

Program FLUX

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1 Preliminary Analysis

1.1 Coordinate Systems

- Cartesian coordinates: X, Y, Z .
- Cylindrical coordinates: $R \equiv (X^2 + Y^2)^{1/2}$, $\phi \equiv \tan^{-1}(Y/X)$, Z , so

$$|\nabla\phi| = \frac{1}{R}. \quad (1)$$

- Flux coordinates: $r(R, Z)$, $\theta(R, Z)$, ϕ , where

$$(\nabla r \cdot \nabla\theta \times \nabla\phi)^{-1} = \mathcal{J}, \quad (2)$$

$$\mathcal{J} = \frac{r R^2}{R_0}, \quad (3)$$

where R_0 is a convenient scale major radius, flux-surface label r has units of length, and θ is a poloidal angle.

1.2 Useful Identities

Easily demonstrated that:

$$\mathbf{A} = A^r \mathcal{J} \nabla\theta \times \nabla\phi + A^\theta \mathcal{J} \nabla\phi \times \nabla r + A^\phi \mathcal{J} \nabla r \times \nabla\theta, \quad (4)$$

$$\mathbf{A} = A_r \nabla r + A_\theta \nabla\theta + A_\phi \nabla\phi, \quad (5)$$

$$\mathbf{A} \cdot \mathbf{B} = A_r B^r + A_\theta B^\theta + A_\phi B^\phi = A^r B_r + A^\theta B_\theta + A^\phi B_\phi, \quad (6)$$

$$(\mathbf{A} \times \mathbf{B})_r = \mathcal{J} (A^\theta B^\phi - A^\phi B^\theta), \quad (7)$$

$$(\mathbf{A} \times \mathbf{B})_\theta = \mathcal{J} (A^\phi B^r - A^r B^\phi), \quad (8)$$

$$(\mathbf{A} \times \mathbf{B})_\phi = \mathcal{J} (A^r B^\theta - A^\theta B^r), \quad (9)$$

$$\mathcal{J} (\mathbf{A} \times \mathbf{B})^r = A_\theta B_\phi - A_\phi B_\theta, \quad (10)$$

$$\mathcal{J}(\mathbf{A} \times \mathbf{B})^\theta = A_\phi B_r - A_r B_\phi, \quad (11)$$

$$\mathcal{J}(\mathbf{A} \times \mathbf{B})^\phi = A_r B_\theta - A_\theta B_r, \quad (12)$$

$$\mathcal{J} \nabla \cdot \mathbf{A} = \frac{\partial(\mathcal{J} A^r)}{\partial r} + \frac{\partial(\mathcal{J} A^\theta)}{\partial \theta} + \frac{\partial(\mathcal{J} A^\phi)}{\partial \phi}, \quad (13)$$

$$\mathcal{J}(\nabla \times \mathbf{A})^r = \frac{\partial A_\phi}{\partial \theta} - \frac{\partial A_\theta}{\partial \phi}, \quad (14)$$

$$\mathcal{J}(\nabla \times \mathbf{A})^\theta = \frac{\partial A_r}{\partial \phi} - \frac{\partial A_\phi}{\partial r}, \quad (15)$$

$$\mathcal{J}(\nabla \times \mathbf{A})^\phi = \frac{\partial A_\theta}{\partial r} - \frac{\partial A_r}{\partial \theta}. \quad (16)$$

1.3 Equilibrium Magnetic Field

Equilibrium magnetic field:

$$\mathbf{B} = B_0 R_0 [f(r) \nabla \phi \times \nabla r + g(r) \nabla \phi], \quad (17)$$

$$q(r) = \frac{r g}{R_0 f}, \quad (18)$$

where B_0 is a convenient scale magnetic field strength, and g and f are dimensionless. So,

$$B^r = 0, \quad (19)$$

$$B^\theta = B_0 R_0^2 \frac{f}{r R^2}, \quad (20)$$

$$B^\phi = B_0 R_0^2 \frac{q f}{r R^2} = B_0 B_0 \frac{g}{R^2}, \quad (21)$$

$$B_r = -B_0 r f \nabla r \cdot \nabla \theta, \quad (22)$$

$$B_\theta = B_0 r f |\nabla r|^2, \quad (23)$$

$$B_\phi = B_0 R_0^2 \frac{q f}{r} = B_0 R_0 g, \quad (24)$$

and

$$\mathbf{B} \cdot \nabla = B_0 R_0^2 \frac{f}{r R^2} \left(\frac{\partial}{\partial \theta} + q \frac{\partial}{\partial \phi} \right). \quad (25)$$

1.4 Equilibrium Plasma Current

The relation $\mu_0 \mathbf{J} = \nabla \times \mathbf{B}$ yields

$$\mu_0 \mathcal{J} J^r = 0, \quad (26)$$

$$\mu_0 \mathcal{J} J^\theta = -B_0 R_0 \frac{dg}{dr}, \quad (27)$$

$$\mu_0 \mathcal{J} J^\phi = B_0 \frac{\partial}{\partial r} (r f |\nabla r|^2) + B_0 \frac{\partial}{\partial \theta} (r f \nabla r \cdot \nabla \theta). \quad (28)$$

1.5 Grad-Shafranov Equation

The equilibrium force balance relation

$$\mathbf{J} \times \mathbf{B} = \nabla P \quad (29)$$

yields

$$\mathcal{J} (J^\theta B^\phi - J^\phi B^\theta) = \frac{dP}{dr}, \quad (30)$$

which reduces to the Grad-Shafranov equation,

$$\frac{f}{r} \frac{\partial}{\partial r} (r f |\nabla r|^2) + \frac{f}{r} \frac{\partial}{\partial \theta} (r f \nabla r \cdot \nabla \theta) + g \frac{dg}{dr} + \frac{\mu_0}{B_0^2} \left(\frac{R}{R_0} \right)^2 \frac{dP}{dr} = 0. \quad (31)$$

1.6 Inhomogeneous Tearing Mode Dispersion Relation

1.7 High- q Limit

Now,

$$\nabla \cdot \delta \mathbf{B} = 0 \quad (32)$$

yields

$$\frac{\partial(\mathcal{J} \delta B^r)}{\partial r} + \frac{\partial(\mathcal{J} \delta B^\theta)}{\partial \theta} + \frac{\partial(\mathcal{J} \delta B^\phi)}{\partial \phi} = 0. \quad (33)$$

