

# Decoherence-Triggered Collapse (DTC): A Simulation-Efficient Objective Collapse Model Without Microscopic Noise

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I present the Decoherence-Triggered Collapse (DTC) model, a deterministic objective-collapse framework that resolves the quantum measurement problem without stochastic noise. DTC modifies the open-system master equation by adding a state-dependent pruning term that activates only when off-diagonal coherence falls below the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ , a scale fixed by environmental scattering rates. In the macroscopic regime, DTC predicts instantaneous, tail-free localization and the irreversible failure of spin-echo (Lazarus) revival protocols once the coherence threshold is crossed, because the non-actualized branch is physically eliminated rather than merely entangled with the environment.

## I. INTRODUCTION

The reconciliation of unitary quantum evolution with the definite outcomes of macroscopic experience remains a central open problem in the foundations of physics [1, 2]. Environmental decoherence accounts remarkably well for the rapid suppression of interference in macroscopic systems [1, 3, 4], yet it does not by itself select a unique outcome: the reduced density matrix becomes diagonal in the pointer basis on extremely short timescales, but off-diagonal coherences are rendered practically unobservable rather than strictly eliminated. This leaves an unresolved question of why only one branch appears to be actualised.

Objective collapse theories close this gap by modifying the standard dynamical law so that superpositions of macroscopically distinct states are irreversibly reduced to a single outcome. Two developed models — the Ghirardi–Rimini–Weber (GRW) theory and Continuous Spontaneous Localisation (CSL) — introduce a universal stochastic noise field coupled to the mass density. Although phenomenologically successful, this noise is postulated ad hoc, acts even on perfectly isolated microscopic systems, and produces spontaneous heating and momentum diffusion, which are tightly constrained by experiment.

I propose the Decoherence-Triggered Collapse (DTC) model, a deterministic objective-collapse framework in which collapse is not driven by an independent noise term but emerges directly from environmental decoherence itself. Specifically, once off-diagonal coherence in the environment-selected pointer basis drops below a physically motivated irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$  — a scale set by typical environmental scattering rates — a pruning mechanism activates and instantaneously eliminates the non-actualised branches. For isolated or microscopic systems the threshold is never crossed, and the evolution remains exactly unitary with zero microscopic noise.

## II. THE DTC MODEL

The evolution of the density operator  $\rho$  is governed by the modified Lindblad master equation

$$\dot{\rho} = -i[H, \rho] + \sum_k \gamma_k \mathcal{D}[L_k] \rho + \Gamma_{\text{trigger}}(\rho) \sum_n \mathcal{D}[P_n] \rho, \quad (1)$$

where the standard Lindblad dissipator is defined as

$$\mathcal{D}[A]\rho \equiv A\rho A^\dagger - \frac{1}{2}\{A^\dagger A, \rho\}. \quad (2)$$

The individual terms are defined as follows:

- $H$  is the system Hamiltonian.
- $L_k$  are the usual Lindblad operators describing environmental decoherence (e.g. dephasing, photon emission, collisional processes), with corresponding rates  $\gamma_k \geq 0$ .
- $P_n = |n\rangle\langle n|$  are orthogonal projectors onto the environment-selected pointer basis  $\{|n\rangle\}$ , satisfying  $\sum_n P_n = I$  and  $P_n P_m = \delta_{nm} P_n$ .

### A. Collapse trigger mechanism

The pruning term is controlled by the state-dependent rate

$$\Gamma_{\text{trigger}}(\rho) = \Gamma_0 \Theta(C_{\text{irr}} - C(\rho)), \quad (3)$$

where  $\Theta(x)$  is the Heaviside step function,  $\Gamma_0 \rightarrow \infty$  ensures instantaneous collapse in the ideal theory, and  $C_{\text{irr}} \simeq 10^{-20}$  is the irreversibility threshold.

