

Öving 6, Matte 4K

Rendell Cale, grupp 2

12, 11

OK

$$3) u(x,t) = \cos(4t) \sin(2x)$$

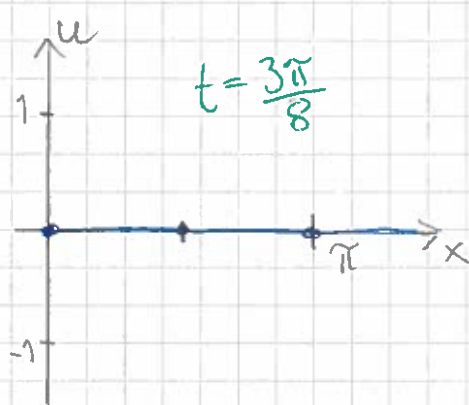
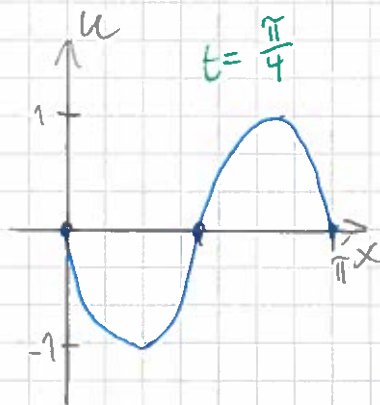
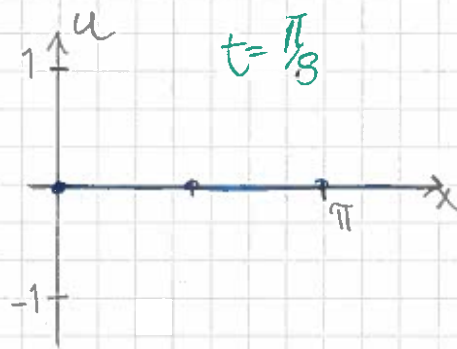
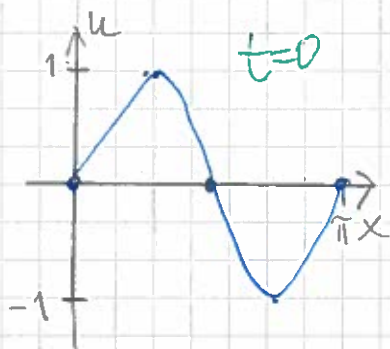
We verify that u solves the wave equation by solving for c .

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

$$-16 \cos(4t) \sin(2x) = c^2 (-4 \cos(4t) \sin(2x))$$

$$\Leftrightarrow c^2 = 4$$

$$\Leftrightarrow |c| = 2 \quad \checkmark$$



$$1) u(x,t) = e^{-\pi^2 t} \cos(25x)$$

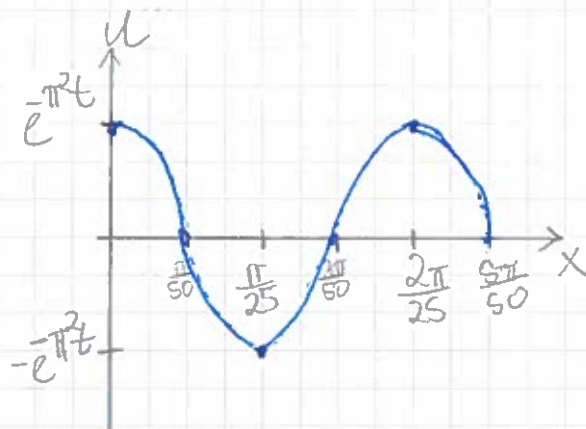
$$\frac{\partial u}{\partial t} = -\pi^2 e^{-\pi^2 t} \cos(25x)$$

$$\frac{\partial u}{\partial x^2} = -25^2 e^{-\pi^2 t} \cos(25x)$$

Heat equation:

$$-\pi^2 e^{-\pi^2 t} \cos(25x) = -25^2 e^{-\pi^2 t} \cos(25x) c^2$$

$$\Leftrightarrow |c| = \frac{\pi}{25} \quad \checkmark$$



$$15) \quad u(x,y) = a \ln(x^2+y^2) + b$$

$$\begin{aligned} (u_x)_x &= \left(\frac{2ax}{x^2+y^2} \right)_x \\ &= \frac{(x^2+y^2)2a - 2x \cdot 2ax}{(x^2+y^2)^2} \\ &= \frac{2ax^2 + 2ay^2 - 4ax^2}{(x^2+y^2)^2} \\ &= 2a \frac{y^2 - x^2}{(x^2+y^2)^2} \end{aligned}$$

Due to symmetry, we get

$$\begin{aligned} u_{yy} &= 2a \frac{x^2 - y^2}{x^2 + y^2} \\ \hookrightarrow &= -u_{xx} \end{aligned}$$

$$\Leftrightarrow u_{yy} + u_{xx} = 0$$

So u satisfies Laplace's equation.

