

Coherent State Path Integral Monte Carlo for ${}^4\text{He}$

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I. INTRODUCTION

In the publication “Coherent State Path Integral Monte Carlo” we performed several Path Integral Monte Carlo (PIMC) simulations for a two dimensional, $d = 2$, ${}^4\text{He}$ liquid [1] with either Boltzmann or Bose statistics. The liquid has a surface density $\rho = N/\Omega$ where N is the number of helium atoms in an area Ω of a flat surface, at an inverse temperature $\beta = 1/k_B T$ with k_B Boltzmann constant. The Hamiltonian of the fluid is $\hat{H} = \hat{T} + v\hat{V} = \hat{P}^2/2m + vV(Q)$ with $Q = (\mathbf{q}_1, \mathbf{q}_2, \dots, \mathbf{q}_N)$ the particles positions and momenta $P = (\hat{\mathbf{p}}_1, \hat{\mathbf{p}}_2, \dots, \hat{\mathbf{p}}_N) = -i(\nabla_{\mathbf{q}_1}, \nabla_{\mathbf{q}_2}, \dots, \nabla_{\mathbf{q}_N})$. We compared two different PIMC versions: The conventional Plane Waves PIMC (**PWPIMC**) [2] and our Coherent States PIMC (**CSPIMC**) new algorithm. The two PIMC differ for the expression of the hot kinetic density matrix at an imaginary timestep $\tau = \beta/M$ with M a large number of timeslices. In our quantum simulations we have to take care of two limiting procedures: The continuum limit where the ultraviolet cutoff $\tau \rightarrow 0$ or $M \rightarrow \infty$

at fixed absolute temperature T and the thermodynamic limit where $N \rightarrow \infty$ and $\Omega \rightarrow \infty$ at fixed density ρ . We will explicitly worry about the continuum limit in the last two Table VII and VIII for Boltzmann and Bose statistics respectively. To mimic the thermodynamic limit we will simply use a periodic square cell which permeates the whole infinite space. This will give rise to spurious finite size effects unavoidable on a computer.

The PWPIMC requires a multidimensional integral over dNM coordinates whereas our CSPIMC requires $5dMN$ integrations, the usual dNM real particles coordinates, $2dNM$ ghost particles coordinates, and $2dNM$ ghost particles momenta. The coherent states are generated by an Harmonic Oscillator (HO) with mass $m_{h.o.}$ and elastic constant $k = m_{h.o.}\omega^2$. We will introduce the parameter $\xi \equiv m_{h.o.}\omega/2$ and the adimensional one $\varphi \equiv \xi\tau/m = (\sigma_{p.w.}/\sigma_{c.s.})^2/2$ where $\sigma_{p.w.} \equiv \sqrt{2\lambda\tau}$ is the standard deviation in the plane wave scheme, with $\lambda = 1/2m$, and $\sigma_{c.s.} \equiv \sqrt{1/m_{h.o.}\omega}$ is the standard deviation in the coherent state scheme. We will work in units where $\hbar = k_B = 1$ so that the imaginary time has dimensions of temperature.

II. THE SHORT TIME DENSITY MATRIX FOR CSPIMC

The short time τ Green function for CSPIMC can be written as follows:

$$\rho(Q, Q'; \tau) \approx e^{-\tau v V(Q')} \frac{1}{N!} \sum_{\mathcal{P}} (-1)^{\mathcal{P}} \prod_{\alpha=1}^N \zeta_{\alpha} [\mathbf{q}_{\mathcal{P}\alpha} | \mathbf{q}'_{\alpha}; \tau, m, \xi], \quad (2.1)$$

with \mathcal{P} a permutation of the N **real** identical particles and

$$\zeta_{\alpha} [\mathbf{q} | \mathbf{q}'; \tau, m, \xi] \approx \int \frac{d\mathbf{q}_a d\mathbf{p}_a}{(2\pi)^d} \frac{d\mathbf{q}_b d\mathbf{p}_b}{(2\pi)^d} \psi_{\alpha}^{\mathbf{a}}(\mathbf{q}) \psi_{\alpha}^{\mathbf{b}*}(\mathbf{q}') G_{\mathbf{a}, \mathbf{b}} \exp [-\tau(\mathbf{p}_a^2 + \mathbf{p}_b^2)/4m],$$

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where $\psi_{\alpha}^{\mathbf{a}}(\mathbf{q}) \equiv \psi^{\mathbf{a}}(\mathbf{q}_{\alpha})$ is the CS wave function with ghost canonical variables \mathbf{q}_a and \mathbf{p}_a , the **ghost** of the real particle α , namely

$$\begin{aligned}\psi^{\mathbf{a}}(\mathbf{q}_{\alpha}) &\equiv \langle \mathbf{q}_{\alpha} | \mathbf{q}_a, \mathbf{p}_a \rangle \\ &= \left(\frac{m_{h.o.}\omega}{\pi} \right)^{d/4} \exp \left\{ -\frac{m_{h.o.}\omega}{2} \left[\mathbf{q}_{\alpha} - \sqrt{\frac{2}{m_{h.o.}\omega}} \text{Re}(\mathbf{a}) \right]^2 + i\mathbf{q}_{\alpha} \cdot \sqrt{2m_{h.o.}\omega} \text{Im}(\mathbf{a}) - i2\text{Re}(\mathbf{a}) \cdot \text{Im}(\mathbf{a}) \right\},\end{aligned}\quad (2.2)$$

$$G_{\mathbf{a},\mathbf{b}} = \exp \left[-\frac{1}{2}(|\mathbf{a}|^2 + |\mathbf{b}|^2) + \mathbf{a}^* \cdot \mathbf{b} + \frac{i}{2}(\mathbf{q}_a \cdot \mathbf{p}_a - \mathbf{q}_b \cdot \mathbf{p}_b) \right], \quad (2.3)$$

$$\mathbf{a} = \frac{1}{\sqrt{2m_{h.o.}\omega}}(m_{h.o.}\omega\mathbf{q}_a + i\mathbf{p}_a), \quad (2.4)$$

$$\mathbf{b} = \frac{1}{\sqrt{2m_{h.o.}\omega}}(m_{h.o.}\omega\mathbf{q}_b + i\mathbf{p}_b), \quad (2.5)$$

where $G_{\mathbf{a},\mathbf{b}}$ is the scalar product of the two ghosts “a” and “b”. All this is used to find the acceptance probability in subroutine `cs` in the code listing at the end of this material.

III. RESULTS

We compare conventional PWPIMC with our CSPIMC algorithm on specific simulations. We performed computer experiments for $N = 16$ ${}^4\text{He}$ atoms ($m = 0.0830594 \approx 1/12 \text{ \AA}^{-2}\text{K}^{-1}$) in two dimensions $d = 2$, in a square periodic cell of area $\Omega = L^2$, interacting with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556 \text{ \AA}$, $\varepsilon = 10.22 \text{ K}$ and a cutoff distance $r_{\text{cut}} = 2.5 \text{ \AA}$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $\rho = N/L^2 = 0.05 \text{ \AA}^{-2}$ and various temperatures T or timeslices number M , either with Boltzmann and Bose statistics. As we can see from the phase diagram of this fluid [1] at this density it undergoes a phase transition from a fluid phase at high temperature to a superfluid phase at low temperatures.

We take as initial configuration each atom sitting at all timeslices on a random position choosen so to avoid overlaps with the other atoms.