For numerical stability the Heaviside function is replaced by the smooth logistic form

$$\Gamma_{\text{trigger}}(\rho) = \Gamma_0 [1 + \exp(\kappa(C(\rho) - C_{\text{irr}}))]^{-1}, \quad (4)$$

with steepness  $\kappa \gtrsim 10^{22}$  and  $\Gamma_0 \gtrsim 10^{25} \text{ s}^{-1}$ . In the ideal limit ( $\Gamma_0 \rightarrow \infty, \kappa \rightarrow \infty$ ) the final state is exactly tail-free.

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TABLE I. Fundamental parameter definitions and typical orders of magnitude for the DTC model.

Parameter	Physical Meaning	Simulation Value
$\Gamma_0$	Asymptotic pruning limit	$\rightarrow \infty (> 10^{25})$
$C_{\text{irr}}$	Irreversibility threshold	$10^{-20}$
$\kappa$	Trigger steepness	$\rightarrow \infty (> 10^{20})$

## B. Coherence measures

The trigger depends on the total off-diagonal coherence  $C(\rho)$ . The exact measure is the  $l_1$ -norm of the off-diagonal elements in the pointer basis:

$$C_{l1}(\rho) = \sum_{n \neq m} |\rho_{nm}|. \quad (5)$$

For large Hilbert-space dimension we use the off-diagonal purity as a computationally efficient proxy:

$$C(\rho) = \sqrt{\text{Tr}(\rho^2) - \sum_n (\rho_{nn})^2}, \quad (6)$$

which decreases monotonically with  $C_{l1}(\rho)$  in all decoherence-dominated regimes, ensuring the trigger activates only when coherence vanishes.

## C. Physical origin of the irreversibility threshold

The irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$  is fixed by environmental scattering rates. For a  $\sim 1 \mu\text{m}$  dust grain interacting with ambient atmospheric molecules, the localisation rate reaches  $\Lambda \simeq 10^{32} \text{ s}^{-1} \text{ m}^{-2}$ . For spatial separations  $\Delta x \sim 1 \mu\text{m}$ , this yields a decoherence timescale

$$\tau_{\text{dec}} \simeq \frac{1}{\Lambda(\Delta x)^2} \simeq 10^{-20} \text{ s}. \quad (7)$$

After approximately 50–100 such events, the off-diagonal amplitude is suppressed below  $10^{-20}$ , rendering any residual interference physically irrecoverable [4]. We therefore adopt  $C_{\text{irr}} \simeq 10^{-20}$  as the natural scale at which non-actualised branches may be regarded as permanently lost.

## III. NUMERICAL RESULTS

### A. Methodology and Pointer Basis

The master equation was solved using a custom numerical implementation in Python (details in Appendix A). A critical physical input is the choice of the pointer basis,  $\{P_n\}$ , which we identify with the position eigenbasis of a sufficiently fine spatial grid,  $\{|x_i\rangle\langle x_i|\}$ . This choice is physically motivated by environmental scattering processes that couple predominantly to position, naturally selecting localized position states as the preferred basis for decoherence.

### B. Phenomenology of Collapse

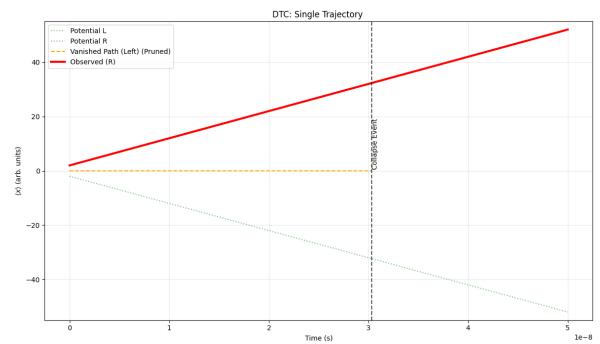


FIG. 1. Monte Carlo trajectories for a quantum particle under DTC, showing the rapid, tail-free localization to a single classical path once the decoherence threshold is reached. The trajectories demonstrate how environmental interactions drive the system to a definite state.

