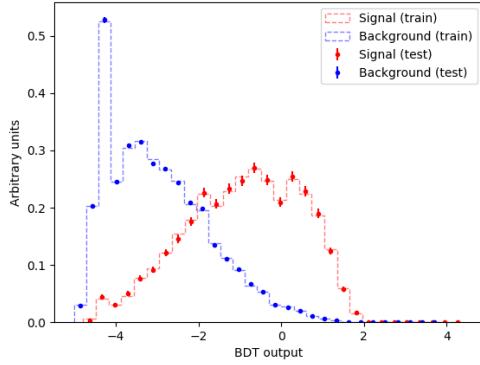


Figure 38: Normalised histograms of the object-level BDT discriminator output from the signal and background classes for the training and test sets from the 5th fold in a 5 fold kfold cross validation is shown. The output of the object-level BDT is shown on the x-axis and the relative number of events (in arbitrary units) is shown on the y-axis. The training set for the signal class is shown by the red dotted histogram. The test set for the signal class is shown by the red points, with the total uncertainty represented by the vertical error bars. The training set for the background class is shown by the blue dotted histogram. The test set for the background class is shown by the blue points, with the total uncertainty represented by the vertical error bars.

866 that no over-training occurred.

867
868 The output of the object-level BDT is converted to an event-level variable to be used in the event-level BDT. This
869 variable, $BDTScore(\frac{\text{Best}}{\text{2nd Best}})$, takes the ratio of the scores of the top scoring ℓb system to the 2nd best scoring ℓb
870 system. The 2nd best scoring ℓb system in a tWZ event is expected to be low, since there is only one ℓb system
871 originating from a top. Thus $BDTScore(\frac{\text{Best}}{\text{2nd Best}})$ is expected to be large for tWZ events and closer to one for ttZ

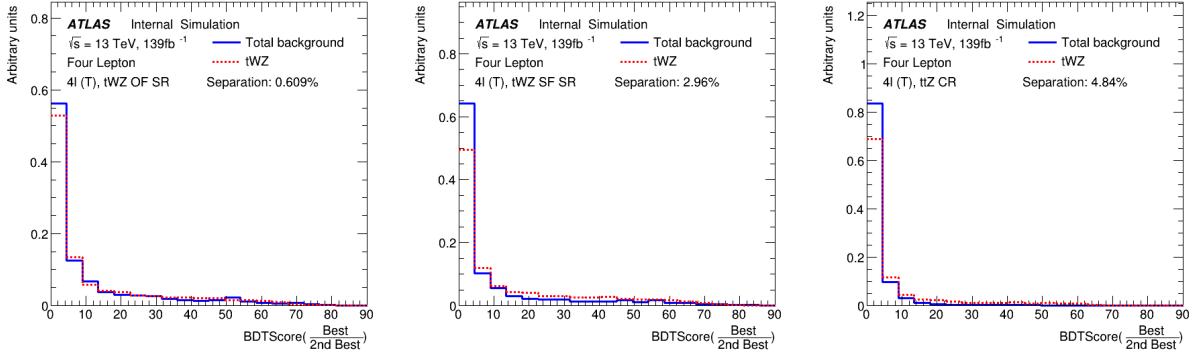


Figure 39: Normalised distributions of the signal and total background of the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable in the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR are shown (left to right). The dotted red and solid blue lines represent the distributions of the signal and total background events respectively. These histograms are normalised to an area of 1. The x-axis shows the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ and the y-axis show the relative number of events (in arbitrary units).

events, therefore providing discrimination between tWZ and $t\bar{t}Z$. In Figure 39, normalised distributions of the signal and total background of the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable in the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR are shown.

The amount of discrimination can be quantified by the separation metric, which gives the percentage of the total area of the distributions which do not overlap. A value of 1 indicates that the distributions are fully separated (no overlap) and a value of 0 indicates that the distributions have no separation (fully overlapped). The separation between signal and background for $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ in the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR are 0.609%, 2.96% and 4.84% respectively. The larger separation in the $t\bar{t}Z$ CR, compared to the tWZ SRs, can be explained since there is a larger proportion of $t\bar{t}Z$ events (events with two ℓb systems) in this region, due to the baseline selection requirement of exactly two b -tagged jets. In a similar way, the smaller separation in the two tWZ SRs can be explained by the tighter selection on the number of b -tagged jets (exactly one) leading to regions which are enriched in only one ℓb system which originates from a top quark. Using the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable in training in the event-level BDT (see Section 4.7.3) improves the mean accuracy of the BDT. This tells us that the event-level BDT is taking advantage of the discrimination between signal and background present in the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable.

In an attempt to optimise the performance of the object-level BDT, we aim to train on signal events which are as pure in ℓb systems originating from top quarks, as possible. Similarly, we aim to train on background events which are as pure in ℓb systems not originating from top quarks, as possible. This is done by studying the distribution of ΔR between the reconstructed ℓb system and the truth ℓb system ($\Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}})$), and excluding ℓb systems from training which are moderately matched in ΔR to their truth counterparts, leaving well matched ℓb systems being labelled as signal and badly matched ℓb systems labelled as background. We call the ΔR range where ℓb systems are excluded from training, the exclusion region. In Figure 40, the distribution of ΔR between the reconstructed ℓb system and the truth ℓb system ($\Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}})$) in the alternative $t\bar{t}$ sample, along with the exclusion region, is shown.

A large number of reconstructed ℓb systems have $\Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}})$ at values near 0. These are matched (in ΔR) extremely well to truth ℓb systems originating from top quarks. We therefore define our exclusion region to be between $0.05 < \Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}}) < 3.0$, such that all reconstructed ℓb systems with $\Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}}) \leq 0.05$ are labelled as signal and reconstructed ℓb systems with $\Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}}) \geq 3.0$ are labelled as background. All reconstructed ℓb systems with $0.05 < \Delta R((lb)_{\text{reco}}, (lb)_{\text{truth}}) < 3.0$ are excluded from training.

The performance of the object-level BDT with and without the exclusion region can be compared by studying the discrimination between signal and background events in the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable (object-level output converted to an event-level variable to be used in the event-level BDT) for both object-level BDTs. In Figure 41, normalised distributions of $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ using the object-level BDT without the exclusion region (see Figure 40) for the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR are shown.

We can compare the separation metrics between the $\text{BDTScore}(\frac{\text{Best}}{\text{2nd Best}})$ variable in the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR for the ℓb classifier with (Figure 39) and without (Figure 41) the exclusion region by taking the absolute difference between the two values in each region. The differences are 0.31%, 0.37% and 0.36% for the tWZ

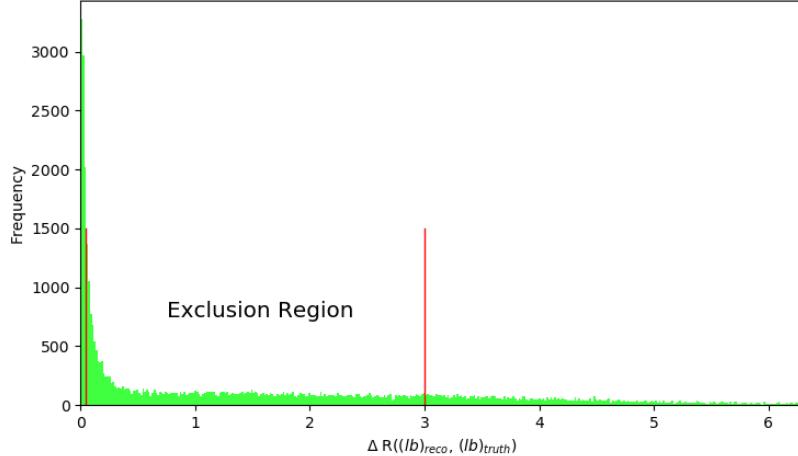


Figure 40: The distribution of ΔR between the reconstructed ℓb system and the truth ℓb system ($\Delta R((lb)_{reco}, (lb)_{truth})$) in the alternative $t\bar{t}$ sample, along with the exclusion region, is shown. The ΔR distribution is shown in green. ΔR between the reconstructed ℓb system and the truth ℓb system ($\Delta R((lb)_{reco}, (lb)_{truth})$) is shown on the x-axis. The bin frequency is shown on the y-axis. The exclusion region is shown between the vertical red lines situated at $\Delta R((lb)_{reco}, (lb)_{truth}) = 0.05$ and $\Delta R((lb)_{reco}, (lb)_{truth}) = 3.0$. Reconstructed ℓb systems with $\Delta R((lb)_{reco}, (lb)_{truth}) \leq 0.05$ are labelled as signal and reconstructed ℓb systems with $\Delta R((lb)_{reco}, (lb)_{truth}) \geq 3.0$ are labelled as background.

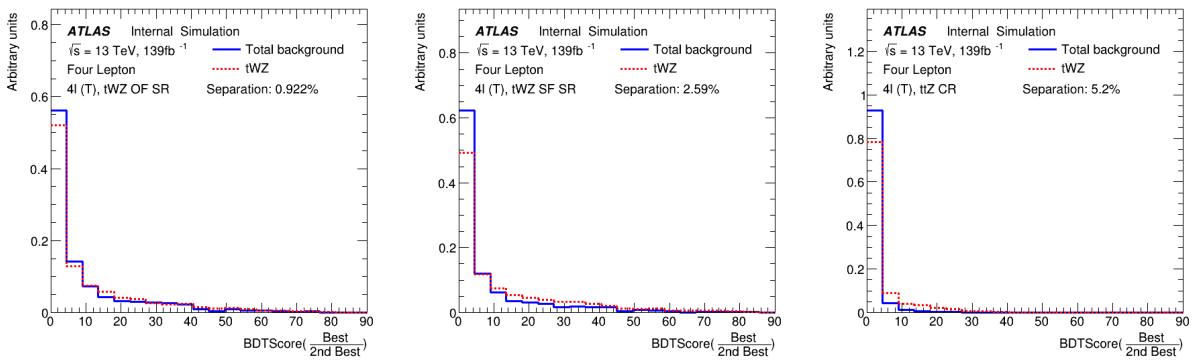


Figure 41: Normalised distributions of $BDTScore(\frac{\text{Best}}{\text{2nd Best}})$ using the object-level BDT without the exclusion region (see Figure 40) for the tWZ OF SR, tWZ SF SR and ttZ CR are shown (left to right). The dotted red and solid blue lines represent the distributions of the signal and total background events respectively. These histograms are normalised to an area of 1. The x-axis shows the $BDTScore(\frac{\text{Best}}{\text{2nd Best}})$ and the y-axis show the relative number of events (in arbitrary units).

910 OF SR, tWZ SF SR and $t\bar{t}Z$ CR, respectively. These differences are minimal and the object-level BDT with the
 911 exclusion region outperforms the object-level BDT without the exclusion region in the tWZ SF SR. Due to the
 912 small differences in performance between the two BDTs, we chose to keep the BDT with the exclusion region.

913 4.7.2 Two Neutrino Scanning Method (2ν SM) Algorithm

914 The Two Neutrino Scanning Method (2ν SM) algorithm¹ [60, 59] aims to reconstruct $t\bar{t}$ systems in the 2ℓ , 3ℓ and
 915 4ℓ final states (e.g. 2ℓ case: $t\bar{t} \rightarrow \ell^+ \nu_\ell b \ell^- \bar{\nu}_\ell \bar{b}$). This was initially designed to suppress the $t\bar{t}$ background in the $t\bar{t}Z$
 916 analysis. We can re-purpose this algorithm to distinguish between tWZ and $t\bar{t}Z$ by removing the easily-identifiable
 917 Z boson.

918 The 2ν SM algorithm reconstructs a $t\bar{t}$ system by scanning through the components of two possible neutrino 4-
 919 vectors (ν_1 and ν_2). It then aims to determine which ν_1 and ν_2 correspond to the two neutrinos which originate
 920 from the decay of a $t\bar{t}$ system the best (quantified by an output weight, $w_{2\nu SM}$). $w_{2\nu SM}$ is the likelihood under
 921 the $t\bar{t}$ dilipeton final state hypothesis. We are able to use this algorithm in our analysis to discriminate between
 922 tWZ and $t\bar{t}Z$ since we can easily reconstruct the OSSF leptons which decay from the Z boson and remove it before
 923 inputting the event into the algorithm. We would then expect that the 2ν SM algorithm returns a higher score from
 924 a $t\bar{t}Z$ event (~ 1 , i.e. it looks like a $t\bar{t}$ event after removal of the Z boson) and a lower score from a tWZ event (\sim
 925 0, i.e. it does not look like a $t\bar{t}$ event after removal of the Z boson).

927 4.7.2.1 The algorithm

928 The 2ν SM algorithm starts off by writing down four equations which correspond to the invariant masses of the top
 929 quark ($m(t)$) and W boson ($m(W)$) for the two top decays (i.e. $t \rightarrow W^+ b \rightarrow \ell^+ \nu_\ell$) in a dileptonic $t\bar{t}$ event. These
 930 can be written as,

$$(927) (\ell_1 + \nu_1)^2 = m(W)^2 = (80.385 \text{ GeV})^2 \quad (4.2)$$

$$(927) (\ell_1 + \nu_1 + b_{1,2})^2 = m(t)^2 = (172.5 \text{ GeV})^2 \quad (4.3)$$

$$(927) (\ell_2 + \nu_2)^2 = m(W)^2 = (80.385 \text{ GeV})^2 \quad (4.4)$$

$$(927) (\ell_2 + \nu_2 + b_{2,1})^2 = m(t)^2 = (172.5 \text{ GeV})^2 \quad (4.5)$$

931 where the subscripts indicate that these particles originate from the decay of two different top quarks in a $t\bar{t}$
 932 system. We assume that the mass of the neutrinos (ν_1 and ν_2) are close to zero, which leaves us with 6 unknowns,
 933 $p_{T\nu_1}$, ϕ_{ν_1} , η_{ν_1} , $p_{T\nu_2}$, ϕ_{ν_2} and η_{ν_2} (components of the two neutrino's 4-vectors).