Furthermore,

$$\delta \mathbf{B} = \delta B_r \nabla r + \delta B_\theta \nabla \theta + \delta B_\phi \nabla \phi. \quad (34)$$

So,

$$\delta B^r = \delta \mathbf{B} \cdot \nabla r = |\nabla r|^2 \delta B_r + (\nabla r \cdot \nabla \theta) \delta B_\theta, \quad (35)$$

$$\delta B^\theta = \delta \mathbf{B} \cdot \nabla \theta = (\nabla r \cdot \nabla \theta) \delta B_r + |\nabla \theta|^2 \delta B_\theta, \quad (36)$$

$$\delta B^\phi = \delta \mathbf{B} \cdot \nabla \phi = \frac{\delta B_\phi}{R^2}. \quad (37)$$

Follows that

$$\delta B_r = \left(\frac{1}{|\nabla r|^2} \right) \delta B^r - \left(\frac{\nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) \delta B_\theta, \quad (38)$$

and

$$\delta B^\theta = \left(\frac{\nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) \delta B^r + \left[|\nabla \theta|^2 - \frac{(\nabla r \cdot \nabla \theta)^2}{|\nabla r|^2} \right] \delta B_\theta. \quad (39)$$

But, Eqs. (2) and (3) imply that

$$|\nabla r|^2 |\nabla \theta|^2 - (\nabla r \cdot \nabla \theta)^2 = \frac{R_0^2}{r^2 R^2}. \quad (40)$$

Hence,

$$\delta B^\theta = \left(\frac{\nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) \delta B^r + \left(\frac{R_0^2}{r^2 R^2 |\nabla r|^2} \right) \delta B_\theta. \quad (41)$$

Assume that high- q limit is equivalent to $\delta \mathbf{B}$ being curl-free. Follows that

$$\frac{\partial \delta B_\phi}{\partial \theta} = \frac{\partial \delta B_\theta}{\partial \phi}, \quad (42)$$

$$\frac{\partial \delta B_r}{\partial \phi} = \frac{\partial \delta B_\phi}{\partial r}, \quad (43)$$

$$\frac{\partial \delta B_\theta}{\partial r} = \frac{\partial \delta B_r}{\partial \theta}. \quad (44)$$

Previous three equations imply that

$$\delta B_\phi \sim \frac{n}{m} \delta B_\theta, \quad \frac{n}{m} r \delta B_r \quad (45)$$

and, hence, that

$$\delta B^\phi \sim \frac{n}{m} \left(\frac{r}{R} \right)^2 \delta B^\theta, \quad \frac{n}{m} \left(\frac{r}{R} \right)^2 \frac{\delta B^r}{r}, \quad (46)$$

Consequently, final term on right-hand side of (33) is of order $(n/m)^2 (r/R)^2$ smaller than other two terms, and, therefore, negligible. Thus, we get

$$\frac{\partial(\mathcal{J} \delta B^r)}{\partial r} + \frac{\partial(\mathcal{J} \delta B^\theta)}{\partial \theta} = 0. \quad (47)$$

Follows from (41) and previous equation that

$$r \frac{\partial}{\partial r} \left(\frac{r R^2 \delta B^r}{R_0^2} \right) = - \frac{\partial}{\partial \theta} \left[\left(\frac{r \nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) \left(\frac{r R^2 \delta B^r}{R_0^2} \right) + \left(\frac{1}{|\nabla r|^2} \right) \delta B_\theta \right]. \quad (48)$$

Follows from (3), (38) and (44) that

$$r \frac{\partial \delta B_\theta}{\partial r} = \frac{\partial}{\partial \theta} \left[\left(\frac{R_0^2}{R^2 |\nabla r|^2} \right) \left(\frac{r R^2 \delta B^r}{R_0^2} \right) - \left(\frac{r \nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) \delta B_\theta \right]. \quad (49)$$

Let

$$\frac{r R^2 \delta B^r(r, \theta, \phi)}{R_0^2} = \sum_j \delta \hat{B}_j^r(r) e^{i(m_j \theta - n \phi)}, \quad (50)$$

$$\delta B_\theta(r, \theta, \phi) = \sum_j \delta \hat{B}_{\theta j}(r) e^{i(m_j \theta - n \phi)}. \quad (51)$$

Follows that

$$\sum_{j'} r \frac{d \delta \hat{B}_{j'}^r}{dr} e^{i m_{j'} \theta} = - \frac{\partial}{\partial \theta} \sum_{j'} \left[\left(\frac{r \nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) e^{i m_{j'} \theta} \delta \hat{B}_{j'}^r + \left(\frac{1}{|\nabla r|^2} \right) e^{i m_{j'} \theta} \delta \hat{B}_{\theta j'} \right], \quad (52)$$

$$\sum_{j'} r \frac{d \delta \hat{B}_{\theta j'}}{dr} e^{i m_{j'} \theta} = \frac{\partial}{\partial \theta} \sum_{j'} \left[\left(\frac{R_0^2}{R^2 |\nabla r|^2} \right) e^{i m_{j'} \theta} \delta \hat{B}_{j'}^r - \left(\frac{r \nabla r \cdot \nabla \theta}{|\nabla r|^2} \right) e^{i m_{j'} \theta} \delta \hat{B}_{\theta j'} \right]. \quad (53)$$

Operating with $\oint (\dots) e^{-i m_j \theta} d\theta / (2\pi)$, we get

$$r \frac{d \delta \hat{B}_j^r}{dr} = -i m_j \sum_{j'} \left(-i c_{jj'} \delta \hat{B}_{j'}^r + a_{jj'} \delta \hat{B}_{\theta j'} \right), \quad (54)$$

$$r \frac{d \delta \hat{B}_{\theta j}}{dr} = -i m_j \sum_{j'} \left(-i c_{jj'} \delta \hat{B}_{\theta j'} - b_{jj'} \delta \hat{B}_{j'}^r \right), \quad (55)$$

where

$$a_{jj'} = \oint \frac{1}{|\nabla r|^2} e^{-i(m_j - m_{j'})\theta} \frac{d\theta}{2\pi}, \quad (56)$$

$$b_{jj'} = \oint \frac{R_0^2}{R^2 |\nabla r|^2} e^{-i(m_j - m_{j'})\theta} \frac{d\theta}{2\pi}, \quad (57)$$

$$c_{jj'} = \oint \frac{i r \nabla r \cdot \nabla \theta}{|\nabla r|^2} e^{-i(m_j - m_{j'})\theta} \frac{d\theta}{2\pi}. \quad (58)$$