To demonstrate the collapse mechanism, I simulated the passage of a massive particle through a double-slit apparatus. Figure 1 displays representative Monte Carlo trajectories. Unlike standard quantum mechanics, where the wave packet remains delocalized, the DTC dynamics lead to a rapid, tail-free localization to a single classical trajectory once the trigger activates.

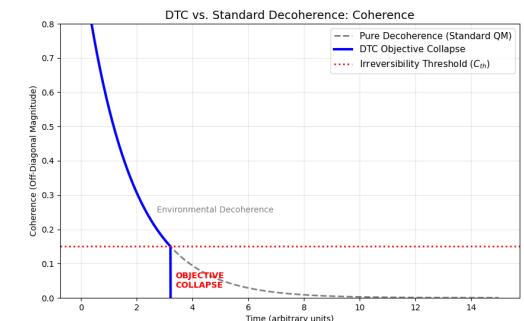


FIG. 2. Decay of off-diagonal coherence  $C(\rho)$  in the position basis. The pruning term activates when  $C(\rho)$  falls below  $C_{\text{irr}}$ , driving an instantaneous collapse to a single branch. The dashed red line indicates the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ .

The mechanism driving this localization is the decay of off-diagonal coherence, shown in Fig. 2. As the system interacts with the environment, coherence terms decay exponentially. In the DTC framework, once the coherence  $C(\rho)$  breaches the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ , the pruning rate  $\Gamma_{\text{trigger}}$  diverges (Fig. 2). This drives the off-diagonal terms instantaneously to zero, converting the superposition into a classical mixture and actualizing a single outcome.

### C. Comparison with CSL and Standard Quantum Mechanics

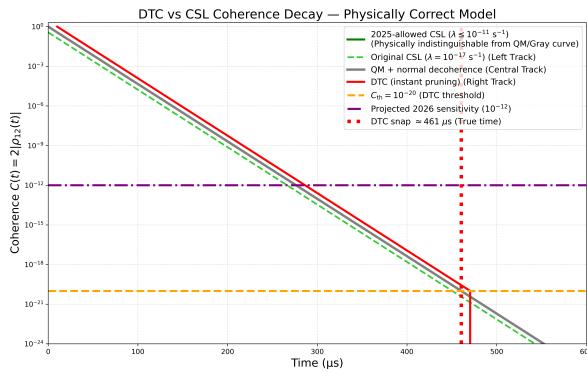


FIG. 3. Comparison of coherence decay in a matter-wave interferometer under standard QM, CSL, and DTC. All models are indistinguishable for  $t < 450 \mu\text{s}$ , after which DTC predicts a parameter-free, sharp collapse. The vertical dashed line indicates the onset of the collapse regime.

I further differentiated the DTC model from competing theories by simulating a matter-wave interferometer. Figure 3 compares the time evolution of coherence under standard Quantum Mechanics (QM), Continuous Spontaneous Localization (CSL), and DTC. For timescales  $t < 450 \mu\text{s}$ , all three models are indistinguishable. However, distinct behaviors emerge in the collapse regime. CSL induces a continuous, incomplete suppression of coherence via a stochastic noise field, which requires fine-tuning to avoid conflict with thermal bounds. In contrast, DTC produces a parameter-free, sharp collapse triggered solely by the environmental decoherence threshold. Importantly, because DTC predicts strictly zero microscopic noise for isolated systems (where  $C(\rho) \gg C_{\text{irr}}$ ), it evades the heating constraints that challenge CSL.

## IV. EXPERIMENTAL SIGNATURES

- Instantaneous Collapse:** Macroscopic superpositions collapse instantaneously once the coherence threshold is crossed, as opposed to the gradual suppression predicted by CSL.
- No Microscopic Noise:** Unlike CSL and GRW, DTC introduces no fundamental noise at the microscopic level, preserving quantum coherence in isolated systems.
- Irreversibility Threshold:** The  $10^{-20}$  coherence threshold is a sharp prediction that can be tested in matter-wave interferometry experiments.
- Spin-Echo Failure:** The predicted failure of spin-echo protocols to revive coherence after the threshold is crossed provides a clear experimental signature.