The Monte Carlo used is the standard Metropolis algorithm [4, 5]. In our simulation we choose as transition move a uniform displacement of each of the dMN real path coordinate $\mathbf{q}_{\alpha}(i\tau) \rightarrow \mathbf{q}_{\alpha}(i\tau) + (1/2 - \eta)\Delta$ for $\alpha = 1, \dots, N$ and $i = 1, \dots, M$, where η is a uniform pseudo-random number in $[0, 1]$ and Δ a fixed 3-dimensional vector whose magnitude is chosen so to have acceptance ratios close to 1/2. And of each of the $4dMN$ ghost path canonical variables $\mathbf{q}_{a_{\alpha}}(i\tau) \rightarrow \mathbf{q}_{a_{\alpha}}(i\tau) + (1/2 - \eta)\Delta$, $\mathbf{p}_{a_{\alpha}}(i\tau) \rightarrow \mathbf{p}_{a_{\alpha}}(i\tau) + (1/2 - \eta)\Delta$, and $\mathbf{q}_{b_{\alpha}}(i\tau) \rightarrow \mathbf{q}_{b_{\alpha}}(i\tau) + (1/2 - \eta)\Delta$, $\mathbf{p}_{b_{\alpha}}(i\tau) \rightarrow \mathbf{p}_{b_{\alpha}}(i\tau) + (1/2 - \eta)\Delta$. So that the transition probability density is just a constant and drops out of the acceptance probability. A MC step consists of displacing the M timeslices of a randomly chosen ghost and real particle one by one (in order to activate just this move in the CSPIMC simulation it is necessary to set TRUE the entry 24 of the input data file `data-cs-2.in`) and of a displacement of a fixed random number $\leq M$ of timeslices, connecting two randomly chosen real particles, all at once, in a sort of “brownian bridge” (for de-

tails look at the subroutine `bridge` in the code listing at the end of this material), so to realize an exchange, `swap`, of two real particles. Here it is also necessary to bridge the “a” and “b” ghosts of particle α so that $\mathbf{q}_{a_{\alpha}}^{new} = \mathbf{q}_{b_{\alpha}}^{new} = \mathbf{q}_{\alpha}^{new}$ (where $\mathbf{q}_{\alpha}^{new}$ is the new position of the real particle α after the brownian bridge), and swap the real particles together with their two ghosts each (in order to activate just this move in the CSPIMC simulation it is necessary to set TRUE the entry 25 of the input data file `data-cs-2.in`).¹

Our PWPIMC simulations (see Tables I, II and III, IV) confirm that at high temperature (classical limit) the nature of the statistics is not important. Whereas it becomes important at low temperatures (quantum regime). The zero temperature (ground state) limit can only be reached through an extrapolation of the PIMC results.

In Tables I and II we use PWPIMC at constant $\tau = 0.025 \text{ K}^{-1}$ (same as Ref. [1]) and various temperatures.

In Tables III and V we compare the PWPIMC with the CSPIMC for Boltzmann statistics at fixed $M = 250$ and various temperatures (In order to use Boltzmann statistics with our CSPIMC code listed in the Appendix we chose FALSE in entry 24 of the input data file `data-cs-2.in`). In Tables IV and VI we do the same for Bose statistics (In order to use Bose statistics with our CSPIMC code listed in the Appendix we chose TRUE in entry 2 and FALSE in entries 24 and 25 of the input data file `data-cs-2.in`).

In all our tables we denote with $\mathcal{E}_p = v(V)$ and

$$\mathcal{E}_k = \begin{cases} \frac{dN}{2\tau} - \frac{\langle (Q_i - Q_{i-1})^2 \rangle}{4\lambda\tau^2} & \text{PWPIMC} \\ \frac{dN}{2\tau} \frac{\varphi}{1+\varphi} \phi - \frac{\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle}{4m} & \text{CSPIMC} \end{cases} \quad (3.1)$$

¹ We also tried two independent brownian bridges one between two random real particles and one between two random ghost particle but this does not work because the ghost particle $\mathbf{q}_{a_{\alpha}}$ may freely end up very close to a real particle \mathbf{q}_{β} . And due to the harmonic coupling the ghost will carry in its neighborhood a real particle \mathbf{q}_{α} which may then overlap with the real particle β producing an explosion of the potential energy.

where $Q_i = Q(i\tau)$ are the N real particles coordinates at the i th timestep and P_a, P_b are the N momenta of the two ghosts. This corresponds to the *thermodynamic* estimator of Ref. [2]. The parameter ϕ , in the CSPIMC case, is necessary in order to find agreement with the kinetic energy of the PWPIPMC ,as explained in the main article. Note that the kinetic energy from Eq. (3.1) is the small difference between two large quantities, infinite in the continuum $\tau \rightarrow 0$ limit.

TABLE I. Results from PWPIPMC for $N = 16$ ^4He atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $\tau = 0.025$, with **Boltzmann statistics**. In the Table \mathcal{E}_k and \mathcal{E}_p are the total kinetic and potential energies respectively. $\mathcal{E} = \mathcal{E}_k + \mathcal{E}_p$ is the total internal energy (note that unlike Ref. [1] we fix M and not τ).

T (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	$-\mathcal{E}$ (K)
1.0	66.6(2)	94.4(1)	27.8
0.5	60.8(2)	92.9(1)	32.1
0.2	59.0(2)	93.00(8)	34.0
0.1	58.2(2)	93.23(8)	35.0

TABLE II. Results from PWPIPMC for $N = 16$ ^4He atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $\tau = 0.025$, with **Bose statistics**. In the Table \mathcal{E}_k and \mathcal{E}_p are the total kinetic and potential energies respectively. $\mathcal{E} = \mathcal{E}_k + \mathcal{E}_p$ is the total internal energy.

T (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	$-\mathcal{E}$ (K)
1.0	57.1(1)	94.8(3)	37.7
0.5	55.0(5)	93.4(2)	38.4
0.2	55.8(1)	93.55(5)	37.7
0.1	55.9(1)	93.46(5)	37.6

In Tables VII and VIII we performed simulations at fixed $\varphi = \pi/2$, and fixed temperatures $T = 1$ (K) and $T = 0.2$ (K) respectively, at increasing numbers M of timeslices. We see that it is necessary to keep φ constant upon taking the continuum limit in order to find agreement between the results for the CSPIMC and the ones for the PWPIPMC for both the kinetic and potential energies. The convergence for the kinetic energy measured through the estimator of Eq. (3.1) is rather slow for CSPIMC and it would be desirable to try also other alternative estimators. Also the equilibration time starting from a random configuration of the atoms is rather

TABLE III. Results from PWPIPMC for $N = 16$ ^4He atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $M = 250$, with **Boltzmann statistics**. The runs are 2×10^6 steps long. In the Table \mathcal{E}_k and \mathcal{E}_p are the total kinetic and potential energies respectively. $\mathcal{E} = \mathcal{E}_k + \mathcal{E}_p$ is the total internal energy (note that unlike Ref. [1] we fix M and not τ).

T (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	$-\mathcal{E}$ (K)
1.0	82.3(6)	83.6(1)	1.3(6)
0.5	74.6(4)	84.8(1)	10.2(4)
0.2	62.5(4)	90.5(1)	28.0(4)
0.1	49.0(2)	101.1(1)	52.1(2)

TABLE IV. Results from PWPIPMC for $N = 16$ ^4He atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $M = 250$, with **Bose statistics**. The runs are 2×10^6 steps long. In the Table \mathcal{E}_k and \mathcal{E}_p are the total kinetic and potential energies respectively. $\mathcal{E} = \mathcal{E}_k + \mathcal{E}_p$ is the total internal energy (note that unlike Ref. [1] we fix M and not τ).

T (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	$-\mathcal{E}$ (K)
1.0	77(1)	84.1(4)	7.1(4)
0.5	69.7(7)	84.6(4)	14.9(7)
0.2	60.0(5)	90.6(1)	30.6(5)
0.1	49.3(3)	101.1(2)	51.8(3)

long, about 6×10^6 MC steps. We find that, in the continuum $\tau \rightarrow 0$ limit, ϕ tends to $\approx \sqrt{2}$, instead of $\approx \sqrt{3}$ as happens in the non interacting case (this is shown in the main article in Table II there), and this is due to the fact that for the interacting system there are additional terms in \mathcal{E}_k due to the interaction as shown by the *direct* estimator of Ref. [2].