Let

$$\delta \hat{B}_j^r(r) = i \psi_j(r), \quad (59)$$

$$\delta \hat{B}_{\theta j}(r) = -\chi_j(r). \quad (60)$$

It follows that

$$r \frac{d \psi_j}{dr} = m_j \sum_{j'} (-c_{jj'} \psi_{j'} + a_{jj'} \chi_{j'}), \quad (61)$$

$$r \frac{d \chi_j}{dr} = m_j \sum_{j'} (-c_{jj'} \chi_{j'} + b_{jj'} \psi_{j'}). \quad (62)$$

Hence,

$$\frac{r R^2 \delta B^r(r, \theta, \phi)}{R_0^2} = i \sum_j \psi_j(r) e^{i(m_j \theta - n \phi)}, \quad (63)$$

$$\frac{r R^2 \delta B^\theta(r, \theta, \phi)}{R_0^2} = - \sum_j \frac{1}{m_j} \frac{d \psi_j}{dr} e^{i(m_j \theta - n \phi)}, \quad (64)$$

$$R^2 \delta B^\phi(r, \theta, \phi) = n \sum_j \frac{\chi_j(r)}{m_j} e^{i(m_j \theta - n \phi)}. \quad (65)$$

where use has been made of (42). Also have

$$\delta B_r(r, \theta, \phi) = i \sum_j \frac{1}{m_j} \frac{d\chi_j}{dr} e^{i(m_j \theta - n \phi)}, \quad (66)$$

$$\delta B_\theta(r, \theta, \phi) = - \sum_j \chi_j(r) e^{i(m_j \theta - n \phi)}, \quad (67)$$

$$\delta B_\phi(r, \theta, \phi) = n \sum_j \frac{\chi_j(r)}{m_j} e^{i(m_j \theta - n \phi)}, \quad (68)$$

We can write

$$\mathcal{J} \mu_0 \delta J^r = \frac{\partial \delta B_\phi}{\partial \theta} - \frac{\partial \delta B_\theta}{\partial \phi}, \quad (69)$$

$$\mathcal{J} \mu_0 \delta J^\theta = \frac{\partial \delta B_r}{\partial \phi} - \frac{\partial \delta B_\phi}{\partial r}, \quad (70)$$

$$\mathcal{J} \mu_0 \delta J^\phi = \frac{\partial \delta B_\theta}{\partial r} - \frac{\partial \delta B_r}{\partial \theta}. \quad (71)$$

Normally, all three components are zero. However, at k th resonant surface (at which $r = r_k$, where $q(r_k) = m_k/n$) ψ_k , $\psi_{j \neq k}$, and $\chi_{j \neq k}$ are continuous, whereas χ_k is discontinuous. Hence,

$$\mathcal{J} \mu_0 \delta J^r(r, \theta, \phi) = 0, \quad (72)$$

$$\mathcal{J} \mu_0 \delta J^\theta(r, \theta, \phi) = - \sum_k \frac{n}{m_k} [\chi_k]_{r_{k-}}^{r_{k+}} \delta(r - r_k) e^{i(m_k \theta - n \phi)}, \quad (73)$$

$$\mathcal{J} \mu_0 \delta J^\phi(r, \theta, \phi) = - \sum_k [\chi_k]_{r_{k-}}^{r_{k+}} \delta(r - r_k) e^{i(m_k \theta - n \phi)}. \quad (74)$$

Now,

$$(\mu_0 \delta \mathbf{J} \times \delta \mathbf{B})_\theta = \frac{1}{4} (\mathcal{J} \mu_0 \delta J^\phi \delta B^{r*} + \text{c.c.}), \quad (75)$$

$$(\mu_0 \delta \mathbf{J} \times \delta \mathbf{B})_\phi = \frac{1}{4} (-\mathcal{J} \mu_0 \delta J^\theta \delta B^{r*} + \text{c.c.}), \quad (76)$$

which implies that

$$\mathcal{J} (\mu_0 \delta \mathbf{J} \times \delta \mathbf{B})_\theta = \frac{R_0}{2} \sum_j \sum_k \text{Re} \left\{ i [\chi_k]_{r_{k-}}^{r_{k+}} \psi_j^*(r_k) e^{i(m_k - m_j) \theta} \right\} \delta(r - r_k), \quad (77)$$

$$\mathcal{J} (\mu_0 \delta \mathbf{J} \times \delta \mathbf{B})_\phi = -\frac{R_0}{2} \sum_j \sum_k \frac{n}{m_k} \text{Re} \left\{ i [\chi_k]_{r_{k-}}^{r_{k+}} \psi_j^*(r_k) e^{i(m_k - m_j) \theta} \right\} \delta(r - r_k). \quad (78)$$

Let

$$\Psi_k = \frac{\psi_k(r_k)}{m_k}, \quad (79)$$

$$\Delta\Psi_k = [\chi_k]_{r_{k-}}^{r_{k+}}, \quad (80)$$

$$\delta T_{\theta k} \equiv \delta \mathbf{T} \cdot \mathbf{e}_\theta = \int_{r_{k-}}^{r_{k+}} \oint \oint (\delta \mathbf{J} \times \delta \mathbf{B})_\theta \mathcal{J} dr d\theta d\phi, \quad (81)$$

$$\delta T_{\phi k} \equiv \delta \mathbf{T} \cdot \mathbf{e}_\phi = \int_{r_{k-}}^{r_{k+}} \oint \oint (\delta \mathbf{J} \times \delta \mathbf{B})_\phi \mathcal{J} dr d\theta d\phi, \quad (82)$$

where $\delta \mathbf{T}(r) dr$ is the net electromagnetic torque acting on the plasma between r and $r + dr$. Here, $\mathbf{e}_\theta = (R^2/R_0) \nabla \phi \times \nabla r$ and $\mathbf{e}_\phi = R \nabla \phi$. Follows that

$$\delta T_{\theta k} = -\frac{2\pi^2 R_0 m_k}{\mu_0} \text{Im}(\Psi_k^* \Delta\Psi_k), \quad (83)$$

$$\delta T_{\phi k} = \frac{2\pi^2 R_0 n}{\mu_0} \text{Im}(\Psi_k^* \Delta\Psi_k). \quad (84)$$