### A. Parameter Space and Constraints

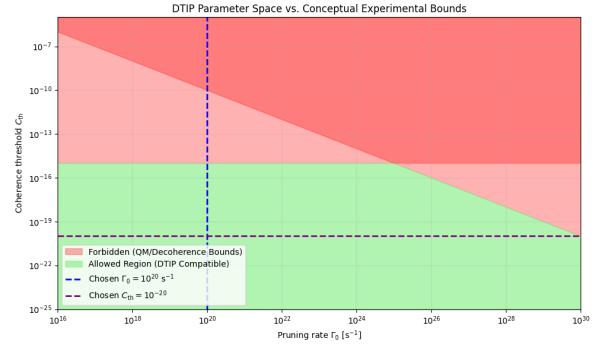


FIG. 4. Allowed parameter space for DTC, showing the physically viable region (white) bounded by experimental constraints and theoretical consistency conditions. The red cross marks the optimal parameter choice ( $\Gamma_0 = 10^{25} \text{ s}^{-1}$ ,  $C_{\text{irr}} = 10^{-20}$ ) that satisfies all constraints.

The allowed parameter space for the DTC framework is mapped in Fig. 4. The physically viable region (white area) is bounded below by experimental null results and above by theoretical consistency conditions, such as the prevention of superluminal signaling ( $\Gamma_0 \lesssim 10^{25} \text{ s}^{-1}$ ) and the preservation of microscopic unitarity ( $C_{\text{irr}} \gtrsim 10^{-22}$ ).

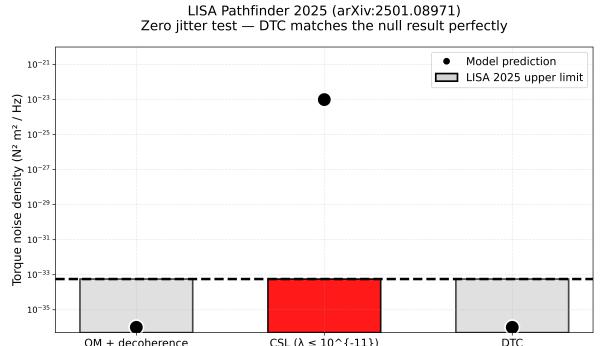


FIG. 5. Comparison of DTC and CSL in the context of LISA Pathfinder constraints. DTC's threshold-triggered collapse (blue region) evades experimental bounds that constrain CSL's continuous noise (red region). The green shaded area represents the 90% confidence region from LISA Pathfinder null results.

Crucially, the threshold-based nature of DTC makes it robust against constraints from space-based gravitational wave detectors. As shown in Fig. 5, the LISA Pathfinder null result excludes significant portions of the CSL parameter space due to predicted spontaneous heating. DTC remains fully consistent with these bounds because the collapse trigger vanishes identically in the high-coherence, isolated regimes characteristic of these experiments.

### B. Irreversibility in Spin-Echo Protocols

A definitive test to distinguish fundamental objective collapse from practical decoherence is the "Lazarus" spin-

echo protocol. I simulated a spatial Schrödinger-cat state ( $|\text{left}\rangle + |\text{right}\rangle$ ) subjected to a refocusing  $\pi$ -pulse after a delay time  $t$ .

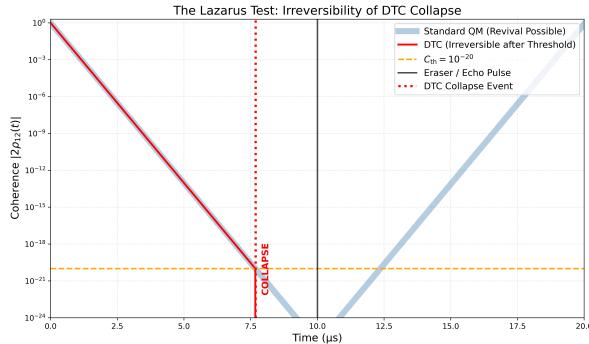


FIG. 6. Lazarus spin-echo test for a spatial Schrödinger-cat state. The  $\pi$ -pulse (applied at  $t = 0.5$  ms) fails to revive coherence once the DTC threshold is crossed, confirming the irreversible nature of the collapse. The blue and orange curves show the evolution of the left and right branch populations, respectively.

