Path Integral Monte Carlo on a Sphere

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CONTENTS

I. INTRODUCTION

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II. MANY BODY PATH INTEGRAL ON A RIEMANNIAN MANIFOLD

A many body system is composed of N *distinguishable* particles of mass m with positions in $R = (\boldsymbol{r}_1, \boldsymbol{r}_2, \ldots, \boldsymbol{r}_N) = (\{\boldsymbol{r}_i\})$ where each position vector $\boldsymbol{r}_i = (r_i^1, r_i^2, \ldots, r_i^d) = (\{r_i^\alpha\})$ in d dimensions. On a Riemannian manifold of dimension d and metric tensor $g_{\alpha\beta}(\boldsymbol{r})$, the geodesic distance between two infinitesimally close points R and R' is $d\tilde{s}^2(R, R') = \sum_{i=1}^N ds^2(\boldsymbol{r}_i, \boldsymbol{r}_i')$ where $ds^2(\boldsymbol{r}, \boldsymbol{r}') = g_{\alpha\beta}(\boldsymbol{r} - \boldsymbol{r}')^\alpha(\boldsymbol{r} - \boldsymbol{r}')^\beta$. Moreover,

$$\tilde{g}_{\mu\nu}(R) = g_{\alpha_1\beta_1}(\boldsymbol{r}_1) \otimes \ldots \otimes g_{\alpha_N\beta_N}(\boldsymbol{r}_N), \tag{2.1}$$

$$\tilde{g}(R) = \prod_{i=1}^N \det \|g_{\alpha_i\beta_i}(\boldsymbol{r}_i)\|, \tag{2.2}$$

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where $||\tilde{g}_{\mu\nu}||$ is a matrix made of N diagonal blocks $||g_{\alpha_i\beta_i}||$ with $i = 1, 2, \ldots, N$. The Laplace-Beltrami operator on the manifold of dimension dN is

$$\Delta_R = \tilde{g}^{-1/2}\nabla_\mu(\tilde{g}^{1/2}\tilde{g}^{\mu\nu}\nabla_\nu), \tag{2.3}$$

where $\nabla = \nabla_R$, $\tilde{g}^{\gamma\nu}$ is the inverse of $\tilde{g}_{\gamma\nu}$, i.e. $\tilde{g}_{\mu\gamma}\tilde{g}^{\gamma\nu} = \delta_\mu{}^\nu$ the Kronecker delta, and a sum over repeated indexes is tacitly assumed.

We will first assume *free*, non interacting bodies, with an Hamiltonian $\mathscr{H}$ that reduces to the one of the free gas in flat space. For the sake of simplicity [1] we will choose

$$\mathscr{H} = -\lambda\Delta_R, \tag{2.4}$$

with $\lambda = \hbar^2/2m$.

For *interacting* bodies we will then have more generally

$$\mathscr{H} = -\lambda\Delta_R + V(R), \tag{2.5}$$

where V is the potential energy of the system, that we here assume only a function of the particles positions and bounded from below.

The density matrix ρ of the system obeys Bloch equation

$$\frac{\partial\rho(t)}{\partial t} = -\mathscr{H}\rho(t), \tag{2.6}$$

$$\rho(0) = \mathbb{1}, \tag{2.7}$$

where t is the imaginary time with the dimensions of an energy and $\mathbb{1}$ the identity matrix. The position representation of the density matrix is then obtained from $\rho(R, R'; t) = \langle R|\rho(t)|R'\rangle$ with $\langle R|R'\rangle = \delta(R - R')/\sqrt{\tilde{g}(R)}$ where δ is a dN dimensional Dirac delta function. In the small imaginary time τ limit the position representation of the density matrix is

$$\rho(R, R'; \tau) \propto \tilde{g}(R)^{-1/4}\sqrt{\mathscr{D}(R, R'; \tau)}\,\tilde{g}(R')^{-1/4}e^{\lambda\tau\mathscr{R}(R)/6}e^{-\mathscr{S}(R,R';\tau)}, \tag{2.8}$$

