Neoclassical Physics

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1 Plasma Equilibrium

Let

$$\vec{B} = I(\psi) \, \vec{\nabla}\zeta + \vec{\nabla}\zeta \times \vec{\nabla}\psi, \tag{1}$$

$$\vec{E} = -\vec{\nabla}\Phi(\psi) + E_{\parallel} \vec{b},\tag{2}$$

where $\vec{b} = \vec{B}/B$. Expect $n_a = n_a(\psi)$, $p_a = p_a(\psi)$, and $T_a = T_a(\psi)$. Follows that

$$\mu_0 \vec{j} \equiv \vec{\nabla} \times \vec{B} = \vec{\nabla} I \times \vec{\nabla} \zeta + R^2 \vec{\nabla} \cdot \left(\frac{\vec{\nabla} \psi}{R^2}\right) \nabla \zeta. \tag{3}$$

Force balance implies that

$$\vec{j} \times \vec{B} = \vec{\nabla}p,\tag{4}$$

which yields the Grad-Shafranov equation:

$$R^{2} \vec{\nabla} \cdot \left(\frac{\vec{\nabla} \psi}{R^{2}} \right) = -I \frac{dI}{d\psi} - \mu_{0} R^{2} \frac{dp}{d\psi}.$$
 (5)

Finally,

$$\mu_0 \vec{j} \cdot \vec{B} = -\left(\frac{I^2}{R^2} + \frac{|\vec{\nabla}\psi|^2}{R^2}\right) \frac{dI}{d\psi} - \mu_0 I \frac{dp}{d\psi}.$$
 (6)

2 Fluid Velocities and Heat Fluxes

We can write

$$\vec{u}_{j}^{a} = u_{\parallel j}^{a} \vec{b} + \vec{u}_{\perp j}^{a}, \tag{7}$$

and

$$\vec{u}_0^a = \vec{V}_a,\tag{8}$$

$$\vec{u}_1^a = -\frac{2}{5} \frac{\vec{q}_a}{p_a},\tag{9}$$

with

$$\vec{u}_{\perp 0}^{a} = \frac{\vec{E} \times \vec{B}}{B^{2}} + \frac{\vec{B} \times \vec{\nabla} p_{a}}{n_{a} e_{a} B^{2}}, \tag{10}$$

$$\vec{u}_{\perp 1}^a = -\frac{\vec{B} \times \vec{\nabla} T_a}{e_a B^2}.\tag{11}$$

It follows that

$$u_{\theta j}^{a} = \frac{u_{\parallel j}^{a}}{B} - \frac{V_{\psi j}^{a}}{B^{2}},\tag{12}$$

where

$$u_{\theta j}^{a}(\psi) = \frac{\vec{u}_{j}^{a} \cdot \vec{\nabla} \theta}{\vec{B} \cdot \vec{\nabla} \theta},\tag{13}$$

$$V_{\psi \, 0}^{a}(\psi) = -\left(\frac{\vec{u}_{\perp \, 0} \cdot \vec{\nabla} \theta}{\vec{B} \cdot \vec{\nabla} \theta}\right) B^{\, 2} = -\frac{I \, T_{a}}{e_{a}} \left(\frac{d \ln p_{a}}{d \psi} + \frac{e_{a}}{T_{a}} \frac{d\Phi}{d \psi}\right),\tag{14}$$

$$V_{\psi 1}^{a}(\psi) = -\left(\frac{\vec{u}_{\perp 1} \cdot \vec{\nabla}\theta}{\vec{B} \cdot \vec{\nabla}\theta}\right) B^{2} = \frac{I T_{a}}{e_{a}} \frac{d \ln T_{a}}{d\psi}.$$
 (15)