In order to reproduce the PWPIPMC results for the total kinetic and potential energies it is necessary to fix in the CSPIMC $\varphi = m_{h.o.}\omega\tau/2m = \pi/2$ and take the continuum limit $\xi = m_{h.o.}\omega/2 \rightarrow \infty$ or $\tau \rightarrow 0$. From Table VII we see that at fixed $\varphi = \pi/2$, ϕ remains close to $\sqrt{2}$ at small τ . In order to find reasonable results for the kinetic energy in the small τ cases of Tables VI and VIII it is necessary to approach equilibrium “adiabatically” (fast) by keeping the acceptance ratios for the single slice displacement move below or much below $1/2$ ². This is also

² This suggests that it would be more convenient to distinguish

TABLE V. Results from **CSPIMC** with $\varphi = \pi/2$ for $N = 16$ ${}^4\text{He}$ atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $M = 250$, with **Boltzmann statistics**. The runs are 2×10^6 steps long. In the Table $\mathcal{E}_k = \frac{N}{\tau} \frac{\pi}{2+\pi} \phi - \frac{\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle}{4m}$ is taken from Table III and \mathcal{E}_p is the total potential energy.

T (K)	$\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle / 4m$ (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	ϕ
1.0	3383(5)	82.3	85.3(4)	1.42
0.5	1735(3)	74.6	85.9(2)	1.48
0.2	668.1(9)	62.5	90.4(2)	1.49
0.1	317.6(4)	49.0	99.3(2)	1.50

TABLE VI. Results from **CSPIMC** with $\varphi = \pi/2$ for $N = 16$ ${}^4\text{He}$ atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and various temperatures T with $M = 250$, with **Bose statistics**. The runs are 2×10^6 steps long. In the Table $\mathcal{E}_k = \frac{N}{\tau} \frac{\pi}{2+\pi} \phi - \frac{\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle}{4m}$ is taken from Table IV and \mathcal{E}_p is the total potential energy. We counted very few atoms swaps. For example at the lower temperature of the table, we counted just 13 swaps over 10^6 MC steps.

T (K)	$\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle / 4m$ (K)	\mathcal{E}_k (K)	$-\mathcal{E}_p$ (K)	ϕ
1.0	3868(6)	77	85.8(4)	1.61
0.5	1722(3)	69.7	85.6(2)	1.47
0.2	643.9(8)	60.0	94.3(2)	1.44
0.1	328.0(4)	49.3	100.4(1)	1.54

important for the measure of the potential energy since it doesn't cost anything for two ghosts to approach each other, then it may happen that in correspondence of the binding of the ghosts of two different particles these also bind producing an artificial positive jump in the potential energy. It is therefore necessary to give a kick Δ larger than the extent $\sigma_{c.s.}$ of the cloud of ghosts around their particle in order to unbind the pair.

The slight disagreement in the potential energy between PWPIMC and CSPIMC is due to the different approach to the continuum $\tau \rightarrow 0$ limit for the two algorithms, where in the CSPIMC this is affected also by the other available parameter ξ .

between a displacement Δ for the positions and one for the momenta.

TABLE VII. Results from **CSPIMC** with $\varphi = \pi/2$ for $N = 16$ ${}^4\text{He}$ atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and a temperature $T = 1$ with various values of M , with **Boltzmann statistics**. In the Table $\mathcal{E}_k = \frac{N}{\tau} \frac{\pi}{2+\pi} \phi - \frac{\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle}{4m} \approx 82.3$ and \mathcal{E}_p are the total kinetic and potential energies.

M	$\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle / 4m$ (K)	ϕ	$-\mathcal{E}_p$ (K)
250	3383(5)	1.42	85.3(4)
500	6818(4)	1.41	80.6(3)
750	10232(5)	1.41	80.1(4)
1000	13696(5)	1.41	79.9(4)

TABLE VIII. Results from **CSPIMC** with $\varphi = \pi/2$ for $N = 16$ ${}^4\text{He}$ atoms ($m = 0.0830594$) in two dimensions, with a Lennard-Jones pair potential with parameters [3] $\sigma = 2.556, \varepsilon = 10.22$ and a cutoff distance $r_{\text{cut}} = 2.5$ (so that $v(r) = 0$ for $r > r_{\text{cut}}\sigma$ without long range corrections), at a density $N/L^2 = 0.05$ and a temperature $T = 0.2$ with various values of M , with **Bose statistics**. In the Table $\mathcal{E}_k = \frac{N}{\tau} \frac{\pi}{2+\pi} \phi - \frac{\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle}{4m} \approx 60.0$ and \mathcal{E}_p are the total kinetic and potential energies. In order to find reasonable results for the small τ entries in the table is important to have acceptance ratios for the single slice displacement move less or much less than 1/2.

M	$\langle P_a^2(i\tau) + P_b^2(i\tau) \rangle / 4m$ (K)	ϕ	$-\mathcal{E}_p$ (K)
250	643.9(8)	1.44	94.3(2)
500	1371(2)	1.46	85.2(2)
750	2514(4)	1.75	83.5(1)
1000	5163(3)	2.67	79.3(1)

IV. CONCLUSIONS

In this supplementary material we proved the statement offered in the publication “Coherent State Path Integral Monte Carlo” that in the CSPIMC algorithm it is necessary to take the continuum limit at fixed $\varphi = \xi\tau/m = \pi/2$. We also presented various simulation results supporting the findings of the main article. It remains an open problem to theoretically asses the, for the time being, empiric values for the two parameters φ and ϕ .

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the development of the brownian bridge and the consequent particles permutation sampling.