where $\mathscr{R}$ is the scalar curvature of the manifold [2], $\mathscr{S}$ the action, and $\mathscr{D}$ the van Vleck's determinant [4, 5]

$$\mathscr{D}_{\mu\nu} = \nabla_\mu\nabla'_\nu\mathscr{S}(R, R'; \tau), \tag{2.9}$$

$$\det||\mathscr{D}_{\mu\nu}|| = \mathscr{D}(R, R'; \tau), \tag{2.10}$$

where $\nabla = \nabla_R$ and $\nabla' = \nabla_{R'}$. This determinant is the Jacobian of the transformation from the initial conditions given by fixing the pair of momentum and coordinate to the boundary conditions given by specifying the pair of initial and final coordinates needed in the path integral formulation. For the density matrix (2.8) the volume element for integration is $\sqrt{\tilde{g}(R)}\,dR$. The two factors $\tilde{g}^{-1/4}$ are needed in order to have for the density matrix a bidensity for which the boundary condition to Bloch equation is simply a Dirac delta function $\rho(R, R'; 0) = \delta(R - R')$. The square root of the van Vleck determinant factor takes into account the density of paths among the minimum extremal region for the action (see Chapter 12 of Ref. [5]).

For the *action* $\mathscr{S}$, the *kinetic-action* $\mathscr{K}$, and the *inter-action* $\mathscr{U}$ we have [3]

$$\mathscr{S}(R, R'; \tau) = \mathscr{K}(R, R'; \tau) + \mathscr{U}(R, R'; \tau), \tag{2.11}$$

$$\mathscr{K}(R, R'; \tau) = \frac{dN}{2}\ln(4\pi\lambda\tau) + \frac{d\tilde{s}^2(R, R')}{4\lambda\tau}. \tag{2.12}$$

In particular the kinetic-action is responsible for a diffusion of the random walk with a single particle variance on the α, β components equal to $\sigma^2_{\alpha\beta}(\boldsymbol{r}) = 2\lambda\tau/g_{\alpha\beta}(\boldsymbol{r})$. The inter-action is defined as $\mathscr{U} = \mathscr{S} - \mathscr{K}$ and for potential energies bounded from below one can resort to Trotter formula [6] to reach the *primitive approximation* [4]

$$\mathscr{U}(R, R'; \tau) = \tau[V(R) + V(R')]/2. \tag{2.13}$$

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[1] This is a delicate point and should be studied more carefully [2]. Especially for what concerns ordering ambiguities. We here appeal to simplicity.

[2] The factor depending on the curvature of the manifold is due to Bryce DeWitt [3]. For a space of constant curvature there is clearly no effect, as the term due to the curvature just leads to a constant multiplicative factor that has no influence on the measure of the various observables.

[3] The expression for $\mathscr{K}$ is the one of Eq. (24.16) of Ref. [5] to lowest order in $R - R'$.

[4] See Ref. [7] for a numerical analysis of the accuracy of this approximation and for possible its refinements.

For non interacting bodies $\mathscr{U} = 0$. Note that, even to lowest order in $R - R'$ [5], the path integral in the curved manifold for the non interacting system will not coincide with the one in flat space since it is not possible with a change of coordinates to simply remove the metric factor from both $d\tilde{s}^2$ and the volume element of integration, if not only locally. In fact this would require a *non coordinate basis* [9].

Given then an observable $\mathscr{O}$ we can determine its thermal average at an absolute temperature T from

$$\langle \mathscr{O} \rangle = \mathrm{tr}\{\rho(\beta)\mathscr{O}\}/Z_N, \tag{2.14}$$

$$Z_N = \mathrm{tr}\{\rho(\beta)\}, \tag{2.15}$$

where $\beta = 1/k_B T$ with k_B Boltzmann constant and Z_N the canonical partition function.