934 The 2ν SM algorithm takes the 4-vectors of the two reconstructed leptons (not from the Z boson) and the two jets
 935 with the highest DL1r b -tagger score as input. For each neutrino (ν_1 and ν_2), we scan over a range of possible η and
 936 ϕ values. These values were chosen to be $\phi_{\nu_1}, \phi_{\nu_2} \in [-\pi, \pi]$ with a step size of ≈ 0.25 and $\eta_{\nu_1}, \eta_{\nu_2} \in [-5, 5]$ with a
 937 step size of ≈ 0.31 . These ranges were chosen to maximize accuracy and minimize computation time. For each of
 938 these possible η and ϕ values, we calculate the corresponding p_T for each neutrino. The transverse momentum of
 939 a neutrino, $p_{T\nu}$, can be calculated via (**recne somewhere here**),

$$(934) p_{T\nu} = \frac{\frac{1}{2}(m(W)^2 - m(\ell)^2)}{E_\ell \cosh \eta_\nu - p_{\ell,z} \sinh \eta_\nu - p_{\ell,x} \cos \phi_\nu - p_{\ell,y} \sin \phi_\nu} \quad (4.6)$$

941 where E_ℓ is the energy of the lepton and $p_{\ell,z}$, $p_{\ell,x}$, $p_{\ell,y}$ are the z , x and y components of lepton's momentum. At
 942 this stage, we have possible 4-vectors for ν_1 and ν_2 . Using these possible neutrino 4-vectors, we reconstruct the two
 943 possible $t\bar{t}$ systems,

$$(944) t_1 = \ell_1 + b_1 + \nu_1 \text{ and } t_2 = \ell_2 + b_2 + \nu_2 \quad (4.7)$$

OR

$$(944) t_1 = \ell_1 + b_2 + \nu_1 \text{ and } t_2 = \ell_2 + b_1 + \nu_2 \quad (4.8)$$

¹software tool and weights provided by Thomas McCarthy ($t\bar{t}Z$ analysis group - Max Planck Institute)

These reconstructed $t\bar{t}$ systems are then used to calculate a weight, $w_{2\nu SM}$. The $w_{2\nu SM}$ weight (a value ranging from 0 to 1) is defined as a product of four probabilities (described below) and can be written as,

$$w_{2\nu SM} = P_{m_{t_1}} \times P_{m_{t_2}} \times P_{\Delta E_x} \times P_{\Delta E_y} \quad (4.9)$$

The $w_{2\nu SM}$ is calculated for each pair of reconstructed neutrinos (or reconstructed $t\bar{t}$ systems), with the maximum value being chosen as the final value for the event.

4.7.2.2 Calculating $w_{2\nu SM}$

We use distributions of well modelled observables ($m_{b\ell\nu}$ and ΔE_x) from simulated $t\bar{t}$ events in order to determine how well our reconstructed neutrinos (and in turn top quarks) resemble neutrinos (and top quarks) present in a $t\bar{t}$ event.

$P_{m_{t_1}}$ and $P_{m_{t_2}}$

A normalised distribution of the mass of reconstructed top quarks ($m_{b\ell\nu}$) from a $t\bar{t}$ sample is generated to determine the probabilities $P_{m_{t_1}}$ and $P_{m_{t_2}}$. The distribution is generated from reco-level leptons, generator-level neutrinos and reoc-level jets matched in ΔR to generator-level b -quarks, therefore only filling the distribution with correct detector-level objects. We then use the distribution to interpolate our two reconstructed top quarks, which returns a weight value from 0 to 1, with higher values corresponding to a reconstructed top quark which has a mass close to that of a top quark from a $t\bar{t}$ system. This interpolation is done for both reconstructed tops, t_1 and t_2 , corresponding to probabilities $P_{m_{t_1}}$ and $P_{m_{t_2}}$. The distribution used is shown in Figure 42.

In Figure 42, the $m_{b\ell\nu}$ distribution (generated from simulated $t\bar{t}$ events), used to calculate $P_{m_{t_1}}$ and $P_{m_{t_2}}$ is shown.

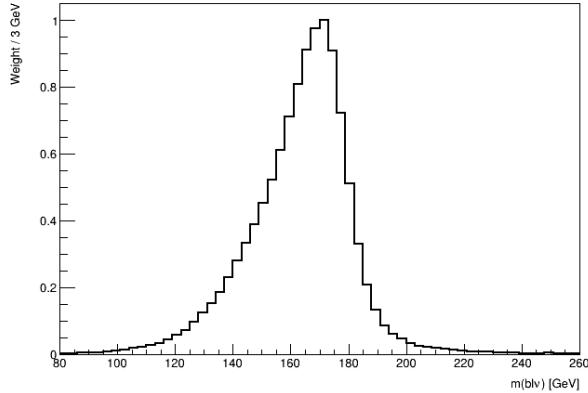


Figure 42: $m_{b\ell\nu}$ distribution generated from simulated $t\bar{t}$ events, used to calculate $P_{m_{t_1}}$ and $P_{m_{t_2}}$ is shown. The $m_{b\ell\nu}$ distribution is shown by the black lined histogram. The mass of the $b\ell\nu$ system is shown on the x-axis. The corresponding weight of the $m_{b\ell\nu}$ distribution is shown on the y-axis.

$P_{\Delta E_x}$ and $P_{\Delta E_y}$

A similar method is used to determine $P_{\Delta E_x}$ and $P_{\Delta E_y}$. In this case we generate a weight distribution of $\Delta E_x = (p_{T,\nu_1})_x + (p_{T,\nu_2})_x - (E_T^{\text{miss}})_x$ based off simulated $t\bar{t}$ events. In particular, this distribution is generated using reco-level E_T^{miss} and generator-level neutrinos. The use of this distribution lies under the assumption that neutrinos are the dominant source of E_T^{miss} , and therefore, $(E_T^{\text{miss}})_x \approx (p_{T,\nu_1})_x + (p_{T,\nu_2})_x$ and $(E_T^{\text{miss}})_y \approx (p_{T,\nu_1})_y + (p_{T,\nu_2})_y$. We then use the distribution to interpolate the value of ΔE_x and ΔE_y from our reconstructed neutrinos. This returns a weight value from 0 to 1, with higher values corresponding to ΔE_x and ΔE_y (and in turn our reconstructed neutrino's p_T) closer to those observed in a $t\bar{t}$ event. We expect the ΔE_x and ΔE_y distributions to have the same shapes, therefore we only need to generate one (we have chosen ΔE_x). In Figure 43, the $m_{b\ell\nu}$ distribution (generated from simulated $t\bar{t}$ events), used to calculate $P_{m_{t_1}}$ and $P_{m_{t_2}}$ is shown.

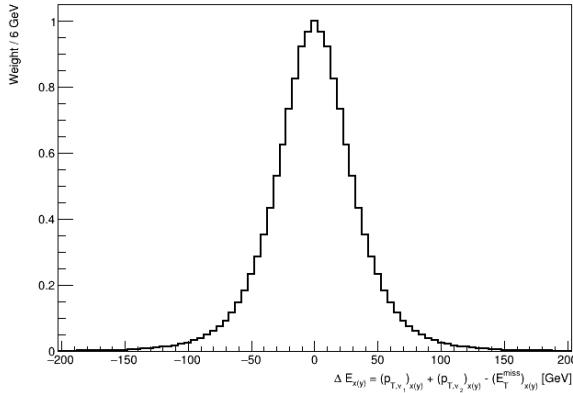


Figure 43: ΔE_x distribution generated from simulated $t\bar{t}$ events, used to calculate $P_{\Delta E_x}$ and $P_{\Delta E_y}$ is shown. The ΔE_x distribution is shown by the black lined histogram. ΔE_x is shown on the x-axis. The corresponding weight of ΔE_x distribution is shown on the y-axis.

974 4.7.2.3 Kinematic Vetoess

975 The 2ν SM algorithm is extremely computationally intensive. The computation time depends on the number step
 976 size of the ϕ and η ranges which we scan over to reconstruct the neutrinos. For example, consider the step sizes
 977 chosen in this analysis, $\Delta\eta \approx 0.31$ and $\Delta\phi \approx 0.25$ which corresponds to 32 values for η and 25 values for ϕ . There
 978 will be $(32)(32)(25)(25) = 640\,000$ possible pairs of neutrinos (ν_1 and ν_2) to consider **per event**. Since we have to
 979 consider two possible $t\bar{t}$ systems (See Equations 4.7 and 4.8), this number effectively increases to $(2)(640\,000) = 128\,000$
 980 iterations **per event**. In order to reduce the number of $t\bar{t}$ systems we need to consider, therefore decreasing
 981 computation time, we look at distributions of well modelled observables from $t\bar{t}$ events and veto (discard) a possible
 982 reconstructed $t\bar{t}$ system if the observable in question is improbable or unlikely to be observed in a $t\bar{t}$ event. To
 983 achieve this, we define a threshold range for these observables (See Figure 46 and Figure 48), and if the possible
 984 reconstructed $t\bar{t}$ system's corresponding value for this observable lies outside this range, it is vetoed and the algorithm
 985 continues with the next iteration.

986 $\Delta\langle m(\ell b) \rangle$

987 The first observable which we consider is the difference between average mass of the two possible ℓb system combi-
 988 nations, $\Delta\langle m(\ell b) \rangle$. The two possible ℓb system combinations are,

$$(1) (\ell b)_1 = \ell_1 + b_1 \text{ and } (\ell b)_2 = \ell_2 + b_2 \quad (4.10)$$

OR

$$(2) (\ell b)_1 = \ell_1 + b_2 \text{ and } (\ell b)_2 = \ell_2 + b_1 \quad (4.11)$$

$$(4.12)$$

990 $\Delta\langle m(\ell b) \rangle$ is therefore defined as,

$$\Delta\langle m(\ell b) \rangle = \frac{1}{2} |[(m(\ell_1 b_1) + m(\ell_1 b_1)) - (m(\ell_2 b_2) + m(\ell_2 b_1))]| \quad (4.13)$$

991 The idea here is that, if $\Delta\langle m(\ell b) \rangle$ is large, it's more likely that we can simply select the ℓb combination with the
 992 smaller (minimum) average mass. To illustrate this, we look at the distribution (constructed from $t\bar{t}$ events) of
 993 $P(\text{Correct combination of } \ell b \text{ systems} | \text{minimum}\langle m(\ell b) \rangle)$ vs $\Delta\langle m(\ell b) \rangle$ for b -tagged jets in the same ($\eta(b_1) \times \eta(b_2) \geq 0$)
 994 and opposite hemispheres ($\eta(b_1) \times \eta(b_2) < 0$).

995 In Figure 44 the $P(\text{Correct combination of } \ell b \text{ systems} | \text{minimum}\langle m(\ell b) \rangle)$ vs $\Delta\langle m(\ell b) \rangle$, for b -tagged jets in the same
 996 and opposite hemispheres, constructed from $t\bar{t}$ events is shown.

997 From Figure 44, for both cases where the b -tagged jets are in the same and opposite hemispheres, the
 998 probability for a correct ℓb system being chosen given that we are considering the ℓb system with mini-
 999 mum average mass is an increasing function which plateaus to 1 at ~ 90 GeV. We use these two distribu-
 1000 tions to interpolate the $P(\text{Correct combination of } \ell b \text{ systems} | \text{minimum}\langle m(\ell b) \rangle)$ from $\Delta\langle m(\ell b) \rangle$. We require that

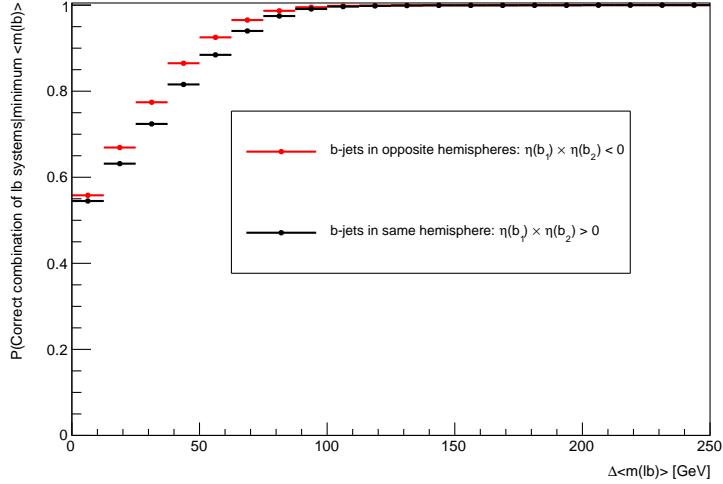


Figure 44: $P(\text{Correct combination of } lb \text{ systems}|\text{minimum}\langle m(lb) \rangle)$ vs $\Delta\langle m(lb) \rangle$, for b -tagged jets in the same and opposite hemispheres, constructed from $t\bar{t}$ events is shown. The horizontal red lines represent the distribution in the case when the two b -jets are in opposite hemispheres. The dot in the middle of the line represents the midpoint of the line. The horizontal black lines represent the distribution in the case when the two b -jets are in the same hemispheres. The dot in the middle of the line represents the midpoint of the line. The average $m(lb)$ is shown on the x-axis. The $P(\text{Correct combination of } lb \text{ systems}|\text{minimum}\langle m(lb) \rangle)$ is shown on the y-axis.