1.8 Homogeneous Solution

We can write

$$\delta \mathbf{B} = \nabla \times \delta \mathbf{A}. \quad (85)$$

Suppose that $r \delta A_r$, δA_θ are negligible with respect to δA_ϕ . In fact, it is easily demonstrated from $\nabla \cdot \delta \mathbf{A} = 0$ that $r \delta A_r$, $\delta A_\theta \sim (n/m)(r/R)^2 \delta A_\phi$. It follows that

$$\mathcal{J} \delta B^r \simeq \frac{\partial \delta A_\phi}{\partial \theta}, \quad (86)$$

$$\mathcal{J} \delta B^\theta \simeq -\frac{\partial \delta A_\phi}{\partial r}, \quad (87)$$

$$\mathcal{J} \delta B^\phi \simeq 0. \quad (88)$$

The neglected terms in the previous three equations are $(n/m)^2 (r/R)^2$ smaller than the dominant terms. The previous three expressions are consistent with (63), (64), and (65) provided

$$\delta A_\phi(r, \theta, \phi) \simeq R_0 \sum_j \frac{\psi_j(r)}{m_j} e^{i(m_j \theta - n \phi)}. \quad (89)$$

According to the Biot-Savart law:

$$\delta A_\phi(\mathbf{x}) = \frac{1}{4\pi} \int \frac{R R' \mu_0 \delta \mathbf{J}(\mathbf{x}') \cdot \nabla \phi}{|\mathbf{x} - \mathbf{x}'|} d^3 \mathbf{x}'. \quad (90)$$

Assume that

$$\delta A_\phi(R, \phi, Z) = \delta A_\phi(R, 0, Z) e^{-i n \phi}. \quad (91)$$

Hence, can evaluate integral at $\phi = 0$ without loss of generality. Follows that

$$\mathbf{x} = (R, 0, Z), \quad (92)$$

$$\mathbf{x}' = (R' \cos \phi', R' \sin \phi', Z'), \quad (93)$$

and

$$|\mathbf{x} - \mathbf{x}'| = [R^2 + R'^2 + (Z - Z')^2 - 2 R R' \cos \phi']^{1/2}. \quad (94)$$

Now,

$$\delta \mathbf{J}(R', \phi', Z') \cdot \nabla \phi = \delta J^\phi(R', 0, Z') e^{-i n \phi'} \cos \phi', \quad (95)$$

so

$$\delta A_\phi(R, 0, Z) = \frac{1}{4\pi} \int_0^\infty \oint \oint \frac{R R' \mu_0 \delta J^\phi(R', 0, Z') e^{-i n \phi'} \cos \phi' \mathcal{J}' dr' d\theta' d\phi'}{[R^2 + R'^2 + (Z - Z')^2 - 2 R R' \cos \phi']^{1/2}}, \quad (96)$$

which can be written

$$\delta A_\phi(R, 0, Z) = \frac{1}{4\pi} \int_0^\infty \oint R R' \mu_0 \delta J^\phi(R', 0, Z') G(R, Z; R', Z') \mathcal{J}' dr' d\theta', \quad (97)$$

where

$$G(R, Z; R', Z') = \oint \frac{\cos \phi' \cos(n \phi') d\phi'}{[R^2 + R'^2 + (Z - Z')^2 - 2 R R' \cos \phi']^{1/2}}, \quad (98)$$

or

$$G(R, Z; R', Z') = \frac{1}{2} \oint \frac{(\cos[(n-1)\phi'] + \cos[(n+1)\phi']) d\phi'}{[R^2 + R'^2 + (Z - Z')^2 - 2 R R' \cos \phi']^{1/2}}. \quad (99)$$

Now,

$$\begin{aligned} P_{-1/2}^n(\cosh \eta) &= \frac{(-1)^n}{2\pi} \frac{\Gamma(1/2)}{\Gamma(1/2 - n)} \oint \frac{\cos(n \varphi) d\varphi}{(\cosh \eta + \sinh \eta \cos \varphi)^{1/2}} \\ &= \frac{1}{2\pi} \frac{\Gamma(1/2)}{\Gamma(1/2 - n)} \oint \frac{\cos(n \varphi) d\varphi}{(\cosh \eta - \sinh \eta \cos \varphi)^{1/2}} \\ &= \frac{(-1)^n \Gamma(1/2) \Gamma(1/2 + n)}{2\pi^2} \oint \frac{\cos(n \varphi) d\varphi}{(\cosh \eta - \sinh \eta \cos \varphi)^{1/2}}. \end{aligned} \quad (100)$$

Let

$$\alpha \cosh \eta = R^2 + R'^2 + (Z - Z')^2, \quad (101)$$

$$\alpha \sinh \eta = 2 R R'. \quad (102)$$

Hence,

$$\eta = \tanh^{-1} \left[\frac{2 R R'}{R^2 + R'^2 + (Z - Z')^2} \right], \quad (103)$$

and

$$\alpha = \frac{R^2 + R'^2 + (Z - Z')^2}{\cosh \eta}. \quad (104)$$

It follows that

$$G(R, Z; R', Z') = \pi \left[\frac{\cosh \eta}{R^2 + R'^2 + (Z - Z')^2} \right]^{1/2} \times \left[\frac{\Gamma(3/2 - n)}{\Gamma(1/2)} P_{-1/2}^{n-1}(\cosh \eta) + \frac{\Gamma(-1/2 - n)}{\Gamma(1/2)} P_{-1/2}^{n+1}(\cosh \eta) \right]. \quad (105)$$