The position representation of the density matrix at an imaginary time $t = \beta$ is obtained through a path integral

$$\rho(R, R'; \beta) = \langle R|\rho(\beta)|R'\rangle = \int \prod_{k=0}^{M-1} [\rho(R_k, R_{k+1}; \tau)\, dR_k]\, \delta(R_0 - R)\delta(R_M - R')\, dR_M, \tag{2.16}$$

where we have discretized the imaginary time β into M *timeslices* with a small *timestep* $\tau = \beta/M$, a *bead* $R_k = (\{\boldsymbol{r}_{i,k}\}) = (\{r_{i,k}^\alpha\})$ at each timeslice $k = 1, 2, \ldots, M$. We will also call *link* a pair of contiguous beads.

For *identical* bodies if they satisfy to the Bose-Einstein statistics one needs to symmetrize the distinguishable density matrix, if they satisfy to the Fermi-Dirac statistics one needs to antisymmetrize it [10]. In these cases we can then write [6]

$$\rho_\pm(R, R'; \beta) = \frac{1}{N!} \sum_{\mathscr{P}} \mathrm{sgn}(\mathscr{P})\rho(\mathscr{P}R, R'; \beta), \tag{2.17}$$

$$\mathrm{sgn}(\mathscr{P}) = (\pm 1)^{\sum_{\nu=1}^{N}(\nu-1)C_\nu}, \tag{2.18}$$

where $\mathscr{P}$ is any permutation of the N particles such that $\mathscr{P}R = (\boldsymbol{r}_{\mathscr{P}1}, \boldsymbol{r}_{\mathscr{P}2}, \ldots, \boldsymbol{r}_{\mathscr{P}N})$, with sign $\mathrm{sgn}(\mathscr{P})$. Any permutations can be broken into cycles $\mathscr{P} = \{C_\nu\}$ where C_ν is the number of cycles of length ν in $\mathscr{P}$. In the sum over the permutation one should use a $+1$ for the symmetrization necessary for bosons and -1 for the antisymmetrization necessary for fermions, in $\mathrm{sgn}(\mathscr{P})$.

On a surface, $d = 2$, for *impenetrable* identical bodies, one can also have anyonic statistics [11]. In this case it is necessary to consider, more generally,

$$\rho_\nu(R, R'; \beta) = \sum_{\alpha \in B_N} \mathrm{Re}[\chi(\alpha)]\rho_\alpha(R, R'; \beta), \tag{2.19}$$

$$\chi(\text{paths } R(t) \text{ with } n \text{ braids among the pairs of single particle paths } \boldsymbol{r}(t)) = e^{-i\nu n\pi}, \tag{2.20}$$

where B_N is the infinite braid group which admits an infinite number of unitary one dimensional representations χ parametrized by an arbitrary number ν which determines the statistics and ρ_α is the distinguishable density matrix obtained from paths of kind α only. Clearly for $\nu = 0$ we recover the Bose-Einstein statistics and for $\nu = 1$ the Fermi-Dirac statistics. So we will be interested in values of $0 < \nu < 1$.

The braid group is the fundamental group of the quotient space $(S^{2N} - \Delta)/S_N$ where S^2 is the (two) sphere, $\Delta = \{R \mid \boldsymbol{r}_i = \boldsymbol{r}_j \text{ for some } i \neq j\}$, and S_N is the group of permutation of N bodies. We then see how paths $R(t)$ with different numbers of crossings between single particle paths $\boldsymbol{r}(t)$ belong to different homotopy classes and one cannot be deformed continuously into the other. Therefore in order to take care of the density matrix of identical impenetrable bodies it is necessary to sum over all the topologically disjoint homotopy classes as is done in Eq. (2.19).

III. THE SPHERE

A sphere of radius a is the surface, $d = 2$ [7], with metric $ds^2 = g_{\alpha\beta}dr^\alpha dr^\beta = s^2(d\theta^2 + \sin^2\theta\, d\varphi^2)$, of constant positive scalar curvature $2/a^2$ so that $\mathscr{R} = 2N/a^2$. The polar angle $r^1 = \theta \in]0, \pi]$ and the azimuthal angle $r^2 = \varphi \in]-\pi, \pi]$ are the contravariant coordinates of the position vector $\boldsymbol{r} \in \mathscr{C}$ with $\mathscr{C} =]0, \pi] \times]-\pi, \pi]$ the single particle consfiguration space. On the sphere $g(\boldsymbol{r}) = |\sin\theta|$ and in the small $\tau \to 0$ limit $\tilde{g}(R)^{-1/4}\sqrt{D(R, R'; \tau)}\,\tilde{g}(R')^{-1/4} \to (1/2\lambda\tau)^N$. So

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[5] For next orders corrections see for example Ref. [8].