Appendix A: The code

This is the code used for the CSPIMC computer experiment. We list here the main FORTRAN code `cspimc.f` with its included `mc-CS.par` parameters file. And the input data file `data-CS-2.in` that is read at the beginning of the run.


```

J=NEXT(I)
PREV(J)=I
ENDDO

NIP = NEXT(IP)
PIP = PREV(IP)

DO 90 C1 = 1, FTNO

  DO IDIM = 1, DIM
    RXIOLD(IDIM) = RX(IDIM,C1,IP)
    raold(idim) = ra(idim,c1,ip)
    rbold(idim) = rb(idim,c1,ip)
    paold(idim) = pa(idim,c1,ip)
    pbold(idim) = pb(idim,c1,ip)
  ENDDO

C ** CALCULATE THE ENERGY OF I IN THE OLD CONFIGURATION **

  CALL PENERGY ( POTK, RXIOLD, IP, C1,
                 VOLD, WOLD )

  CALL CSKENERGY ( PAOLD, PBOLD, IP, C1,
                   KEOLD )

  call cs(raold,paold,rb(:,c1+1,ip),pbold,
         rxioold,rx(:,c1+1,ip),zold)
  call cs(ra(:,c1-1,ip),paold,rbold,pbold,
         rx(:,c1-1,ip),rxioold,zo)
  zold=zold*zo

C ** INSTANTANEOUS VALUE OF THE ACTION **

  ACTOLD = KEOLD + VOLD

C ** MOVE I AND PICKUP THE CENTRAL IMAGE **

  DO IDIM = 1, DIM
    RXINNEW(IDIM) = RXIOLD(IDIM) +
      ( 2.0 * RANF ( DUMMY ) - 1.0 )*DRMAX
    RXINNEW(IDIM) = RXINNEW(IDIM) -
      DRINT ( RXINNEW(IDIM)/SIGMA )*SIGMA
    ranew(idim) = raold(idim) +
      ( 2.0 * rafn ( dummy ) - 1.0 )*drmax
    ranew(idim) = ranew(idim) -
      drint ( ranew(idim)/sigma )*sigma
    rbnew(idim) = rbold(idim) +
      ( 2.0 * rafn ( dummy ) - 1.0 )*drmax
    rbnew(idim) = rbnew(idim) -
      drint ( rbnew(idim)/sigma )*sigma
    panew(idim) = paold(idim) +
      ( 2.0 * rafn ( dummy ) - 1.0 )*drmax
    pbnew(idim) = pbold(idim) +
      ( 2.0 * rafn ( dummy ) - 1.0 )*drmax
  ENDDO

C ** CALCULATE THE ENERGY OF I IN THE NEW CONFIGURATION **

  CALL PENERGY ( POTK, RXINNEW, IP, C1,
                 VNEW, WNEW )

  CALL CSKENERGY ( PANEW, PBNEW, IP, C1,
                   KENEW )

  call cs(ranew,panew,rb(:,c1+1,ip),pbnew,
         rxinew,rx(:,c1+1,ip),znew)
  call cs(ra(:,c1-1,ip),panew,rbnew,pbnew,
         rx(:,c1-1,ip),rxinew,zn)
  znew=znew*zn

C ** INSTANTANEOUS VALUE OF THE ACTION **

  ACTNEW = KENEW + VNEW

  DELTV = VNEW - VOLD
  DELTKE = KENEW - KEOLD

C ** CHECK FOR ACCEPTANCE **

  DELTACTB = dble(znew)/dble(zold)*exp(-LS*(VNEW-VOLD))

  IF ( DELTACTB .GT. RANF ( DUMMY ) ) THEN
    V      = V + DELTV
    KE     = KE + DELTKE
    ACATMA = ACATMA + 1.0
    DO IDIM = 1, DIM
      RX(IDIM,C1,IP) = RXINNEW(IDIM)
      ra(idim,c1,ip) = ranew(idim)
      rb(idim,c1,ip) = rbnew(idim)
      pa(idim,c1,ip) = panew(idim)
      pb(idim,c1,ip) = pbnew(idim)
    ENDDO
    IF ( C1.EQ.1 ) THEN
      RX(IDIM,FTNO+1,PIP)=RX(IDIM,C1,IP)
      RA(IDIM,FTNO+1,PIP)=RA(IDIM,C1,IP)
      RB(IDIM,FTNO+1,PIP)=RB(IDIM,C1,IP)
    ENDIF
    ENDDO
    ENDIF
    IF ( STEP.GT.IEQUI ) THEN
      ACM = ACM + 1.0
    ENDIF

C ** CALCULATE INSTANTANEOUS VALUES **

  IF (POTK .EQ. 'LJ') THEN
    VN = ( V + VLRC )
    ELSE
      VN = V
    ENDIF

    C   ** ACCUMULATE AVERAGES **

      ACV     = ACV     + VN
      ACVSQ  = ACVSQ  + VN*VN
      ACKE   = ACKE   + KE
      ACKESQ = ACKESQ + KE*KE
    ENDIF

    C ****END DISPLACEMENT MOVE ****
    C ****END LOOP OVER TIME SLICES ****

  90  CONTINUE

  C   **** WRITE OUT THE INSTANTANEOUS VALUES ON FORT.9 ***
  IF ( MOD ( STEP, IPRINT ) .EQ. 0 ) THEN
    WRITE(9,*), STEP, DRMAX
  ENDIF

  C   ** CREATE POSITION DISTRIBUTION **
  CALL DISTR(30_8,STEP*FTNO*NP/30)

  C   ** PERFORM PERIODIC OPERATIONS **

  IF ( MOD ( STEP, IRATIO ) .EQ. 0 ) THEN
    C   ** ADJUST MAXIMUM DISPLACEMENT **

      RATIO = ACATMA / DBLE ( FTNO * IRATIO )
      RATIOB = ACATMAB / DBLE ( IRATIO )
      RATIOS = ACATMAS / DBLE ( IRATIO )

      IF ( RATIO .GT. 0.5 ) THEN
        DRMAX = DRMAX * 1.05
      ELSE
        DRMAX = DRMAX * 0.95
      ENDIF

      ACATMA = 0.0
      ACATMAB = 0.0
      ACATMAS = 0.0

    ENDIF

    IF ( MOD ( STEP, IPRINT ) .EQ. 0 ) THEN
      WRITE(*,*), step, drmax
      WRITE(*,'(I8, 2XE10.4)'), STEP, DRMAX
      WRITE(*,*), acm, ratio, ratiob, ratios
      :   ke, v
      :   WRITE(*,'(8,5(2XE10.4))'), INT(ACM),
      :   :   RATIO, RATIOB, RATIOS,
      :   :   KE/FTNO, V/FTNO
      :   :   WRITE(*,*), <ke>, <v>
      :   :   WRITE(*,'(2(2XE13.7))'), ACKE/ACM/FTNO, ACM/ACM/FTNO
    ENDIF

    IF ( MOD ( STEP, ISAVE ) .EQ. 0 ) THEN
      C   ** WRITE OUT THE CONFIGURATION AT INTERVALS **

        CALL WRITCN ( CNFILE )
        ENDIF

      100  CONTINUE

      C ****ENDS THE LOOP OVER CYCLES ****
      C **** CHECKS FINAL VALUE OF THE POTENTIAL ENERGY IS CONSISTENT **

        CALL CSSUMUP ( POTK, OVRALP, KEEND, VEND )
        IF ( ABS(VEND - V) .GT. 1.0d-03 ) THEN
          WRITE(*,'(>> PROBLEM WITH V ENERGY !!!!<)')
          WRITE(*,'(>> VEND = ', E20.6'), VEND
          WRITE(*,'(>> V      = ', E20.6'), V
        ENDIF

        IF ( ABS(KEEND - KE) .GT. 1.0d-03 ) THEN
          WRITE(*,'(>> PROBLEM WITH KE ENERGY !!!!<)')
          WRITE(*,'(>> KEEND = ', E20.6'), KEEND
          WRITE(*,'(>> KE      = ', E20.6'), KE
        ENDIF

      C   ** WRITE OUT THE FINAL CONFIGURATION FROM THE RUN **

    ENDIF
  ENDIF

```

```

CALL WRITCN ( CNFILE )

C ** CALCULATE AND WRITE OUT RUNNING AVERAGES **

      AVV = ACV / ACM
      ACVSQ = ( ACVSQ / ACM ) - AVV ** 2
      AVKE = ACKE / ACM
      ACKESQ = ( ACKESQ / ACM ) - AVKE ** 2

C ** CALCULATE FLUCTUATIONS **

      IF ( ACVSQ .GT. 0.0 ) FLV = SQRT ( ACVSQ/ACM )/FTNO
      IF ( ACKESQ .GT. 0.0 ) FLKE = SQRT ( ACKESQ/ACM )/FTNO

      WRITE(*,'(//> AVERAGES //>)')
      WRITE(*,'(<V/N> = '' ,E12.6)') AVV/FTNO
      WRITE(*,'(<KE/N> = '' ,E12.6)') AVKE/FTNO

      WRITE(*,'(//> FLUCTUATIONS //>')

      WRITE(*,'(<V/N> FLUCTUATION IN <V/N> = '' ,E12.6)') FLV
      WRITE(*,'(<KE/N> FLUCTUATION IN <KE/N> = '' ,E12.6)') FLKE
      WRITE(*,'(//> END OF SIMULATION //>')

      STOP
      END

      subroutine switch(i,j)
!      switch i and j
!      implicit none
!      integer*8 i,j,k
!      k=j
!      i=j
!      j=k
!      return
!      end

      subroutine update(j,rxp,ip,nip)
!      implicit none
!      updates a portion of the current path x using the proposed path xp
!      INCLUDE 'mc-cs.par'

      integer*8 ip,nip,j,l,k,idim
      integer*8 mcm
      real*8 rxp(mdim,0:n)

      mcm = j+mbm-floor((j+mbm-1)/ftn0)*ftn0

      l=0
      if(j+mbm.le.ftn0)then
        do k=j+1,mcm-1
          l=l+1
          do idim=1,dim
            rx(idim,k,ip)=rxp(idim,l)
            ra(idim,k,ip)=rxp(idim,l)
            rb(idim,k,ip)=rxp(idim,l)
          enddo
        enddo
      else
        do k=j+1,ftn0
          l=l+1
          do idim=1,dim
            rx(idim,k,ip)=rxp(idim,l)
            ra(idim,k,ip)=rxp(idim,l)
            rb(idim,k,ip)=rxp(idim,l)
          enddo
        enddo
        do k=1,mcm-1
          l=l+1
          do idim=1,dim
            rx(idim,k,nip)=rxp(idim,l)
            ra(idim,k,nip)=rxp(idim,l)
            rb(idim,k,nip)=rxp(idim,l)
          enddo
        enddo
      endif
      do idim=1,dim
        ra(idim,j,ip)=rx(idim,j,ip)
        rb(idim,j,ip)=rx(idim,j,ip)
        ra(idim,mcm,ip)=rx(idim,mcm,ip)
        rb(idim,mcm,ip)=rx(idim,mcm,ip)
        rx(idim,ftn0+1,ip)=rx(idim,1,nip)
        rx(idim,0,nip)=rx(idim,ftn0,ip)
        ra(idim,ftn0+1,ip)=ra(idim,1,nip)
        ra(idim,0,nip)=ra(idim,ftn0,ip)
        rb(idim,ftn0+1,ip)=rb(idim,1,nip)
        rb(idim,0,nip)=rb(idim,ftn0,ip)
      enddo

      return
    end

    subroutine swap(j,ip,nip,kp,nkp)
!      implicit none
!      updates a portion of the current path x using the proposed path xp
!      INCLUDE 'mc-cs.par'

      integer*8 ip,nip,kp,nkp,j,k,idim
      integer*8 mcm
      real*8 rr

      mcm = j+mbm-floor((j+mbm-1)/ftn0)*ftn0

      if(j+mbm.le.ftn0)then
        do k=mcm,ftn0
          do idim=1,dim
            rr=rx(idim,k,ip)