3 Friction Forces

$$\begin{pmatrix} \vec{F}_{a0} \\ \vec{F}_{a1} \end{pmatrix} = \sum_{b} \begin{pmatrix} l_{00}^{ab}, & l_{01}^{ab} \\ l_{10}^{ab}, & l_{11}^{ab} \end{pmatrix} \begin{pmatrix} \vec{u}_{0}^{b} \\ \vec{u}_{1}^{b} \end{pmatrix}, \tag{16}$$

$$l_{jk}^{ab} = \left(\sum_{c} \frac{n_a \, m_a}{\tau_{ac}} \, M_{ac}^{jk}\right) \delta_{ab} + \frac{n_a \, m_a}{\tau_{ab}} \, N_{ab}^{jk},\tag{17}$$

$$\frac{1}{\tau_{ab}} = \frac{4}{3\sqrt{\pi}} \frac{4\pi \, n_b \, e_a^2 \, e_b^2 \, \ln \Lambda}{(4\pi \, \epsilon_0)^2 \, m_a^2 \, v_{Ta}^3},\tag{18}$$

$$v_{Ta} = \sqrt{\frac{2T_a}{m_a}},\tag{19}$$

$$M_{ab}^{00} = -\frac{1 + m_a/m_b}{(1 + x_{ab}^2)^{3/2}},\tag{20}$$

$$M_{ab}^{01} = -\frac{3}{2} \frac{1 + m_a/m_b}{(1 + x_{ab}^2)^{5/2}} = M_{ab}^{10}, \tag{21}$$

$$M_{ab}^{11} = -\frac{13/4 + 4x_{ab}^2 + (15/2)x_{ab}^4}{(1 + x_{ab}^2)^{5/2}},$$
(22)

$$N_{ab}^{00} = \frac{1 + m_a/m_b}{(1 + x_{ab}^2)^{3/2}},\tag{23}$$

$$N_{ab}^{01} = \frac{3}{2} \frac{T_a}{T_b} \frac{1 + m_b/m_a}{x_{ab} (1 + x_{ba}^2)^{5/2}},$$
(24)

$$N_{ab}^{10} = \frac{3}{2} \frac{1 + m_a/m_b}{(1 + x_{ab}^2)^{5/2}},\tag{25}$$

$$N_{ab}^{11} = \frac{27}{4} \frac{T_a}{T_b} \frac{x_{ab}^2}{(1 + x_{ab}^2)^{5/2}},\tag{26}$$

where

$$x_{ab} = \frac{v_{Tb}}{v_{Ta}}. (27)$$

Species are e, i, and I. Let

$$Z_{\text{eff}} = \frac{n_i \, e_i^2 + n_I \, e_I^2}{n_e \, e_e^2}.$$
 (28)

Follows that

$$\frac{n_e}{n_i} = \frac{Z_I - 1}{Z_I - Z_{\text{eff}}},\tag{29}$$

$$\frac{n_I}{n_i} = \frac{Z_{\text{eff}} - 1}{Z_I \left(Z_I - Z_{\text{eff}} \right)},\tag{30}$$

$$\frac{n_I}{n_e} = \frac{Z_{\text{eff}} - 1}{Z_I (Z_I - 1)}. (31)$$

Thus,

$$M_{aa}^{00} = -\frac{1}{\sqrt{2}},\tag{32}$$

$$M_{aa}^{01} = -\frac{3}{4\sqrt{2}} = M_{aa}^{10},\tag{33}$$

$$M_{aa}^{11} = -\frac{59}{16\sqrt{2}},\tag{34}$$

$$N_{aa}^{00} = \frac{1}{\sqrt{2}},\tag{35}$$

$$N_{aa}^{01} = \frac{3}{4\sqrt{2}},\tag{36}$$

$$N_{aa}^{10} = \frac{3}{4\sqrt{2}},\tag{37}$$

$$N_{aa}^{11} = \frac{27}{16\sqrt{2}},\tag{38}$$

for a = e, i, I. Now,

$$M_{iI}^{00} = -\frac{1 + m_i/m_I}{(1 + x_{iI}^2)^{3/2}},\tag{39}$$

$$M_{iI}^{01} = -\frac{3}{2} \frac{1 + m_i/m_I}{(1 + x_{iI}^2)^{5/2}} = M_{iI}^{10}, \tag{40}$$

$$M_{iI}^{11} = -\frac{13/4 + 4x_{iI}^2 + (15/2)x_{iI}^4}{(1 + x_{iI}^2)^{5/2}},$$
(41)

$$N_{iI}^{00} = \frac{1 + m_i/m_I}{(1 + x_{iI}^2)^{3/2}},\tag{42}$$

$$N_{iI}^{01} = \frac{3}{2} \frac{T_i}{T_I} \frac{1 + m_I/m_i}{x_{iI} (1 + x_{Ii}^2)^{5/2}},$$
(43)

$$N_{iI}^{10} = \frac{3}{2} \frac{1 + m_i/m_I}{(1 + x_{iI}^2)^{5/2}},\tag{44}$$

$$N_{iI}^{11} = \frac{27}{4} \frac{T_i}{T_I} \frac{x_{iI}^2}{(1 + x_{iI}^2)^{5/2}}.$$
 (45)