1001 $P(\text{Correct combination of } lb \text{ systems}|\text{minimum}\langle m(lb) \rangle) > 0.8$, before vetoing any lb combination, such that we
 1002 have are at least 80% certain that we know the correct lb combination. In this case, the lb combination with the
 1003 maximum $\Delta\langle m(lb) \rangle$ is vetoed. If $P(\text{Correct combination of } lb \text{ systems}|\text{minimum}\langle m(lb) \rangle) < 0.8$ we need to consider
 1004 both possible lb system combinations.

1005 $\eta(b\bar{b}\ell\ell)$

1006 We consider η of the $b\bar{b}\ell\ell$ system, $\eta(b\bar{b}\ell\ell)$ to veto improbable $\eta(\nu_1)$ and $\eta(\nu_2)$ values.

1007 In the same way as for $\Delta\langle m(lb) \rangle$, we generate a distribution to determine values $\eta(\nu)$ which are improbable for a $t\bar{t}$
 1008 event. In this case, we generate a 2D histogram from simulated $t\bar{t}$ events (dileptonic final state) at generator-level
 1009 of $\eta(\nu)$ vs $\eta(b\bar{b}\ell\ell)$.

1010 In Figure 45, a heatmap of occupancy for $\eta(\nu)$ vs $\eta(b\bar{b}\ell\ell)$ (produced from simulated $t\bar{t}$ events) is shown.
 1011 Using the above heatmap, we define a veto region (where a $t\bar{t}$ event is extremely unlikely to occur) based off
 1012 double-sided 95% limits (**something here on confidence limit??**). We apply a veto if either possible neutrino
 1013 lies within this region. The veto region is shown in Figure 46.

1014 In Figure 46, the veto region (extracted from Figure 45) for vetoing improbable neutrinos is shown.

1015 L_T

1016 The final kinematic constraint which we consider is the scalar sum of lepton p_T , $L_T = p_T(\ell_1) + p_T(\ell_2)$ which we
 1017 use to veto certain possible neutrinos, ν_1 and ν_2 .

1018 Again, we generate a distribution to determine (and veto) improbable possible neutrinos in simulated $t\bar{t}$ events
 1019 (dilepton final state).

1020 In Figure 47, a heatmap of occupancy for $\Delta R(\ell, \nu)$ vs L_T (produced from simulated $t\bar{t}$ events) is shown.
 1021 Using the same method as described in Section 4.7.2.3, we define a veto region where a veto is applied if either
 1022 possible neutrino lies within this region. In Figure 46, the veto region (extracted from Figure 47) for vetoing
 1023 improbable neutrinos values is shown.

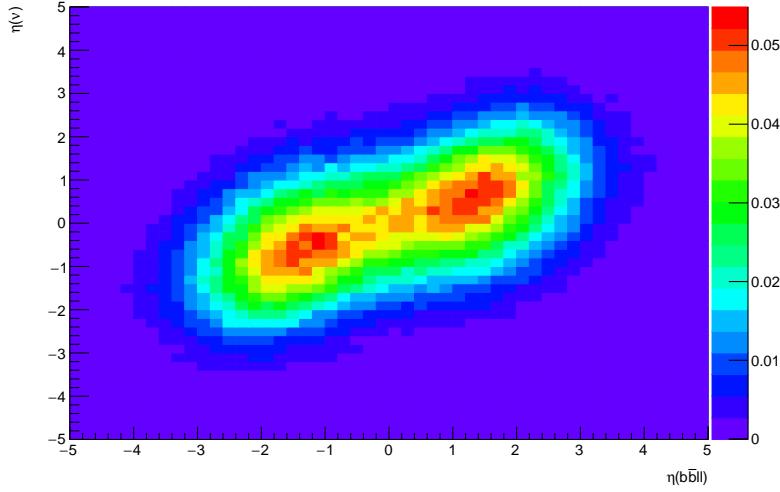


Figure 45: Heatmap of occupancy for $\eta(\nu)$ vs $\eta(b\bar{b}\ell\ell)$ produced from simulated $t\bar{t}$ events (dileptonic final state) at generator-level is shown. η of the $b\bar{b}\ell\ell$ system is shown on the x-axis. η of the neutrino is shown on the y-axis. The colorbar on the right represents the occupancy (normalised) in the phase space.

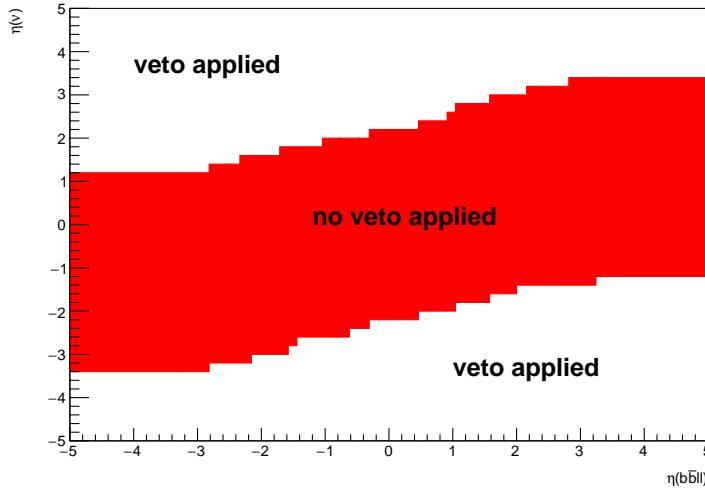


Figure 46: The regions where vetoes are applied for the $\eta(b_1 b_2 \ell_1 \ell_2)$ constraint is shown. η of the $b\bar{b}\ell\ell$ system is shown on the x-axis. η of the neutrino is shown on the y-axis. The red band shows the region where the neutrino would not be vetoed. The white areas (above and below the red band) are regions where the neutrino is vetoed.

1030 4.7.3 Event-level BDT

1031 The event-level BDT is used to distinguish between signal and its major background events, $t\bar{t}Z$ and ZZ . The key
 1032 difference between the object-level BDT and the event-level BDT is that while the former exploits information
 1033 associated with ℓb systems and thus distinguishes between ℓb systems, the event-level BDT exploits information
 1034 based on the entirety of the event and thus distinguishes between events.

1035
 1036 The event-level BDT was trained on 50% of the tWZ MC sample's events for the signal class and similarly,
 1037 50% of the $t\bar{t}Z$ and ZZ MC sample's events were used for the background class. The samples we train on are
 1038 individual events, with the features being carefully chosen observables. These observables are chosen on the basis
 1039 that they are somewhat uncorrelated from one another and show a relatively large amount of separation power
 1040 between tWZ and $t\bar{t}Z$. Similarly to the object-level BDT, the optimum values for the hyper-parameters used
 1041 were determined via hyper-parameter optimisation. After hyperparameter optimisation, the mean accuracy of

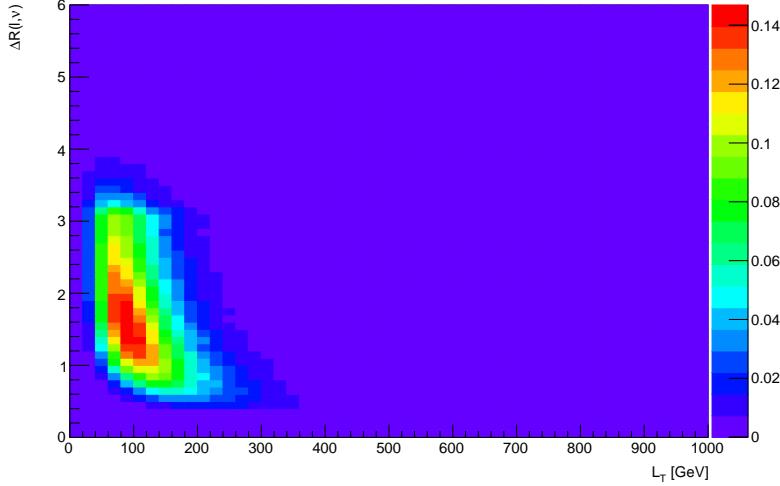


Figure 47: A heatmap of occupancy for $\Delta R(\ell, \nu)$ vs L_T produced from simulated $t\bar{t}$ events (dileptonic final state) at generator-level is shown. ΔR between leptons and neutrinos is shown on the x-axis. L_T (scalar sum of lepton p_T) is shown on the y-axis. The colorbar on the right represents the occupancy (normalised) in the phase space.

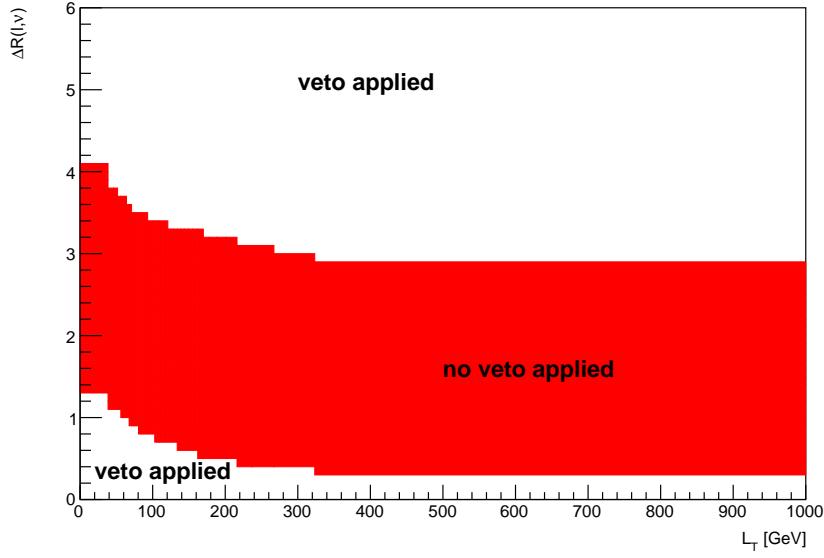


Figure 48: The regions where vetoes are applied for the L_T constraint is shown. ΔR between leptons and neutrinos is shown on the x-axis. L_T (scalar sum of lepton p_T) is shown on the y-axis. The red band shows the region where the neutrino would not be vetoed. The white areas (above and below the red band) are regions where the neutrino is vetoed.

1042 each fold (determined from 5 fold kfold cross validation) increased from 0.72 to 0.74 ($\sim 3\%$ increase). The variable
1043 importance of each variable was computed in the same way as described for the object-level BDT (See Section 4.7.1).

1044 In Table 10, the variables used in training the event-level BDT are shown.

1045 In Figure 49, normalised distributions of the variables used in the event-level BDT, for the signal and background
1046 classes are shown.

1047 Overall the variables used in the BDT show a reasonable amount of discrimination. In particular the output weight
1048 from the 2ν SM algorithm shows the most discrimination. When determining which variables to use in training the
1049 event-level BDT, the output weight from 2ν SM was shown to provide the most sizeable boost in performance of

Observable	Description
$2\nu\text{SM}$	Maximum weight from the $2\nu\text{SM}$ algorithm (See Section 4.7.2)
HT	Scalar sum of jet p_T
LT	Scalar sum of lepton p_T
$\sum p_T(b - \text{jet})$	Scalar sum of b -tagged jet p_T
$\text{BDTScore}_{(2^{\text{nd}} \text{ Best})}^{(\text{Best})}$	Ratio of the top scoring ℓb system to the 2nd best scoring ℓb system from the output of the object-level BDT (ℓb classifier)
$\Delta\eta(\ell_1, \text{non-}Z, \ell_2, \text{non-}Z)$	$\Delta\eta$ between the two leptons, not coming from a Z candidate

Table 10: A list of the observables used in the event-level BDT, ordered by variable importance (descending, top to bottom) is shown.

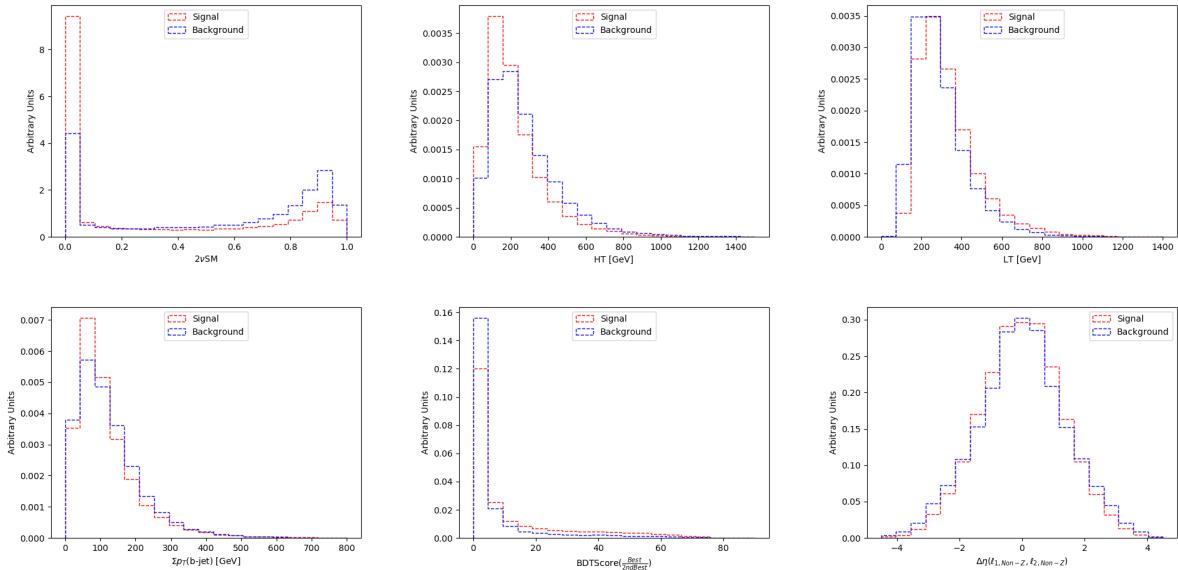


Figure 49: Normalised distributions of the variables used in the event-level BDT (ordered from top left to bottom right via decreasing variable importance), for the signal and background classes are shown. **From top left to bottom right:** Output weight from the $2\nu\text{SM}$ algorithm (See Section 4.7.2). Scalar sum of jet p_T . Scalar sum of lepton p_T . Sum of b -tagged jet p_T . Ratio of the top scoring ℓb system to the 2nd best scoring ℓb system from the output of the object-level BDT (See Section 4.7.1). $\Delta\eta$ between the two leptons, not coming from a Z candidate. The red and blue dotted lined histograms represent the signal and background classes events, respectively. These histograms are normalised to an area of 1. The variable used in training is shown on the x-axis. The y-axis shows the relative number of events for the signal and background classes (in arbitrary units).

the BDT. Surprisingly, the least important variable, $\Delta\phi$ between the non- Z lepton system (leptons not originating from a Z -candidate) and the leading b -tagged jet, seem to discriminate well between signal and background. A possible explanation for its low ranking variable importance is due to it being relatively highly correlated with many of the other input variables.