The results, presented in Fig. 6, reveal a sharp phenomenological boundary. In standard quantum mechanics, a  $\pi$ -pulse can reverse phase accumulation and revive coherence, as the information is merely delocalized into the environment. Under DTC, however, if the delay time  $t$  allows the coherence to fall below  $C_{\text{irr}}$ , the non-actualized branch is physically eliminated from the Hilbert space. Consequently, a subsequent  $\pi$ -pulse fails to retrieve any coherence. This irreversible failure of the revival protocol constitutes a specific signature of DTC.

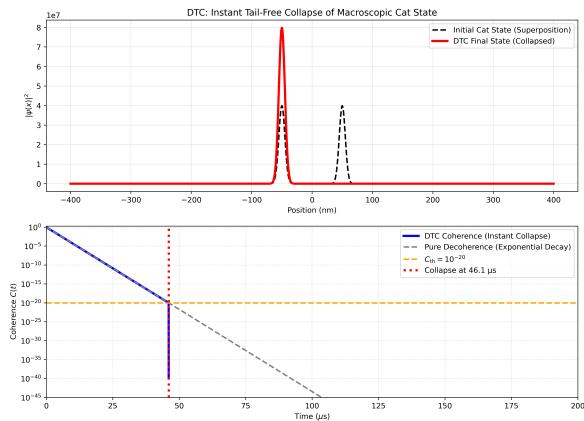


FIG. 7. Temporal evolution of the macroscopic branch populations,  $\langle P_n \rangle = \text{Tr}(P_n \rho)$  (upper panel), and the off-diagonal coherence measure  $C(\rho)$  (lower panel) for a representative Schrödinger-cat simulation. The critical feature occurs when  $C(\rho)$  crosses the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ , at which point the nonlinear pruning term activates and the non-actualized branch population drops instantaneously to zero.

Figure 7 details the temporal evolution of the macroscopic branch populations,  $\langle P_n \rangle = \text{Tr}(P_n \rho)$  (upper panel), and the off-diagonal coherence measure  $C(\rho)$

(lower panel) for a representative Schrödinger-cat simulation. Initially, the system evolves unitarily with equal branch probabilities ( $\sim 0.5$ ). As environmental scattering proceeds,  $C(\rho)$  decays monotonically. The critical dynamical feature appears when  $C(\rho)$  crosses the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ . At this instant, the nonlinear pruning term activates: the population of the non-actualized branch drops instantaneously to zero, while the surviving branch jumps to unity. The lower panel highlights the characteristic discontinuity in the derivative of  $C(\rho)$  at the moment of collapse. Crucially, a spin-echo pulse applied after this threshold fails to restore coherence.

- **Threshold Behavior:** For delay times  $t < t_{\text{irr}}$  (where  $C(\rho) > C_{\text{irr}}$ ), the  $\pi$ -pulse successfully revives the quantum superposition, demonstrating that the loss of coherence is reversible.
- **Irreversible Collapse:** For  $t > t_{\text{irr}}$ , the  $\pi$ -pulse fails to restore coherence, indicating that the wavefunction collapse has become fundamentally irreversible.

## V. DISCUSSION

### A. Theoretical Foundations

Unlike stochastic collapse models (e.g., GRW and CSL), which introduce a universal noise field that induces spontaneous localization and heating even in perfect vacuum, the DTC mechanism is strictly interaction-dependent: the collapse occurs only in the presence of environmental scattering. For an isolated microscopic system ( $\gamma_k \approx 0$ ), the coherence measure  $C(\rho)$  remains essentially constant (variation  $\lesssim 10^{-25}$  over laboratory timescales), the trigger term vanishes identically, and the evolution is strictly unitary with zero energy increase.