[6] One can symmetrize or antisymmetrize respect to the first, the second or both the arguments of the distinguishable density matrix. We here choose the first case.

[7] So it is conformally flat as any Riemannian manifold of dimension $d \leq 3$.

we see how both the curvature term and the van Vleck factor, being constant, simply drop off from the measure of the various observables of Eq. (2.14).

The position of a particle on the sphere in the three dimensional Euclidean space embedding the sphere is

$$\begin{cases} x = a\sin\theta\cos\varphi \\ y = a\sin\theta\sin\varphi \\ z = a\cos\theta \end{cases} \tag{3.1}$$

and the particle path in it is $\boldsymbol{q}(t) = (x(t), y(t), z(t))$.

The geodesic distance between particles $\boldsymbol{r}_i$ and $\boldsymbol{r}_j$ is

$$s_{ij} = s(\boldsymbol{r}_i, \boldsymbol{r}_j) = a\arccos\left[\cos(r_i^1)\cos(r_j^1) + \sin(r_i^1)\sin(r_j^1)\cos(r_i^2 - r_j^2)\right], \tag{3.2}$$

whereas the Euclidean distance is

$$d_{ij} = d(\boldsymbol{r}_i, \boldsymbol{r}_j) = 2a\sin[\arccos(\hat{\boldsymbol{q}}_i \cdot \hat{\boldsymbol{q}}_j)/2]., \tag{3.3}$$

where $\hat{\boldsymbol{q}}_i = \boldsymbol{q}_i/a$ is the versor that from the center of the sphere points towards the center of the ith particle.

We use the Metropolis algorithm [12, 13] to evaluate the average of Eq. (2.14). In order to explore ergodically the configuration space $\boldsymbol{r} = (\theta, \varphi) \in \mathscr{R}$ to sample the distinguishable density matrix we use the transition *displacement* move described in Appendix A. In order to sample the permutation sum of Eq. (2.17) needed for identical bodies we use a transition move combination of 2 Brownian *bridges* between unlike bodies as described in Appendix B. And in order to sample the sum over the homotopy classes of Eq. (2.19) needed for identical impenetrable bodies we use a combination of bridge and displacement transition moves as described in Appendix C. Note that the displacement moves can be freely substituted by moves of bridges connecting only like bodies. But we found them still useful from a purely formal point of view where one starts from simple single bead moves and only later builds more complex many beads moves. Moreover since the single bead displacement move is simple to construct it can serve as a test upon the employment of many beads moves.

We will work in the canonical ensemble with fixed number of particles N, surface area $A = 4\pi a^2$, surface density $\sigma = N/4\pi a^2$, and absolute temperature $T = 1/k_B\beta$. The many body system *degeneracy parameter* is $\Theta = T/T_D$ where the degeneracy temperature $T_D = \sigma\hbar^2/mk_B$. For temperatures higher than T_D, $\Theta \gg 1$, quantum effects are less relevant. We will treat both the non interacting fluid $V = 0$ and the Coulomb fluid

$$V(R) = \sum_{i<j} \frac{e^2}{d_{ij}}, \tag{3.4}$$

where e is the unit of charge and we are assuming that the particles, moving on the sphere, interact with the three dimensional Coulomb potential [8]. The *Coulomb coupling constant* is $\Gamma = \beta e^2/a_0 r_s$ with $a_0 = \hbar^2/me^2$ the Bohr radius and the Wigner-Seitz radius $r_s = (4\pi\sigma)^{-1/2}/a_0$. At weak coupling, $\Gamma \ll 1$, the plasma becomes weakly correlated and approach the ideal gas limit. This will occur at high temperature and/or low density.