Want to find a, b such that

$$110 = a \ln(1) + b \quad (1)$$

$$0 = a \ln(100) + b \quad (2)$$

$$(1) \Rightarrow b = 110$$

$$(2) \Rightarrow a = \frac{-b}{\ln(100)} = \frac{-110}{\ln(100)} \approx -23.89$$

$$a = \frac{-110}{\ln(100)} \text{ and } b = 110 \text{ are the}$$

desired constants.

12, 3:

1) The fundamental mode is given by

$$u_1(x,t) = (B_1 \cos(\lambda_1 t) + B_1^* \sin(\lambda_1 t)) \sin\left(\frac{\pi}{L}x\right)$$

$$\text{where } \lambda_1 = \frac{c\pi}{L}, \quad c^2 = \frac{T}{\rho}$$

The freq. of the fundamental mode decreases when L increases because λ_1 decreases.

When the mass increases (per unit length); c decreases so the frequency also decreases.

If the tension is doubled, then c is increased by a factor of $\sqrt{2}$, and thus the frequency is increased by the same ($\sqrt{2}$) factor.

The contrabass has to produce lower notes (lower freq.) so it has thicker and longer strings.

$$7) u(x, 0) = Kx(1-x)$$

$$L=1, c^2=1, u_x(x, 0) = 0$$

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2}$$

Since the string is fastened at the end-points, we have

$$u(0, t) = 0, u(1, t) = 0$$

Using separation of variables we have

$$u_n(x, t) = F_n(x) G_n(t)$$

where

$$F_n(x) = \sin(n\pi x)$$

$$\text{and } G_n(t) = B_n \cos(n\pi t) + B_n^* \sin(n\pi t)$$

$$\text{where } B_n = 2 \int_0^1 Kx(1-x) \sin(n\pi x) dx$$

$$\text{and } B_n^* = 2 \int_0^1 u_x(x, 0) \sin(n\pi x) dx = 0$$

$$B_n = \frac{-2K}{\pi n} \left[\frac{x(1-x)}{L} \cos(n\pi x) \right]_0^1 + \frac{2K}{\pi n} \int_0^1 (-2x+1) \cos(n\pi x) dx = 0$$

$$\begin{aligned}
 \Rightarrow B_n &= \frac{2K(-2x+1)\sin(\pi x)}{\pi^2} \Big|_0^1 + \frac{4K}{\pi^2} \int_0^1 \sin(\pi x) dx \\
 &= -\frac{4K}{\pi^3} \left[\cos(\pi x) \right]_0^1 \\
 &= -\frac{4K}{\pi^3} ((-1)^n - 1)
 \end{aligned}$$

n even gives $B_n = 0$

n odd gives $B_n = \frac{8K}{n^3\pi^3}$

$$\text{So } G_{2m+1}(t) = \frac{8K}{(2m+1)^3\pi^3} \cos(\pi t), \quad m=0,1,2,\dots$$

This gives

$$u_{2m+1}(x,t) = F_{2m+1}(x) G_{2m+1}(t)$$

$$= \sin((2m+1)\pi x) \cdot \frac{8K}{(2m+1)^3\pi^3} \cos((2m+1)\pi t)$$

which means

$$\begin{aligned}
 u(x,t) &= \frac{8K}{\pi^3} \left[\frac{\sin(\pi x)\cos(\pi t)}{1^3} + \frac{\sin(3\pi x)\cos(3\pi t)}{3^3} \right. \\
 &\quad \left. + \frac{\sin(5\pi x)\cos(5\pi t)}{5^3} \dots \right]
 \end{aligned}$$

15) ~~non-homogeneous~~

$$\frac{\partial^2 u}{\partial t^2} = -c^2 \frac{\partial^4 u}{\partial x^4}, \quad c^2 = \frac{EI}{\rho A} \quad (21)$$

Assume $u(x,t) = F(x) \cdot G(t)$, then

$$\frac{\partial^2 u}{\partial t^2} = F(x) \ddot{G}(t)$$

$$\frac{\partial^4 u}{\partial x^4} = F^{(4)}(x) G(t)$$

Subst. into (21) gives

$$F(x) \ddot{G}(t) = -c^2 F^{(4)}(x) G(t)$$

$$\Leftrightarrow \frac{F^{(4)}(x)}{F(x)} = -\frac{\ddot{G}(t)}{c^2 G(t)}$$

Since $\frac{F^{(4)}(x)}{F(x)}$ depends only on x and $-\frac{\ddot{G}(t)}{c^2 G(t)}$

depends only on t , both sides must be the same constant, call it K .

$$\frac{F^{(4)}(x)}{F(x)} = -\frac{\ddot{G}(t)}{c^2 G(t)} = K.$$

Want to show $K \geq 0$.