However,

$$\Gamma(3/2 - n) = \frac{(-1)^{n+1} \pi (n - 1/2)}{\Gamma(n + 1/2)}, \quad (106)$$

$$\Gamma(-1/2 - n) = \frac{(-1)^{n+1} \pi}{\Gamma(n + 1/2) (n + 1/2)}, \quad (107)$$

so

$$G(R, Z; R', Z') = \frac{(-1)^{n+1} \pi^2}{\Gamma(1/2) \Gamma(n + 1/2)} \left[\frac{\cosh \eta}{R^2 + R'^2 + (Z - Z')^2} \right]^{1/2} \times \left[(n - 1/2) P_{-1/2}^{n-1}(\cosh \eta) + \frac{P_{-1/2}^{n+1}(\cosh \eta)}{n + 1/2} \right]. \quad (108)$$

Now, from (74) and (80),

$$\mathcal{J} \mu_0 \delta J^\phi(r, \theta, 0) = - \sum_k \Delta \Psi_k \delta(r - r_k) e^{i m_k \theta}. \quad (109)$$

Furthermore, from (79) and (89),

$$\Psi_k = \frac{1}{R_0} \oint \delta A_\phi(r_k, \theta, 0) e^{-i m_k \theta} \frac{d\theta}{2\pi}. \quad (110)$$

Hence, we obtain the homogeneous tearing mode dispersion relation,

$$\Psi_k = \sum_{k'} F_{kk'} \Delta \Psi_{k'}, \quad (111)$$

where

$$F_{kk'} = \oint \oint \mathcal{G}(R_k, Z_k; R'_k, Z'_k) e^{-i(m_k \theta - m_{k'} \theta')} \frac{d\theta}{2\pi} \frac{d\theta'}{2\pi}, \quad (112)$$

and

$$\mathcal{G}(R, Z; R', Z') = \frac{(-1)^n \pi^2 R R' / R_0}{2 \Gamma(1/2) \Gamma(n + 1/2)} \left[\frac{\cosh \eta}{R^2 + R'^2 + (Z - Z')^2} \right]^{1/2} \times \left[(n - 1/2) P_{-1/2}^{n-1}(\cosh \eta) + \frac{P_{-1/2}^{n+1}(\cosh \eta)}{n + 1/2} \right]. \quad (113)$$

Here, R_k, Z_k are the R, Z coordinates of the k th resonant surface in the plane $\phi = 0$. Note that

$$\mathcal{G}(R', Z'; R, Z) = \mathcal{G}(R, Z; R', Z'), \quad (114)$$

which implies that

$$F_{k'k} = F_{kk'}^*. \quad (115)$$

1.9 Inhomogeneous Solution

Suppose that there are currents flowing in a number of external poloidal field coils. Let I_l, R_l , and Z_l be the peak current, and coordinates of the l th field coil. The currents are assumed to modulate like $e^{-in\phi}$. It follows that

$$\delta J_{\text{ext}}^\phi(R, 0, Z) = \sum_j \frac{I_l}{R_l} \delta(R - R_l) \delta(Z - Z_l). \quad (116)$$

Hence,

$$\Psi_k = \sum_{k'} F_{kk'} \Delta \Psi_k - \sum_l g_{kl} \mu_0 I_l, \quad (117)$$

where

$$g_{kj} = \frac{1}{2\pi} \oint \mathcal{G}(r_k, \theta; R_j, Z_j) e^{-im_k \theta} \frac{d\theta}{2\pi}. \quad (118)$$

Thus, we obtain the inhomogeneous tearing mode dispersion relation,

$$\Delta \Psi_k = \sum_{k'} E_{kk'} \Psi_{k'} + |E_{kk}| \chi_k, \quad (119)$$

where $E_{kk'}$ is the inverse of the $F_{kk'}$ matrix, and

$$\chi_k = \frac{1}{|E_{kk}|} \sum_j h_{kl} \mu_0 I_l, \quad (120)$$

$$h_{kl} = \sum_{k'} E_{kk'} g_{k'l}. \quad (121)$$

1.10 Perturbed Poloidal Magnetic Field

The perturbed poloidal magnetic field generated by the current sheets inside the plasma and the currents flowing in the external field coils can be expressed as

$$\delta \mathbf{B}_p = \nabla \delta A_\phi \times \nabla \phi, \quad (122)$$

where

$$\frac{\delta A_\phi(R, \phi, Z)}{R_0} = \text{Re} \left[\sum_{k=1, K} e_k(R, Z) \Delta \Psi_k e^{-in\phi} - \sum_{l=1, L} f_l(R, Z) \mu_0 I_l e^{-in\phi} \right], \quad (123)$$

and

$$e_k(R, Z) = \oint \mathcal{G}(R, Z; R_k; Z_k) e^{im_k \theta_k} \frac{d\theta_k}{2\pi}, \quad (124)$$

$$f_l(R, Z) = \frac{1}{2\pi} \mathcal{G}(R, Z; R_l, Z_l). \quad (125)$$

As before, k indexes the resonant surfaces within the plasma, whereas l indexes the external current filaments. Moreover, the integral in Eq. (124) is taken around the k th resonant surface.

1.11 Electromagnetic Torques in Presence of External Currents

Suppose that

$$\Psi_k = B_0 R_0 \hat{\Psi}_k e^{-i\varphi_k}, \quad (126)$$

$$\chi_k = B_0 R_0 \hat{\chi}_k e^{-i\zeta_k}, \quad (127)$$

$$E_{kk'} = \hat{E}_{kk'} e^{-i\xi_{kk'}}, \quad (128)$$

where $\hat{\Psi}_k$, $\hat{\chi}_k$, and $\hat{E}_{kk'}$ are real and positive, whereas φ_k , ζ_k , and $\xi_{kk'}$ are real. (Note that all hatted quantities in this report are dimensionless.) It follows from Eqs. (83), (84), and (119) that

$$\delta T_{\theta k} = - \left(\frac{2\pi^2 B_0^2 R_0^3}{\mu_0} \right) m_k \delta \hat{T}_k, \quad (129)$$

$$\delta T_{\phi k} = \left(\frac{2\pi^2 B_0^2 R_0^3}{\mu_0} \right) n \delta \hat{T}_k, \quad (130)$$

where

$$\delta \hat{T}_k = \sum_{k'=1, K} \hat{E}_{kk'} \hat{\Psi}_k \hat{\Psi}_{k'} \sin(\varphi_k - \varphi_{k'} - \xi_{kk'}) + \hat{E}_{kk} \hat{\Psi}_k \hat{\chi}_k \sin(\varphi_k - \zeta_k). \quad (131)$$