Choosing length in units of Wigner-Seitz radius, $a_0 r_s = a/\sqrt{N}$, and energy in units of Rydberg, $\mathrm{Ry} = \hbar^2/2ma_0^2$, we have $\lambda = \mathrm{Ry}/r_s^2$, $\Gamma = \beta(2/r_s)$, and $\Theta = (2\pi r_s^2)/\beta$.

TABLE I. Thermodynamic states treated in our simulations: N number of particles, β (Ry^{-1}) inverse temperature, r_s Wigner-Seitz radius, e_K (Ry) kinetic energy per particle from the thermodynamic estimator as explained in Ref. [7], and e_V (Ry) potential energy per particle. The other quantities were introduced in the main text. We chose length in units of Wigner-Seitz radius and energy in units of Rydberg. We chose M such as to have $\tau = 0.01$ in all cases.

case	M	N	a	$a/\sqrt{N}$	β	Γ	Θ	e_K	e_V
A									
B									
C									
...									

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[8] Note that this is not the only possible choice since we could as well choose particles "living in" [14] the surface of the sphere as done for example in Ref. [15–17].

IV. DISTINGUISHABLE BODIES

Here we use the Metropolis algorithm with the transition displacement moves of Eqs. (A1) and (A2) or equivalently bridge moves of Eqs. (B1) and (B2) but between like particles.

Apart from the *kinetic energy* per particle e_K and the *potential energy* per particle e_V [9] we will also measure the *radial distribution function*, $g(r) = \langle \mathcal{O} \rangle$ for which we may use the following histogram estimator,

$$O(R; r) = \sum_{i \neq j} \frac{1_{]r-\Delta/2, r+\Delta/2]}(d_{ij})}{N n_{id}(r)},$$
(4.1)

where Δ is the histogram bin, $1_{]a,b]}(x) = 1$ if $x \in]a, b]$ and 0 otherwise, and

$$n_{id}(r) = N \left[\left(\frac{r + \Delta/2}{2a} \right)^2 - \left(\frac{r - \Delta/2}{2a} \right)^2 \right],$$
(4.2)

is the average number of particles on the spherical crown $]r - \Delta/2, r + \Delta/2]$ for the ideal gas of density σ. We have that $\sigma^2 g(r)$ gives the probability, that sitting on a particle at $\boldsymbol{r}$, one finds another particle at $\boldsymbol{r}'$ with $r = d(\boldsymbol{r}, \boldsymbol{r}')$.

FIG. 1. The radial distribution function for the Boltzmann plasma on the sphere of radius $a = 5a_0$ at an inverse temperature β I.

V. IDENTICAL BODIES

Here we use the Metropolis algorithm with the transition displacement moves of Eqs. (A1) and (A2) and bridge moves of Eqs. (B1) and (B2) between unlike particles as described in Appendix B to produce the necessary particles exchanges.

Bosons

Given the *superfluid* density σ_s and the *normalfluid* density $\sigma_n = 1 - \sigma_s$, the area estimator for the superfluid fraction is given by [7, 18]

$$f_s = \frac{\sigma_s}{\sigma} = 1 - \frac{\sigma_n}{\sigma} = \frac{2m \langle \mathscr{A}^2 \rangle}{\beta \lambda I_c},$$
(5.1)

where, if ϵ is the Levi-Civita antisymmetric symbol,

$$\mathscr{A} = \frac{1}{2} \sum_{i,k} \epsilon_{\alpha\beta} r_{i,k}^\alpha r_{i,k+1}^\beta \sqrt{g(\boldsymbol{r}_{i,k})},$$
(5.2)

is the area occupied by all the single particle paths and

$$I_c = m \left\langle \sum_{i,k} g_{\alpha\beta}(\boldsymbol{r}_{i,k}) r_{i,k}^\alpha r_{i,k+1}^\beta \right\rangle,$$
(5.3)

is the classical moment of inertia of all the single particle paths.

FIG. 2. The supefluid fraction for the spinless boson plasma on the sphere of radius $a = 5a_0$ at an inverse temperature β I.

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[9] The estimators for these observables are carefully described in Ref. [7].