If $K < 0$ then $K = -\beta^4$ for some $\beta \in \mathbb{R}$.

$$\Rightarrow \ddot{G}(t) = \beta^4 c^2 G(t) = (\beta c)^2 G(t)$$

$$\Rightarrow G(t) = e^{\beta^2 c^2 t} + b e^{-\beta^2 c^2 t}$$

$$\Rightarrow |G| \rightarrow \infty \text{ as } |t| \rightarrow \infty$$

so G is unbounded (above).

$$\frac{F^{(4)}}{F} = -\beta^4$$

$$\Leftrightarrow F^{(4)} = -\beta^4 F$$

$$\Rightarrow \mathcal{L}\{F^{(4)}\} = -\beta^4 \mathcal{L}\{F\} \quad , \quad V(s) = \mathcal{L}\{F\}(s)$$

$$s^4 V - s^3 F(0) - s^2 F'(0) - s F''(0) - F'''(0) = -\beta^4 V$$

Assume $k > 0$, then $k = \beta^4$ for some β .

This gives two equations:

$$(1) \ddot{G}(t) = -\beta^4 c^2 G(t)$$

$$(2) F^{(4)}(x) = \beta^4 F(x)$$

$$(1) \text{ has solution } G(t) = a \cos(\beta^2 c t) + b \sin(\beta^2 c t)$$

(2) is an ODE of order 4, so we need four linearly independent solutions.

It is quite clear that $\cos \beta x$, $\sin \beta x$, $\cosh \beta x$ and $\sinh \beta x$ are four candidates and since they are linearly indep. we get

$$F(x) = A \cos(\beta x) + B \sin(\beta x) + C \cosh(\beta x) + D \sinh(\beta x)$$

16)

$$\frac{\partial^2 u}{\partial t^2} = -c^2 \frac{\partial^2 u}{\partial x^2} \quad (*)$$

$$\begin{cases} u_t(x, 0) = 0 & (1) \end{cases}$$

$$\begin{cases} u(0, t) = 0 & (2) \end{cases}$$

$$\begin{cases} u(L, t) = 0 & (3) \end{cases}$$

$$\begin{cases} u_{xx}(0, t) = 0 & (4) \end{cases}$$

$$\begin{cases} u_{xx}(L, t) = 0 & (5) \end{cases}$$

$$\begin{cases} u_n = F_n(x) G_n(t) \\ F_n(x) = A_n \cos \beta_n x + B_n \sin \beta_n x + C_n \cosh(\beta_n x) + D_n \sinh(\beta_n x) \\ G_n(t) = a_n \cos(c \beta_n^2 t) + b_n \sin(c \beta_n^2 t) \end{cases}$$

$$(1): u_t(x, 0) = 0 \Leftrightarrow F_n(x) \dot{G}_n(0) = 0$$

$$\Rightarrow \dot{G}_n(0) = 0$$

$$\Leftrightarrow \left[-a_n c \beta_n^2 \sin(c \beta_n^2 t) + b_n c \beta_n^2 \cos(c \beta_n^2 t) \right]_{t=0} = 0$$

$$\Leftrightarrow b_n = 0$$

$$\text{So } G_n(t) = a_n \cos(c \beta_n^2 t)$$

$$(2): u(0, t) = 0 \Leftrightarrow F_n(0) G(t) = 0$$

$$\Rightarrow F_n(0) = 0$$

$$\Leftrightarrow A_n + C_n = 0$$

$$(4): u_{xx}(0, t) = 0 \Leftrightarrow F_n''(0) G(t) = 0$$

$$\Rightarrow F_n''(0) = 0$$

$$\Leftrightarrow -A_n + C_n = 0$$

$$\text{So } A_n = C_n = 0$$

$$\text{So } F_n(x) = B_n \sin(\beta_n x) + D_n \sinh(\beta_n x)$$

$$(3): u(L, t) = F(L) G(t) = 0$$

$$\Rightarrow F(L) = 0$$

$$\Leftrightarrow B_n \sin(\beta_n L) + D_n \sinh(\beta_n L) = 0 \quad (3*)$$

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$$(5): u_{xx}(L, t) = F''(L) G(t) = 0$$