```

```

IMPLICIT NONE
C FACTORIAL FUNCTION
INTEGER*8 FACT,N,P,I
P=1
DO I=1,N
  P=P*I
ENDDO
FACT=P
END

SUBROUTINE CSSUMUP (POTK, OVLAP, KE, V)
IMPLICIT NONE
C *****
C ** CALCULATES THE TOTAL ACTION
C ***
C ** USAGE:
C ***
C ** THE SUBROUTINE RETURNS THE TOTAL ACTION AT THE
C ** BEGINNING AND END OF THE RUN.
C *****
C INCLUDE      'mc-cs.par'
REAL*8      V, KE, VV, KK
LOGICAL      OVLAP
CHARACTER    POTK*(*)
REAL*8      RXII, RXIJ
REAL*8      VIJ, WIJ, RIJSQ, W, WW
INTEGER*8    TAU, I, J, IDIM
C *****
C POTENTIAL ACTION
VV      = 0.0
C ** LOOP OVER ALL THE PAIRS IN THE LIQUID **
DO TAU = 1, FTNO
  DO 100 I = 1, NP - 1
    DO 99 J = I + 1, NP
      RIJSQ = 0.0
      DO IDIM = 1, DIM
        RXII = RX(IDIM,TAU,I) - RX(IDIM,TAU,J)
      C ** MINIMUM IMAGE THE PAIR SEPARATIONS **
        RXII = RXII -
          :      DNINT ( RXII/SIGMA )*SIGMA
        RIJSQ = RIJSQ + RXII * RXII
      ENDDO
      CALL POT (RIJSQ, VIJ, WIJ, POTK)
      VV      = VV + VIJ
      WW      = WW + WIJ
      99  CONTINUE
    100 CONTINUE
    V=VV
  ENDDO
C KINETIC ACTION
KK      = 0.0
DO TAU = 1, FTNO
  DO I = 1, NP
    DO IDIM = 1, DIM
      KK=KK+(PA(IDIM,TAU,I)**2.+PB(IDIM,TAU,I)**2.)/4/FTM
    ENDDO
    ENDDO
    KE=KK
  ENDDO
RETURN
END

SUBROUTINE PENERGY ( POTK, RXI, I, TAU,
:      V, W )
IMPLICIT NONE
C *****
C ** RETURNS THE POTENTIAL ENERGY OF ATOM I WITH ALL OTHER ATOMS. **
C ***
C ** USAGE:
C ***
C ** THIS SUBROUTINE IS USED TO CALCULATE THE CHANGE OF ENERGY
C ** DURING A TRIAL MOVE OF ATOM I. IT IS CALLED BEFORE AND
C ** AFTER THE RANDOM DISPLACEMENT OF I.
C *****
C INCLUDE      'mc-cs.par'
REAL*8      RXI(MDIM), V, W
INTEGER*8    I, J, TAU, IDIM
CHARACTER    POTK*(*)
REAL*8      RXIJ, RIJSQ, VIJ, WIJ
C *****
C V      = 0.0
C W      = 0.0
C ** LOOP OVER ALL MOLECULES EXCEPT I **
DO 100 J = 1, NP
  IF ( I .NE. J ) THEN
    RIJSQ = 0.0
    DO IDIM = 1, DIM
      RXIJ = RXI(IDIM) - RX(IDIM,TAU,J)
      RXIJ = RXIJ -
        :      DNINT ( RXIJ/SIGMA )*SIGMA
      RIJSQ = RIJSQ + RXIJ * RXIJ
    ENDDO
    CALL POT (RIJSQ, VIJ, WIJ, POTK)
    V      = V + VIJ
    W      = W + WIJ
  ENDIF
  100 CONTINUE
RETURN
END

SUBROUTINE PENERGY ( POTK, RXI, I, TAU,
:      V, W )
IMPLICIT NONE
C *****
C ** RETURNS THE POTENTIAL ENERGY OF ATOM I WITH ALL OTHER ATOMS. **
C ***
C ** USAGE:
C ***
C ** THIS SUBROUTINE IS USED TO CALCULATE THE CHANGE OF ENERGY
C ** DURING A TRIAL MOVE OF ATOM I. IT IS CALLED BEFORE AND
C ** AFTER THE RANDOM DISPLACEMENT OF I.
C *****
C INCLUDE      'mc-cs.par'
REAL*8      RXI(MDIM), V, W
INTEGER*8    I, J, TAU, IDIM
CHARACTER    POTK*(*)
REAL*8      RXIJ, RIJSQ, VIJ, WIJ
C *****
C V      = 0.0
C W      = 0.0
C ** LOOP OVER ALL MOLECULES EXCEPT I **
DO 100 J = 1, NP
  IF ( I .NE. J ) THEN
    RIJSQ = 0.0
    DO IDIM = 1, DIM
      RXIJ = RXI(IDIM) - RX(IDIM,TAU,J)
      RXIJ = RXIJ -
        :      DNINT ( RXIJ/SIGMA )*SIGMA
      RIJSQ = RIJSQ + RXIJ * RXIJ
    ENDDO
    CALL POT (RIJSQ, VIJ, WIJ, POTK)
    V      = V + VIJ
    W      = W + WIJ
  ENDIF
  100 CONTINUE
RETURN
END

SUBROUTINE CSKENERGY ( PAI, PBI, I, TAU, KE )
IMPLICIT NONE
C *****
C ** RETURNS THE KINETIC ENERGY OF GHOST I.
C ***
C ** USAGE:
C ***
C ** THIS SUBROUTINE IS USED TO CALCULATE THE CHANGE OF ENERGY
C ** DURING A TRIAL MOVE OF ATOM I. IT IS CALLED BEFORE AND
C ** AFTER THE RANDOM DISPLACEMENT OF I.
C *****