Furthermore,

$$M_{Ii}^{00} = -\frac{1 + m_I/m_i}{(1 + x_{Ii}^2)^{3/2}},\tag{46}$$

$$M_{Ii}^{01} = -\frac{3}{2} \frac{1 + m_I/m_i}{(1 + x_{Ii}^2)^{5/2}} = M_{Ii}^{10}, \tag{47}$$

$$M_{Ii}^{11} = -\frac{13/4 + 4x_{Ii}^2 + (15/2)x_{Ii}^4}{(1 + x_{Ii}^2)^{5/2}},$$
(48)

$$N_{Ii}^{00} = \frac{1 + m_I/m_i}{(1 + x_{Ii}^2)^{3/2}},\tag{49}$$

$$N_{Ii}^{01} = \frac{3}{2} \frac{T_I}{T_i} \frac{1 + m_i/m_I}{x_{Ii} (1 + x_{iI}^2)^{5/2}},$$
(50)

$$N_{Ii}^{10} = \frac{3}{2} \frac{1 + m_I/m_i}{(1 + x_{Ii}^2)^{5/2}},\tag{51}$$

$$N_{Ii}^{11} = \frac{27}{4} \frac{T_I}{T_i} \frac{x_{Ii}^2}{(1 + x_{Ii}^2)^{5/2}}.$$
 (52)

Also,

$$M_{eq}^{00} = -1, (53)$$

$$M_{ea}^{01} = -\frac{3}{2} = M_{ea}^{10}, (54)$$

$$M_{ea}^{11} = -\frac{13}{4},\tag{55}$$

$$N_{ea}^{00} = 1, (56)$$

$$N_{ea}^{01} = 0, (57)$$

$$N_{ea}^{10} = \frac{3}{2},\tag{58}$$

$$N_{ea}^{11} = 0, (59)$$

for a = i, I, where we have neglected terms of order $(m_e/m_i)^{1/2}$. Finally,

$$M_{ae}^{00} = 0, (60)$$

$$M_{ae}^{01} = 0 = M_{ae}^{10}, (61)$$

$$M_{ae}^{11} = 0, (62)$$

$$N_{ae}^{00} = 0, (63)$$

$$N_{ae}^{01} = 0, (64)$$

$$N_{ae}^{10} = 0 (65)$$

$$N_{ae}^{11} = 0, (66)$$

for a = i, I. Now,

$$l_{jk}^{ii} = \frac{n_i \, m_i}{\tau_{ii}} \left[\left(\sum_a \frac{n_a \, e_a^2}{n_i \, e_i^2} \, M_{ia}^{jk} \right) + N_{ii}^{jk} \right] \tag{67}$$

$$= \frac{n_i \, m_i}{\tau_{ii}} \left(M_{ii}^{jk} + N_{ii}^{jk} + \alpha \, M_{iI}^{jk} \right), \tag{68}$$

where

$$\alpha = \frac{n_I e_I^2}{n_i e_i^2} = (Z_{\text{eff}} - 1) \left(\frac{Z_I}{Z_I - Z_{\text{eff}}} \right).$$
 (69)

Thus,

$$l_{00}^{ii} = -\frac{n_i m_i}{\tau_{ii}} \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{3/2}},$$
(70)

$$l_{01}^{ii} = -\frac{n_i m_i}{\tau_{ii}} \frac{3}{2} \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{5/2}}, \tag{71}$$

$$l_{10}^{ii} = -\frac{n_i m_i}{\tau_{ii}} \frac{3}{2} \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{5/2}}, \tag{72}$$

$$l_{11}^{ii} = -\frac{n_i m_i}{\tau_{ii}} \left\{ \sqrt{2} + \frac{\alpha \left[13/4 + 4 x_{iI}^2 + (15/2) x_{iI}^4 \right]}{(1 + x_{iI}^2)^{5/2}} \right\}.$$
 (73)