Fermions

Here we use the Metropolis algorithm within the *restricted path integral* strategy [1, 10] which uses the density matrix of the non interacting system to guide the *nodes* [19] of the interacting system in order to avoid the *fermions sign problem* [10].
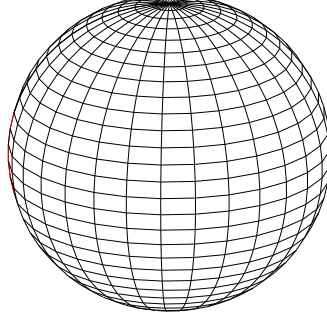

FIG. 3. Snapshot of the macroscopic path during the simulation. The different paths have different colors.

VI.   ANYONIC BODIES

Here we use the Metropolis algorithm with the transition displacement moves of Eqs. (A1) and (A2) and bridge moves of Eqs. (B1) and (B2) between unlike particles as described in Appendix C counting at each 2 bridges transition the number of single particles crossings in the newly generated R.

Appendix A: The transition displacement move

In order to explore the θ and φ configuration space $\mathscr{C} =]0,\pi]\times]-\pi,\pi]$ on the sphere it is convenient to propose the following transition move for each particle in a randomly chosen bead

$$\theta_{\text{new}} = \theta_{\text{old}} + \Delta_\theta(\eta - 1/2), \tag{A1}$$
$$\varphi_{\text{new}} = \varphi_{\text{old}} + \Delta_\varphi(\eta - 1/2), \tag{A2}$$

where $\eta \in [0,1]$ is a uniform pseudo random number and Δ_θ and Δ_φ are two positive quantities measuring the θ-displacement and the φ-displacement respectively.

This transition move can bring $\boldsymbol{r}_{\text{new}}$ out of $\mathscr{C}$ so it is also necessary to bring it back into $\mathscr{C}$ with the following subsequent chain of transformations

$$\begin{cases} \theta_{\text{new}} \to 2\pi \, \texttt{ABS}[\theta_{\text{new}}/2\pi - \texttt{FLOOR}(\theta_{\text{new}}/2\pi + 1/2)], \\ \varphi_{\text{new}} \to \varphi_{\text{new}} - \pi \left\{ \texttt{INT}(\theta_{\text{new}}/\pi) + [\texttt{SGN}(\theta_{\text{new}}) - 1]/2 \right\}, \\ \varphi_{\text{new}} \to \varphi_{\text{new}} - 2\pi \, \texttt{NINT}(\varphi_{\text{new}} - 2\pi \, \texttt{NINT}(\varphi_{\text{new}}/2\pi)/2\pi), \end{cases} \tag{A3}$$

where $\texttt{ABS}$ is the function absolute value, $\texttt{SGN}$ is the sign function, $\texttt{INT}$ is the function integer part, $\texttt{NINT}$ is the nearest integer function and $\texttt{FLOOR}$ is the floor function. This transformations chain is such that when the path

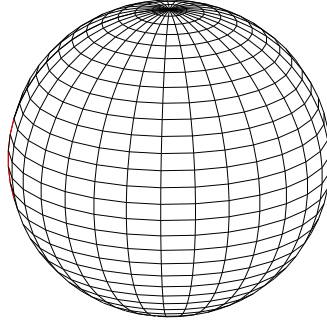FIG. 4. Snapshot of the macroscopic path during the simulation. Paths belonging to different homotopy classes have different colors.

$\boldsymbol{r}(t) = (\theta(t), \varphi(t))$ crosses a pole φ changes by π. One can easily convince himself that this chain does not alter the uniformity of the probability distribution of $\boldsymbol{r}_{\rm new}$ in $\mathscr{C}$.

In the simulation we choose Δ_θ and Δ_φ so to have acceptance ratios as close as possible to $1/2$ in the acceptance/rejection choices for the random walk transition displacement moves of the Metropolis algorithm. The transition probability distribution function for the displacement move of the Metropolis algorithm will be uniform so it will drop out of the acceptance probability distribution function.