```

```

INCLUDE      'mc-cs.par'

REAL*8       PAI(MDIM), PBI(MDIM), KE
INTEGER*8    I, TAU, IDIM

C ****
KE = 0.d0

DO IDIM = 1, DIM
  KE=KE+(PAI(IDIM)**2.+PBI(IDIM)**2.)/4/FTM
ENDDO

RETURN
END

SUBROUTINE CSKKNERGY ( I, TAU, KE )
IMPLICIT NONE
C ****
C ** RETURNS THE KINETIC ENERGY OF GHOST I.          **
C **                                                 **
C ** USAGE:                                           **
C **                                                 **
C ** THIS SUBROUTINE IS USED TO CALCULATE THE CHANGE OF ENERGY   **
C ** DURING A TRIAL MOVE OF ATOM I. IT IS CALLED BEFORE AND     **
C ** AFTER THE RANDOM DISPLACEMENT OF I.                   **
C ****
C INCLUDE      'mc-cs.par'

REAL*8       KE
INTEGER*8    I, NI, TAU, IDIM

C ****
KE = 0.d0

DO IDIM = 1, DIM
  KE=KE+(PA(IDIM,TAU,I)**2.+PB(IDIM,TAU,I)**2.)/4/FTM
ENDDO

RETURN
END

REAL*8 FUNCTION RANF ( DUMMY )

C ****
C ** RETURNS A UNIFORM RANDOM VARIATE IN THE RANGE 0 TO 1.      **
C **                                                 **
C **                                                 **
C **                                                 **
C ** GOOD RANDOM NUMBER GENERATORS ARE MACHINE SPECIFIC.        **
C ** PLEASE USE THE ONE RECOMMENDED FOR YOUR MACHINE.           **
C **                                                 **
C **                                                 **
C ** RAND(FLAG) returns a pseudo-random number from a uniform   **
C ** distribution between 0 and 1. If FLAG is 0, the next number   **
C ** in the current sequence is returned; if FLAG is 1, the       **
C ** generator is restarted by CALL SRAND(0); if FLAG has any    **
C ** other value, it is used as a new seed with SRAND.           **
C ****
C REAL*8       DUMMY

RANF = RAND ( )

RETURN
END

SUBROUTINE READCN ( CNFILE )
IMPLICIT NONE
C ****
C ** SUBROUTINE TO READ IN THE CONFIGURATION FROM UNIT 10        **
C ****
C INCLUDE      'mc-cs.par'

CHARACTER    CNFILE*(*), HASH*1, MYFMT*77

INTEGER*8    CNUUNIT, I, J, NNP, NI, IDIM
PARAMETER ( CNUUNIT = 10 )

C ****
OPEN ( UNIT = CNUUNIT, FILE = CNFILE, STATUS = 'OLD' )
WRITE(MYFMT,'(A,I10,A)') '(I3,3X,',5*DIM,'(F12.6,3X))'

READ ( CNUUNIT,* ) HASH, FTNO, NNP
IF ( NNP .NE. NP ) STOP 'N ERROR IN READCN'

DO 100 I = 1, NNP
  READ ( CNUUNIT,* ) HASH, NEXT(I)
  NI = NEXT(I)
  DO 90 J = 1, FTNO
    READ (CNUUNIT, MYFMT) LEXT(I), (RX(IDIM,J,I),IDIM=1,DIM),
    : (RA(IDIM,J,I),IDIM=1,DIM), (RB(IDIM,J,I),IDIM=1,DIM),
    : (PA(IDIM,J,I),IDIM=1,DIM), (PB(IDIM,J,I),IDIM=1,DIM)
  90 ENDDO
  READ (CNUUNIT, MYFMT) LEXT(I), (RX(IDIM,FTNO+1,I),IDIM=1,DIM),
  : (RX(IDIM,FTNO+1,I),IDIM=1,DIM), (RX(IDIM,FTNO+1,I),IDIM=1,DIM),
  : (RX(IDIM,FTNO+1,I),IDIM=1,DIM), (RX(IDIM,FTNO+1,I),IDIM=1,DIM)
  READ ( CNUUNIT,13 )
  READ ( CNUUNIT,13 )
  DO IDIM = 1, DIM
    RX(IDIM,0,NI) = RX(IDIM,FTNO,I)
  ENDDO
  RETURN
END

SUBROUTINE WRITCN ( CNFILE )
IMPLICIT NONE
C ****
C ** SUBROUTINE TO WRITE OUT THE CONFIGURATION TO UNIT 10        **
C ****
C INCLUDE      'mc-cs.par'

CHARACTER    CNFILE*(*), MYFMT*77

REAL*8       RXI(DIM), RAI(DIM), RBI(DIM), PAI(DIM), PBI(DIM)
INTEGER*8    CNUUNIT, I, J, IDIM
PARAMETER ( CNUUNIT = 10 )

C ****
OPEN ( UNIT = CNUUNIT, FILE = CNFILE, STATUS = 'UNKNOWN' )
WRITE(MYFMT,'(A,I10,A)') '(I3,3X,',5*DIM,'(F12.6,3X))'

WRITE(*,*) 'output to file -----'
WRITE ( CNUUNIT,* ) '#',FTNO,NP

DO 100 I = 1, NP
  WRITE ( CNUUNIT, MYFMT ) LEXT(I),(RX(IDIM,1,DIM),
  : (RA(IDIM,1,DIM), (RB(IDIM,1,DIM),
  : (PA(IDIM,1,DIM), (PB(IDIM,1,DIM)
  ENDDO
  WRITE ( CNUUNIT, MYFMT ) LEXT(I),(RX(IDIM,1,DIM),
  : (RA(IDIM,1,DIM), (RB(IDIM,1,DIM),
  : (PA(IDIM,1,DIM), (PB(IDIM,1,DIM)
  ENDDO
  DO IDIM = 1, DIM
    RXI(IDIM) = RX(IDIM,FTNO+1,I)
  ENDDO
  WRITE ( CNUUNIT, MYFMT ) LEXT(I),(RX(IDIM,1,DIM),
  : (RA(IDIM,1,DIM), (RB(IDIM,1,DIM),
  : (PA(IDIM,1,DIM), (PB(IDIM,1,DIM)
  ENDDO
  WRITE ( CNUUNIT,13 )
  WRITE ( CNUUNIT,13 )
  ENDDO
  RETURN
END

SUBROUTINE INITCN ( CNFILE )
IMPLICIT NONE
C ****
C ** SUBROUTINE TO INITIALIZE THE CONFIGURATION TO UNIT 10       **
C ****
C INCLUDE      'mc-cs.par'

REAL*8       RANF
CHARACTER    CNFILE*(*), MYFMT*77

INTEGER*8    CNUUNIT, I, J, IDIM, II
REAL*8       RXI(DIM), RXII(DIM, NP), RIJSQ
PARAMETER ( CNUUNIT = 10 )

C ****
OPEN ( UNIT = CNUUNIT, FILE = CNFILE, STATUS = 'UNKNOWN' )
WRITE(MYFMT,'(A,I10,A)') '(I3,3X,',5*DIM,'(F12.6,3X))'

WRITE ( CNUUNIT,* ) '#',FTNO,NP
RXII = 0.d0

DO 100 I = 1, NP
  WRITE ( CNUUNIT,* ) '#', NEXT(I)
  DO IDIM = 1, DIM
    RXI(IDIM) = SIGMA*(RANF(0.d0)-.5)
  ENDDO
  DO II = 1, NP
    RIJSQ = 0.d0
    DO IDIM = 1, DIM
      RXIJ = RXI(IDIM)-RXII(IDIM,II)
      RXIJ = RXIJ -
      DIINT ( RXIJ/SIGMA ) *SIGMA
      RIJSQ = RIJSQ + RXIJ * RXIJ
    ENDDO
    IF (RIJSQ.LE.EPSR*EPSR) THEN
      GOTO 10
    ENDIF
  DO IDIM = 1, DIM
    RXII(IDIM,II) = RXI(IDIM)
  ENDDO
  DO J = 1, FTNO+1
    WRITE ( CNUUNIT, MYFMT ) I, (RXI(IDIM,1,DIM),
    : (RXI(IDIM,1,DIM), (RXI(IDIM,1,DIM),
    : (RXI(IDIM,1,DIM), (RXI(IDIM,1,DIM)
    ENDDO
  ENDIF
  WRITE ( CNUUNIT,13 )
  WRITE ( CNUUNIT,13 )
  ENDDO
  RETURN
END

