

N 6.1

Запишем уравнения:

$$\frac{\partial \rho}{\partial t} + \operatorname{div} \bar{j} = 0; \quad \bar{j} = \sigma \bar{E}; \quad \bar{D} = \epsilon \bar{E}, \quad \operatorname{div} \bar{D} = 4\pi \rho \Rightarrow$$

$$\bar{j} = \frac{\sigma}{\epsilon} \bar{D}; \quad \operatorname{div} \bar{j} = \frac{\sigma}{\epsilon} \operatorname{div} \bar{D} = \frac{4\pi \rho \sigma}{\epsilon}$$

$$\frac{\partial \rho}{\partial t} + \frac{4\pi \sigma}{\epsilon} \rho = 0 \Rightarrow \frac{\partial \rho}{\rho} = - \frac{4\pi \sigma}{\epsilon} dt \Rightarrow \rho = C \exp \left\{ - \frac{4\pi \sigma}{\epsilon} t \right\}$$

Из начальных условий:  $\rho(\bar{r}, t=0) = \rho_0(\bar{r}) \Rightarrow$

$$\rho(\bar{r}, t) = \rho_0(\bar{r}) \exp \left\{ - \frac{4\pi \sigma}{\epsilon} t \right\}$$

Найдем изменение поля  $\bar{E}$ :

$$\rho = \frac{\operatorname{div} \bar{D}}{4\pi} = \frac{\epsilon \operatorname{div} \bar{E}}{4\pi}$$

$$\frac{\partial}{\partial t} \frac{\epsilon \operatorname{div} \bar{E}}{4\pi} + \operatorname{div} \bar{j} = \frac{\partial}{\partial t} \frac{\epsilon \operatorname{div} \bar{E}}{4\pi} + \sigma \operatorname{div} \bar{E} = \operatorname{div} \left( \frac{\epsilon}{4\pi} \frac{\partial \bar{E}}{\partial t} + \sigma \bar{E} \right) = 0 \Rightarrow$$

$$\Rightarrow \frac{\epsilon}{4\pi} \frac{\partial \bar{E}}{\partial t} + \sigma \bar{E} = 0; \quad \bar{E}(\bar{r}, t) = \bar{E}_0(\bar{r}) \exp \left\{ - \frac{4\pi \sigma}{\epsilon} t \right\}$$

$$\operatorname{rot} \bar{H} = \frac{4\pi}{c} \bar{j} + \frac{1}{c} \frac{\partial \bar{D}}{\partial t} = \frac{1}{c} \left( 4\pi \sigma \bar{E} + \epsilon \frac{\partial \bar{E}}{\partial t} \right)$$

$$\epsilon \frac{\partial \bar{E}}{\partial t} = - 4\pi \sigma \bar{E} \Rightarrow \operatorname{rot} \bar{H} = 0 \Rightarrow |\bar{H}| = 0$$

N 6.8

$$\operatorname{rot} \bar{H} = \frac{4\pi}{c} \bar{j} + \frac{1}{c} \frac{\partial \bar{D}}{\partial t}; \quad \frac{4\pi}{c} \sigma \bar{E}$$



Для комплексного амплитуды

$$\text{rot } \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{i\omega}{c} \vec{D} = \frac{4\pi}{c} \vec{E} + \frac{i\omega}{c} \epsilon \vec{E} =$$

$$= \frac{i\omega}{c} \left( \frac{4\pi\sigma}{i\omega} + \epsilon \right) \vec{E} = \frac{i\omega}{c} \left( \epsilon - i \frac{4\pi\sigma}{\omega} \right) \vec{E} = \frac{i\omega}{c} \epsilon_k \vec{E}$$

Учитывая, что  $\sigma \gg \omega$

$$\epsilon_k \approx -i \frac{4\pi\sigma}{\omega}$$

$$\begin{cases} \text{rot } \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} i\omega \vec{D} \\ \text{rot } \vec{E} = -\frac{1}{c} i\omega \vec{H} \end{cases} \quad \text{rot} = \frac{i\omega}{c} \epsilon_k \vec{E}$$

$$\text{rot rot } \vec{H} = \frac{4\pi}{c} \epsilon_k \text{rot } \vec{E} \quad \text{rot rot } \vec{E} = -\mu \frac{i\omega}{c} \text{rot } \vec{H}$$

$$\text{grad div } \vec{H} - \Delta \vec{H} = -\mu \frac{i\omega}{c} \text{rot } \vec{H}$$

$$-\Delta \vec{E} = -\frac{i\omega}{c} \mu \frac{i\omega}{c} \epsilon_k \vec{E}, \quad \Delta \vec{E} + \frac{\omega^2}{c^2} \mu \epsilon_k \vec{E} = 0$$

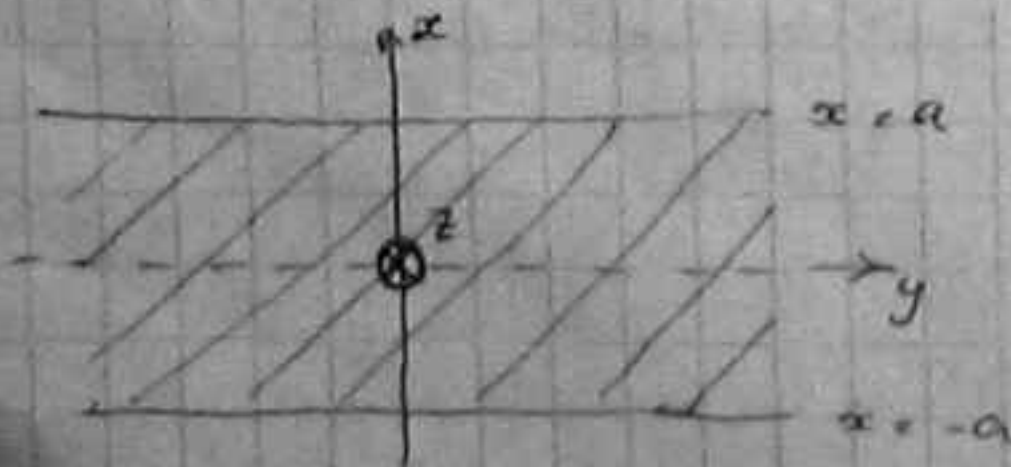
$$\vec{E} = \vec{E}(x)$$

$$\frac{d^2 \vec{E}}{dx^2} + \frac{\omega^2}{c^2} \mu \epsilon_k \vec{E} = 0$$

$$\text{div } \vec{D} = 0 \Rightarrow \text{div } \vec{E} = 0 \Rightarrow \frac{dE_x}{dx} = 0 \Rightarrow E_x = \text{const}$$

$$\frac{d^2 E_x}{dx^2} + \frac{\omega^2}{c^2} \mu \epsilon_k E_x = 0 \Rightarrow E_x = 0$$

Тангенциальная составляющая задана как  $E_y \Rightarrow E_z = 0$



$$\frac{d^2 E_y}{dx^2} + \frac{\omega^2}{c^2} \mu \epsilon_k E_y = 0$$

$$E_y = C_1 \exp \left\{ \frac{\omega}{c} \sqrt{\mu \epsilon_k} x \right\} + C_2 \exp \left\{ -\frac{\omega}{c} \sqrt{\mu \epsilon_k} x \right\}$$

$$\sqrt{\mu \epsilon_k} = \sqrt{\mu \left( -i \frac{4\pi\sigma}{\omega} \right)} = \pm \sqrt{\frac{2\pi\sigma\mu}{\omega}} (1+i)$$

$$E_y = C_1 \text{sh} \left\{ \frac{\omega}{c} \sqrt{\frac{2\pi\sigma\mu}{\omega}} (1+i) x \right\} + C_2 \text{ch} \left\{ \frac{\omega}{c} \sqrt{\frac{2\pi\sigma\mu}{\omega}} (1+i) x \right\}$$

$$E_y(a) = E_y(-a) = E_0 - \text{четная функция}$$

$$E_y = C_2 \text{ch} \left\{ \frac{\omega}{c} \sqrt{\frac{2\pi\sigma\mu}{\omega}} (1+i) x \right\} = C_2 \text{ch} \left\{ \frac{1}{c} \sqrt{2\pi\sigma\mu\omega} (1+i) x \right\} =$$

$$= C_2 \text{ch} \left\{ \frac{x}{\delta} (1+i) \right\}$$

$$E_y(a) = C_2 \text{ch} \left\{ \frac{a}{\delta} (1+i) \right\} = E_0, \quad C_2 = E_0 \text{ch}^{-1} \left\{ \frac{a}{\delta} (1+i) \right\}$$

$$E_y = E_0 \text{ch}^{-1} \left\{ \frac{a}{\delta} (1+i) \right\} \text{ch} \left\{ \frac{x}{\delta} (1+i) \right\} =$$

$$= E_0 \frac{e^{\frac{x}{\delta}} e^{\frac{ix}{\delta}} + e^{-\frac{x}{\delta}} e^{-\frac{ix}{\delta}}}{e^{\frac{a}{\delta}} e^{\frac{ia}{\delta}} + e^{-\frac{a}{\delta}} e^{-\frac{ia}{\delta}}} =$$

$$= E_0 \frac{e^{\frac{x}{\delta}} (\cos \frac{x}{\delta} + i \sin \frac{x}{\delta}) + e^{-\frac{x}{\delta}} (\cos \frac{x}{\delta} - i \sin \frac{x}{\delta})}{e^{\frac{a}{\delta}} (\cos \frac{a}{\delta} + i \sin \frac{a}{\delta}) + e^{-\frac{a}{\delta}} (\cos \frac{a}{\delta} - i \sin \frac{a}{\delta})} =$$

$$= E_0 \frac{\text{ch} \frac{x}{\delta} \cos \frac{x}{\delta} + i \text{sh} \frac{x}{\delta} \sin \frac{x}{\delta}}{\text{ch} \frac{a}{\delta} \cos \frac{a}{\delta} + i \text{sh} \frac{a}{\delta} \sin \frac{a}{\delta}}$$

$$= E_0 \frac{\text{ch}^2 \frac{x}{\delta} \cos^2 \frac{x}{\delta} - \text{sh}^2 \frac{x}{\delta} \sin^2 \frac{x}{\delta}}{\text{ch}^2 \frac{a}{\delta} \cos^2 \frac{a}{\delta} - \text{sh}^2 \frac{a}{\delta} \sin^2 \frac{a}{\delta}}$$

$$a) \quad a \gg \delta \Rightarrow \frac{a}{\delta} \gg 1$$



$$E_y = |Re\{E_y\}| e^{i\omega t}$$

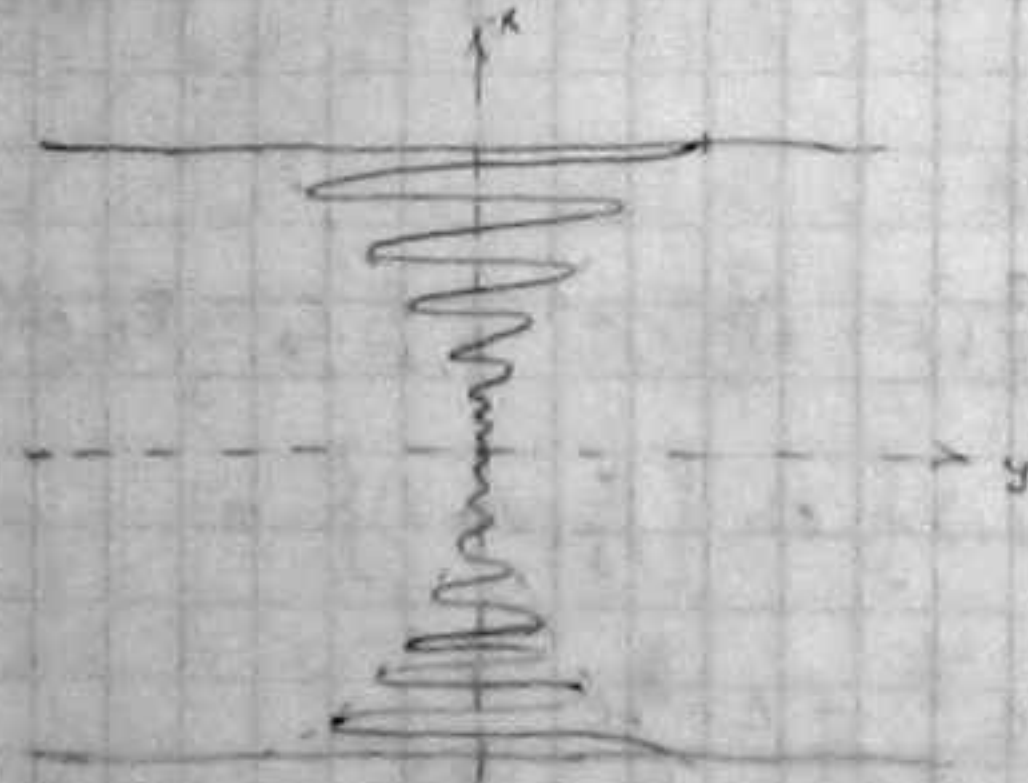
$$a) a \gg \delta, \frac{a}{\delta} \gg 1$$

$$E_y = E_0 \frac{e^{\frac{x}{\delta}} e^{\frac{ix}{\delta}} + e^{-\frac{x}{\delta}} e^{-\frac{ix}{\delta}}}{e^{\frac{a}{\delta}} e^{\frac{ia}{\delta}}}$$

$$E_y e^{i\omega t} = E_0 e^{\frac{x}{\delta} - \frac{a}{\delta}} e^{i(\frac{x}{\delta} - \frac{a}{\delta} + \omega t)} + E_0 e^{-\frac{x}{\delta} - \frac{a}{\delta}} e^{i(\frac{x}{\delta} - \frac{a}{\delta} + \omega t)} =$$

$$\approx E_0 e^{i\omega t} \cos(\omega t + \frac{x}{\delta} - \frac{a}{\delta}) + E_0 e^{-\frac{x}{\delta} - \frac{a}{\delta}} \cos(\omega t - \frac{x}{\delta} - \frac{a}{\delta})$$