We can check the modelling of the input variables to the event-level BDT by referring to the expected number of events of data and simulation in control regions where they are defined. Note that certain variables which are ill-defined in certain regions (e.g. $\Delta\eta(\ell_{1,\text{non-}Z}, \ell_{2,\text{non-}Z})$ in the ZZb CR, as all leptons originate from a b -jet in this region) will not be shown.

In Figure 50, MC predictions for the input variables to the event-level BDT in the $t\bar{t}Z$ CR are shown.

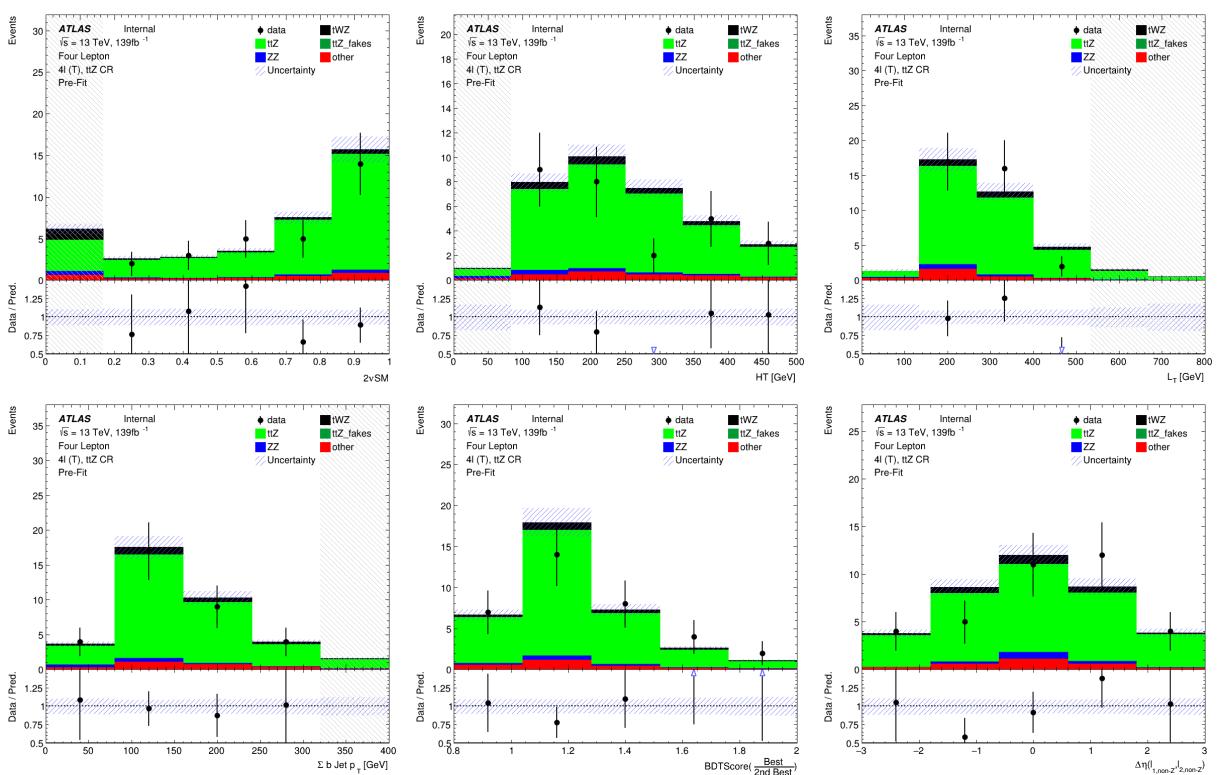


Figure 50: Pre-fit distributions of variables used as input to the event-level BDT (ordered from top left to bottom right via decreasing variable importance), in the $t\bar{t}Z$ CR, are shown. **From top left to bottom right:** Output weight from the 2 ν SM algorithm (See Section 4.7.2). Scalar sum of jet p_T . Scalar sum of lepton p_T . Sum of b -tagged jet p_T . Ratio of the top scoring llb system to the 2nd best scoring llb system from the output of the object-level BDT (See Section 4.7.1). $\Delta\eta$ between the two leptons, not coming from a Z candidate. The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the blue diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. Bins with $\frac{\text{signal}}{\text{background}}$ greater than 0.1 are kept blinded. Blinded bins are shaded with black diagonal lines and their data points are omitted.

In Figure 51, MC predictions for the input variables to the event-level BDT in the ZZb CR are shown.

In Figure 52, MC predictions for the input variables to the event-level BDT in the $(tWZ)_{\text{fake}}$ CR are shown.

Overall, the agreement between data and simulation for the input variables to the event-level BDT is good. Therefore, these variables are well-modelled and reasonable to include as inputs to the event-level BDT.

In Table 11, the hyper-parameters used in the event-level BDT are shown.

Since we are training on $t\bar{t}Z$ and ZZ events for the background class, we ensure that the relative weighting of these events are such that they mimic the amount of $t\bar{t}Z$ and ZZ expected to be present in the regions where we aim to use the BDT discriminator (tWZ SRs and $t\bar{t}Z$ CR). This is done by applying normalization weights to each event,

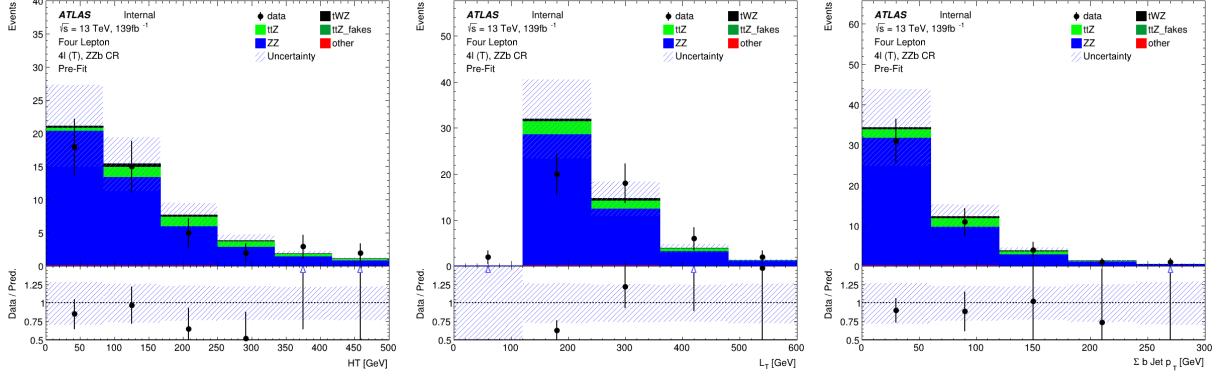


Figure 51: Pre-fit distributions of variables used as input to the event-level BDT (ordered from top left to bottom right via decreasing variable importance), in the ZZb CR, are shown. **From left to right:** Scalar sum of jet p_T . Scalar sum of lepton p_T . Sum of b -tagged jet p_T . The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the blue diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. Bins with $\frac{\text{signal}}{\text{background}}$ greater than 0.1 are kept blinded. Blinded bins are shaded with black diagonal lines and their data points are omitted.

Hyper-parameter	Value	Description
<code>loss</code>	<code>deviance</code>	The loss function to be optimised
<code>criterion</code>	<code>friedman_mse</code>	The function used to measure the quality of a split
<code>n_estimators</code>	200	The number of boosting stages to perform
<code>learning_rate</code>	0.1	The step size at each iteration during optimisation
<code>max_depth</code>	6	The maximum depth of the individual regression estimators
<code>min_samples_split</code>	2	The minimum number of samples (events) required to split an internal node
<code>min_samples_leaf</code>	1	The minimum number of samples (events) required to be at a leaf node
<code>validation_fraction</code>	0.1	The proportion of training data to set aside as validation set for early stopping
<code>n_iter_no_change</code>	20	Training terminates when the validation score (determined by the validation set) does not improve in all of the previous

Table 11: A list of the hyper-parameters used in the event-level BDT is shown. Hyperparameters not listed in this table use the default values as stated in the Scikit-learn Documentation[72].

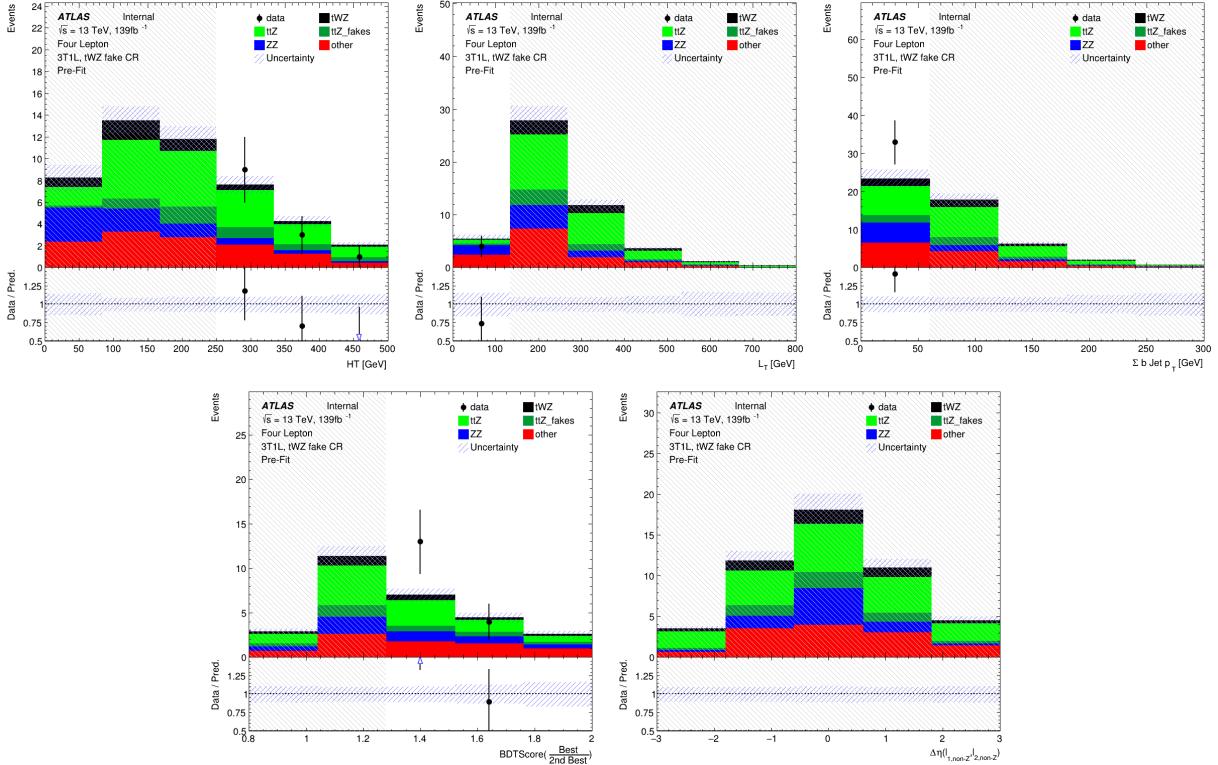


Figure 52: Pre-fit distributions of variables used as input to the event-level BDT (ordered from top left to bottom right via decreasing variable importance), in the (tWZ)_{fake} CR, are shown. **From top left to bottom right:** Output weight from the 2ν SM algorithm (See Section 4.7.2). Scalar sum of jet p_T . Scalar sum of lepton p_T . Sum of b -tagged jet p_T . Ratio of the top scoring ℓb system to the 2nd best scoring ℓb system from the output of the object-level BDT (See Section 4.7.1). $\Delta\eta$ between the two leptons, not coming from a Z candidate. The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the blue diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. Bins with $\frac{\text{signal}}{\text{background}} > 0.1$ are kept blinded. Blinded bins are shaded with black diagonal lines and their data points are omitted.

1071 defined as,

$$W = \frac{\sigma \mathcal{L} \text{weight(MC)}}{\text{totalWeight(MC)}} \quad (4.14)$$

1072 where σ is the cross section of the process, \mathcal{L} is the integrated luminosity, weight(MC) is the weight assigned to
1073 the event by the MC generator and totalWeight(MC) is the sum of those weights for all the generated events.