$$\Rightarrow F''(L) = 0$$

$$\Leftrightarrow -\beta_n^2 B_n \sin(\beta_n L) + \beta_n^2 D_n \sinh(\beta_n L) = 0$$

$$\Leftrightarrow -B_n \sin(\beta_n L) + D_n \sinh(\beta_n L) = 0 \quad (5*)$$

Assume $B_n \neq 0$, then

$$(3*) - (5*):$$

$$2B_n \sin(\beta_n L) = 0$$

$$\Rightarrow \beta_n = \frac{n\pi}{L}, \quad n = 0, 1, 2, \dots$$

$$(3*) + (5*): D_n = 0$$

$$\text{So } F(x) = B_n \sin\left(\frac{n\pi}{L}x\right)$$

$$\text{and } u_n(x, t) = B_n \sin\left(\frac{n\pi}{L}x\right) a_n \cos\left(\frac{cn\pi}{L}t\right)$$

$$= c_n \sin\left(\frac{n\pi}{L}x\right) \cos\left(\frac{cn\pi}{L}t\right)$$

$$, \quad c_n = B_n a_n$$

$$17) u_n(x,t) = c_n \sin\left(\frac{n\pi}{L}x\right) \cos\left(c\frac{n\pi}{L}t\right)$$

$$\Rightarrow u(x,t) = \sum_{n=0}^{\infty} c_n \sin\left(\frac{n\pi}{L}x\right) \cos\left(c\frac{n\pi}{L}t\right)$$

$$u(x,0) = f(x) = x(L-x)$$

$$= \sum_{n=0}^{\infty} c_n \sin\left(\frac{n\pi}{L}x\right)$$

$$\Rightarrow c_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L}x\right) dx$$

$$= \frac{2}{L} \int_0^L x(L-x) \sin\left(\frac{n\pi}{L}x\right) dx$$

$$= \frac{2}{L} \cdot \frac{L}{n\pi} \int_0^L (L-2x) \cos\left(\frac{n\pi}{L}x\right) dx$$

$$= \frac{4L}{n^2\pi^2} \int_0^L \sin\left(\frac{n\pi}{L}x\right) dx$$

$$= \frac{4L^2}{n^3\pi^3} (\cos(n\pi) - 1)$$

$$= \frac{4L^2}{n^3\pi^3} ((-1)^n - 1)$$

$$= \frac{-8L^2}{(2m+1)^3\pi^3}, \quad m=0,1,2,\dots$$

$$\Rightarrow u(x,t) = -\frac{8L^2}{\pi^3} \sum_{m=0}^{\infty} \frac{1}{(2m+1)^3} \cdot \sin\left(\frac{(2m+1)\pi}{L}x\right) \cos\left(c\frac{(2m+1)\pi}{L}t\right)$$

12,6:

g) $L = 10 \text{ cm}$

$A = 1 \text{ cm}^2$

$\rho = 10,6 \text{ g/cm}^3$

$\sigma = 0,056 \text{ cal/(g}^\circ\text{C)}$

$K = 1,04 \text{ cal/(cm sec }^\circ\text{C)}$

$$C^2 = \frac{K}{\sigma \rho} = \frac{1,04}{0,056 \cdot 10,6} \left[\frac{\text{cal/(cm sec }^\circ\text{C)}}{\frac{\text{cal/(g }^\circ\text{C)}}{\text{g/cm}^3}} \right]$$
$$= 1,75 \frac{1}{\text{cm}^2 \text{ sec}}$$

$\Rightarrow C = 1,32$

$u(x, 0) = f(x) = \sin(0,1\pi x)$

$$u(x, t) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi}{L}x\right) e^{-\lambda_n^2 t}$$

$t=0: u(x, 0) = f(x) \Rightarrow \frac{n}{L} = 0,1$
 $\Leftrightarrow n = 1$

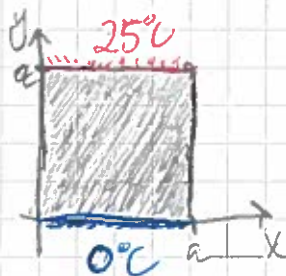
So $B_1 = 1$ and all other $B_n = 0$,

$\lambda_1^2 = \left(\frac{0,1\pi}{L}\right)^2 = (0,41)^2 = 0,17$

So

$$u(x, t) = \sin(0,1\pi x) e^{-0,17t}$$

21)



- $u(x, y) = \lim_{t \rightarrow \infty} u(x, y, t)$
- $u(x, 0, t) = 0$
- $u(x, a, t) = 25$