$\approx 0$



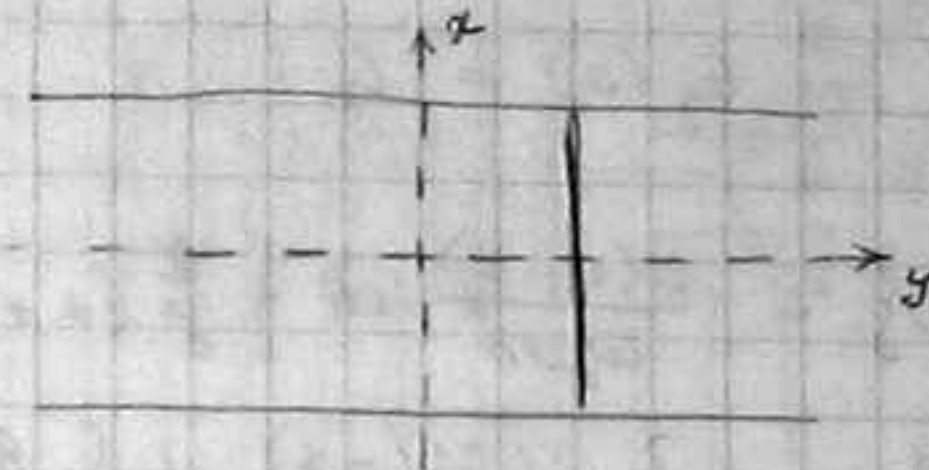
$$b) a \ll \delta, \frac{a}{\delta} \ll 1$$

$$E_y \approx E_0 \frac{(1 + \frac{x}{\delta}) e^{\frac{ix}{\delta}} + (1 - \frac{x}{\delta}) e^{-\frac{ix}{\delta}}}{(1 + \frac{a}{\delta}) e^{\frac{ia}{\delta}} + (1 - \frac{a}{\delta}) e^{-\frac{ia}{\delta}}}$$

$$= E_0 \frac{(1 + \frac{x}{\delta})(1 + \frac{ix}{\delta}) + (1 - \frac{x}{\delta})(1 - \frac{ix}{\delta})}{(1 + \frac{a}{\delta})(1 + \frac{ia}{\delta}) + (1 - \frac{a}{\delta})(1 - \frac{ia}{\delta})}$$

$$\approx E_0 \frac{1 + \frac{x}{\delta} + \frac{ix}{\delta} + i\frac{x^2}{\delta^2} + 1 - \frac{x}{\delta} - \frac{ix}{\delta} + i\frac{x^2}{\delta^2}}{1 + \frac{a}{\delta} + \frac{ia}{\delta} + i\frac{a^2}{\delta^2} + 1 - \frac{a}{\delta} - \frac{ia}{\delta} + i\frac{a^2}{\delta^2}} = \frac{2 + i\frac{x^2}{\delta^2}}{2 + i\frac{a^2}{\delta^2}} \approx E_0$$

$$Re\{E_y e^{i\omega t}\} = E_0 \frac{2 \cos \omega t - \frac{x^2}{\delta^2} \sin \omega t}{2 \cos \omega t - \frac{a^2}{\delta^2} \sin \omega t} = E_0 \cos \omega t$$



нб. 10

Из соображений непрерывности грани  $E_x = E_z = 0$

$$a) E_y(-a) = E_0$$

$$E_y(0) = 0$$

$$E_y = C_1 \operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\} + C_2 \operatorname{ch} \left\{ \frac{x}{\delta} (1+i) \right\}$$

$$E_y(x=0) = C_2 = 0;$$

$$E_y(-a) = -C_1 \operatorname{sh} \left\{ \frac{a}{\delta} (1+i) \right\} = E_0 \rightarrow$$

$$C_1 = -E_0 \operatorname{sh}^{-1} \left\{ \frac{a}{\delta} (1+i) \right\}$$

$$E_y = -E_0 \frac{\operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\}}{\operatorname{ch} \left\{ \frac{a}{\delta} (1+i) \right\}}$$

$$b) E_y(-a) = E_0$$

$$E_y = C_1 \operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\} + C_2 \operatorname{ch} \left\{ \frac{x}{\delta} (1+i) \right\}$$

$$\operatorname{rot} \vec{E} = -\frac{i\omega}{c} \mu \vec{H}$$

$$\operatorname{rot} \vec{E} = \begin{vmatrix} \vec{x}_0 & \vec{y}_0 & \vec{z}_0 \\ \frac{\partial}{\partial x} & 0 & 0 \\ 0 & E_y & 0 \end{vmatrix} = \vec{z}_0 \frac{\partial E_y}{\partial x}, \quad \vec{H} = \frac{c}{\omega \mu} \vec{z}_0 \frac{\partial E_y}{\partial x}$$



$$H_z = \frac{C_1}{\omega \mu} \frac{i+1}{\delta} \left[ C_1 \operatorname{ch} \left\{ \frac{x}{\delta} (1+i) \right\} + C_2 \operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\} \right]$$

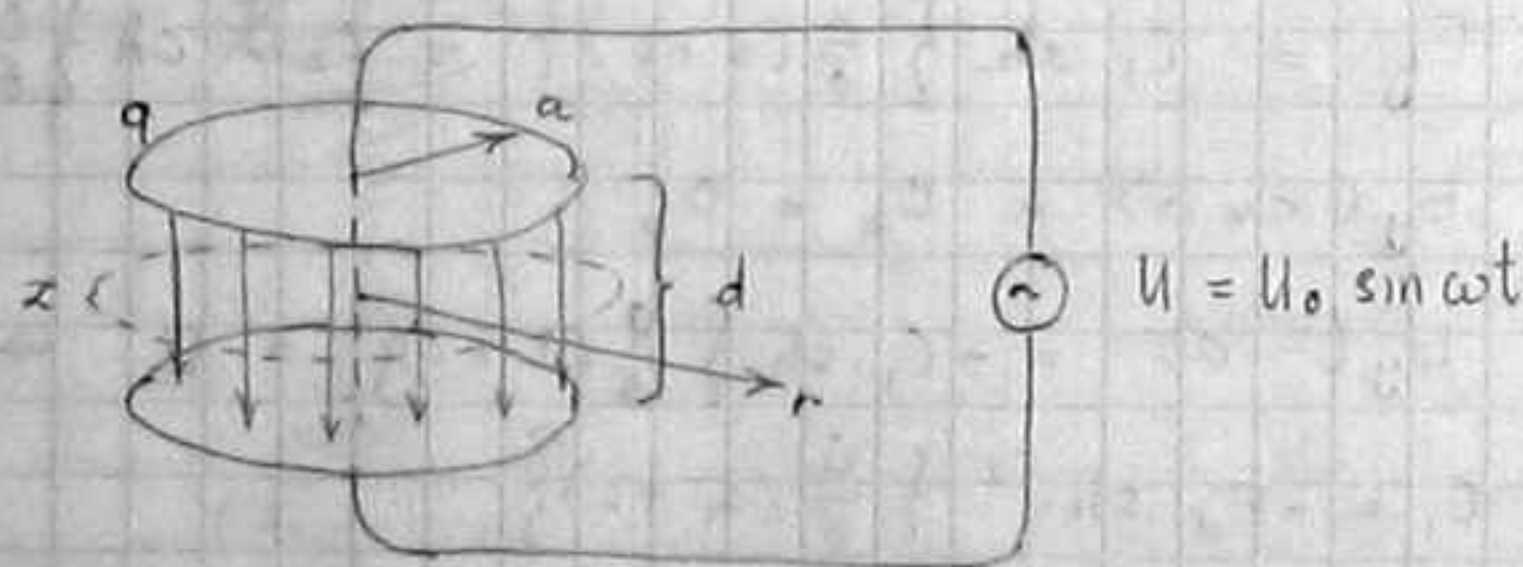
$$H_z(0) = \frac{C(i-1)}{\omega \mu \delta} C_1 = E_y(0) = C_2$$

$$E_y = C_1 \left[ \operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\} + \frac{C(i-1)}{\omega \mu \delta} \operatorname{ch} \left\{ \frac{x}{\delta} (1+i) \right\} \right]$$

$$E_y(-a) = C_1 \left[ -\operatorname{sh} \left\{ \frac{a}{\delta} (1+i) \right\} + \frac{C(i-1)}{\omega \mu \delta} \operatorname{ch} \left\{ \frac{a}{\delta} (1+i) \right\} \right] = E_0$$

$$E_y = E_0 \cdot \frac{\frac{C(i-1)}{\omega \mu \delta} \operatorname{ch} \left\{ \frac{x}{\delta} (1+i) \right\} + \operatorname{sh} \left\{ \frac{x}{\delta} (1+i) \right\}}{\frac{C(i-1)}{\omega \mu \delta} \operatorname{ch} \left\{ \frac{a}{\delta} (1+i) \right\} - \operatorname{sh} \left\{ \frac{a}{\delta} (1+i) \right\}}$$

N6, 12



$$q = CV = CU_0 \sin \omega t$$

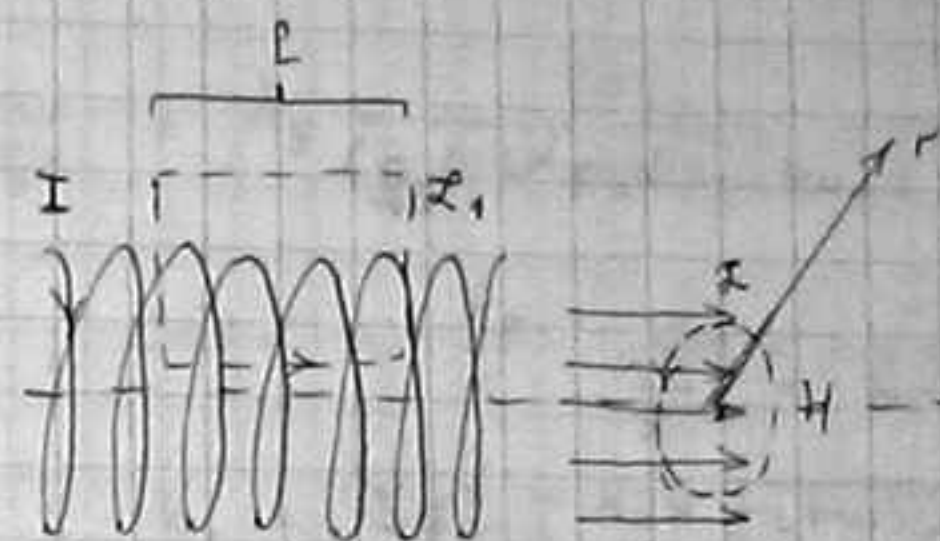
$$Ed = U_0 \sin \omega t, \quad E = \frac{U_0 \sin \omega t}{d}$$

$$\int_x \bar{H} d\bar{l} = \frac{1}{c} \frac{\partial}{\partial t} \int_s \bar{D} d\bar{s} = \frac{E}{c} \frac{\partial}{\partial t} \int_s \bar{E} d\bar{s} = \frac{E}{c} \frac{\partial}{\partial t} (E s) =$$

$$= \frac{E}{c} \frac{\partial}{\partial t} \left( \frac{U_0 \sin \omega t}{d} \pi r^2 \right) = \frac{\omega \epsilon U_0 \pi r^2 \cos \omega t}{cd}$$

$$\int_x \bar{H} d\bar{l} = 2\pi r H, \quad H = \frac{\omega \epsilon U_0 r \cos \omega t}{2cd}$$

N6, 13



$$\oint_{x_1} \bar{H} d\bar{l} = \frac{4\pi}{c} I_z$$

$$\oint_{x_1} \bar{H} d\bar{l} = lH; \quad I_z = nl I_0 \sin \omega t; \quad H = \frac{4\pi n I_0}{c} \sin \omega t$$