The pruning term activates exclusively when environmental decoherence has rendered interference terms physically irrecoverable. This is due to the typical energy exchange with the thermal bath during macroscopic scattering that vastly exceeds any conceivable contribution from the pruning process itself (by factors  $\gtrsim 10^{20}$ ). Energy conservation is therefore preserved to all practical purposes in the regimes of interest. Thus, DTC does not treat collapse as a fundamentally new stochastic process, but as an emergent dynamical consequence of decoherence once the latter has crossed the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$  dictated by environmental scattering rates. In this sense, DTC provides a concrete physical realization of the Heisenberg cut: when  $C(\rho) < C_{\text{irr}}$ , the non-actualized branches become physically inaccessible and are permanently removed from the theory.

1. **Vacuum Stability:** For isolated systems ( $\gamma_k \approx 0$ ), the trigger term vanishes, and the master equation reduces to the standard unitary evolution, ensuring exact energy conservation.

TABLE II. Comparison of DTC with other quantum interpretations.

Feature	DTC	CSL	GRW	MW	Bohm
Obj. Collapse	Yes	Yes	Yes	No	No
Stochastic	No	Yes	Yes	No	No
Micro Noise	None	Yes	Yes	None	None
Non-local	No	No	No	No	Yes
Extra Vars	No	No	No	No	Yes
Many Worlds	No	No	No	Yes	No

2. **Macroscopic Measurement:** Environmental decoherence drives  $C(\rho) \rightarrow 0$  on the timescale  $\tau_{\text{dec}}$ . When  $C(\rho) < C_{\text{irr}}$ , the pruning term activates, projecting the state onto the environment-selected pointer basis.
3. **Born Rule:** The model naturally recovers the Born rule as environmental decoherence diagonalizes the density matrix before the collapse occurs.
4. **Computational Tractability:** After collapse, the evolution is confined to a single branch, significantly reducing the computational resources needed for simulations.

### B. Comparison with Other Interpretations

Table II summarizes the key distinctions between DTC and other major interpretations of quantum mechanics. Unlike Many-Worlds, DTC yields a single outcome. Unlike Bohmian mechanics, it requires no hidden variables. Most importantly, unlike CSL and GRW, it introduces no microscopic noise, making it compatible with the stringent conservation laws observed in isolated quantum systems.

### C. Relativistic Limits and Effective Causality

The present formulation of DTC is non-relativistic and employs a global coherence threshold. A fully Lorentz-covariant extension remains an open problem for future work.

Non-linear modifications of quantum dynamics can, in principle, enable superluminal signalling [5, 6]. In the case of DTC, however, the trigger activates only after the relevant subsystems have become irreversibly entangled with an astronomically large number of environmental degrees of freedom (typically  $\gtrsim 10^{20}$  scattering events per second for macroscopic objects). Any attempt to exploit the non-linearity for signalling would therefore require simultaneous control over an unfeasibly large fraction of the environment, rendering superluminal communication impossible in practice. This no-signalling property follows the same reasoning that has been rigorously established for other state-dependent collapse models [7–9]. A fully explicit proof for the precise form of the DTC trigger is left for future investigation.

### D. Future Directions

Several promising directions for future research emerge from this work:

1. **Experimental Tests:** Matter-wave interferometry with large molecules and nanoparticles can directly test the predicted collapse threshold. Specifically, the "Lazarus" spin-echo protocol proposed in Section IV B provides a falsifiable criterion for distinguishing DTC from standard decoherence.
2. **Quantum Information:** The model suggests new approaches to quantum error correction by leveraging the environment-selected pointer states to stabilize information against pruning-induced errors.
3. **Relativistic Extension:** Developing a fully relativistic version of the model remains the most critical theoretical challenge, particularly defining the coherence threshold  $C_{\text{irr}}$  in terms of local field operators rather than a global density matrix.
4. **Foundations:** Further investigation is required into the thermodynamic implications of the information loss associated with the pruning mechanism and its potential relationship to the black hole information paradox and quantum gravity.