```

```

100    CONTINUE
13    FORMAT(82)
      CLOSE ( UNIT = CNUNIT )
      RETURN
END

      SUBROUTINE POT (RIJSQ, VIJ, WIJ, POK)
      IMPLICIT NONE
C **** SUBROUTINE FOR THE PAIR POTENTIAL ****
C ****
C INCLUDE      'mc-cs.par'

      CHARACTER POK*(*) 
      REAL*8   RIJ, RIJSQ, RMIN, RMSQ, RCSQ, SR2, SR6, VIJ, WIJ
      REAL*8   COU3, F
      PARAMETER ( COU3 = 4.76015472795910701328763470057d0 )

C     PI = ACOS(-1.d0)

      VIJ = 0.0
      WIJ = 0.0
      IF (DIM .GT. 3) RETURN
      IF (POK .EQ. 'FREE') RETURN

      IF (POK .EQ. 'LJ') THEN
        RMIN = 2.40**(.1/.6)
        RMIN = TINY(PI)
        RMIN = EPSA
        RMIN = 0.d0
        RIJSQ = RIJSQ / ( SIG * SIG)
        RMSQ = RMIN*RMIN
        RCSQ = RCUT*RCUT
        IF (RIJSQ .LT. RMSQ) THEN
          RIJSQ = RMSQ
          SR2 = 1.d0 / RIJSQ
          SR6 = SR2 * SR2
          VIJ = SR6 * ( SR6 - 1.d0 )
          WIJ = SR6 * ( SR6 - 5.d-1 )
          VIJ = 4.d0 * EPS * VIJ
          WIJ = 48.d0 * EPS * WIJ / 3.d0
        ELSEIF (RIJSQ .GT. RMSQ .AND. RIJSQ .LT. RCSQ) THEN
          SR2 = 1.d0 / RIJSQ
          SR6 = SR2 * SR2
          VIJ = SR6 * ( SR6 - 1.d0 )
          WIJ = SR6 * ( SR6 - 5.d-1 )
          VIJ = 4.d0 * EPS * VIJ
          WIJ = 48.d0 * EPS * WIJ / 3.d0
        ENDIF
      ELSEIF (POK .EQ. 'BUMP') THEN
        RIJ = SQRT( RIJSQ )
        IF (RIJ .LE. SIG) THEN
          VIJ = EPSR
        ENDIF
      ELSEIF (POK .EQ. 'PSW') THEN
        RIJ = SQRT( RIJSQ )
        IF (RIJ .LE. SIG) THEN
          VIJ = EPSR
        ELSEIF (RIJ .LE. SIG + RCUT .AND. RIJ .GT. SIG) THEN
          VIJ = -EPSA
        ENDIF
      ELSEIF (POK .EQ. 'COULOMB') THEN
        F = NP/(NP-1)
        RIJ = SQRT( RIJSQ )
        IF (RIJ .GT. EPS) THEN
          IF (DIM .EQ. 3) THEN
            VIJ = -F*COUS*SIG/SIGMA/2.
            VIJ = VIJ + SIG/RIJ
          ELSEIF (DIM .EQ. 2) THEN
            VIJ = -F*(6.-PI+LOG(4.)-4.*LOG(SIGMA/SIG))/4.
            VIJ = VIJ - LOG(RIJ/SIG)
          ELSEIF (DIM .EQ. 1) THEN
            VIJ = F*SIGMA/SIG/4.
            VIJ = VIJ - RIJ/SIG
          ENDIF
        ELSE
          RIJ = EPS
          IF (DIM .EQ. 3) THEN
            VIJ = -F*COUS*SIG/SIGMA/2.
            VIJ = VIJ + SIG/RIJ
          ELSEIF (DIM .EQ. 2) THEN
            VIJ = -F*(6.-PI+LOG(4.)-4.*LOG(SIGMA/SIG))/4.
            VIJ = VIJ - LOG(RIJ/SIG)
          ELSEIF (DIM .EQ. 1) THEN
            VIJ = F*SIGMA/SIG/4.
            VIJ = VIJ - RIJ/SIG
          ENDIF
        ENDIF
      ELSEIF (POK .EQ. 'HARMONIC') THEN
        VIJ = 0.5*RIJSQ/SIG**2.
      ENDIF

      RETURN
END

      subroutine cs(qqa,ppa,qqb,ppb,q,qp,zeta)
      implicit none
C **** THE CHOERENT STATE ****
C ****
C INCLUDE      'mc-cs.par'

      complex*16 ii,aa(mdim),bb(mdim),psia,psib,gab,zeta
      real*8   mo, qqa(mdim), ppa(mdim), qqb(mdim), ppb(mdim)
      real*8   q(mdim), qp(mdim), aux
      real*8   rgab,rzeta
      integer*8 idim

      parameter ( ii = cmplx(0.d0,1.d0) )
      harmonic oscillator
      mo = moho

      do idim=1,dim
        aa(idim) = (mo*qqa(idim)+ii*ppa(idim))/sqrt(2*mo)
        bb(idim) = (mo*qqb(idim)+ii*ppb(idim))/sqrt(2*mo)
      enddo

      coherent states
      psia=0.d0
      psib=0.d0
      do idim=1,dim
        aux=q(idim)-sqrt(2/mo)*dble(aa(idim))
        aux=aux-dnint(aux/sigma)*sigma
        psia = psia - aux**2.*mo/2
        psia = psia + ii*aux*sqrt(2*mo)*dimag(aa(idim))
        aux=qp(idim)-sqrt(2/mo)*dble(bb(idim))
        aux=aux-dnint(aux/sigma)*sigma
        psib = psib - aux**2.*mo/2
        psib = psib + ii*aux*sqrt(2*mo)*dimag(bb(idim))
      enddo
      psia = exp(psia)
      psia = (mo/pi)**(dim/4.)*psia
      psib = exp(psib)
      psib = (mo/pi)**(dim/4.)*psib

      normalization
      gab = 0.d0
      do idim=1,dim
        gab = gab - (cdabs(aa(idim))**2.+cdabs(bb(idim))**2.)/2
        gab = gab + dconjg(aa(idim))*bb(idim)
        gab = gab + (qqa(idim)*ppa(idim)-qqb(idim)*ppb(idim))*ii/2
      enddo
      gab = exp(gab)
      rgab = dble(gab)

      zeta
      zeta = 0.d0
      do idim=1,dim
        zeta = zeta - ls*(ppb(idim)**2.+ppa(idim)**2.)/4/ftm
      enddo
      zeta = psia*dconjg(psib)*gab*exp(zeta)
      rzeta = dble(zeta)

      print *, zeta
      return
end

      subroutine acc_pcs(p,ip,kp,j)
      implicit none
! complete acceptance probability
      INCLUDE      'mc-cs.par'

      integer*8 ip,kp,j,idim
      real*8 p, rho
      real*8 rxink,rxnik
      real*8 rxini,rxknk
      integer*8 mcm

      p*exp(-ls*p) ! contribution from the pair potential
      if (ip.eq.kp) return

      mcm = j+mbm-floor((j+mbm-1)/ftn0)*ftn0
      rho=0.d0
      do idim=1,dim
        if(j+mbm.le.ftn0)then
          rxink=rx(idim,j,ip)-rx(idim,mcm,kp)
          rxnik=rx(idim,j,kp)-rx(idim,mcm,ip)
          rxini=rx(idim,j,ip)-rx(idim,mcm,ip)
          rxknk=rx(idim,j,kp)-rx(idim,mcm,kp)
          rxink=rxink-dnint(rxink/sigma)*sigma
          rxnik=rxnik-dnint(rxnik/sigma)*sigma
          rxini=rxini-dnint(rxini/sigma)*sigma
          rxknk=rxknk-dnint(rxknk/sigma)*sigma
          rho=rho+rxink**2+rxnik**2-rxini**2-rxknk**2
        else
          rxink=rx(idim,j,ip)-rx(idim,mcm,next(kp))
          rxnik=rx(idim,j,kp)-rx(idim,mcm,next(ip))
          rxini=rx(idim,j,ip)-rx(idim,mcm,next(ip))
          rxknk=rx(idim,j,kp)-rx(idim,mcm,next(kp))
          rxink=rxink-dnint(rxink/sigma)*sigma
          rxnik=rxnik-dnint(rxnik/sigma)*sigma
          rxini=rxini-dnint(rxini/sigma)*sigma
          rxknk=rxknk-dnint(rxknk/sigma)*sigma
          rho=rho+rxink**2+rxnik**2-rxini**2-rxknk**2
        endif
      enddo
      rho=ftm*rho/(2.*mbm*ls)
      p*p*exp(-3*rho)
      return
end

      subroutine acccs(p,ip,kp,j,rxp,rxpp)
      implicit none
! complete acceptance probability
      INCLUDE      'mc-cs.par'

      integer*8 ip,kp,j,idim
      real*8 p,rxp(mdim,0:n),rxpp(mdim,0:n)
      real*8 rho,rxr,rxpp,rxxp
      integer*8 mcm,i,ll,nip,nkp

      mcm = j+mbm-floor((j+mbm-1)/ftn0)*ftn0
      nip=next(ip)