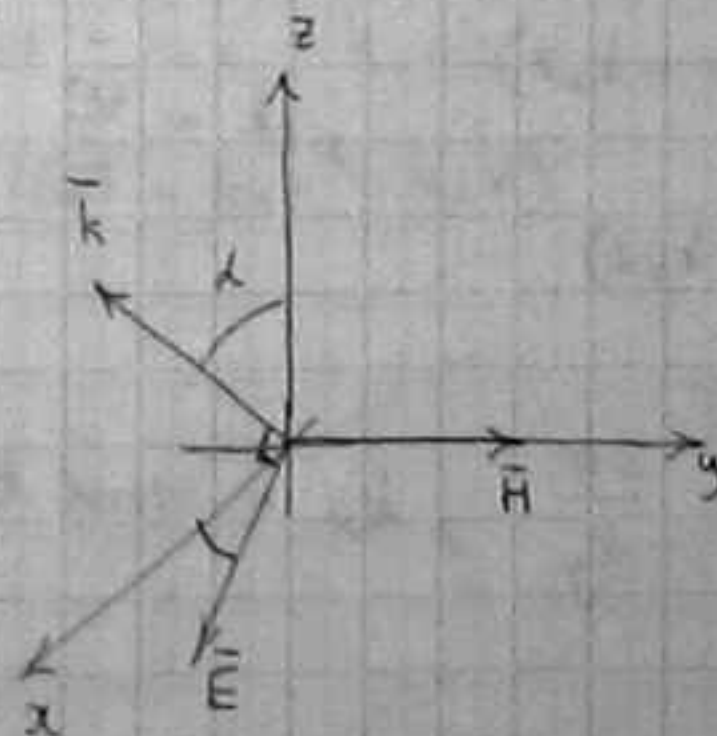
$$\oint_x \bar{E} d\bar{l} = -\frac{1}{c} \frac{\partial}{\partial t} \int_s \bar{B} d\bar{s}$$

$$\oint_x \bar{E} d\bar{l} = 2\pi r E; \quad -\frac{1}{c} \frac{\partial}{\partial t} \int_s \bar{B} d\bar{s} = -\frac{\mu}{c} \frac{\partial}{\partial t} (H \pi r^2) =$$

$$= -\frac{\mu}{c} \frac{4\pi}{c} n I_0 \pi r^2 \omega \cos \omega t$$

$$E = -\frac{4\pi \mu n I_0 r \omega}{2c^2} \cos \omega t = -\frac{2\pi \mu n I_0 r \omega}{c^2} \cos \omega t$$

N7, 1





$$a) \vec{E} = \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$$

$$\vec{E}_x = \cos \alpha |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$\cos \alpha = \frac{k_z}{k}; \quad \sin \alpha = \frac{k_x}{k}$$

$$E_x = \frac{k_z}{k} |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$E_z = -\frac{k_x}{k} |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$E_y = 0$$

$$H_x = 0, H_z = 0$$

$$\text{rot } \vec{E} = -\frac{i\omega}{c} \vec{H}; \quad \vec{H} = \frac{ic}{\omega} \text{rot } \vec{E}$$

$$\text{rot } \vec{E} = \begin{vmatrix} \vec{x}_0 & \vec{y}_0 & \vec{z}_0 \\ \frac{\partial}{\partial x} & 0 & \frac{\partial}{\partial z} \\ E_x & 0 & E_z \end{vmatrix} = \vec{y}_0 \left( \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \right) =$$

$$= \vec{y}_0 \left( -\frac{ik_z^2}{k} |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)} + \frac{ik_x^2}{k} |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)} \right) =$$

$$= \vec{y}_0 (-ik) |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$H_y = \frac{c}{\omega} k |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}; \quad k = \frac{\omega}{c}$$

$$H_y = |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$b) \lambda_x = \frac{2\pi}{k_x}; \quad \lambda_y = \frac{2\pi}{k_y}; \quad k_x = \sin \alpha k; \quad k_z = \cos \alpha k$$

$$k = \frac{\omega}{c} \sqrt{\epsilon_{\mu}}$$

$$\lambda_x = \frac{2\pi c}{\sin \alpha \omega \sqrt{\epsilon_{\mu}}}; \quad \lambda_y = \frac{2\pi c}{\cos \alpha \omega \sqrt{\epsilon_{\mu}}}$$

$$b) \lambda_x = \frac{2\pi}{k_x}; \quad \cos \alpha = \frac{k_z}{k}; \quad \sin \alpha = \sqrt{1 - \cos^2 \alpha} = \sqrt{1 - \frac{k_z^2}{k^2}}$$

$$k = \frac{\omega}{c} \sqrt{\epsilon_{\mu}}; \quad k_x = k \sin \alpha = \sqrt{k^2 - k_z^2}; \quad \lambda_z = \frac{2\pi}{k_z}$$

$$k_z = \frac{2\pi}{\lambda_z}; \quad k_x = \sqrt{\frac{\omega^2}{c^2} \epsilon_{\mu} - \frac{4\pi^2}{\lambda_z^2}}$$

$$\lambda_x = \frac{2\pi}{\sqrt{\frac{\omega^2}{c^2} \epsilon_{\mu} - \frac{4\pi^2}{\lambda_z^2}}} = \frac{1}{\sqrt{\frac{\omega^2 \epsilon_{\mu}}{4\pi^2 c^2} - \frac{1}{\lambda_z^2}}} = \frac{\lambda_z}{\sqrt{\frac{\lambda_z^2 \omega^2 \epsilon_{\mu}}{4\pi^2 c^2} - 1}}$$

$$2) \vec{v}^{(z)} = \vec{v} \cos \alpha; \quad v = \frac{c}{\sqrt{\epsilon_{\mu}}}; \quad \lambda_x = \frac{\lambda}{\sin \alpha}; \quad \lambda = \frac{c}{\sqrt{\epsilon_{\mu}}} \frac{2\pi}{\omega}$$

$$\lambda_x = \frac{c}{\sqrt{\epsilon_{\mu}}} \cdot \frac{2\pi}{\omega} \cdot \frac{1}{\sin \alpha}; \quad \cos \alpha = \frac{v^{(z)}}{v} = v^{(z)} \frac{\sqrt{\epsilon_{\mu}}}{c}$$

$$\sin \alpha = \sqrt{1 - \frac{v^{(z)2} \epsilon_{\mu}}{c^2}}$$

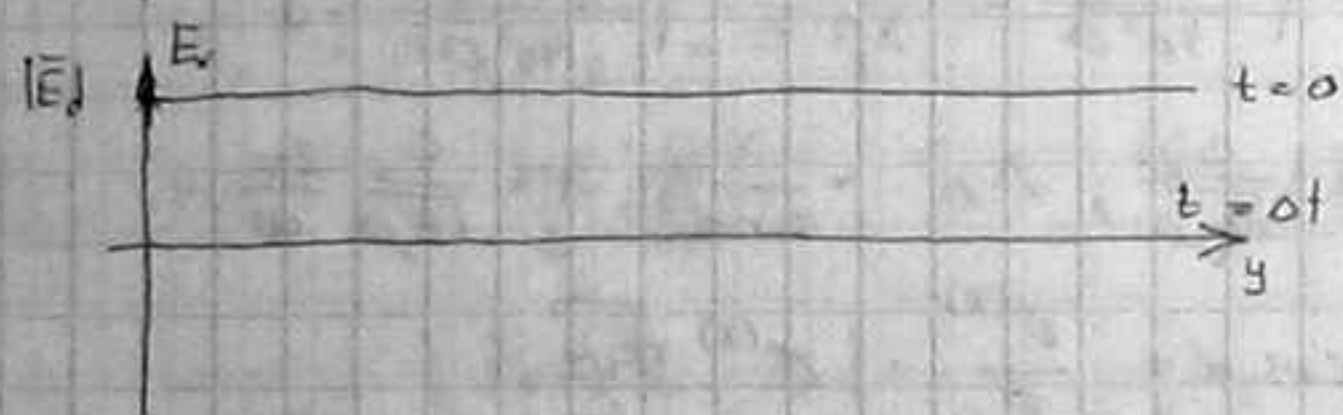
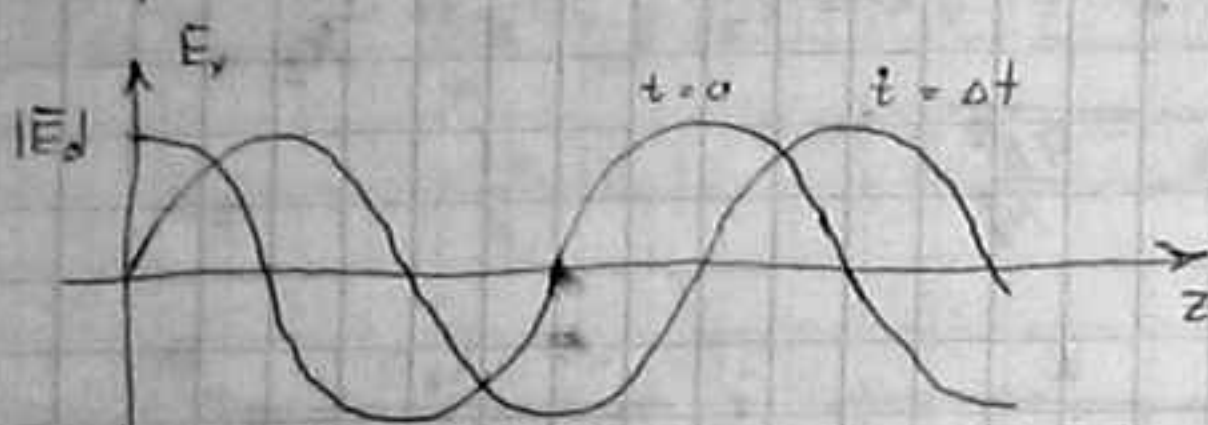
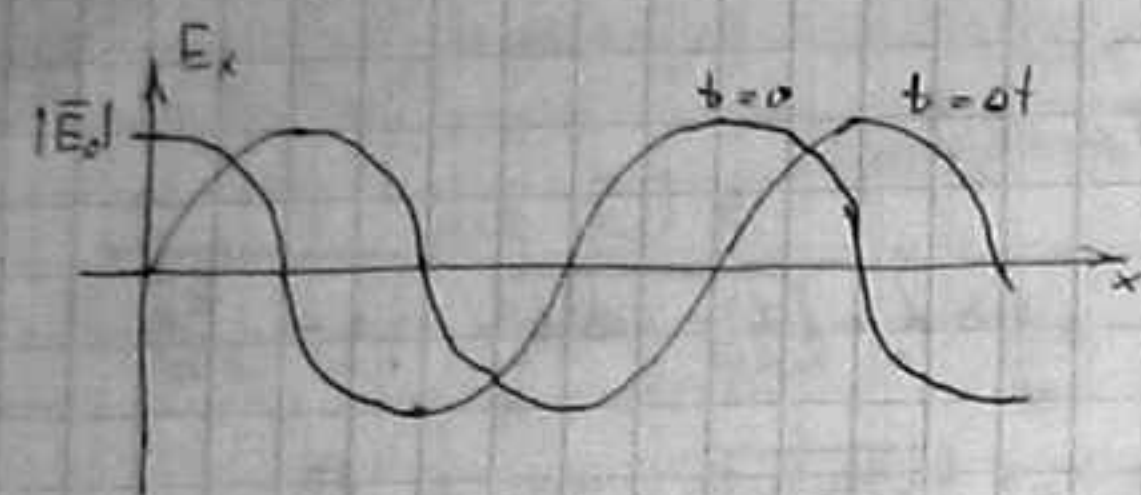
$$\omega = \frac{2\pi c}{\lambda_x \sqrt{\epsilon_{\mu}} \sin \alpha} = \frac{2\pi c}{\lambda_x \sqrt{\epsilon_{\mu}}} \left[ 1 - \frac{v^{(z)2} \epsilon_{\mu}}{c^2} \right]^{-1/2}$$

$$g) \text{ В случае } \vec{k} \parallel \vec{z}_0; \quad k_z = k; \quad k_x = 0$$