1074 The number of events used in training for the signal and background classes were 41066 and 22608 respectively.
1075 Similarly to the object-level BDT, there is a dataset imbalance. We correct this imbalance (in the same way as
1076 before with the object-level BDT) by ensuring that the relative weighting of each event is such that the sum of the
1077 signal weights is equal to the sum of the background weights.

1078 In Figure 53 the normalised histograms of the training and test sets (extracted from fold 5 from a 5 fold kfold cross
1079 validation) for signal and background is shown.

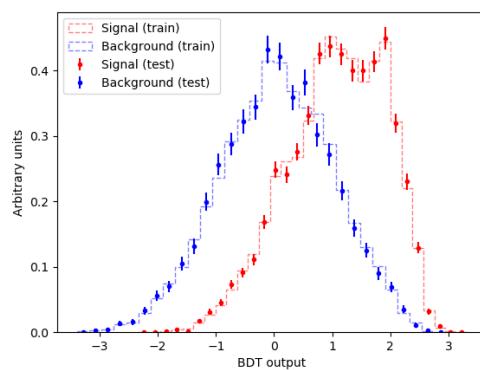


Figure 53: Normalised histograms of the event-level BDT discriminator output from the signal and background classes for the training and test sets from the 5th fold in a 5 fold kfold cross validation are shown. The output of the event-level BDT is shown on the x-axis and the relative number of events (normalised to have an area of 1, in arbitrary units) is shown on the y-axis. The training set for the signal class is shown by the red dotted histogram. The test set for the signal class is shown by the red points, with the total uncertainty represented by the vertical error bars. The training set for the background class is shown by the blue dotted histogram. The test set for the background class is shown by the blue points, with the total uncertainty represented by the vertical error bars.

1081 We can see that the shapes of the training and test sets for both signal and background are very similar. This is
1082 a good indicator that no over-training occurred. As with the object-level BDT, we perform another over-training
1083 check, by ensuring that the variance of the mean accuracy of each folds' test set in a 5 fold kfold cross validation is
1084 sufficiently small. This indicates that fluctuations in features from different training sets are not learnt by the BDT.
1085 For the event-level classifier, a variance of 0.00026 was calculated for the mean accuracies of each folds' test set in
1086 cross validation, providing further evidence that no over-training occurred. In Figure 54, normalised distributions
1087 of the signal and total background of the event-level BDT discriminator output in the tWZ OF SR, tWZ SF SR
1088 and $t\bar{t}Z$ CR, are shown.
1089 The event-level BDT discriminates well between signal and background events in the tWZ OF SR, tWZ SF SR
1090 and $t\bar{t}Z$ CR, with separations of 8.98%, 10.6% and 20.6%, respectively.

1092 4.8 Systematics

1093 The systematic uncertainties can be separated into experimental (detector) systematics, which are related to the
1094 reconstruction of physics objects in the detector and theoretical uncertainties related to the modelling of the different
1095 processes background.

1096 4.8.1 Experimental uncertainties

1097 In this section, the experimental systematics are outlined.

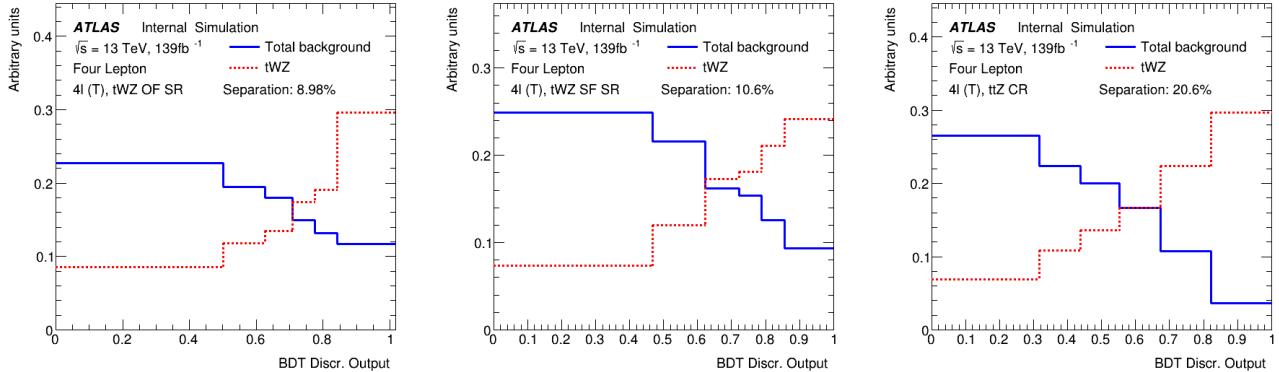


Figure 54: Normalised distributions of the signal and total background of the event-level BDT discriminator output in the tWZ OF SR, tWZ SF SR and $t\bar{t}Z$ CR are shown (left to right). The dotted red and solid blue lines represent the distributions of the signal and total background events respectively. These histograms are normalised to an area of 1. The x-axis shows the event-level BDT discriminator output and the y-axis shows the relative number of events (in arbitrary units).

• Luminosity:

The 2015–2018 luminosity estimate of 139fb^{-1} has a relative uncertainty of 3%. This uncertainty is obtained using the LUCID-2 detector [23] for the primary luminosity measurements. This systematic uncertainty affects all processes modelled using MC simulations.

• Pile-up reweighting:

An uncertainty related to the SFs used for MC to account for differences in pile-up distributions between MC and data is applied. This uncertainty is obtained by re-scaling the $\langle \mu \rangle$ value in data by 1.00 and 1/1.18 corrections are only applied to MC.

• Jet vertex tagger:

Uncertainties associated to the JVT are applied via the `JetJvtEfficiency` package [21] which account for the residual contamination from pile-up jets after pile-up suppression and the MC generator choice [14].

• Heavy- and light-flavor tagging:

The efficiency of the flavour-tagging algorithm is measured for each jet flavour using control samples in data and in simulation. From these measurements, correction factors are derived to correct the tagging rates in the simulation. In the case of b -tagged jets, the correction factors and their uncertainties are estimated from data using dileptonic $t\bar{t}$ events [16, 11]. In the case of c -jets, they are derived from jets arising from W boson decays in $t\bar{t}$ events [15]. In the case of light-flavour jets, the correction factors are derived using dijet events [12]. Sources of uncertainty affecting the b - and c -tagging efficiencies are evaluated as a function of jet p_T , including bin-to-bin correlations. The uncertainties in the efficiency for tagging light-flavour jets depend on the jet p_T and on η . An additional uncertainty is assigned to account for the extrapolation of the b -tagging efficiency measurement from the p_T region used to determine the correction factors to regions with higher p_T .

• Electron efficiency:

Uncertainties associated with the electron efficiency SFs are provided by the egamma CP group [19] and arise from the reconstruction, ID, isolation and trigger efficiencies. They correct for the efficiency difference between data and MC [13] and are measured with a “tag-and-probe” method in $Z \rightarrow e^+e^-$ and $J/\psi \rightarrow e^+e^-$ events. The information on the correlation of the different components of the systematic uncertainties are provided for all efficiency measurements. The default correlation model for the uncertainties is used, which provides one up/down variation for each of the SF components separately [19, 20].

• Muon efficiency:

As for electrons, SFs obtained from $Z \rightarrow \mu^+\mu^-$ and $J/\psi \rightarrow \mu^+\mu^-$ events are applied to correct for the differences between data and MC in the muon ID, isolation and trigger efficiencies [17]. Uncertainties on these

1129 SFs are provided by the muon CP group [22] and applied as up/down variations of the nominal SFs for each
 1130 component.

1131 4.8.2 Theoretical uncertainties

1132 In this section, the theoretical systematics are outlined.

1133 • **$t\bar{t}Z$ background:**

1134 An overall normalization uncertainty of 10% is considered for the $t\bar{t}Z$ background. Two generic shape sys-
 1135 tematics are considered for the $t\bar{t}Z$ background. They are constructed (see Section 4.8.3) by either applying a
 1136 linear or triangular interpolation to up and down variations which are defined to be $\pm 20\%$ from the nominal
 1137 $t\bar{t}Z$ background.

1138 • **ZZ background:**

1139 An overall normalization uncertainty of 30% is considered for the ZZ background.

1140 • **$t\bar{t}H$ background:**

1141 An overall normalization uncertainty of 20% is considered for the $t\bar{t}H$ background.

1142 • **tZq background:**

1143 An overall normalization uncertainty of 14% is considered for the tZq background.

1144 • **$t\bar{t}Z$ fake background:**

1145 An overall normalization uncertainty of 50% is considered for the $t\bar{t}Z$ fake background.

1146 • **other background processes:**

1147 The 'other' background consists of many processes which have minimal but non-negligible contribution in
 1148 the signal regions (See Table 6). An overall normalization uncertainty of 30% is considered for the 'other'
 1149 background processes.

1150 • **tWZ :** A modelling uncertainty on tWZ is considered by comparing the nominal sample (using the DR1
 1151 scheme) and a minimal DR2 sample.

1152 Two generic shape systematics are considered for the tWZ background. They are constructed (see Section
 1153 4.8.3) by either applying a linear or triangular interpolation to up and down variations which are defined to
 1154 be $\pm 20\%$ from the nominal tWZ background.

1155 4.8.3 Generic shape systematics

1156 It is evident that the tetralepton channel is statistically limited. We therefore expect that the uncertainty on u_{tWZ}
 1157 is dominated by statistical uncertainty and that the impact of shape systematics will be negligible in comparison.

1158 In order to include shape uncertainties related to the modelling of our samples, we construct generic shape
 1159 systematics for any given sample process. Given that we choose a sufficiently large set of values for which the
 1160 systematics can take in the fit, the constructed systematics could represent many shape systematics which we have
 1161 not yet considered to include in the fitting procedure.

1162 We start by constructing an envelope (error bars) consisting of two MC templates. One with the nominal MC
 1163 template increased by 20% on its normalisation and the other with the nominal MC template decreased by 20% on
 1164 its normalisation. The templates are then modified from their original shape either by doing linear interpolation
 1165 (from the leftmost-up variation to the rightmost-down variation) or triangular interpolation (shape is set to zero at
 1166 the rightmost and leftmost parts and reaches the envelope in the middle). The linear and triangular interpolation is
 1167 done using TRF's `ForceShape` option [68], which alters the original templates (as described above). This envelope
 1168 now represents the bounds which the systematic can vary in the fit.

1169 In Figure 55 the envelope before and after the shape change, for both the linear and triangular interpolations, for
 1170 the $t\bar{t}Z$ background in the $t\bar{t}Z$ CR is shown.

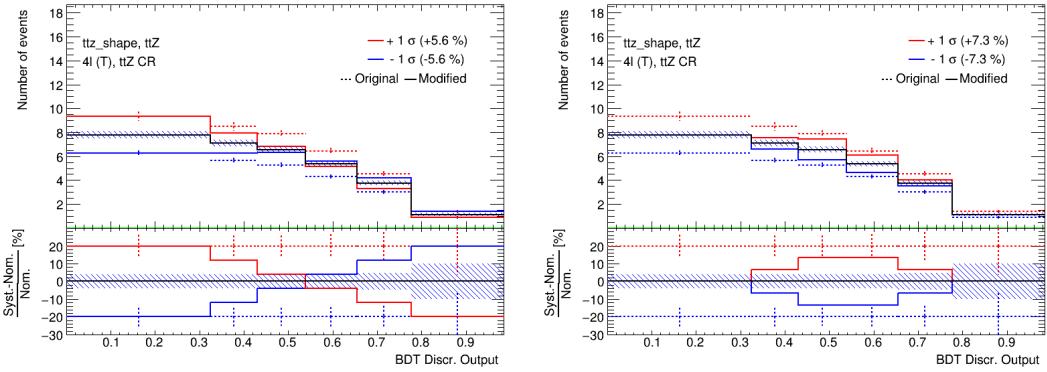


Figure 55: $t\bar{t}Z$ generic shape systematic before (original) and after (modified) linear (left) and triangular (right) interpolation in the $t\bar{t}Z$ CR is shown. The output from the event-level BDT shown on the x-axis. In the upper panel, the number of events is shown on the y-axis. In the lower panel, the difference between the systematic variation (the envelope's templates) and the nominal template, divided by the nominal template, is shown on the y-axis. The nominal $t\bar{t}Z$ template is shown by the solid black lined histogram, with the diagonal lined bands representing its total uncertainty. The templates of the upper and lower envelopes (before modification) is given by the dotted red and blue lined histograms respectively. The templates of the upper and lower envelopes (after modification) is given by the solid red and blue lined histograms respectively, with the vertical dotted lines representing its total uncertainty.

1174 We consider two shape systematics (linear and triangular shapes) for tWZ and the most dominant background
 1175 processes across both channels. In particular, we consider these shape systematics for $t\bar{t}Z$ for both the tri- and
 1176 tetralepton channels. Additionally, we consider these shape systematics for the $WZ + b$ and $WZ + c$ backgrounds
 1177 in the trilepton channel.

1178 4.9 Analysis Pipeline and **TRExFitter**

1179 We make use of industry standard **ROOT**² wrappers in this analysis, namely, **PyROOT** and **TRExFitter**.

1180
 1181 Python is used extensively in many fields of science (not limited to physics and data science) due to its simplicity
 1182 and ongoing support by the communities which utilize it. **PyROOT** allows users to access the full **ROOT** functionality
 1183 within Python. More specifically, **PyROOT** provides Python bindings for **ROOT**.

1184
 1185 **TRExFitter** is a framework for binned template profile likelihood fits[79]. In this analysis, we used **TRExFitter**
 1186 (tag: **TRExFitter-00-04-13**) to produce all plots (including fit statistics, e.g. limit, significance, $\mu_{best-fit}$).