1.12 GPEC Coupling

1.12.1 PRL Derivation

We have

$$\frac{d\psi_p}{dr} = B_0 R_0 f(r), \quad (132)$$

$$\mathbf{B} \cdot \nabla \phi = \frac{B_0 R_0}{R^2} g(r), \quad (133)$$

$$\delta B^r = \frac{i}{r} \left(\frac{R_0}{R} \right)^2 \sum_j \psi_j e^{i(m_j \theta - n \phi)}. \quad (134)$$

According to PRL **99**, 195003 (2007),

$$\Delta_j e^{i(m_j \theta - n \phi)} = \left[\frac{\partial}{\partial \psi_p} \frac{\delta \mathbf{B} \cdot \nabla \psi_p}{\mathbf{B} \cdot \nabla \phi} \right]_{r_j} = \left[\frac{\partial}{\partial r} \frac{\delta \mathbf{B} \cdot \nabla r}{\mathbf{B} \cdot \nabla \phi} \right]_{r_j} = \left[\frac{\partial}{\partial r} \frac{\delta B^r}{\mathbf{B} \cdot \nabla \phi} \right]_{r_j}. \quad (135)$$

It follows that

$$\Delta_j = \frac{i}{r_j} \left(\frac{R_0}{R} \right)^2 \frac{R^2}{B_0 R_0 g_j} \left[\frac{d\psi_j}{dr} \right]_{r_j}, \quad (136)$$

But,

$$\left[r \frac{d\psi_j}{dr} \right]_{r_j} = m_j a_{jj} [\chi_j]_{r_j} = m_j a_{jj} \Delta \Psi_j, \quad (137)$$

and

$$\chi_j = \frac{\Delta \Psi_j}{|E_{jj}|}, \quad (138)$$

which implies that

$$\Delta_j = i \left(\frac{R_0}{r_j} \right)^2 \frac{m_j a_{jj}}{g_j} \frac{\Delta \Psi_j}{B_0 R_0}, \quad (139)$$

or

$$\frac{\chi_j}{R_0 B_0} = -i \left(\frac{r_j}{R_0} \right)^2 \frac{g_j}{m_j a_{jj}} \frac{\Delta_j}{|E_{jj}|}. \quad (140)$$

Here, the Δ_j values can be determined from the GPEC code.

1.12.2 PoP Derivation

Now, $d\psi_p/dr = B_0 R_0 f$. It follows that

$$(\nabla \psi_p \times \nabla \theta \cdot \nabla \phi)^{-1} = \frac{R^2 q}{B_0 R_0 g}. \quad (141)$$

Hence,

$$\mathbf{B} = \nabla \phi \times \nabla \psi_p + B_0 R_0 g \nabla \phi = \nabla \phi \times \nabla \psi_p + q \nabla \psi_p \times \nabla \theta. \quad (142)$$

Let $\Psi_p = 2\pi \psi_p$ and $d\Psi_p/d\Psi_t = 1/q$. Hence,

$$\frac{d\Psi_p}{dr} = 2\pi B_0 R_0 f, \quad (143)$$

$$\frac{d\Psi_t}{dr} = 2\pi B_0 r g, \quad (144)$$

and [cf. PoP **13**, 102501 (2006), Eq. (41)]

$$2\pi \mathbf{B} = q^{-1} \nabla \phi \times \nabla \Psi_t + \nabla \Psi_t \times \nabla \theta. \quad (145)$$

We have

$$\delta J^r = 0, \quad (146)$$

$$\mathcal{J} \mu_0 \delta J^\theta = - \sum_j \frac{\Delta \Psi_j}{q_j} e^{i(m_j \theta - n \phi)} \delta(r - r_j), \quad (147)$$

$$\mathcal{J} \mu_0 \delta J^\phi = - \sum_j \Delta \Psi_k e^{i(m_j \theta - n \phi)} \delta(r - r_j), \quad (148)$$

which implies that

$$\mu_0 \delta \mathbf{J} = -2\pi \sum_j \Delta \Psi_j e^{i(m_j \theta - n \phi)} \delta(\psi_t - \psi_{tj}) \mathbf{B}. \quad (149)$$

However, according to PoP **13**, 102501 (2006),

$$\Delta \Psi_j = -i \frac{\mu_0 J_c \Delta_j}{2\pi m_j}, \quad (150)$$

where

$$\frac{1}{\mu_0 J_c} = \left(\oint \frac{B^2}{|\nabla \psi_t|^2} \frac{d\theta d\phi}{2\pi \mathbf{B} \cdot \nabla \phi} \right)_{r_j}. \quad (151)$$

It is easily demonstrated that

$$\frac{1}{\mu_0 J_c} = \left(\frac{R_0}{r_j} \right)^2 \frac{1}{2\pi B_0 R_0 g_j} \left[a_{jj} + \left(\frac{r_k}{R_0 q_j} \right)^2 \right]. \quad (152)$$

Hence,

$$\frac{\Delta \Psi_j}{B_0 R_0} = -i \Delta_j \left(\frac{r_j}{R_0} \right)^2 \frac{g_j}{m_j [a_{jj} + (r_k/R_0 q_j)^2]}, \quad (153)$$

which implies that

$$\frac{\chi_j}{B_0 R_0} = -i \frac{\Delta_j}{|E_{jj}|} \left(\frac{r_j}{R_0} \right)^2 \frac{g_j}{m_j [a_{jj} + (r_k/R_0 q_j)^2]}. \quad (154)$$

(Note: This is what is actually implemented in EPEC.) As before, the Δ_j values can be determined from the GPEC code.