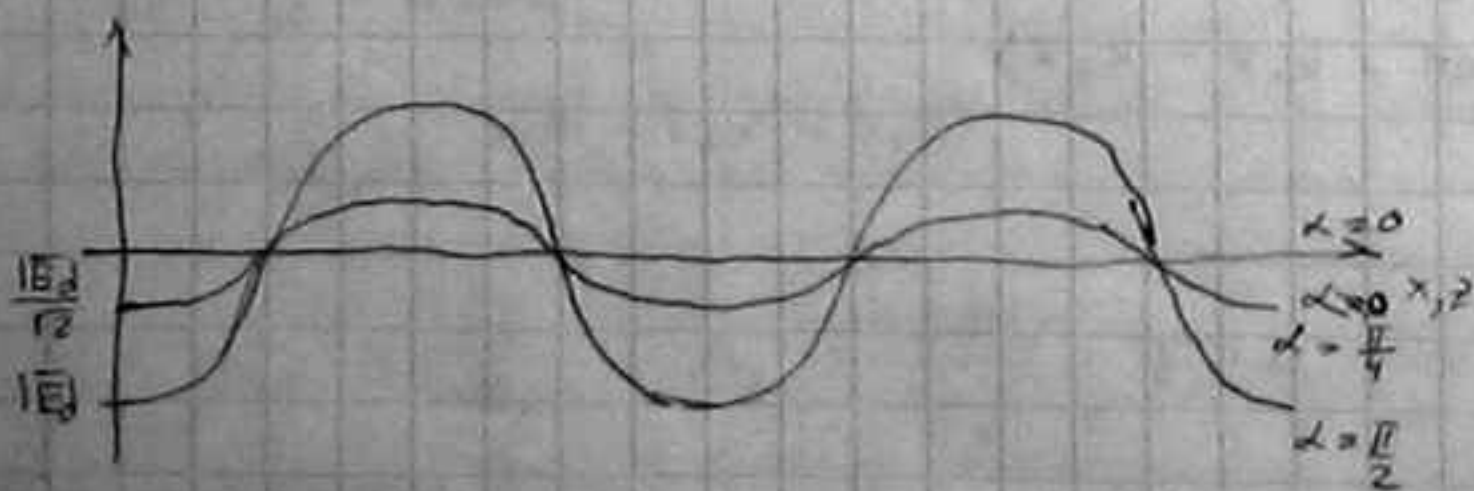
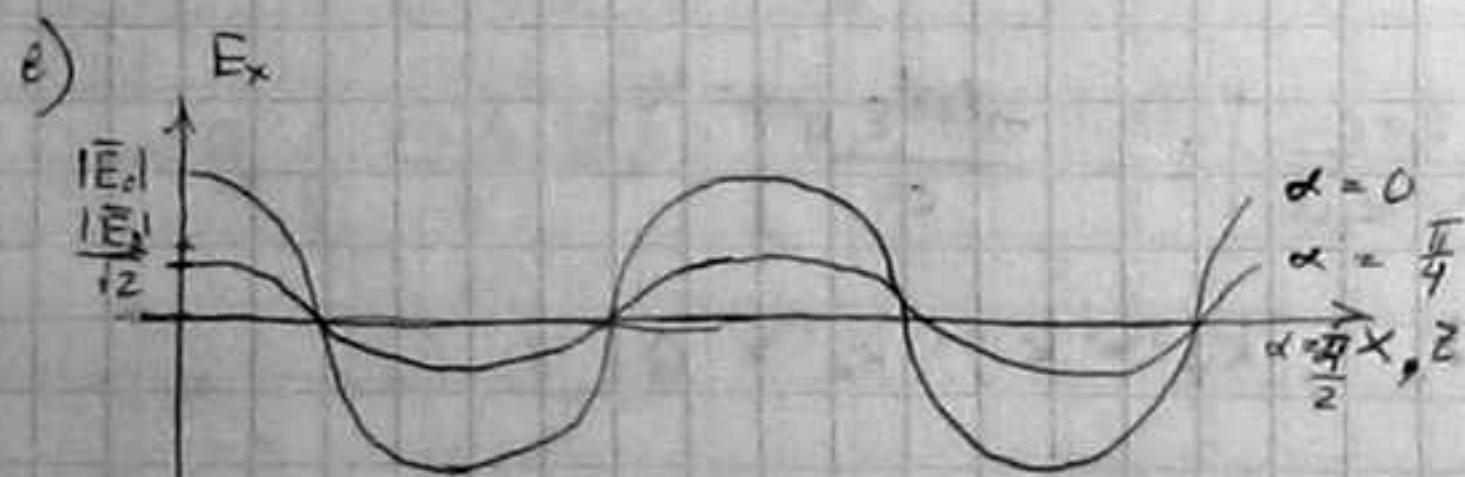
$$E_x = |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$

$$\neq H_y = |\vec{E}_0| e^{i(\omega t - k_x x - k_z z)}$$



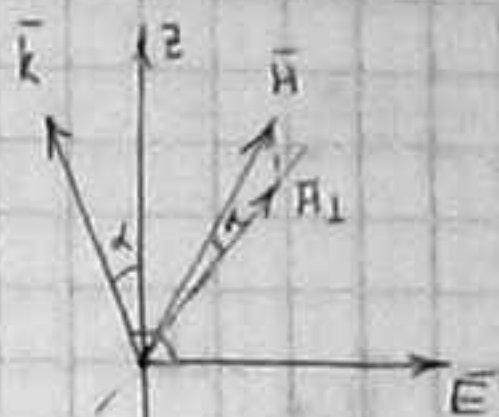


Аналогично и для  $H_y$



17,2

a) TE

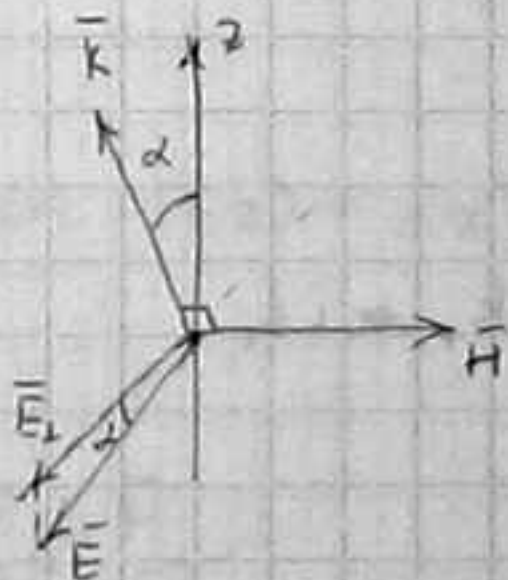


$$E = \eta H, \text{ где } \eta = \sqrt{\frac{\mu}{\epsilon}}$$

$$E_{\perp} = E, \quad H_{\perp} = H \cos \alpha; \quad \eta \frac{H_{\perp}}{\cos \alpha} = E_{\perp},$$

$$\eta_{\perp} = \frac{\eta}{\cos \alpha} = \sqrt{\frac{\mu}{\epsilon}} \frac{1}{\cos \alpha}$$

б) TM



$$E_{\perp} = E \cos \alpha; \quad H_{\perp} = H, \quad \eta H_{\perp} = \frac{E_{\perp}}{\cos \alpha}, \quad E_{\perp} = \eta \cos \alpha H_{\perp}$$

$$\eta_{\perp} = \sqrt{\frac{\mu}{\epsilon}} \cos \alpha$$

в) TEM

$$E_{\perp} = E, \quad H_{\perp} = H; \quad E_{\perp} = \eta H_{\perp}, \quad \eta_{\perp} = \sqrt{\frac{\mu}{\epsilon}}$$



N 7.3

$$S = \operatorname{Re} \frac{c}{8\pi} [\vec{E} \times \vec{H}^*]$$

$$\vec{E} = \vec{E}_0 e^{-ikr}$$

$$\vec{H} = \frac{\vec{E}}{\eta} = \sqrt{\frac{\epsilon}{\mu}} \vec{E}_0 e^{-ikr} = \vec{H}_0 e^{-ikr}$$

$$S = \operatorname{Re} \frac{c}{8\pi} [\vec{E}_0 e^{-ikr} \times \vec{H}_0 e^{ikr}] = \operatorname{Re} \frac{c}{8\pi} |\vec{E}_0|^2 \sqrt{\frac{\epsilon}{\mu}} =$$

$$= \frac{c}{8\pi} |\vec{E}_0|^2 \sqrt{\frac{\epsilon}{\mu}}, \quad |\vec{E}_0| = \left( \frac{8\pi}{c} \sqrt{\frac{\mu}{\epsilon}} \right)^{1/2}$$

$$|\vec{H}_0| = \sqrt{\frac{\epsilon}{\mu}} \left( \frac{8\pi}{c} \sqrt{\frac{\mu}{\epsilon}} \right)^{1/2} = \left( \frac{8\pi}{c} \sqrt{\frac{\epsilon}{\mu}} \right)^{1/2}$$

N 7.5

$$a) \vec{E} = \vec{E}_0 e^{-ikr}$$

$$k = \frac{\omega}{c} \sqrt{\epsilon\mu} = \frac{\omega}{c} \sqrt{\epsilon} = k_r - ik_i$$

$$\vec{E} = \vec{E}_0 e^{-k_i r} e^{-ik_r r}$$

Для скорости волнового фронта  $v = \frac{\omega}{k_r}; k_r = \frac{\omega}{v}$

Для затухания  $k_i h = 1, k_i = \frac{1}{h}$

$$\frac{\omega^2}{c^2} (\epsilon_r + i\epsilon_i) = k_r^2 - 2ik_r k_i - k_i^2$$

$$\frac{\omega^2}{c^2} \epsilon_r = k_r^2 - k_i^2, \quad \epsilon_r = \frac{c^2}{\omega^2} (k_r^2 - k_i^2)$$

$$\epsilon_r = \frac{c^2}{\omega^2} \left( \frac{\omega^2}{v^2} - \frac{1}{h^2} \right) = \frac{c^2}{v^2} - \frac{c^2}{\omega^2 h^2}$$

$$\frac{\omega^2}{c^2} \epsilon_i = -2k_r k_i; \quad \frac{\omega^2}{c^2} \epsilon_i = -2 \frac{\omega}{v} \frac{1}{h};$$

$$\epsilon_i = -2 \frac{c^2}{\omega^2} \frac{\omega}{v} \frac{1}{h} = -\frac{2c^2}{\omega v h}$$

N 7.6)

$$E_0 = \eta H_0 = \sqrt{\frac{\mu}{\epsilon}} H_0 = \frac{1}{\sqrt{\epsilon}} H_0$$

$$E_0 = H_0 e^{i\varphi} \quad \# \quad \frac{H_0}{p} e^{i\varphi}$$

$$\frac{1}{\sqrt{\epsilon}} = p e^{i\varphi}; \quad \frac{1}{\epsilon} = p^2 e^{2i\varphi}; \quad \epsilon = \frac{1}{p^2} e^{-2i\varphi}$$

$$\epsilon_r + i\epsilon_i = \frac{1}{p^2} (\cos 2\varphi - i \sin 2\varphi)$$

$$\epsilon_r = \frac{\cos 2\varphi}{p^2}, \quad \epsilon_i = -\frac{\sin 2\varphi}{p^2}$$

N 7.7

$$E_y = E_0 \exp \{ i(\omega t - hz) - kx \}$$

$$E_x = E_z = 0$$

$$\operatorname{rot} \vec{E} = -\frac{i\omega}{c} \mu \vec{H}$$

$$\vec{H} = \frac{ic}{\mu\omega} \operatorname{rot} \vec{E}$$

$$\operatorname{rot} \vec{E} = \begin{vmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ \frac{\partial}{\partial x} & 0 & \frac{\partial}{\partial z} \\ 0 & E_y & 0 \end{vmatrix} =$$

$$= -\vec{e}_x \frac{\partial E_y}{\partial z} + \vec{e}_z \frac{\partial E_y}{\partial x}$$

$$H_x = -\frac{ic}{\mu\omega} (-ih) E_0 \exp \{ -ihz - kx \} =$$



$$= -\frac{ch}{\omega\mu} E_0 \exp\{-i\hbar z - kx\} = \frac{ch}{\omega\mu} E_0 \exp\{-i\hbar z - kx + i\pi\}$$

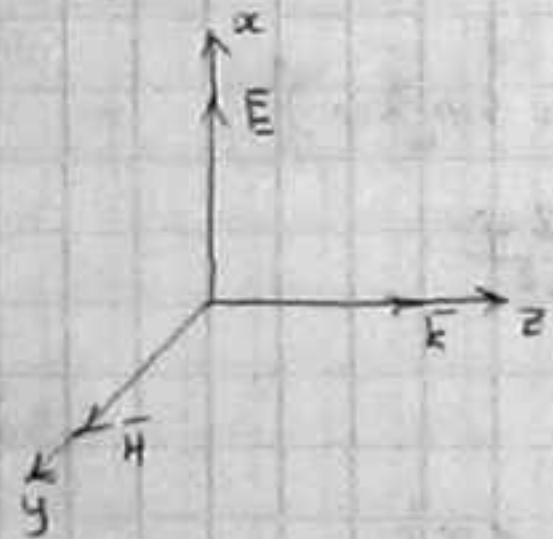
$$\mathcal{H}_2 = -\frac{ic}{\mu\omega} (-k) \exp\{-i\hbar z - kx\} =$$

$$= \frac{ick}{\mu\omega} \exp\{-i\hbar z - kx\} = \frac{ck}{\mu\omega} \exp\{-i\hbar z - kx + i\frac{\pi}{2}\}$$

$$k \neq \frac{\omega^2}{c^2} \mu$$

н 7.11

Для поперечной волны



$$\begin{cases} \text{rot } \vec{E} = -\frac{1}{c} i\omega\mu \vec{H} & | \text{rot} \\ \text{rot } \vec{H} = \frac{1}{c} i\omega \vec{D} \end{cases}$$

$$-\Delta \vec{E} = -\frac{1}{c} i\omega\mu \left( \frac{1}{c} i\omega \vec{D} \right)$$

$$\Delta \vec{E} + \frac{\omega^2}{c^2} \mu \vec{D} = 0$$

$$\vec{D} = \epsilon \vec{E} + \delta^2 \text{grad div } \vec{E}$$

$$\epsilon = 1 - \frac{\omega_p^2}{\omega^2}; \quad \delta = \frac{v_T}{\omega}$$

$$\vec{E} = E_x \vec{x}_0 = \vec{x}_0 E_0 e^{-ikz}; \quad \text{div } \vec{E} = 0$$