1187
 1188 The analysis pipeline starts with sample derivations (derived dataset) being submitted to the grid for ntuple
 1189 production. This applies cuts and selections to the already reduced derivations and produces ntuples with trees
 1190 containing variables (e.g. scale factors, observables, MC truth flags) that will be used at future stages in the analysis.
 1191 These ntuples are then read by **PyROOT** where the events are looped over, before being written to **ROOT** files as input
 1192 to **TRExFitter**. The Python scripts are used to define the different regions and apply the final cuts and selections
 1193 outlined in Table 5. In addition to this purpose, they are used to train the two BDTs and to produce the output
 1194 from these trained BDTs. As each event is looped over, the cuts and selection criteria are checked for the given
 1195 event and is either thrown away (if the event does not pass the selection criteria), or gets written to a **ROOT** file (if
 1196 the event passes the selection criteria) corresponding to the MC sample and Run 2 data-set which it belongs to.
 1197 These **ROOT** files contain all observables, weights and scale factors (corresponding to an event) which we wish to use
 1198 in **TRExFitter**. **TRExFitter** then takes these files as input, runs a maximum likelihood fit and produces relevant
 1199 plots (e.g. expected number of events, pull plots) and statistical parameters (e.g. limit, significance, $\mu_{best-fit}$).

²CERN’s HEP data analysis framework (written in C++)

1200 4.9.1 Fitting Procedure

1201 Using the TRexFitter framework, binned profile-likelihood fits are performed to determine the signal strength,
 1202 $\mu_{tWZ} = \frac{\sigma_{obs}(tWZ)}{\sigma_{SM}(tWZ)}$, of tWZ production. A fit across all regions in the tetralepton channel is performed to determine
 1203 the sensitivity tWZ in this channel. In Section 4.10.2, a combined fit is performed across all regions in the trilepton
 1204 and tetralepton channels to take advantage of the sensitivity of tWZ in both channels in order to further boost
 1205 the sensitivity of tWZ .

1206 To characterise the sensitivity and associated uncertainty of our measurement of $\mu(tWZ)$, we compute two metrics:
 1207 the expected significance (Z_μ^{exp}) and the expected upper limit (μ_{up}^{exp}). In this context, the expected significance
 1208 can be interpreted as, the probability that the measured signal is due to a background fluctuation. Larger values
 1209 indicate lower probabilities and smaller values indicate higher probabilities. Particle physicists have adopted a
 1210 standard to define the sensitivity necessary for evidence and discovery of a particular particle or phenomena. A
 1211 3σ (corresponding to a background fluctuation probability of $\approx 10^{-3}$) significance is considered to be evidence
 1212 for observation and a 5σ (corresponding to a background fluctuation probability of $\approx 10^{-7}$) is considered to be
 1213 a discovery. The expected upper limit is a single-sided interval test statistic, associated with the POI in the
 1214 maximum-likelihood fit ($\mu(tWZ)$, in our case). In this context, the expected upper limit can be understood in the
 1215 following way: consider running an ensemble of MC toy experiments, each with their own confidence interval (a
 1216 range of possible values for $\mu(tWZ)$). An expected upper limit, at some fixed percentage $x\%$ (or *confidence level*),
 1217 can be determined from this ensemble. The expected upper limit tells us that, $x\%$ of the toy MC experiment's
 1218 confidence intervals will contain the true value of $\mu(tWZ)$. A commonly used percentage in particle physics is
 1219 95%, which we adopt for this analysis. In particle physics, this is referred to as the *CLs Method* [39]. The CLs
 1220 test statistic can be calculated 'brute force' by running these MC toy experiments, however this is very CPU
 1221 intensive. Asymptotic formulae are able to describe the underlining CLs test statistic distributions under certain
 1222 approximations [9]. Instead of running toy MC experiments, we use asymptotic formulae to perform the CLs
 1223 method, which reduces computation time from several millions of CPU hours to several CPU minutes. The
 1224 significance and upper limits which are calculated in this analysis are given a prefix of 'expected' to indicate that
 1225 these are results from a blinded analysis.

1226 In the separate and combined fits, the *mixed data and MC* fit setup [80] is used. This is done to obtain the
 1227 most accurate prediction of the expected results while keeping the signal regions blinded. For this setup, first a
 1228 background-only fit to the control regions using real data is done to determine estimates of the nuisance parameters.
 1229 Then these estimates are used to construct a modified ASIMOV dataset in the signal regions. Finally, the fit is
 1230 performed using real data in the control regions and the aforementioned modified ASIMOV data-set in the signal
 1231 regions.

1232 In these fits, the parameter of interest (POI) is μ_{tWZ} . The POI is ultimately the quantity which we wish to
 1233 measure and is set as a free parameter (unconstrained; can take any value in the fit). The nuisance parameters are
 1234 assigned to the systematic uncertainties outlined in Section 4.8. Furthermore, a gamma (γ) nuisance parameter
 1235 for a bin is added to the likelihood function if the statistical uncertainty in the bin exceeds 0.1% of its nominal value.

1236 Pruning is done per sample and per region on the shape and normalisation uncertainties for samples. A
 1237 sample's shape and normalisation nuisance parameter is pruned (removed from the limit/fit) if the fraction of
 1238 the expected number of signal events to the expected number of total events (signal + background) is less than 0.01.

1239 An auto-binning algorithm, `TransfoD` [29], was used to define the binning. This aims to maximise $\frac{signal}{background}$ in
 1240 each bin. Furthermore, it aims to avoid defining bins with a low number of events.

1246 4.10 Results

1247 In the section, an expected upper limit and an expected significance are set on the cross section of tWZ . This is
 1248 performed for the current analysis in the tetralepton channel as well as for a combined analysis across the trilepton
 1249 and tetralepton channels. The trilepton analysis was performed as an independent study by Benjamin Warren
 1250 (UCT) [81]. Note that throughout this section, all signal regions remain blinded.

1251 4.10.1 Tetralepton Channel

1252 In Figure 56 expected number of events for the variables used in the likelihood fit in each region are shown.

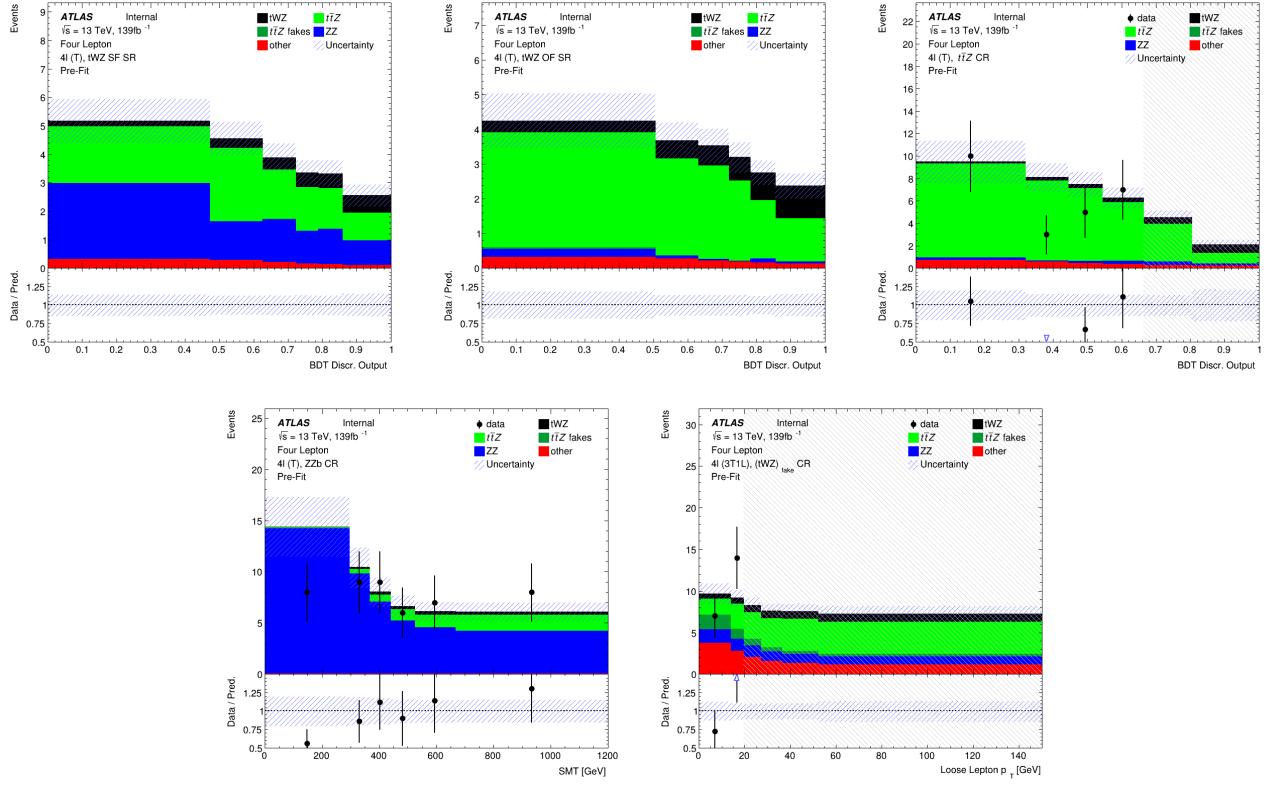


Figure 56: Pre-fit distributions (blinded) of variables used in the fit are shown. **Top left:** The event-level BDT Disc. Output in the tWZ SF (4T) SR region is shown. **Middle Top:** The event-level BDT Disc. Output in the tWZ SF (4T) SR region is shown. **Top right:** The event-level BDT Disc. Output in the $t\bar{Z}$ CR region is shown. **Bottom left:** $SMT = \sum p_T(\ell) + \sum p_T(jet) + E_T^{\text{miss}}$ in the ZZb CR is shown. **Bottom right:** p_T (loose lepton) in the $(tWZ)_{\text{fake}}$ (3T1L) CR is shown. The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. The plots in the tWZ OF SR and tWZ SF SR are kept blinded by omitting the data points.

1253 In Figure 57 expected number of events for the variables used in the likelihood fit in each region are shown.
In Table 12, the expected number of events (after the fit) for each sample in each region is shown.

	tWZ OF SR	tWZ SF SR	$t\bar{Z}$ CR	ZZb CR	$(tWZ)_{\text{fake}}$ CR
$t\bar{Z}$	13.2379 ± 1.52295	9.62061 ± 1.12291	29.9054 ± 3.60908	5.08899 ± 0.620121	18.5393 ± 2.23036
$t\bar{Z}$ fakes	0.0702522 ± 0.0468691	0.0334067 ± 0.0256903	0.0723509 ± 0.0418526	0.0485273 ± 0.029133	5.04378 ± 2.34732
tWZ	7.83414 ± 3.31679	5.33547 ± 2.24801	5.69373 ± 2.58041	2.89889 ± 1.23837	10.278 ± 4.3345
ZZ	0.481776 ± 0.119774	7.72372 ± 1.2351	1.07955 ± 0.182461	40.6067 ± 6.26078	6.86097 ± 1.11443
other	$6.00553e-06 \pm 3.02819e-06$	0.252557 ± 0.442116	0.273507 ± 0.223201	$6.00553e-06 \pm 3.02819e-06$	2.35427 ± 0.917112
	0.0827905 ± 0.0398773	0.0756107 ± 0.0354584	0.063585 ± 0.0293325	0.05884 ± 0.0244084	4.90963 ± 0.745354
	$0.00668643 \pm 0.00792217$	$0.00297948 \pm 0.00287361$	$6.00553e-06 \pm 3.02819e-06$	0.002306 ± 0.00564349	0.943182 ± 0.292148
	0.0442934 ± 0.024156	0.0396511 ± 0.0154282	0.0133471 ± 0.0128199	0.0472562 ± 0.0330315	1.83567 ± 0.392459
	$0.000987164 \pm 0.000766266$	$0.00247481 \pm 0.00136945$	0.0140869 ± 0.00479496	$6.00553e-06 \pm 3.02819e-06$	0.0100558 ± 0.00363157
	0.00934035 ± 0.0080554	0.0107458 ± 0.00849984	0.0571373 ± 0.0204011	$6.00553e-06 \pm 3.02819e-06$	0.0216809 ± 0.00992819
	0.0294618 ± 0.0263174	0.029771 ± 0.0195582	0.264364 ± 0.0926252	0.0129431 ± 0.032803	0.152448 ± 0.058965
	0.280643 ± 0.0853411	0.191433 ± 0.0586778	0.0697266 ± 0.0225059	0.171142 ± 0.0518102	0.266502 ± 0.0810194
$t\bar{H}$	0.846054 ± 0.175495	0.669375 ± 0.140107	1.96662 ± 0.401199	0.150025 ± 0.0353826	2.21264 ± 0.451623
Total	22.9243 ± 2.96284	23.9876 ± 2.11249	39.4734 ± 3.44937	49.0856 ± 6.04162	53.4282 ± 4.31683
data	-	-	36	49	57

Table 12: The expected number of events (after the fit) for each sample in each region is shown.