1.13 Island Width (in r)

We have

$$\mathcal{J} \delta B^r \simeq \frac{\partial \delta A_\phi}{\partial \theta}, \quad (155)$$

$$\mathcal{J} \delta B^\theta \simeq -\frac{\partial \delta A_\phi}{\partial r}, \quad (156)$$

$$\mathcal{J} \delta B^\phi \simeq 0. \quad (157)$$

It follows that

$$\delta \mathbf{B} \cdot \nabla \delta A_\phi = \delta B^r \frac{\partial \delta A_\phi}{\partial r} + \delta B^\theta \frac{\partial \delta A_\phi}{\partial \theta} + \delta B^\phi \frac{\partial \delta A_\phi}{\partial \phi} \simeq 0. \quad (158)$$

Suppose that

$$\delta A_\phi(r, \theta, \phi) \simeq \delta \hat{A}_\phi e^{i(m_k \theta - n \phi)} \quad (159)$$

in the vicinity of the k th resonant surface. It follows that

$$\mathbf{B} \cdot \nabla \delta A_\phi = i B_0 R_0^2 \frac{f}{r R^2} (m_k - n q) \delta \hat{A}_\phi e^{i(m_k \theta - n \phi)}. \quad (160)$$

Let us search for a function,

$$F(r, \theta, \phi) = F_0(r) + \delta A_\phi, \quad (161)$$

which is such that

$$(\mathbf{B} + \delta \mathbf{B}) \cdot \nabla F = 0. \quad (162)$$

It follows that

$$\delta B^r \frac{dF_0}{dr} + i B_0 R_0^2 \frac{f}{r R^2} (m_k - n q) \delta \hat{A}_\phi e^{i(m_k \theta - n \phi)} = 0, \quad (163)$$

or

$$i m_k \frac{R_0}{r R^2} \delta \hat{A}_\phi \frac{dF_0}{dr} + i B_0 R_0^2 \frac{f}{r R^2} (m_k - n q) \delta \hat{A}_\phi = 0, \quad (164)$$

which implies that

$$\frac{dF_0}{dr} = -\frac{B_0 R_0 f}{m_k} (m_k - n q), \quad (165)$$

or

$$F_0(r) \simeq \frac{B_0}{2} \left(\frac{g s}{q} \right)_{r_k} (r - r_k)^2, \quad (166)$$

where $s = r q' / q$. Hence,

$$F(r, \theta, \phi) = \frac{B_0}{2} \left(\frac{g s}{q} \right)_{r_k} (r - r_k)^2 + R_0 \Psi_k \cos(m_k \theta - n \phi) \quad (167)$$

is a flux surface function in the island region. Thus,

$$\frac{F}{R_0 |\Psi_k|} = 2 X^2 + \cos(m_k \theta - n \phi), \quad (168)$$

where

$$X = \frac{2(r - r_k)}{W_k}, \quad (169)$$

and

$$\frac{W_k}{4 R_0} = \left[\left(\frac{q}{g s} \right)_{r_k} \frac{|\Psi_k|}{B_0 R_0} \right]^{1/2}. \quad (170)$$

It follows that W_k is the full radial island width in r . Moreover, W_k has no dependence on θ .

2 Technical Details

2.1 Flux Coordinate System

Let all lengths be normalized to R_0 , and all magnetic field-strengths to B_0 . We have

$$\mathbf{B} = \nabla\phi \times \nabla\psi_p + g(\psi_p) \nabla\phi, \quad (171)$$

and

$$\nabla\psi_p \times \nabla\theta \cdot \nabla\phi = \frac{g}{q R^2}, \quad (172)$$

where $q = q(\psi_p)$.

Let $\Psi = \psi_p/\psi_c = 1 - \Psi_N$, where ψ_c is the value of ψ_p on the magnetic axis. (It is assumed that $\psi_p = 0$ on the plasma boundary.) The previous equation implies that

$$\frac{d\theta}{dl} = \frac{g}{q |\psi_c| R \sqrt{\Psi_R^2 + \Psi_Z^2}}, \quad (173)$$

where dl is an element of poloidal path length on a magnetic flux-surface, and $\Psi_R \equiv \partial\Psi/\partial R$, etc. Furthermore,

$$dR = -\frac{\Psi_Z dl}{\sqrt{\Psi_R^2 + \Psi_Z^2}}, \quad (174)$$

$$dZ = \frac{\Psi_R dl}{\sqrt{\Psi_R^2 + \Psi_Z^2}}. \quad (175)$$

It follows that

$$\frac{q(\Psi)}{g(\Psi)} = \frac{1}{2\pi |\psi_c|} \oint \frac{dl}{R \sqrt{\Psi_R^2 + \Psi_Z^2}}. \quad (176)$$

If we define

$$\tan \zeta = \frac{Z - Z_{\text{axis}}}{R_{\text{axis}} - R} \quad (177)$$

then

$$\frac{dR}{d\zeta} = -\Psi_Z F, \quad (178)$$

$$\frac{dZ}{d\zeta} = \Psi_R F, \quad (179)$$

$$\frac{q(\Psi)}{g(\Psi)} = \frac{1}{2\pi |\psi_c|} \oint \frac{F}{R} d\zeta, \quad (180)$$

where

$$F = \frac{(R_{\text{axis}} - R)^2 + (Z - Z_{\text{axis}})^2}{-(Z - Z_{\text{axis}}) \Psi_Z + (R_{\text{axis}} - R) \Psi_R}. \quad (181)$$

It is helpful to define the length-like flux-surface coordinate r , according to

$$\nabla r \times \nabla \theta \cdot \nabla \phi = \frac{1}{r R^2}. \quad (182)$$

It follows that

$$r(\Psi) = \left[2 |\psi_c| \int_{\Psi}^1 \frac{q(\Psi')}{g(\Psi')} d\Psi' \right]^{1/2}. \quad (183)$$

Let

$$a = r(1). \quad (184)$$

We can calculate $R(r, \theta)$ and $Z(r, \theta)$ from

$$\frac{dR}{d\theta} = -|\psi_c| \frac{q}{g} R \Psi_Z, \quad (185)$$

$$\frac{dZ}{d\theta} = |\psi_c| \frac{q}{g} R \Psi_R. \quad (186)$$

Now,

$$r \frac{dr}{d\Psi} = -|\psi_c| \frac{q(r)}{g(r)}. \quad (187)$$

So

$$\nabla r = \frac{dr}{d\Psi} \nabla \Psi = -|\psi_c| \frac{q(r)}{r g(r)} \nabla \Psi. \quad (188)$$

Hence,

$$a_{jj} = \left(\oint \frac{1}{|\nabla r|^2} \frac{d\theta}{2\pi} \right)_{r_j} = \left(\frac{r g}{|\psi_c| q} \right)_{r_j}^2 \oint \frac{1}{\Psi_R^2 + \Psi_Z^2} \frac{d\theta}{2\pi}. \quad (189)$$

Note that

$$\frac{d\Psi_N}{dr} = \frac{r g(r)}{|\psi_c| q(r)}. \quad (190)$$

Hence, if \overline{W}_k is the full magnetic island width in Ψ_N at the k th resonant surface then

$$\overline{W}_k = \frac{W_k}{R_0} \frac{d\Psi_N(r_k)}{dr}, \quad (191)$$

where W_k is the full island width in r .