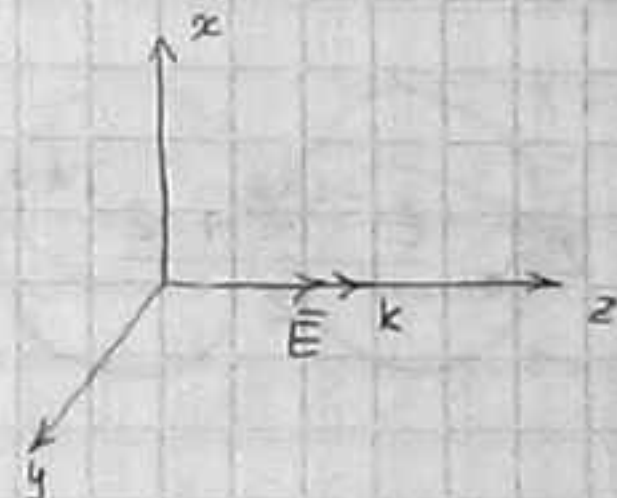
$$\vec{D} = \epsilon \vec{x}_0 E_x$$

$$\Delta E_x + \frac{\omega^2}{c^2} \mu \epsilon E_x = 0$$

$$k^2 = \frac{\omega^2}{c^2} \mu \epsilon = \frac{\omega^2}{c^2} \mu \left( 1 - \frac{\omega_p^2}{\omega^2} \right) = \frac{\omega^2 \mu}{c^2} - \frac{\mu \omega_p^2}{c^2}$$

$$\omega^2 = \frac{k^2 c^2}{\mu} + \omega_p^2$$

Для продольной волны



$$\begin{cases} \text{rot } \vec{E} = -\frac{i\omega}{c} \mu \vec{H} & | \text{rot} \\ \text{rot } \vec{H} = \frac{i\omega}{c} \vec{D} = \frac{i\omega}{c} (\epsilon \vec{E} + \delta^2 \text{grad div } \vec{E}) \end{cases}$$

$$\text{rot rot } \vec{E} = -\frac{i\omega}{c} \mu \text{rot } \vec{H}$$

$$\text{grad div } \vec{E} - \Delta \vec{E} = -\frac{i\omega}{c} \mu \frac{i\omega}{c} (\epsilon \vec{E} + \delta^2 \text{grad div } \vec{E})$$

$$\Delta \vec{E} + \left( \frac{\omega^2}{c^2} \mu \delta^2 - 1 \right) \text{grad div } \vec{E} + \frac{\omega^2}{c^2} \epsilon \mu \vec{E} = 0$$

$$\vec{E} = \vec{x}_0 E_z e^{-ikz}$$

$$-k^2 E_z e^{-ikz} + \left( \frac{\omega^2}{c^2} \mu \delta^2 - 1 \right) (-k^2 E_z e^{-ikz}) + \frac{\omega^2}{c^2} \epsilon \mu E_z e^{-ikz} = 0$$

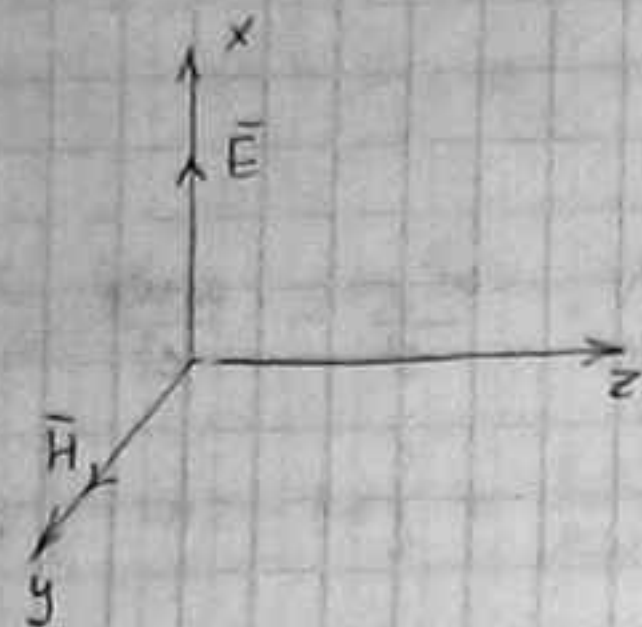
$$-k^2 \frac{\omega^2}{c^2} \mu \delta^2 + \frac{\omega^2}{c^2} \epsilon \mu = 0; \quad k^2 \delta^2 = \epsilon$$

$$k^2 \frac{v_T^2}{\omega^2} = 1 - \frac{\omega_p^2}{\omega^2}$$

$$\omega^2 = \omega_p^2 + k^2 v_T^2$$



N 7.13



Стоящая волна

$$E_x = E_0 e^{-ikz} + E_0 e^{ikz} = 2E_0 \frac{e^{-ikz} + e^{ikz}}{2} =$$

$$= 2E_0 \cos kz$$

$$E_x = 2E_0 \cos kz \cos \omega t - \text{реальная волна}$$

$$\text{rot } \vec{E} = -\frac{i\omega\mu}{c} \vec{H}, \quad \vec{H} = \frac{ic}{\omega\mu} \text{rot } \vec{E}$$

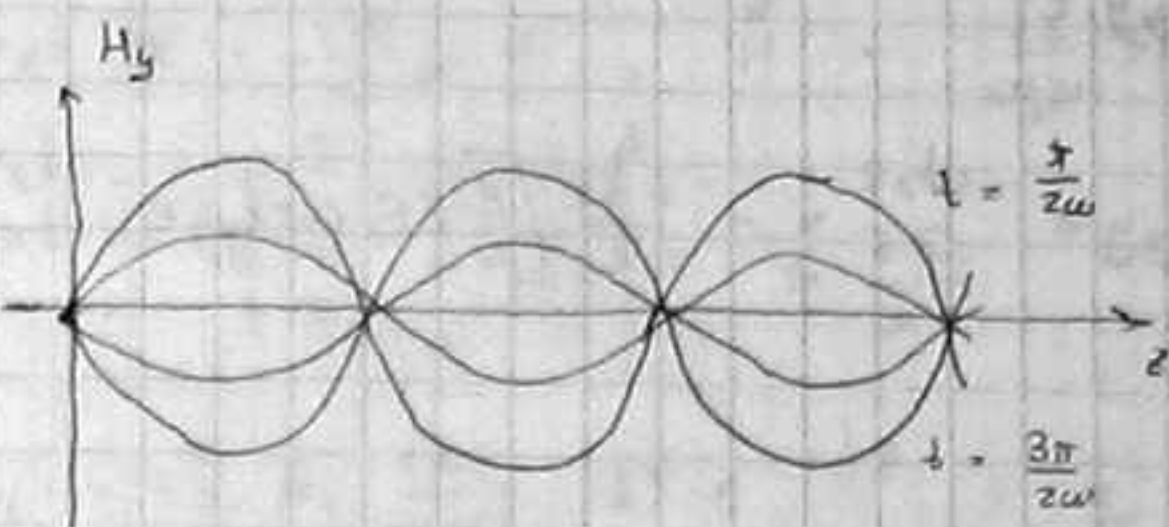
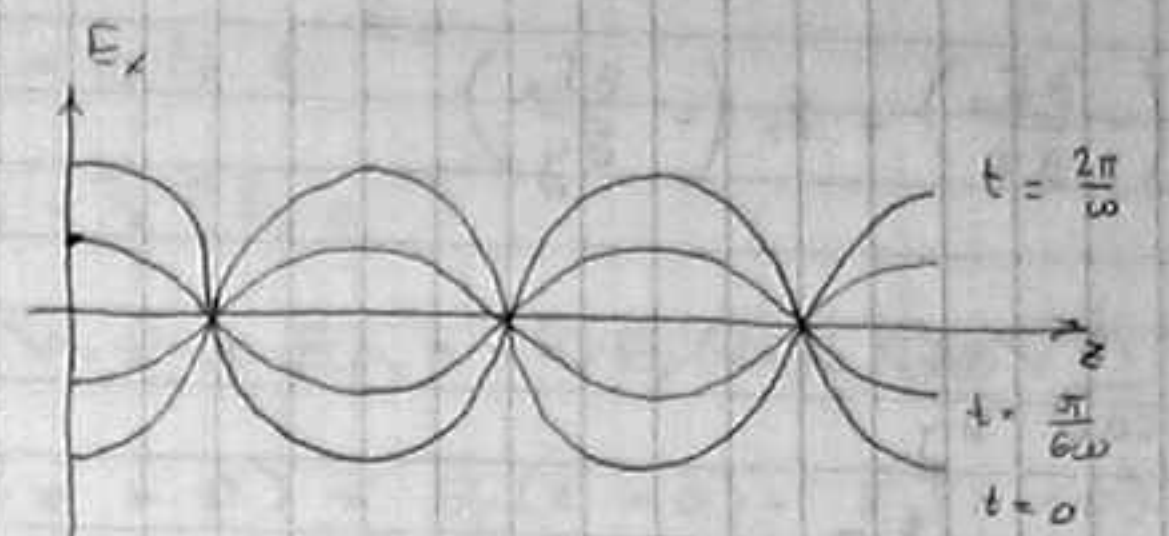
$$\text{rot } \vec{E} = \begin{vmatrix} \vec{x}_0 & \vec{y}_0 & \vec{z}_0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = -\vec{y}_0 \left( -\frac{\partial E_x}{\partial z} \right), \text{ т.е.}$$

$$H_y = \frac{ic}{\omega\mu} \frac{\partial E_x}{\partial z} = -\frac{ic}{\omega\mu} 2E_0 k \sin kz = \frac{2E_0 kc}{\omega\mu} e^{-i\frac{\pi}{2}} \sin kz =$$

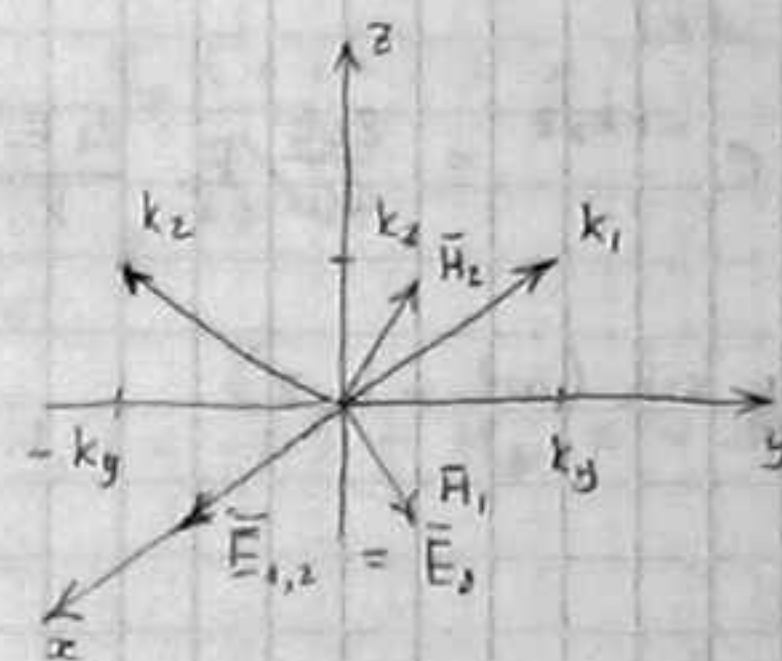
$$= \frac{2E_0}{\mu} \frac{c}{\omega} \frac{\omega}{c} \sqrt{\epsilon\mu} e^{-i\frac{\pi}{2}} \sin kz = 2E_0 \sqrt{\frac{\epsilon}{\mu}} \sin kz e^{-i\frac{\pi}{2}}$$

Две реальных волны

$$H_y = 2E_0 \sqrt{\frac{\epsilon}{\mu}} \sin kz \cos \left( \omega t - \frac{\pi}{2} \right) = 2E_0 \sqrt{\frac{\epsilon}{\mu}} \sin kz \sin \omega t$$