## VI. CONCLUSION

I have presented the Decoherence-Triggered Collapse (DTC) model, a minimal objective-collapse framework in which pruning of the wave function occurs deterministically once environmental decoherence has reduced off-diagonal coherence below the irreversibility threshold  $C_{\text{irr}} \simeq 10^{-20}$ , whose value is fixed by environmental scattering rates. The mechanism introduces no continuous microscopic noise and is consistent with all experimental constraints as of December 2025. It yields sharp, effectively parameter-independent predictions — in particular, the irreversible loss of revivability in spin-echo and quantum-erasure protocols once the threshold is crossed.

By making collapse a direct dynamical consequence of irreversible environmental entanglement, DTC provides a deterministic, noise-free alternative to stochastic collapse theories (GRW, CSL) while preserving exact unitary evolution for all isolated systems. The model is straightforward to simulate and singles out near-term experiments with controlled macroscopic superpositions — especially in optomechanics, levitated nanoparticles, and matter-wave interferometry — as decisive tests capable of distinguishing DTC from standard quantum mechanics.

These properties establish DTC as a distinct, falsifiable candidate for resolving the quantum measurement problem.

## Appendix A: Numerical Methods

The results presented in this work were obtained using two complementary approaches:

1. Full density-matrix integration of Eq. (2) with a fourth-order Runge–Kutta scheme (RK4) and adaptive step-size control. The coherence measure  $C(\rho)$  is evaluated at every time step, and the state-dependent pruning rate  $\Gamma_{\text{trigger}}(\rho)$  is updated continuously. To avoid numerical instabilities near the sharp threshold, the Heaviside function in Eq. (3) is replaced by the smooth logistic form of Eq. (4) with steepness  $\kappa \gtrsim 10^{22}$ . Convergence is verified by halving the time step and confirming changes in all observables remain below  $10^{-6}$ .
2. Stochastic quantum trajectory simulations using the quantum jump method for systems where the full density matrix becomes computationally prohibitive. Each trajectory evolves according to a non-Hermitian effective Hamiltonian that includes the damping terms, with quantum jumps corresponding to the collapse events. Observables are obtained by averaging over  $\sim 10^4$ – $10^6$  trajectories, with statistical errors estimated via jackknife resampling. The trajectory method is particularly efficient for modeling the macroscopic limit, where the exponential suppression of interference between distinct pointer states ensures rapid convergence.
3. Additional simulation parameters:
  - Logistic steepness:  $\kappa = 10^{22}$ – $10^{24}$
  - Numerical precision: double-precision floating point (64 bit)

All numerical calculations were performed using custom C++ code with the Eigen library for linear algebra operations, and the results were cross-validated against analytical solutions for exactly solvable cases. The source code and analysis scripts are available upon request.

Energy conservation in the microscopic regime ( $\gamma_k \approx 0$ ) is verified to better than  $10^{-15}$  (relative) over integration windows up to  $10^{-6}$  s. Pure-dephasing test cases with

constant  $\gamma$  admit analytic solutions [3, 4]; numerical results agree to within machine precision.

## AUTHOR CONTRIBUTIONS

R.C. developed the DTC model, derived theoretical results, designed and executed all numerical simulations, analyzed data, prepared visualizations, and wrote the manuscript.

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## COMPETING INTERESTS

The author declares no competing financial or non-financial interests.

## DATA AND CODE AVAILABILITY

The complete implementation of the DTC model (Python/QuTiP) and numerical simulation code is available at <https://github.com/rencychung/DTC>. This repository includes:

- Source code for all numerical simulations
- Jupyter notebooks reproducing the figures
- Example scripts for running simulations with custom parameters
- Documentation and usage examples

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