2.2 Neoclassical Coordinate System

It is also helpful to define the geometric poloidal angle

$$\mathbf{b} \cdot \nabla \Theta = \gamma(r). \quad (192)$$

It follows that

$$\frac{d\Theta}{dl} = \frac{\gamma B R}{|\psi_c| \sqrt{\Psi_R^2 + \Psi_Z^2}}, \quad (193)$$

where

$$B R = [g^2 + |\psi_c|^2 (\Psi_R^2 + \Psi_Z^2)]^{1/2}. \quad (194)$$

Hence,

$$\frac{1}{\gamma(r)} = \frac{1}{2\pi |\psi_c|} \oint \frac{B R dl}{\sqrt{\Psi_R^2 + \Psi_Z^2}} = \frac{1}{2\pi |\psi_c|} \oint B R F d\zeta. \quad (195)$$

We can calculate $R(r, \Theta)$ and $Z(r, \Theta)$ from

$$\frac{dR}{d\Theta} = -|\psi_c| \frac{\Psi_Z}{\gamma B R}, \quad (196)$$

$$\frac{dZ}{d\Theta} = |\psi_c| \frac{\Psi_R}{\gamma B R}. \quad (197)$$

Note that

$$\frac{d\Theta}{d\theta} = \left(\frac{\gamma q}{g} \right) B R^2. \quad (198)$$

Thus,

$$\frac{1}{\gamma} = \frac{q}{g} \oint B R^2 \frac{d\theta}{2\pi}. \quad (199)$$

2.3 Neoclassical Parameters

The flux-surface average operator has the following properties:

$$\langle 1 \rangle = 1, \quad (200)$$

$$\langle \mathbf{B} \cdot \nabla A \rangle = 0. \quad (201)$$

It follows that

$$\langle A \rangle = \oint R^2 A \frac{d\theta}{2\pi} \bigg/ \oint R^2 \frac{d\theta}{2\pi} = \oint \frac{A}{B} \frac{d\Theta}{2\pi} \bigg/ \oint \frac{1}{B} \frac{d\Theta}{2\pi}. \quad (202)$$

Let

$$I_0 = \oint \frac{1}{B R^2} \frac{d\Theta}{2\pi} = \frac{\gamma q}{g}, \quad (203)$$

$$I_1 = \oint \frac{1}{B} \frac{d\Theta}{2\pi}, \quad (204)$$

$$I_2 = \oint B \frac{d\Theta}{2\pi}, \quad (205)$$

$$I_3 = \oint \left(\frac{\partial B}{\partial \Theta} \right)^2 \frac{1}{B} \frac{d\Theta}{2\pi}, \quad (206)$$

$$I_{4,k} = \sqrt{\frac{2}{k}} \oint \frac{\sin(k\Theta)}{B^2} \frac{\partial B}{\partial \Theta} \frac{d\Theta}{2\pi} = \oint \frac{\sqrt{2k} \cos(k\Theta)}{B} \frac{d\Theta}{2\pi}, \quad (207)$$

$$I_{5,k} = \sqrt{\frac{2}{k}} \oint \frac{\sin(k\Theta)}{B^3} \frac{\partial B}{\partial \Theta} \frac{d\Theta}{2\pi} = \oint \frac{\sqrt{2k} \cos(k\Theta)}{2B^2} \frac{d\Theta}{2\pi}, \quad (208)$$

$$I_6(\lambda) = \oint \frac{\sqrt{1 - \lambda B/B_{\max}}}{B} \frac{d\Theta}{2\pi}. \quad (209)$$

It follows that

$$\langle B \rangle = \frac{1}{I_1}, \quad (210)$$

$$\left\langle \frac{1}{R^2} \right\rangle = \frac{\gamma q}{I_1 g}, \quad (211)$$

$$\langle B^2 \rangle = \frac{I_2}{I_1}, \quad (212)$$

$$\langle (\mathbf{b} \cdot \nabla B)^2 \rangle = \gamma^2 \frac{I_3}{I_1}, \quad (213)$$

$$|\langle \mathbf{B} \cdot \nabla \theta \rangle| = \frac{g}{|q|} \frac{I_0}{I_1} = \frac{|\gamma|}{I_1}, \quad (214)$$

$$\left\langle \sqrt{\frac{2}{k}} \sin(k\Theta) (\mathbf{b} \cdot \nabla \ln B) \right\rangle = \gamma \frac{I_{4,k}}{I_1}, \quad (215)$$

$$\left\langle \sqrt{\frac{2}{k}} \sin(k\Theta) \frac{(\mathbf{b} \cdot \nabla \ln B)}{B} \right\rangle = \gamma \frac{I_{5,k}}{I_1}. \quad (216)$$

Hence,

$$L_c = \frac{1}{|\gamma|} \frac{I_2^2}{I_1^2 I_3} \sum_{k>0} I_{5,k} I_{6,k}, \quad (217)$$

$$\omega_{ta} \equiv \frac{v_{Ta}}{L_c} = K_t |\gamma| v_{Ta}, \quad (218)$$

$$\nu_{*a} \equiv \frac{8}{3\pi} \frac{\langle B^2 \rangle}{\langle (\mathbf{b} \cdot \nabla B)^2 \rangle} \frac{g_t \omega_{ta}}{v_{Ta}^2 \tau_{aa}} = K_* \frac{g_t}{\omega_{ta} \tau_{aa}}, \quad (219)$$

$$f_c = \frac{3}{4} \frac{I_2}{B_{\max}^2} \int_0^1 \frac{\lambda d\lambda}{I_6(\lambda)}, \quad (220)$$

where

$$K_t = \frac{I_1^2 I_3}{I_2^2 \sum_{k>0} I_{4,k} I_{5,k}}, \quad (221)$$

$$K_* = \frac{8}{3\pi} \frac{I_2}{I_3} K_t^2. \quad (222)$$