N 7.14



$$E_x = E_0 e^{-ik_y y - ik_z z} + E_0 e^{ik_y y - ik_z z} =$$

$$= 2E_0 e^{-ik_z z} \frac{e^{-ik_y y} + e^{ik_y y}}{2} = 2E_0 \cos k_y y e^{-ik_z z}$$

$$\text{Re} \{ E_x e^{i\omega t} \} = 2E_0 \cos k_y y \cos (\omega t - k_z z)$$

$$\vec{H} = \frac{ic}{\omega\mu} \text{rot } \vec{E}$$



$$\text{rot } \vec{E} = \begin{vmatrix} \vec{x}_0 & \vec{y}_0 & \vec{z}_0 \\ 0 & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = -\vec{y}_0 \left( -\frac{\partial E_x}{\partial z} \right) + \vec{z}_0 \left( -\frac{\partial E_x}{\partial y} \right)$$

$$H_y = \frac{ic}{\omega \mu} \frac{\partial E_x}{\partial z}, \quad H_z = -\frac{ic}{\omega \mu} \frac{\partial E_x}{\partial y}$$

$$H_y = \frac{ic}{\omega \mu} (-k_z) 2E_0 \cos k_y y e^{-ik_z z} =$$

$$= -i^2 \left( \frac{c}{\omega} \frac{1}{\sqrt{\epsilon \mu}} \right) \sqrt{\frac{\epsilon}{\mu}} k_z 2E_0 \cos k_y y e^{-ik_z z} = -\frac{k_z}{k} \sqrt{\frac{\epsilon}{\mu}} 2E_0 \cos k_y y e^{-ik_z z} =$$

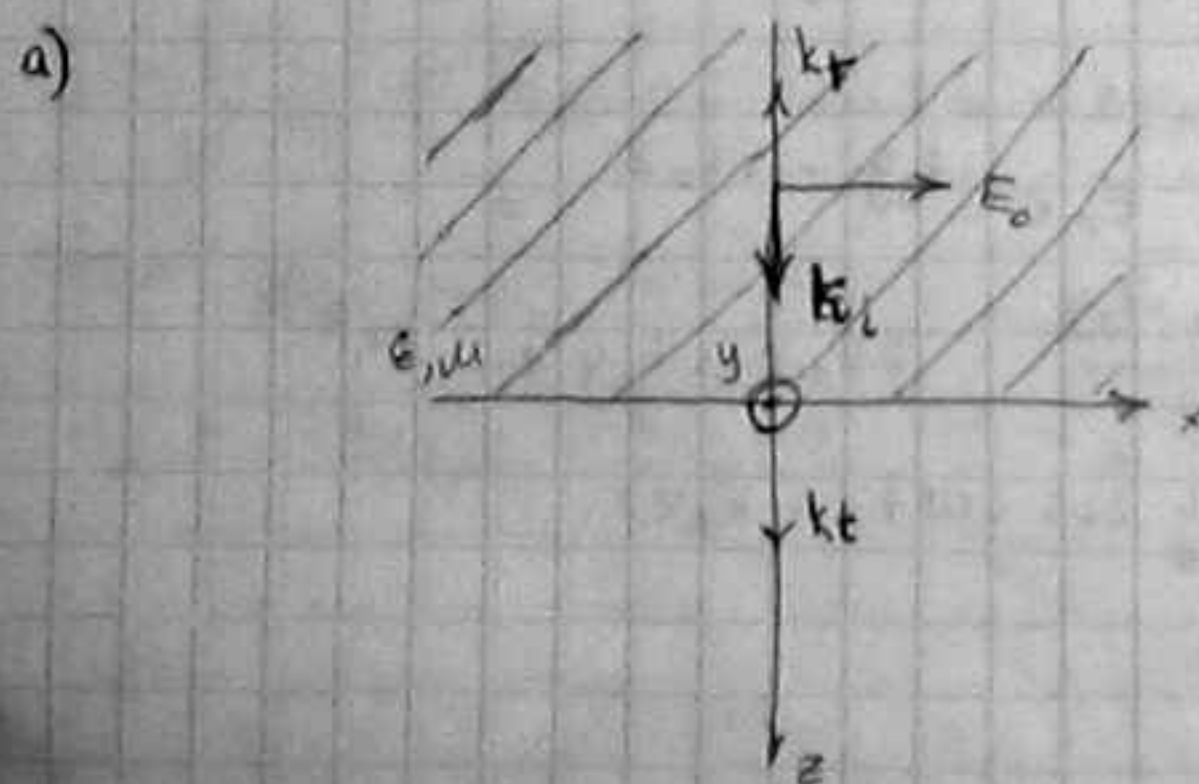
$$= \frac{k_z}{k} \sqrt{\frac{\epsilon}{\mu}} 2E_0 \cos k_y y e^{-ik_z z}$$

$$\text{Re} \{ H_y e^{i\omega t} \} = \frac{k_z}{k} \sqrt{\frac{\epsilon}{\mu}} 2E_0 \cos k_y y \cos(\omega t - k_z z)$$

$$H_z = -\frac{ic}{\omega \mu} (-2E_0 k_y \sin k_y y) e^{-ik_z z} = \frac{2ic}{\omega \mu} E_0 \frac{2iE_0 k_y \sin k_y y}{k} e^{-ik_z z}$$

$$\text{Re} \{ H_z e^{i\omega t} \} = -\frac{2E_0 k_y}{k} \sin k_y y \sin(\omega t - k_z z)$$

N 8.1



$$\vec{E}_i = \vec{x}_0 E_0 e^{-ik_z z}$$

$$\vec{E}_r = \vec{x}_0 \Gamma E_0 e^{+ik_z z}$$

$$\vec{E} = \vec{E}_i + \vec{E}_r = \vec{x}_0 E_0 (1 + \Gamma) e^{-ik_z z}$$

$$\vec{E}(z=0) = \vec{E}_i(z=0) + \vec{E}_r(z=0) = \vec{x}_0 E_0 (1 + \Gamma)$$

$$\text{rot } \vec{E} = -\frac{i\omega}{c} \mu \vec{H}, \quad \vec{H} = \frac{ic}{\omega \mu} \text{rot } \vec{E}$$

$$\text{rot } \vec{E} = \begin{vmatrix} \vec{x}_0 & \vec{y}_0 & \vec{z}_0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = \vec{y}_0 \frac{\partial E_x}{\partial z}$$

$$\vec{H}_i = +\vec{y}_0 E_0 \frac{ic}{\omega \mu} (-ik) e^{-ik_z z} = \vec{y}_0 E_0 \frac{c}{\omega \mu} \frac{\omega}{c} \sqrt{\frac{\epsilon}{\mu}} e^{-ik_z z} =$$

$$= \vec{y}_0 E_0 \sqrt{\frac{\epsilon}{\mu}} e^{-ik_z z}$$

$$\vec{H}_r = \vec{y}_0 \frac{ic}{\omega \mu} E_0 (+ik) \Gamma e^{ik_z z} = -\vec{y}_0 E_0 \frac{c}{\omega \mu} \frac{\omega}{c} \sqrt{\frac{\epsilon}{\mu}} e^{ik_z z} =$$

$$= -\vec{y}_0 \Gamma E_0 \sqrt{\frac{\epsilon}{\mu}} e^{ik_z z}$$

$$\vec{H}(z=0) = \vec{H}_i(z=0) + \vec{H}_r(z=0) = \vec{y}_0 E_0 \sqrt{\frac{\epsilon}{\mu}} (1 - \Gamma)$$

$$\eta_s = \frac{E_x(0)}{H_y(0)}$$

$$\eta_s = \frac{E_0 (1 + \Gamma)}{E_0 \sqrt{\frac{\epsilon}{\mu}} (1 - \Gamma)}; \quad \eta_s \sqrt{\frac{\epsilon}{\mu}} = \frac{1 + \Gamma}{1 - \Gamma}; \quad \eta_\omega = \sqrt{\frac{\mu}{\epsilon}}$$

$$\frac{\eta_s}{\eta_\omega} = \frac{1 + \Gamma}{1 - \Gamma}, \quad \eta_s - \eta_s \Gamma = \eta_\omega + \eta_\omega \Gamma, \quad \eta_s - \eta_\omega = (\eta_s + \eta_\omega) \Gamma$$

$$\Gamma = \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega}$$



$$\begin{aligned} \bar{E}_i &= \bar{x}_0 E_0 e^{-ikz} \\ \bar{E}_r &= \bar{x}_0 \Gamma E_0 e^{ikz} \\ \bar{H}_i &= \bar{y}_0 \frac{E_0}{\eta_\omega} e^{-ikz} \\ \bar{H}_r &= -\bar{y}_0 \frac{\Gamma E_0}{\eta_\omega} e^{ikz} \\ \eta(L) &= \left( \frac{E_x}{E_y} \right)_{z=-L} \end{aligned}$$

$$\begin{aligned} E_x(z = -L) &= E_0 e^{ikL} + \Gamma E_0 e^{-ikL} = E_0 \left( e^{ikL} + \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega} e^{-ikL} \right) = \\ &= E_0 \frac{E_0}{\eta_s + \eta_\omega} \left( \eta_s e^{ikL} + \eta_\omega e^{ikL} + \eta_s e^{-ikL} - \eta_\omega e^{-ikL} \right) = \\ &= \frac{E_0}{\eta_s + \eta_\omega} \left( \eta_s (e^{ikL} + e^{-ikL}) + \eta_\omega (e^{ikL} - e^{-ikL}) \right) = \\ &= \frac{E_0}{\eta_s + \eta_\omega} (2\eta_s \cos kL + 2i\eta_\omega \sin kL) \end{aligned}$$

$$\begin{aligned} H_y(z = -L) &= \frac{1}{\eta_\omega} E_0 e^{ikL} - \frac{1}{\eta_\omega} \Gamma E_0 e^{-ikL} = \\ &= \frac{E_0}{\eta_\omega} e^{ikL} - \frac{E_0}{\eta_\omega} \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega} e^{-ikL} = \frac{E_0}{\eta_\omega (\eta_s + \eta_\omega)} (\eta_s e^{ikL} + \eta_\omega e^{ikL} - \eta_s e^{-ikL} + \eta_\omega e^{-ikL}) = \\ &= \frac{E_0}{\eta_\omega (\eta_s + \eta_\omega)} (2i\eta_s \sin kL + 2\eta_\omega \cos kL) \end{aligned}$$

$$\eta(L) = \eta_\omega \frac{2\eta_s \cos kL + 2i\eta_\omega \sin kL}{2\eta_\omega \cos kL + 2i\eta_s \sin kL} = \eta_\omega \frac{\eta_s + i\eta_\omega \operatorname{tg} kL}{\eta_\omega + i\eta_s \operatorname{tg} kL}$$

$$b) E = E_0 e^{-ikz} + \Gamma E_0 e^{ikz}, \quad \Gamma = |\Gamma| e^{i\varphi}$$

$$|E|^2(z) = E(z) \cdot E^*(z)$$

$$E^* = E_0 e^{ikz} + E_0 |\Gamma| e^{-i\varphi} e^{-ikz}$$

$$\begin{aligned} E \cdot E^* &= E_0^2 + |\Gamma| E_0^2 e^{i\varphi} e^{2ikz} + |\Gamma|^2 E_0^2 e^{-i\varphi} e^{-2ikz} + E_0^2 |\Gamma|^2 = \\ &= E_0^2 + 2|\Gamma| E_0^2 \frac{e^{i(\varphi+2kz)} + e^{-i(\varphi+2kz)}}{2} + E_0^2 |\Gamma|^2 = \\ &= E_0^2 (1 + 2|\Gamma| \cos(\varphi + 2kz) + |\Gamma|^2) \end{aligned}$$

$$|E|_{\max}^2 \leftarrow E_0^2 \text{ так как } \cos(\varphi + 2kz) = 1$$

$$|E|_{\max}^2 = E_0^2 (1 + 2|\Gamma| + |\Gamma|^2) = E_0^2 (1 + |\Gamma|)^2$$

$$|E|_{\min}^2 \text{ так как } \cos(\varphi + 2kz) = -1$$

$$|E|_{\min}^2 = E_0^2 (1 - 2|\Gamma| + |\Gamma|^2) = E_0^2 (1 - |\Gamma|)^2$$

$$KCB = \frac{|E|_{\max}^2}{|E|_{\min}^2} = \left( \frac{1 + |\Gamma|}{1 - |\Gamma|} \right)^2$$

$$e) KCB = 1, \text{ значит } 1 + |\Gamma| = 1 - |\Gamma|, |\Gamma| = -|\Gamma| \Rightarrow$$

$$|\Gamma| = 0$$

$$\Gamma = \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega} \Rightarrow \eta_s = \eta_\omega = \sqrt{\frac{\mu}{\epsilon}}$$