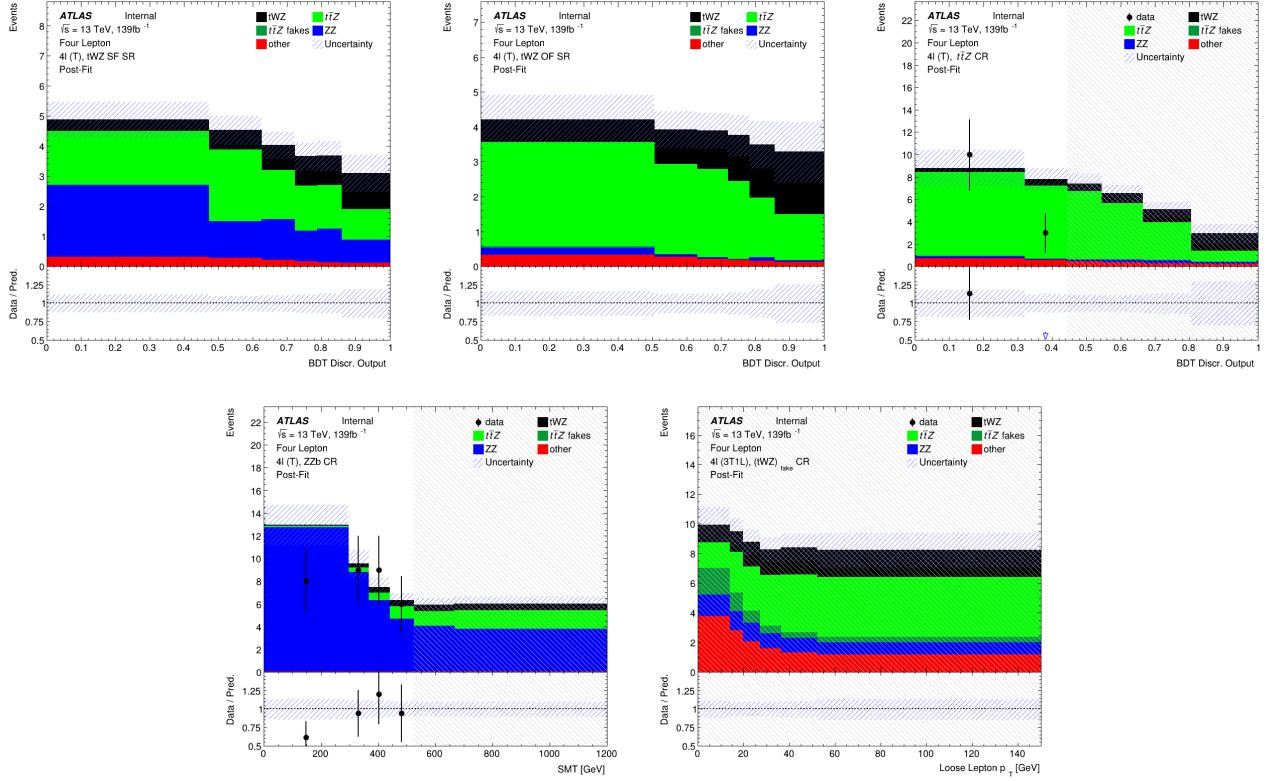


Figure 57: Post-fit distributions (blinded) of variables used in the fit are shown. **Bottom left:** $SMT = \sum p_T(\ell) + \sum p_T(jet) + E_T^{\text{miss}}$ in the ZZb CR is shown. **Bottom right:** p_T (loose lepton) in the $(tWZ)_{\text{fake}}$ (3T1L) CR is shown. The data is given by the black points and the MC predictions for each process are given by the filled histograms. The vertical lines on the data points represent the statistical uncertainty in the data and the diagonally-lined bands represent the total (statistical and systematic added in quadrature) uncertainty. The lower panel in each plot shows the ratios of the data to the theoretical predictions. The plots in the tWZ OF SR and tWZ SF SR are kept blinded by omitting the data points. **Top left:** The event-level BDT Discr. Output in the tWZ SF (4T) SR region is shown. **Middle Top:** The event-level BDT Discr. Output in the tWZ SF (4T) SR region is shown. **Top right:** The event-level BDT Discr. Output in the ttZ CR region is shown.

1255 The expected upper limit of tWZ in the tetralepton channel is measured as,

$$\mu_{up}^{exp} = 1.61^{+2.35}_{-1.16} \quad (4.15)$$

1256 The expected significance of tWZ in the tetralepton channel is measured as,

$$Z_\mu^{exp} = 1.44\sigma \quad (4.16)$$

1257 The best-fit value of the signal strength, $\mu(tWZ) = \frac{\sigma(tWZ)}{\sigma(tWZ)_{SM}}$, from the likelihood fit is measured as,

$$\mu(tWZ) = 1.91^{+0.95}_{-0.82} \quad (4.17)$$

1258 The expected upper limit is in agreement with the extracted best-fit value on the signal strength, therefore no
1259 deviations from the SM cross section of tWZ is observed. Neither the 3σ evidence nor the 5σ discovery standards
1260 are reached for the expected significance. This is not surprising, given the low amount of events present in the
1261 tetralepton channel.

1262 In Figure 58, a ranking plot showing the impact of the systematic uncertainties on the POI, $\mu(tWZ)$ is shown.

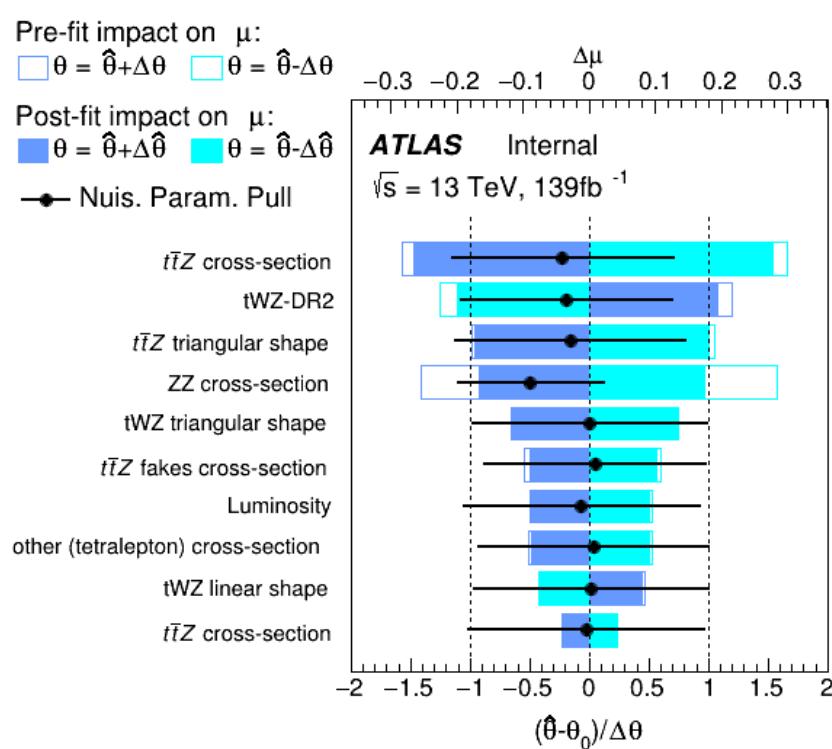


Figure 58: A ranking plot showing the impact (ordered from top to bottom via decreasing impact) of the systematic uncertainties (top 10) on the POI, $\mu(tWZ)$, in the tetralepton channel is shown. $\hat{\theta}$ is the best-fit value of the nuisance parameter. $\Delta\hat{\theta}$ and $\Delta\theta$ are the post-fit and pre-fit uncertainties respectively. The post-fit and pre-fit impact of each nuisance parameter on $\mu(tWZ)$ are shown with the solid and lined rectangles respectively. The empty and solid blue rectangles correspond to the pre-fit and post-fit impacts on $\mu(tWZ)$ respectively. These impacts are shown on the upper axis ($\Delta\mu$). On the lower axis, the nuisance parameter pull, $\frac{\hat{\theta} - \theta_0}{\Delta\theta}$, is shown (θ_0 is the nominal pre-fit value of the nuisance parameter). The nuisance parameter pull is indicated by the black points, with their relative post-fit errors ($\frac{\Delta\hat{\theta}}{\Delta\theta}$) shown by the black horizontal error bars.

1264 The most important systematics are the cross sections of $t\bar{t}Z$ and ZZ , and shape modelling on $t\bar{t}Z$ ($t\bar{t}Z$ triangular
1265 shape) and tWZ (tWZ -DR2 and tWZ triangular shape). The cross section of ZZ is significantly shifted down
1266 in the fit from its nominal value. The $t\bar{t}Z$ cross section, tWZ -DR2 and $t\bar{t}Z$ triangular shape nuisance parameters
1267 are similarly shifted down in the fit from their nominal values, but to a much lesser degree than the cross section
1268 of ZZ . These pulls are all within 1σ uncertainty and are thus relatively small. It is expected that the modelling
1269 uncertainties (shape and normalisations) of the most dominant backgrounds (e.g. $t\bar{t}Z$, ZZ) have relatively large
1270 impacts on $\mu(tWZ)$, since the uncertainty of the analysis is dominated by statistical uncertainty.

Nuisance Parameter	Channel Affected	
	trilepton	tetralepton
$\sigma(t\bar{t}H)$	✓	✓
$\sigma(t\bar{t}Z)$	✓	✓
$\sigma(WZ)$	✓	✓
$\sigma(tZq)$	✓	✓
$\sigma(ZZ)$	✓	✓
$\sigma(\text{other(trilepton)})$	✓	✗
$\sigma(\text{other(tetralepton)})$	✗	✓
$\sigma(t\bar{t}Z)_{\text{fakes}}$	✗	✓
$\sigma(t\bar{t})_{\text{fakes}}$	✓	✗
$\sigma(Z + \text{jets})_{\text{fakes}}$	✓	✗
Luminosity	✓	✓
jvt	✓	✓
pileup	✓	✓
DL1r SF (b jets)	✓	✓
DL1r SF (light jets)	✓	✓
$\sigma(tWZ - DR2)$	✗	✓
lepton SF	✓	✓
$t\bar{t}Z$ triangular shape	✓	✓
$t\bar{t}Z$ linear shape	✓	✓
$WZ + b$ triangular shape	✓	✗
$WZ + b$ linear shape	✓	✗
$WZ + c$ triangular shape	✓	✗
$WZ + c$ linear shape	✓	✗

Table 13: A summary of the nuisance parameters used in the combined fit is shown. The channels which are affected by each nuisance parameter are indicated with a ✓ (is affected) or a ✗ (is not affected).

4.10.2 Trilepton and Tetralepton Channels

In the section, an expected upper limit and an expected significance are set on the cross section of tWZ from the combined fit across all regions of tWZ in the tetralepton and trilepton channels.

The trilepton analysis follows a similar analysis strategy to that of the tetralepton analysis. It includes an event-level BDT which aims to discriminate between the tWZ and all background as well as an object-level BDT which aims to identify hadronically decaying W bosons to discriminate between tWZ and the large WZ background. One tWZ SR is defined and five CRs are defined. WZ and $t\bar{t}Z$ CRs are defined to constrain the dominant WZ and $t\bar{t}Z$ backgrounds. Three CRs which require that one of the three selected leptons are loose, are defined for WZ , $t\bar{t}Z$ and tWZ in order to constraint the fake lepton component (using the MC template method - similar to the method used in Section 4.6 to estimate the fake lepton component).

The expected upper limit of tWZ in the trilepton channel is measured as,

$$\mu_{up}^{exp} = 2.65^{+3.67}_{-1.91} \quad (4.18)$$

The expected significance of tWZ in the trilepton channel is measured as,

$$Z_{\mu}^{exp} = 0.75\sigma \quad (4.19)$$

The best-fit value of the signal strength, $\mu(tWZ) = \frac{\sigma(tWZ)}{\sigma(tWZ)_{SM}}$, from the likelihood fit is measured as,

$$\mu(tWZ) = 1.16^{+1.33}_{-1.30} \quad (4.20)$$

The expected upper limit is in agreement with the extracted best-fit value on the signal strength, therefore no deviations from the SM cross section of tWZ is observed.

In Table 13, the nuisance parameters used in the fit, including which channel's regions are affected by each, are shown.

The expected upper limit of tWZ across both channels is measured as,

$$\mu_{up}^{exp} = 1.43^{+2.04}_{-1.03} \quad (4.21)$$

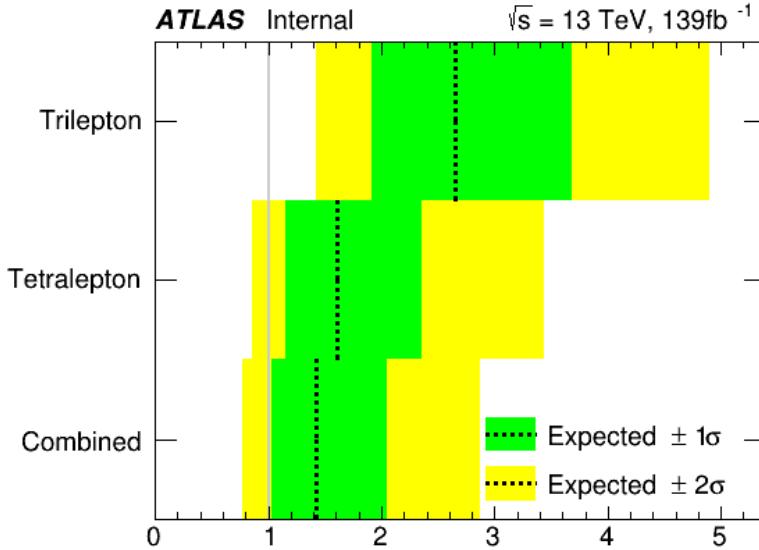


Figure 59: The expected upper limits of the trilepton channel, tetralepton channel and both channels combine are shown. The y-axis shows the channels in which the fitting procedure was performed. The expected limits are represented by the vertical dotted line. One- and two- σ uncertainty bands are shown in green and yellow respectively. The vertical grey line indicates when $\mu(tWZ) = 1$.