Appendix B: The transition bridge move

In order to take into account the particles permutations it is necessary to construct two Brownian bridges between two different [10] randomly chosen particles in two randomly chosen beads to generate an exchange between the two particles. With one bridge we connect particle 1 on bead R_i to particle 2 on bead R_j and with the other we connect particle 2 on bead R_i to particle 1 on bead R_j with $i < j$. This will produce an *exchange* of particles 1 and 2.

The Brownian bridge between particle l at $\boldsymbol{r}_{l,i}$ and particle m at $\boldsymbol{r}_{m,j}$ is built like so [7],

$$\boldsymbol{r}_{\rm new,i} = \boldsymbol{r}_{l,i} \tag{B1}$$

$$\boldsymbol{r}_{\rm new,k} = \boldsymbol{r}_{\rm new,k-1} + \frac{(\boldsymbol{r}_{m,j} - \boldsymbol{r}_{\rm new,k-1})}{j - k + 1} + \boldsymbol{\xi}(\boldsymbol{r}_{\rm new,k-1}) \quad k = i + 1, \dots, j - 1 \tag{B2}$$

where $\xi^\alpha(\boldsymbol{r})$ is a random number with a Gaussian probability distribution with variance $\sigma_{\alpha\alpha}^2(\boldsymbol{r})(j - k)/(j - k + 1)$ where $\sigma_{\alpha\alpha}^2(\boldsymbol{r}) = 2\lambda\tau/g_{\alpha\alpha}(\boldsymbol{r})$ is the diagonal free particle variance. In order to produce an exchange of two *unlike* particles l and m with $l \neq m$ one needs a 2 bridges transition as described above.

At each accepted transition, it is then necessary to bring the $\boldsymbol{r}_{\rm new}$ back into the region $\mathscr{C}$ through the transformations chain of Eq. (A3) in Appendix A.

Note that the metric enters the free particle variance since it is not possible by a change of coordinates to remove it both from the kinetic-action and from the integration measure $\sqrt{g(\boldsymbol{r})}\,d\boldsymbol{r}$, if not only locally. Even for the non interacting system, one would still need to adopt Metropolis algorithm in order to sample additionally the metric factor in the integration measure that can be seen as giving rise to an effective/external single particle potential $\ln \sqrt{g(\boldsymbol{r})}$. The transition probability distribution function of the Metropolis will be Gaussian.

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[10] A bridge between the same particle can still be used to sample the density matrix of distinguishable particles as can be done the displacement move of Appendix A.

Any permutation can be reached through a two particles exchange so the bridge transition move allows to sample the sum in Eq. (2.17).

Appendix C: Combination of bridges and diaplacements

In order to calculate each ρ_α in Eq. (2.19) one needs to combine bridges (described in Appendix B) and displacements (described in Appendix A). In fact one can first use bridge moves to reach the α homotopy class, with a fixed number n of braids, and then use only displacement moves to sample that particular ρ_α.

In order to sample the sum in Eq. (2.19) one needs to explore all possible swaps of two particles between any two timeslices i and j and at the same time count the number n of braids reached in the path $R(t)$ in order to be able to asses which path homotopy class he accessed after each move. This will allow him to determine to which ρ_α he is contributing at each accepted transition move.

If two different particles, say 1 and 2, have k crossings between timeslices i and $j > i$, an acceptance of the move described in Appendix B, where we build a bridge connecting $\boldsymbol{r}_{1,i}$ to $\boldsymbol{r}_{2,j}$ and one connecting $\boldsymbol{r}_{2,i}$ to $\boldsymbol{r}_{1,j}$, will necessarily result in just a single crossing if k is even, and consequently $n \to n - k + 1$, or to no crossings at all, if k is odd, and consequently $n \to n - k$. In order to count the number of crossings between timeslices i and j it is necessary to count the number of times in which $\boldsymbol{r}_{1,k} = \boldsymbol{r}_{2,k}$ for $k \in [i,j]$. In order to take into account of this crossing condition in the discretized imaginary time one can for example determine when both $r_{1,k}^\alpha - r_{2,k}^\alpha$ for $\alpha = 1, 2$ change sign varying k.

AUTHOR DECLARATIONS

Conflicts of interest

None declared.

Data availability

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