```

```

nkp=next(kp)

rho=0.d0
ll=-1
if(j+mbm.le.ftn0)then
  do i=j,mcm-1
    ll=ll+1
    do idim=1,dim
      rrx=rx(idim,i,ip)-rx(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=ra(idim,i,ip)-ra(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rb(idim,i,ip)-rb(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rxp(idim,ll)-rxp(idim,ll+1)
      rrx=rxxp-dnint(rxxp/sigma)*sigma
      rho=rho-3*rrxp**2
    enddo
  enddo
else
  do i=j,ftn0
    ll=ll+1
    do idim=1,dim
      rrx=rx(idim,i,ip)-rx(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=ra(idim,i,ip)-ra(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rb(idim,i,ip)-rb(idim,i+1,ip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rxp(idim,ll)-rxp(idim,ll+1)
      rrx=rxxp-dnint(rxxp/sigma)*sigma
      rho=rho-3*rrxp**2
    enddo
  enddo
  do i=1,mcm-1
    ll=ll+1
    do idim=1,dim
      rrx=rx(idim,i,nip)-rx(idim,i+1,nip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=ra(idim,i,nip)-ra(idim,i+1,nip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rb(idim,i,nip)-rb(idim,i+1,nip)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rxp(idim,ll)-rxp(idim,ll+1)
      rrx=rxxp-dnint(rxxp/sigma)*sigma
      rho=rho-3*rrxp**2
    enddo
  enddo
endif
if (ip.ne.kp) then
  ll=-1
  if(j+mbm.le.ftn0)then
    do i=j,mcm-1
      ll=ll+1
      do idim=1,dim
        rrx=rx(idim,i,kp)-rx(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=ra(idim,i,kp)-ra(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=rb(idim,i,kp)-rb(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=rxp(idim,ll)-rxp(idim,ll+1)
        rrx=rxxp-dnint(rxxp/sigma)*sigma
        rho=rho-3*rrxp**2
      enddo
    enddo
  else
    do i=j,ftn0
      ll=ll+1
      do idim=1,dim
        rrx=rx(idim,i,kp)-rx(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=ra(idim,i,kp)-ra(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=rb(idim,i,kp)-rb(idim,i+1,kp)
        rrx=rrx-dnint(rrx/sigma)*sigma
        rho=rho+rrx**2
        rrx=rxp(idim,ll)-rxp(idim,ll+1)
        rrx=rxxp-dnint(rxxp/sigma)*sigma
        rho=rho-3*rrxp**2
      enddo
    enddo
  do i=1,mcm-1
    ll=ll+1
    do idim=1,dim
      rrx=rx(idim,i,nkp)-rx(idim,i+1,nkp)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=ra(idim,i,nkp)-ra(idim,i+1,nkp)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
      rrx=rb(idim,i,nkp)-rb(idim,i+1,nkp)
      rrx=rrx-dnint(rrx/sigma)*sigma
      rho=rho+rrx**2
    enddo
  enddo
  rrxpp=rrxp(idim,ll)-rrxp(idim,ll+1)
  rrxpp=rrxp-dnint(rrxp/sigma)*sigma
  rho=rho-3*rrxp**2
end
rrxpp=rrxp(idim,ll)-rrxp(idim,ll+1)
rrxpp=rrxp-dnint(rrxp/sigma)*sigma
rho=rho-3*rrxp**2
rho=rho-ftm*rho/(2.*ls)
p=p*exp(-rho)
return
end

mc=cs.par
-----
C ****
C ** MC-CS-PAR **
C ****
C ***** INTEGER*8 MDIM, N, MNP
REAL*8 PI
PARAMETER ( MDIM = 10 ) ! maximum number of dimensions
PARAMETER ( MNP = 1000 ) ! maximum number of particles
PARAMETER ( N = 3000 ) ! maximum number of time slices
PARAMETER ( PI = 3.14159265358979323846264338328d0 )
REAL*8 RX(MDIM,0:N,MNP)
REAL*8 RA(MDIM,0:N,MNP), RB(MDIM,0:N,MNP)
REAL*8 PA(MDIM,0:N,MNP), PB(MDIM,0:N,MNP)
REAL*8 FTN, LS
REAL*8 SIGMA, SIG, EPSA, EPSR, EPS, RCUT
REAL*8 MOHO
INTEGER*8 FTNO, NP, NEXT(MNP), LEXT(MNP), MEXT(N,MNP), MBM
INTEGER*8 DIM

! ghost path
COMMON / BLOCK0 / RA, RB, PA, PB
! real path
COMMON / BLOCK1 / RX, DIM
! pair-potential parameters
COMMON / BLOCK2 / SIG, EPSR, EPSA, EPS, RCUT
! mass, timestep, box edge, # timeslices, # particles
COMMON / BLOCK3 / FTN, LS, SIGMA, FTNO, NP
! permutations
COMMON / BLOCK4 / NEXT, LEXT, MEXT, MBM
! harmonic oscillator for coherent states
COMMON / BLOCK5 / MOHO

----- data-cs-2.in -----
0 number of spatial dimensions (1,2,3,...<= MDIM)
2
1 seed of the random sequence RAND
3
2 if bose (T/F)
T
3 number of particles (<= MNP)
16
4 the potential SUBROUTINE POT
LJ
5 # time slices = 1/temperature/timestep (< N)
250
6 mass (hbar = kb = 1)
0.0830594d0
7 # of cycles (nstep)
5000000000000
8 # of steps between output lines (iprint)
100
9 # of steps between configuration saves (isave)
100
10 # of steps for equilibration (iequi)
0
11 # of steps for acceptance ratios (iratio)
100
12 configuration file name
conf.xyz
13 enter 0 if initialization needed
0
14 density THERMODYNAMICS
.05d0
15 temperature THERMODYNAMICS
1.d0
16 maximum displacement/box edge
.005d0
17 number of bridge timeslices (>= 1; <= #5)
150
18 potential cutoff distance (LJ) SUBROUTINE POT
2.5d0
19 sig (LJ 2.566) SUBROUTINE POT
2.556d0
20 epsr (LJ INIT) SUBROUTINE POT
3.d0
21 epsa SUBROUTINE POT
1.d0
22 eps (LJ 10.22) SUBROUTINE POT
1.022d1
23 moho = mass_ho*omega HO
20.7649d0
24 if only displace (T/F)
F
25 if only bridge (T/F)
F
26 if zero path initially (0 in 13) (T/F)
F

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