$$KCB = \infty, \text{ значит } 1 - |\Gamma| = 0, |\Gamma| = 1$$

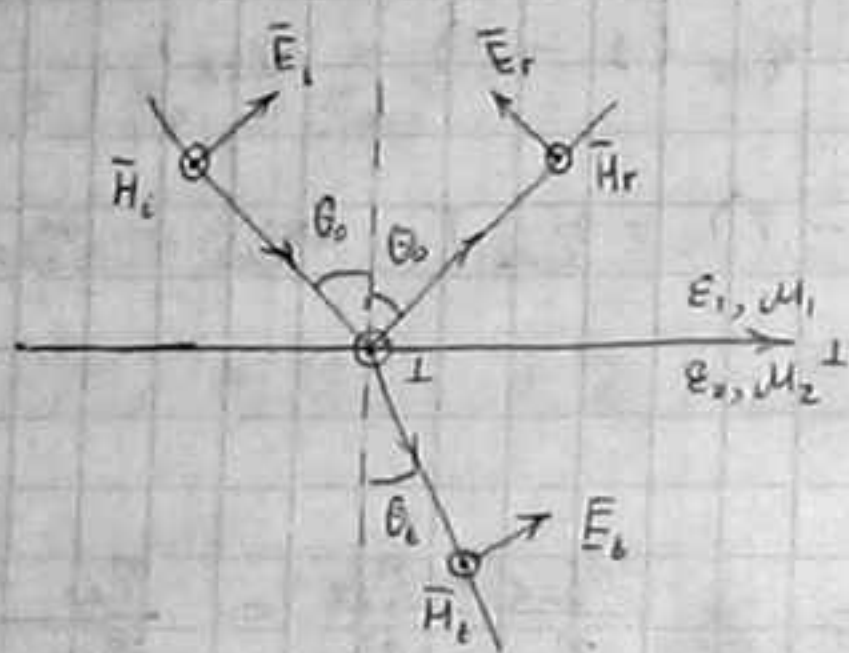
$$\Gamma = \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega}, \text{ при } \eta_s \rightarrow \infty \quad \Gamma = \lim_{\eta_s \rightarrow \infty} \frac{\eta_s - \eta_\omega}{\eta_s + \eta_\omega} =$$

$$= \lim_{\eta_s \rightarrow \infty} \frac{\eta_s}{\eta_s} = 1$$



$$\eta_{12} = \sqrt{\frac{\mu_2}{\epsilon_2}} \left( 1 - \frac{\epsilon_1 \mu_1}{\epsilon_2 \mu_2} \sin^2 \theta_0 \right)^{-1/2}$$

б) ТМ



$$H_{1i} = \frac{E_i}{\eta_{11}}, \quad H_{1r} = -\frac{E_r}{\eta_{11}}, \quad H_{1t} = \frac{E_t}{\eta_{12}}$$

$$E_{1i} = E_{01}; \quad E_{1r} = -\Gamma E_{01}; \quad E_{1t} = T E_{01}$$

$$H_{1i} = H_{12}, \quad \frac{E_{01} + \Gamma E_{01}}{\eta_{11}} = \frac{T E_{01}}{\eta_{12}}$$

$$E_{1i} = E_{12}, \quad E_{01} - \Gamma E_{01} = T E_{01}$$

$$\begin{cases} \eta_{12}(1 + \Gamma) = T \eta_{11} \\ 1 - \Gamma = T \end{cases}$$

$$\Gamma = 1 - T, \quad \eta_{12}(2 - T) = T \eta_{11};$$

$$H_{1i} = H_0, \quad H_{1r} = \Gamma H_0, \quad H_{1t} = T H_0$$

$$E_{1i} = \eta_{11} H_{1i} = \eta_{11} H_0; \quad E_{1r} = -\eta_{11} H_{1r} = -\eta_{11} \Gamma H_0$$

$$E_{1t} = \eta_{12} H_{1t} = \eta_{12} T H_0$$

$$H_{1i} = H_{12}, \quad E_{1i} = E_{12}$$

$$1 + \Gamma = T$$

$$\eta_{11} - \eta_{11} \Gamma = T \eta_{12}$$

$$\Gamma = T - 1; \quad 2\eta_{11} - \eta_{11} T = T \eta_{12};$$

$$T = \frac{2\eta_{11}}{\eta_{11} + \eta_{12}}$$

$$\Gamma = \frac{2\eta_{11}}{\eta_{11} + \eta_{12}} - \frac{\eta_{11} + \eta_{12}}{\eta_{11} + \eta_{12}} = \frac{\eta_{11} - \eta_{12}}{\eta_{11} + \eta_{12}}$$

Для первой среды

$$E_0 = \eta_1 H_0, \quad H_1 = H_0, \quad E_1 = E_0 \cos \theta_0$$

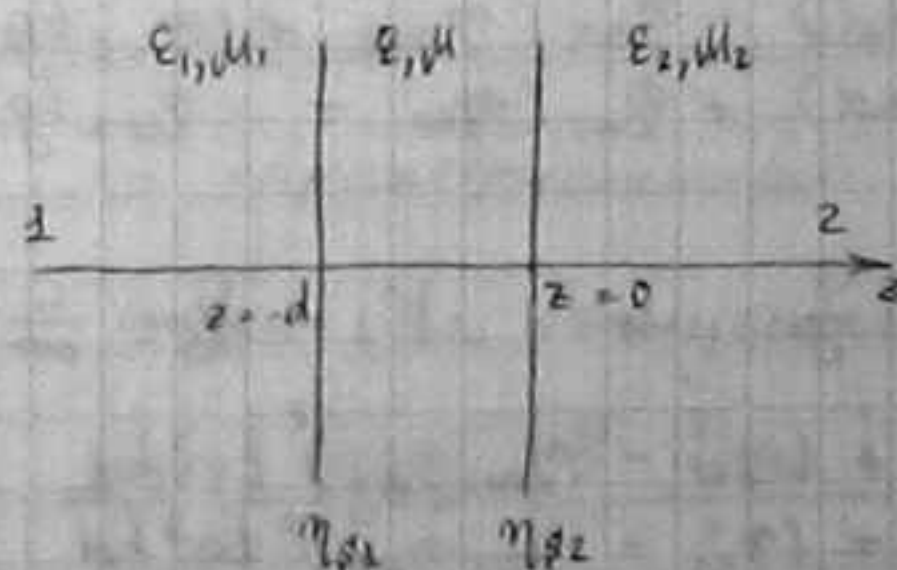
$$\eta_1 = \frac{E_1}{H_1} = \frac{\eta_1 H_0 \cos \theta_0}{H_0} = \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_0$$

Для второй среды

$$E_0 = \eta_2 H_0; \quad H_1 = H_0, \quad E_1 = E_0 \cos \theta_t$$

$$\eta_2 = \sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_t = \sqrt{\frac{\mu_2}{\epsilon_2}} \left( 1 - \frac{\epsilon_1 \mu_1}{\epsilon_2 \mu_2} \sin^2 \theta_0 \right)^{1/2}$$

№ 8,3





Т.к. в области 2 существует плоская волна  
движущаяся по z, где все  $\eta_{xz} = \sqrt{\frac{\mu_2}{\epsilon_2}}$

На границе  $\eta_{f2} = \eta_{w2} = \sqrt{\frac{\mu_2}{\epsilon_2}}$

Формула пересчета импеданса

$$\eta(d) = \eta_w \frac{\eta_{f2} + i \eta_w \operatorname{tg} kd}{\eta_w + i \eta_{f2} \operatorname{tg} kd}, \text{ где}$$

$$\eta_w = \sqrt{\frac{\mu}{\epsilon}}, \quad k = \frac{\omega}{c} \sqrt{\epsilon \mu}$$

На границе слоя со средой 1  $\eta_{f1} = \eta(d)$

$$\eta_{f2} = \eta_w \frac{\eta_{f2} + i \eta_w \operatorname{tg} kd}{\eta_w + i \eta_{f2} \operatorname{tg} kd} = \eta_w \frac{\eta_{w2} + i \eta_w \operatorname{tg} kd}{\eta_w + i \eta_{w2} \operatorname{tg} kd}$$

Для коэффициента  $\Gamma$ :

$$\Gamma = \frac{\eta_{f1} - \eta_{w1}}{\eta_{f1} + \eta_{w1}}, \text{ где } \eta_{w1} = \sqrt{\frac{\mu_1}{\epsilon_1}}$$

$$\Gamma = \frac{\eta_w \frac{\eta_{w2} + i \eta_w \operatorname{tg} kd}{\eta_w + i \eta_{w2} \operatorname{tg} kd} - \eta_{w1}}{\eta_w \frac{\eta_{w2} + i \eta_w \operatorname{tg} kd}{\eta_w + i \eta_{w2} \operatorname{tg} kd} + \eta_{w1}} =$$

$$= \frac{\eta_w (\eta_{w2} + i \eta_w \operatorname{tg} kd) - \eta_{w1} (\eta_{w2} + i \eta_{w2} \operatorname{tg} kd)}{\eta_w (\eta_{w2} + i \eta_w \operatorname{tg} kd) + \eta_{w1} (\eta_w + i \eta_{w2} \operatorname{tg} kd)}$$

$$\Gamma = 0, \text{ если } |\Gamma| = \frac{0}{c} \text{ либо } |\Gamma| = \frac{0}{\infty}$$

$$|\Gamma|^2 = \frac{\eta_w^2 (\eta_{w2} - \eta_{w1})^2 + (\eta_w^2 - \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd}{\eta_w^2 (\eta_{w2} + \eta_{w1})^2 + (\eta_w^2 + \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd}$$

$$\eta_{w2} \neq \eta_{w1} \Rightarrow \eta_w^2 (\eta_{w2} - \eta_{w1})^2 + (\eta_w^2 - \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd \neq 0$$

Отношение  $\Gamma = \frac{0}{c}$  невозможно тогда было возможно

отношение  $\Gamma = \frac{0}{\infty}$ , необходимо  $(\eta_w^2 - \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd = 0$ ,

т.е.  $\eta_w = \sqrt{\eta_{w1} \eta_{w2}}$  и  $\operatorname{tg}^2 kd = \infty$ ,  $kd = \frac{\pi}{2} (2n+1)$ , где

$n = 0, 1, 2, \dots$ , в ином при  $\eta_{w2} \neq \eta_{w1}$

$$kd = \frac{\pi}{2} (2n+1), \quad n = 0, 1, 2, \dots \text{ и } \eta_w = \sqrt{\eta_{w1} \eta_{w2}}$$

Если  $\eta_{w2} = \eta_{w1}$ , то при  $\Gamma = \frac{0}{\infty}$ , невозможно,

$$\text{т.к. } \eta_w^2 (\eta_{w2} - \eta_{w1})^2 + (\eta_w^2 - \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd =$$

$$= (\eta_w^2 - \eta_{w1} \eta_{w2})^2 \operatorname{tg}^2 kd$$

Рассмотрим отношение  $\Gamma = \frac{0}{c}$ , т.е.  $(\eta_w^2 - \eta_{w1} \eta_{w2}) = 0$

либо  $\operatorname{tg} kd = 0$ :

$$\operatorname{tg} kd = 0 \Rightarrow kd = \frac{\pi}{4} (2n+1), \text{ где } n = 0, 1, 2, \dots$$

$$\eta_w^2 - \eta_{w1} \eta_{w2} = 0 \Rightarrow \eta_w = \eta_{w1} \text{ либо}$$