As $t \rightarrow \infty$, the heat distribution should not depend on the x -coordinate, so

$$u(x, y) = G(y)$$

Since $u_{xx} + u_{yy} = 0$, we get

$$\frac{d^2 G}{dy^2} = 0$$

$$\Rightarrow G(y) = my + b$$

$$b = 0 \text{ since } G(0) = b = 0$$

$$m = \frac{25}{a} \text{ since } G(a) - ma = 25$$

So $u(x, y) = \frac{25}{a} y$, and since we have $a = 24$, we get

$$u(x, y) = \frac{25}{24} y$$

Sup. N

$$|e^{z_0}| = 5$$

$$\begin{aligned} |e^{2z_0+3i}| &= |(e^{z_0})^2 e^{3i}| \\ &= |e^{z_0}|^2 \text{ since } |e^{3i}| = 1 \\ &= 5^2 \\ &= \underline{\underline{25}} \end{aligned}$$

Sup. O

a) The superposition principle holds for all linear homogeneous PDE, and since both $(*)$ and $(**)$ are that, it holds for both of them. So $u_1 + u_2$ and $u_1 - u_2$ solves $(*)$ and $(**)$.

b) Let $v(x,t) = u(x,t) - (a + (b-a)x)$ solve $(*)$.

$$\text{Then } \underline{v(0,t)} = u(0,t) - a \stackrel{(*)}{=} \underline{0}$$

$$\text{and } \underline{v(1,t)} = u(1,t) - (a + b - a) = u(1,t) - b \stackrel{(*)}{=} \underline{0}$$

$$\text{So } v(0,t) = 0, \text{ and } v(1,t) = 0$$

Since $v_t - v_{xx} \stackrel{(*)}{=} 0$, $v(x,t)$ satisfies all conditions of $(**)$ so it is a solution.

$$v(x,t) = F(x)G(t)$$

$$\text{Then } v_t - v_{xx} = F(x)\dot{G}(t) - F''(x)G(t) = 0$$

$$\Leftrightarrow \frac{F''}{F} = \frac{\dot{G}}{G} = -p^2 \quad (\text{proof omitted})$$

$$\Leftrightarrow F'' = -p^2 F$$

$$\dot{G} = -p^2 G$$

$$\Rightarrow F(x) = A \cos px + B \sin px$$

$$\frac{dG}{dt} = -p^2 G$$

$$\Leftrightarrow \frac{dG}{G} = -p^2 dt$$

$$\Leftrightarrow \ln|G| = -p^2 t + C$$

$$G(t) = C e^{-p^2 t}$$

Using $v(0,t) = 0$ we get

$$F(0)G(t) = 0$$

$$\Rightarrow F(0) = A = 0$$

$$\Rightarrow F(x) = B \sin(px)$$

Using $v(1,t) = 0$ we get

$$F(1) = 0$$

$$\Rightarrow B \sin p = 0$$

$$\Rightarrow p_n = n\pi, \text{ assuming } B \neq 0$$

$$n = 1, 2, 3, \dots$$

$$\text{So } v_n(x,t) = B_n \sin(p_n x) e^{-p_n^2 t},$$

where the constant C from G is combined into B_n .

The general solution is then

$$\begin{aligned} v(x,t) &= \sum v_n \\ &= \sum_{n=1}^{\infty} C_n \sin(p_n x) e^{-p_n^2 t}, \quad p_n = n\pi \end{aligned}$$

if $v(x,0) = f(x)$ then

$$\sum_{n=1}^{\infty} C_n \sin(p_n x) = f(x) \text{ so the } C_n \text{'s}$$

are given by the Fourier expansion of f .

$$C_n = \int_0^1 f(x) \sin(n\pi x) dx$$

c) If $v(x,t)$ is a solution of $(*)$, then

$$u(x,t) = v(x,t) + (a - (b-a)x), \quad x \in [0,1], t \geq 0$$

is a solution of $(*)$.

We've shown that

$$v(x,t) = \sum_{n=1}^{\infty} C_n \sin(p_n x) e^{-p_n^2 t}, \quad p_n = n\pi$$

$$C_n = \int_0^1 u(x,0) \sin(n\pi x) dx$$

In our case $a=1, b=1$.

Since $u(x,0) = \sin(\pi x)$, all $C_n = 0, n \neq 1$ and

$$C_1 = 1$$

$$\text{So } v(x,t) = u_1(x,t) = \sin(\pi x) e^{-\pi^2 t}$$

$$\text{and } \underline{u(x,t) = \sin(\pi x) e^{-\pi^2 t} - (1+2x)}$$