1292 The expected significance of tWZ across both channels is measured as,

$$Z_{\mu}^{exp} = 1.61\sigma \quad (4.22)$$

1293 The best-fit value of the signal strength, $\mu(tWZ) = \frac{\sigma(tWZ)}{\sigma(tWZ)_{SM}}$, from the likelihood fit is measured as,

$$\mu(tWZ) = 1.80^{+0.70}_{-0.65} \quad (4.23)$$

1294 In Figure 59, the expected upper limits of the trilepton channel, tetralepton channel and both channels combined
1295 are shown.

1296 It can be seen that the sensitivity of tWZ is mostly driven by the tetralepton analysis, with the trilepton analysis
1297 attributing a small decrease in the expected upper limit of the combined analysis, and its associated uncertainty.

1299 In Figure 60, the best-fit values of $\mu(tWZ)$ from the fit for the trilepton channel, tetralepton channel and both
1300 channels combined are shown.

1301 It can be seen from Figures 59 and 60 that the best fit value for the signal strength on tWZ , $\mu(tWZ)$, and the
1302 expected limits for the tri- and tetralepton channels are consistent with one-another (their uncertainties overlap).
1303 Therefore it is appropriate to combine these two analyses.

1304 In Figure 61, a ranking plot showing the impact of the systematic uncertainties on the POI, $\mu(tWZ)$, in the
1306 combined fit across both the tri- and tetralepton channels is shown.

1307 Some nuisance parameters are pulled down from their nominal pre-fit values, however these are all within 1 σ
1308 uncertainty and are thus relatively small. It is expected that the modelling uncertainties (shape and normalisations)
1309 of the most dominant backgrounds (e.g. $t\bar{t}Z$, ZZ , WZ) have relatively large impacts on $\mu(tWZ)$, since the
1310 uncertainty of the analysis is dominated by statistical uncertainty.

1312 The most important systematics are the cross sections of $t\bar{t}Z$, ZZ and $WZ + b$, and shape modelling on $t\bar{t}Z$
1313 ($t\bar{t}Z$ triangular shape). The aforementioned nuisance parameters are significantly shifted down in the fit from its
1314 nominal value. The less important systematics are pulled in the fit from their nominal values, but to a much lesser
1315 degree than the cross section of those mentioned above. It is expected that the modelling uncertainties (shape and
1316 normalisations) of the most dominant backgrounds (e.g. $t\bar{t}Z$, ZZ and $WZ + b$) have relatively large impacts on
1317 $\mu(tWZ)$, since the uncertainty of the analysis is dominated by statistical uncertainty.

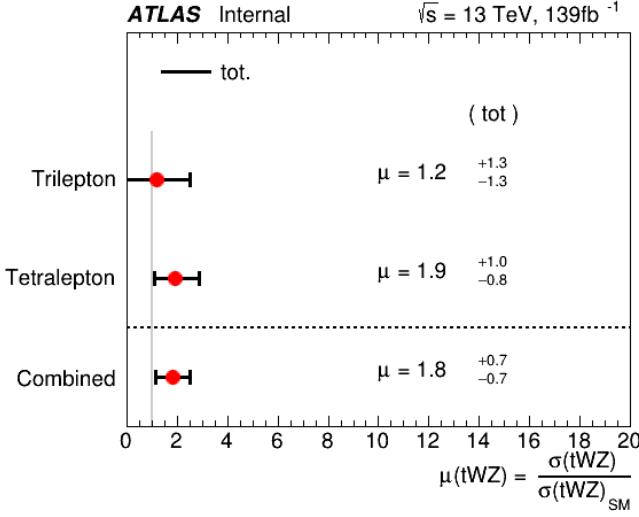


Figure 60: The best-fit values of $\mu(tWZ)$ from the fit for the trilepton channel, tetralepton channel and both channels combined are shown. The y-axis shows the channels in which the fitting procedure was performed. The signal strength $\mu(tWZ)$ is shown on the x-axis. The nominal signal strengths are represented by the red dots. The total uncertainty associated with the best-fit $\mu(tWZ)$ value is shown by the black error bars.

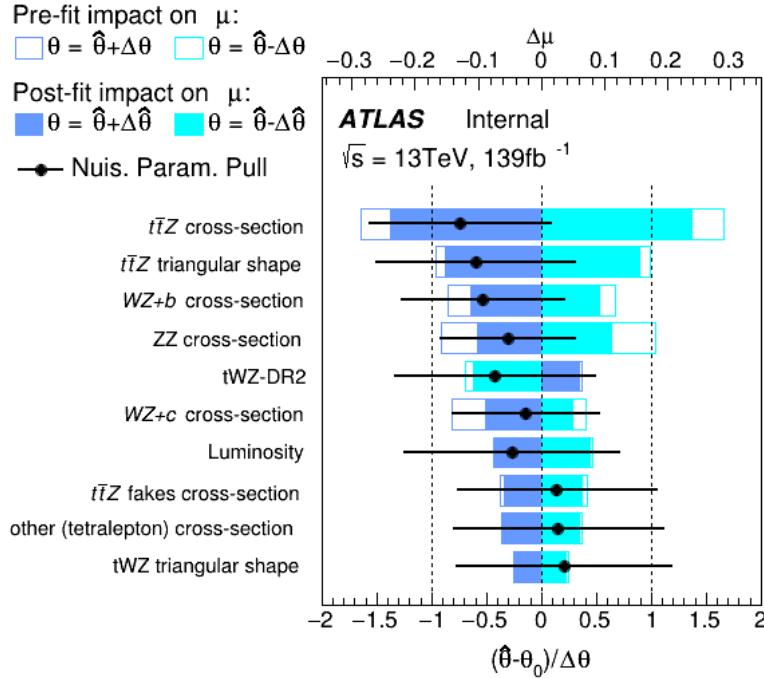


Figure 61: A ranking plot showing the impact of the systematic uncertainties (top 10) on the POI, $\mu(tWZ)$, in the combined fit across both the tri- and tetralepton channels is shown. $\hat{\theta}$ is the best-fit value of the nuisance parameter. $\Delta\hat{\theta}$ and $\Delta\theta$ are the post-fit and pre-fit uncertainties respectively. The post-fit and pre-fit impact of each nuisance parameter on $\mu(tWZ)$ are shown with the solid and lined rectangles respectively. The empty and solid blue rectangles correspond to the pre-fit and post-fit impacts on $\mu(tWZ)$ respectively. These impacts are shown on the upper axis ($\Delta\mu$). On the lower axis, the nuisance parameter pull, $\frac{\hat{\theta} - \theta_0}{\Delta\theta}$, is shown (θ_0 is the nominal pre-fit value of the nuisance parameter). The nuisance parameter pull is indicated by the black points, with their relative post-fit errors ($\frac{\Delta\hat{\theta}}{\Delta\theta}$) shown by the black horizontal error bars.

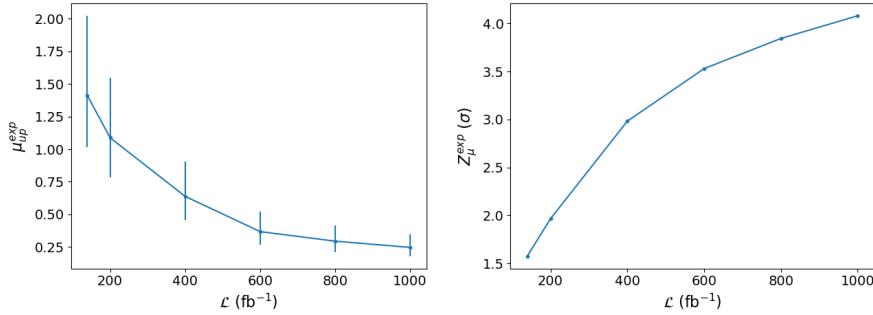


Figure 62: The expected upper limit (left) and significance (right) for combined fully blinded fits to the ASIMOV dataset across both the tri- and tetralepton channels for a range integrated luminosities are shown. **Left:** The integrated luminosity, \mathcal{L} , is shown on the x-axis. The expected upper limit, μ_{up}^{exp} , is shown on the y-axis. The vertical lines represent the total uncertainty ($\pm 1 \sigma$) on the expected upper limit. **Right:** The integrated luminosity, \mathcal{L} , is shown on the x-axis. The expected significance, Z_μ^{exp} , is shown on the y-axis. The vertical lines represent the total uncertainty ($\pm 1 \sigma$) on the expected significance.

1318 4.10.2.1 Projection to Higher Luminosity

1319 In this section we apply a fully blinded fit to the ASIMOV dataset for integrated luminosities larger than the 139fb^{-1}
 1320 currently available from the ATLAS Full Run 2 dataset. This study gives us insight into the sensitivity of tWZ
 1321 which we could expect if we were to replicate the current analysis, given more data. Given the upgrades planned
 1322 for the LHC and the ATLAS detector, we will soon expect a large increase in available pp collision data (reaching
 1323 $\mathcal{L} = 3000\text{fb}^{-1}$ [32]). This study therefore gives us an idea of the sensitivity of tWZ which we can expect in the future.
 1324

1325 In Figure 62, the expected upper limit and significance for combined fully blinded fits to the ASIMOV dataset
 1326 across both the tri- and tetralepton channels for a range integrated luminosities are shown.

1327 As we increase the total integrated luminosity, the sensitivity of tWZ increases. The increase in sensitivity is shown
 1328 by the decrease of the expected upper limit and the increase of the expected significance with increasing integrated
 1329 luminosity. It can be noted that a 3σ expected significance is achieved at $\mathcal{L} \sim 400 \text{ fb}^{-1}$, indicating evidence for
 1330 observation of tWZ . The 5σ significance 'discovery' standard is above our 1000fb^{-1} luminosity range of study.
 1331 Furthermore, these plots show that the sensitivity of tWZ is hindered by the low amount of events we observe,
 1332 rather than the systematic effects. This is evident since increasing the luminosity, therefore increasing the number
 1333 of events we see in the detector, directly causes a steady increase in sensitivity.

1334

Chapter 5

1335

Conclusion

1336 The search for tWZ production using 139 fb^{-1} of pp collision data at a centre-of-mass energy of $\sqrt{s} = 13 \text{ TeV}$,
 1337 recorded by the ATLAS experiment at CERN, has been presented. This thesis targeted the tetralepton final state
 1338 channel. To further increase the sensitivity of tWZ , a combined analysis was done across the tetralepton and
 1339 trilepton (studied in an independent analysis by Benjamin Warren (UCT) [81]) channels.

1340 Two SRs and Three CRs were defined. Two SRs, instead of one, were defined in order to suppress and constrain
 1341 the ZZ background. The definition of the tWZ OF SR and the tWZ SF SR take advantage of the OSSF lepton
 1342 pairs which originate from the decay of a Z boson, by requiring that non- Z leptons in the event have the opposite
 1343 and same flavours, for the tWZ OF SR and tWZ SF SR respectively. The resulting SRs successfully separate the
 1344 ZZ background, with the tWZ OF SR containing around 6% of the total ZZ background yield across both regions
 1345 (implying that the remaining $\sim 94\%$ is contained in the tWZ SF SR). The dominant background processes, $t\bar{t}Z$
 1346 and ZZ were constrained by the definition of $t\bar{t}Z$ and ZZ CRs, respectively. The dominant source of fake leptons,
 1347 originating from the $t\bar{t}Z$ background, was constrained by the $(tWZ)_{fake}$ CR, using the MC template method.
 1348

1349 Two BDTs were implemented: an object-level BDT which aims to classify between ℓb systems coming from top
 1350 quarks and an event-level BDT which aims to discriminate between tWZ and our major backgrounds, $t\bar{t}Z$ and
 1351 ZZ . The output from the object-level BDT was converted to an event-level variable to be used as input to
 1352 the event-level BDT. A kinematic reconstruction algorithm, $2\nu\text{SM}$, was used to reconstruct top quarks in order
 1353 to discriminate between tWZ and $t\bar{t}Z$. The output from this algorithm was used as an input variable to the
 1354 event-level BDT. The trained BDT was shown to discriminate well between signal and background events.
 1355

1356 Using a modified ASIMOV dataset in the SRs and real data in the CRs, a blinded maximum-likelihood fit was
 1357 performed across all regions in the tetralepton channel. The best-fit value of the signal strength in the tetralepton
 1358 channel was,

$$\mu(tWZ) = 1.91^{+0.95}_{-0.82} \quad (5.1)$$

1360 with an expected significance of 1.44σ . The expected upper limit on the signal strength of tWZ in the tetralepton
 1361 channel was,

$$\mu_{up}^{exp} = 1.61^{+2.35}_{-1.16} \quad (5.2)$$

1362 Furthermore, a blinded maximum-likelihood fit was performed across all regions across the trilepton and tetralepton
 1363 channels. The best-fit value of the signal strength across both the trilepton and tetralepton channels were,

$$\mu(tWZ) = 1.80^{+0.70}_{-0.65} \quad (5.3)$$

1364 with an expected significance of 1.61σ . The expected upper limit on the signal strength of tWZ across both the
 1365 trilepton and tetralepton channels were,

$$\mu_{up}^{exp} = 1.43^{+2.04}_{-1.03} \quad (5.4)$$

1366 Although this result does not satisfy the 3σ evidence nor the 5σ discovery standards, this is the tightest ever
 1367 constraint on the tWZ process. The results in this analysis are heavily statistically limited, it is therefore expected
 1368 that future analyses of this process, using larger datasets (such as that from the HL-LHC), would significantly
 1369 improve the results.

1370

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