и/или

$|\Gamma|$  не будет зависеть от толщины слоя d, если

этот слой d по характеристикам является первой

или второй средой, т.е. полукруг нормальное падение

на границу двух сред. В итоге должно быть

$$\eta_w = \eta_{w1}, \text{ либо } \eta_w = \eta_{w2}$$



19.1

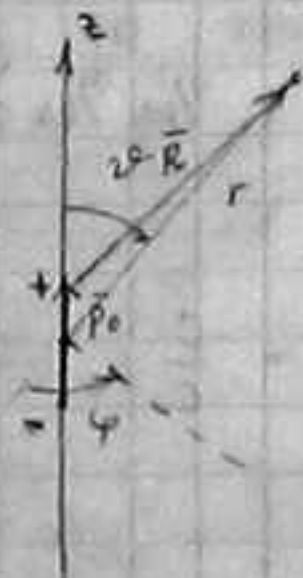
$$a) \bar{p} = \bar{p}_0 \exp \{i\omega t\}; |\bar{p}| = q|\bar{E}|, \quad q\bar{E} = p_0 \exp \{i\omega t\}$$

$$\dot{q} = i\omega \frac{p_0}{l} \exp \{i\omega t\}; \quad kl \ll 1 \Rightarrow$$

$$\bar{j} = \dot{q} \delta(x) \delta(y) = -i\omega \frac{\bar{p}_0}{l} \exp \{i\omega t\} \delta(x) \delta(y)$$

$$\mu = 1; \quad \epsilon = 1$$

$$\bar{A} = \frac{1}{c} \int \frac{\bar{j} e^{-ikR}}{R} dV$$



$$\bar{p}_0 = p_0 \bar{z}_0$$

$$R \gg l \Rightarrow R \approx r \Rightarrow$$

$$\bar{A} = -\frac{1}{c} \int_{-l/2}^{l/2} \frac{i\omega \frac{p_0 \bar{z}_0}{l} e^{-ikr}}{r} dz = -\frac{i\omega p_0 \bar{z}_0}{lc} \frac{e^{-ikr}}{r} \int_{-l/2}^{l/2} dz =$$

$$= -\frac{i\omega p_0 \bar{z}_0}{lc} \frac{e^{-ikr}}{r} l = -\bar{z}_0 \frac{i\omega}{c} p_0 e^{-ikr}$$

$$\bar{H} = \bar{B} = \text{rot } \bar{A}$$

$$\text{rot } \bar{A} = \begin{vmatrix} \bar{x}_0 & \bar{y}_0 & \bar{z}_0 \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & 0 & A_0 \end{vmatrix} = \bar{x}_0 \frac{\partial A_0}{\partial y} - \bar{y}_0 \frac{\partial A_0}{\partial x} =$$

$$= \bar{x}_0 \frac{\partial A_0}{\partial r} \frac{\partial r}{\partial y} - \bar{y}_0 \frac{\partial A_0}{\partial r} \frac{\partial r}{\partial x} = \frac{\partial A_0}{\partial r} \left( \bar{x}_0 \frac{\partial r}{\partial y} - \bar{y}_0 \frac{\partial r}{\partial x} \right)$$

$$\frac{\partial r}{\partial x} = \frac{x}{r}, \quad \frac{\partial r}{\partial y} = \frac{y}{r}, \quad \bar{H} = \frac{\partial A_0}{\partial r} \frac{\bar{x}_0 y - \bar{y}_0 x}{r} =$$

$$= \frac{\partial A_0}{\partial r} \frac{\bar{x}_0 r \sin \vartheta \sin \varphi - \bar{y}_0 r \sin \vartheta \cos \varphi}{r} =$$

$$= \frac{\partial A_0}{\partial r} \sin \vartheta (\bar{x}_0 \sin \varphi - \bar{y}_0 \cos \varphi) = \frac{\partial A_0}{\partial r} \sin \vartheta \varphi_0$$

$$\frac{\partial A_0}{\partial r} = -\frac{i\omega p_0}{c} \left( -\frac{1}{r^2} - \frac{ik}{r} \right) e^{-ikr} = -ik p_0 \left( -\frac{1}{r^2} - \frac{ik}{r} \right) e^{-ikr}$$

$$H_\varphi = +ik p_0 \frac{k}{r} \left( \frac{1}{kr} + i \right) e^{-ikr} \sin \vartheta =$$

$$= +\frac{k^2 p_0}{r} \left( \frac{i}{kr} - 1 \right) e^{-ikr} \sin \vartheta$$

$$\text{rot } \bar{H} = \frac{1}{c} \frac{\partial \bar{E}}{\partial t} = i\omega \bar{D} = \frac{i\omega}{c} \bar{E}, \quad \bar{E} = -\frac{ic}{\omega} \text{rot } \bar{H}$$

$$\text{rot } \bar{H} = \begin{vmatrix} \bar{r}_0/r^2 \sin \vartheta & \bar{\varphi}_0/r & \bar{z}_0/r \sin \vartheta \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \varphi} & \frac{\partial}{\partial \vartheta} \\ 0 & r \sin \vartheta H_\varphi & 0 \end{vmatrix} =$$

$$= -\frac{\bar{r}_0}{r^2 \sin \vartheta} \frac{\partial (r \sin \vartheta H_\varphi)}{\partial \vartheta} + \frac{\bar{z}_0}{r \sin \vartheta} \frac{\partial (r \sin \vartheta H_\varphi)}{\partial r}$$

$$E_r = \frac{ic}{\omega} \frac{1}{r^2 \sin \vartheta} k^2 p_0 \frac{1}{r} \left( \frac{i}{kr} - 1 \right) e^{-ikr} \frac{\partial (\sin^2 \vartheta)}{\partial \vartheta} =$$

$$= \frac{ic}{\omega} \frac{k^2 p_0}{r^2 \sin \vartheta} \left( \frac{i}{kr} - 1 \right) e^{-ikr} 2 \sin \vartheta \cos \vartheta =$$

$$= -\frac{2k p_0}{r^2} \cos \vartheta \left( \frac{1}{kr} + i \right) e^{-ikr}$$



$$E_r = -\frac{ic}{\omega r} \frac{\partial(rH_\varphi)}{\partial r} = -\frac{ic}{\omega r} \frac{\partial}{\partial r} \left( r \frac{k^2 p_0}{r} \left( \frac{i}{kr} - 1 \right) e^{-ikr} \sin \vartheta \right) =$$

$$= -\frac{i}{kr} k^2 p_0 \sin \vartheta \frac{\partial}{\partial r} \left( \left( \frac{i}{kr} - 1 \right) e^{-ikr} \right) =$$

$$= -\frac{i}{r} k p_0 \sin \vartheta \left[ -\frac{i}{kr^2} - ik \left( \frac{i}{kr} - 1 \right) \right] e^{-ikr} =$$

$$= -\frac{k^2 p_0}{r} \sin \vartheta \left( -1 + \frac{i}{kr} + \frac{1}{kr^2} \right) e^{-ikr}$$

Нпу  $kr \ll 1 \Rightarrow \frac{1}{kr} \gg 1 \Rightarrow$

$$H_\varphi = \frac{ik^2 p_0}{r} \frac{1}{kr} e^{-ikr} \sin \vartheta = \frac{ik p_0}{r^2} \sin \vartheta$$

$$E_r = -\frac{2k p_0}{r^2} \cos \vartheta \frac{1}{kr} = -\frac{2p_0}{r^3} \cos \vartheta$$

$$E_\vartheta = -\frac{k^2 p_0}{r} \sin \vartheta \frac{1}{k^2 r^2} = -\frac{p_0}{r^3} \sin \vartheta$$

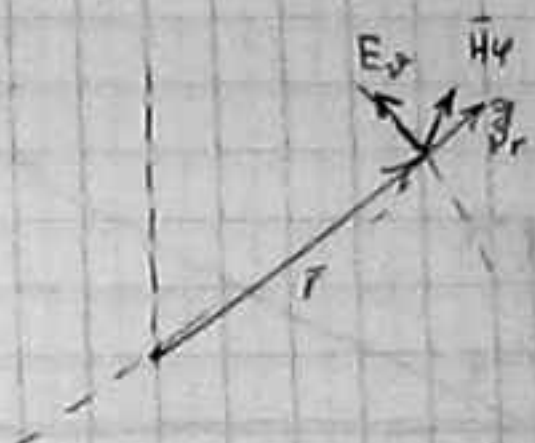
Нпу  $kr \gg 1 \Rightarrow \frac{1}{kr} \ll 1 \Rightarrow$

$$H_\varphi = \frac{k^2 p_0}{r} (-1) e^{-ikr} = -\frac{k^2 p_0}{r} e^{-ikr}$$

$$E_r = -\frac{2k p_0}{r^2} i \cos \vartheta e^{-ikr} = -\frac{2ik p_0}{r^2} \cos \vartheta e^{-ikr}$$

$$E_\vartheta = -\frac{k^2 p_0}{r} \sin \vartheta (-1) e^{-ikr} = \frac{k^2 p_0}{r} \sin \vartheta e^{-ikr}$$

$$5) \quad \tilde{S} = \operatorname{Re} \frac{c}{8\pi} [\vec{E} \times \vec{H}^*]$$



$$\tilde{S}_r = -\operatorname{Re} \frac{c}{8\pi} E_\vartheta H_\varphi^* = +\operatorname{Re} \frac{c}{8\pi} \frac{k^2 p_0}{r} \sin \vartheta \left( -1 + \frac{i}{kr} + \frac{1}{kr^2} \right) e^{-ikr}.$$

$$\cdot \frac{k^2 p_0}{r} \left( -\frac{i}{kr} - 1 \right) e^{ikr} \sin \vartheta =$$

$$= \frac{c}{8\pi} \frac{k^4 p_0^2}{r^2} \sin^2 \vartheta \operatorname{Re} \left\{ \frac{i}{kr} + 1 + \frac{1}{kr^2} - \frac{i}{kr} - \frac{i}{kr^3} - \frac{1}{kr^2} \right\} =$$

$$= \frac{c k^4 p_0^2}{8\pi r^2} \sin^2 \vartheta \operatorname{Re} \left\{ 1 - \frac{i}{kr^3} \right\} = \frac{c k^4 p_0^2}{8\pi r^2} \sin^2 \vartheta$$

$$6) \quad P = \int_S \tilde{S}_r dS = \int_0^{2\pi} d\varphi \int_0^\pi \frac{c k^4 p_0^2}{8\pi r^2} \sin^2 \vartheta r^2 d\vartheta \sin \vartheta d\varphi =$$

$$= \frac{c k^4 p_0^2}{8\pi} \int_0^{2\pi} d\varphi \int_{-\pi/2}^{\pi/2} \sin^2 \vartheta \sin \vartheta d\vartheta = -\frac{c k^4 p_0^2}{4} \int_{-1}^1 (1 - \cos^2 \vartheta) d(\cos \vartheta) =$$

$$= -\frac{c k^4 p_0^2}{4} \left( \cos \vartheta - \frac{\cos^3 \vartheta}{3} \right) \Big|_0^\pi = -\frac{c k^4 p_0^2}{4} \left( -1 - 1 + \frac{1}{3} + \frac{1}{3} \right) =$$

$$= -\frac{c k^4 p_0^2}{4} \left( -\frac{4}{3} \right) = \frac{c k^4 p_0^2}{3} = \frac{c \omega^4 p_0^2}{3 c^4} = \frac{\omega^4 p_0^2}{3 c^3}$$

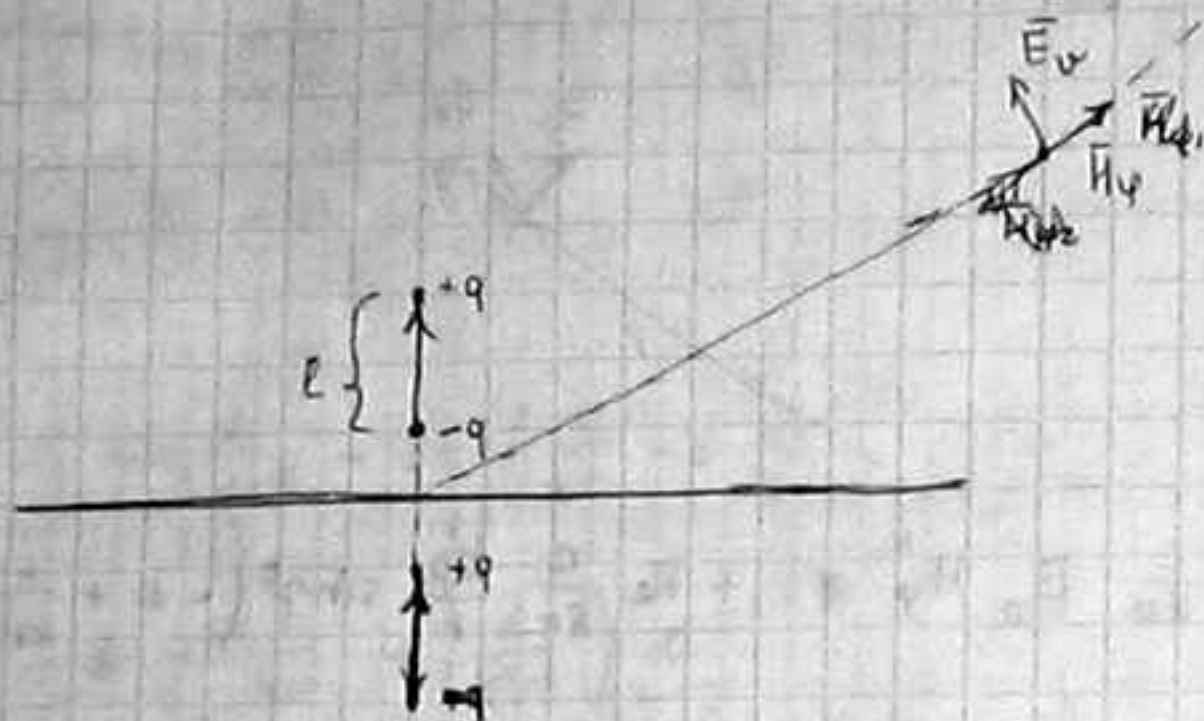
$$I_0 = \frac{\omega p_0}{c}, \quad I_0^2 = \frac{\omega^2 p_0^2}{c^2}$$

$$R_r^{(e)} = \frac{2 \omega^4 p_0^2}{3 c^3} \cdot \frac{l^2}{\omega^2 p_0^2} = \frac{2 \omega^2 l^2}{3 c^3} = \frac{2 (kl)^2}{3 c}$$



№ 9.6

a)



Поле элементарного вибратора в точке  $2l$ .

$$E \sim \sin^2 \alpha$$

b)