Now,

$$l_{jk}^{iI} = \frac{n_i \, m_i}{\tau_{ii}} \, \alpha \, N_{iI}^{jk}, \tag{74}$$

SO

$$l_{00}^{iI} = \frac{n_i m_i}{\tau_{ii}} \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{3/2}},$$
(75)

$$l_{01}^{iI} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{3}{2} \, \frac{T_i}{T_I} \, \frac{\alpha \, (1 + m_I/m_i)}{x_{iI} \, (1 + x_{I_i}^2)^{5/2}},\tag{76}$$

$$l_{10}^{iI} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{3}{2} \, \frac{\alpha \, (1 + m_i / m_I)}{(1 + x_{iI}^2)^{5/2}},\tag{77}$$

$$l_{11}^{iI} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{27}{4} \, \frac{T_i}{T_I} \, \frac{\alpha \, x_{iI}^2}{(1 + x_{iI}^2)^{5/2}}. \tag{78}$$

Now,

$$l_{jk}^{II} = \frac{n_I m_I}{\tau_{II}} \left[\left(\sum_a \frac{n_a e_a^2}{n_I e_I^2} M_{Ia}^{jk} \right) + N_{II}^{jk} \right]$$
 (79)

$$= \frac{n_I \, m_I}{\tau_{II}} \left(M_{II}^{jk} + N_{II}^{jk} + \alpha^{-1} \, M_{Ii}^{jk} \right)$$

$$= \frac{n_i m_i}{\tau_{ii}} \alpha^2 \frac{m_i}{m_I} x_{Ii}^3 \left(M_{II}^{jk} + N_{II}^{jk} + \alpha^{-1} M_{Ii}^{jk} \right), \tag{80}$$

so

$$l_{00}^{II} = -\frac{n_i m_i}{\tau_{ii}} \frac{m_i}{m_I} x_{Ii}^3 \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{3/2}}, \tag{81}$$

$$l_{01}^{II} = -\frac{n_i m_i}{\tau_{ii}} \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}}, \tag{82}$$

$$l_{10}^{II} = -\frac{n_i m_i}{\tau_{ii}} \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}},$$
(83)

$$l_{11}^{II} = -\frac{n_i m_i}{\tau_{ii}} \frac{m_i}{m_I} x_{Ii}^3 \left\{ \sqrt{2} \alpha^2 + \frac{\alpha \left[13/4 + 4 x_{Ii}^2 + (15/2) x_{Ii}^4 \right]}{(1 + x_{Ii}^2)^{5/2}} \right\}.$$
 (84)

Finally,

$$l_{jk}^{Ii} = \frac{n_i m_i}{\tau_{ii}} \alpha \frac{m_i}{m_I} x_{Ii}^3 N_{Ii}^{jk}, \tag{85}$$

SO

$$l_{00}^{Ii} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{m_i}{m_I} \, x_{Ii}^3 \, \frac{\alpha \, (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{3/2}}, \tag{86}$$

$$l_{01}^{Ii} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{m_i}{m_I} \, x_{Ii}^3 \, \frac{3}{2} \, \frac{T_I}{T_i} \, \frac{\alpha \, (1 + m_i/m_I)}{x_{Ii} \, (1 + x_{iI}^2)^{5/2}}, \tag{87}$$

$$l_{10}^{Ii} = \frac{n_i m_i}{\tau_{ii}} \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}},$$
(88)

$$l_{11}^{Ii} = \frac{n_i \, m_i}{\tau_{ii}} \, \frac{m_i}{m_I} \, x_{Ii}^3 \, \frac{27}{4} \, \frac{T_I}{T_i} \, \frac{\alpha \, x_{Ii}^2}{(1 + x_{Ii}^2)^{5/2}}. \tag{89}$$

Thus,

$$(F^{i}) = \frac{n_{i} m_{i}}{\tau_{ii}} \left\{ - \left[F^{ii} \right] (u^{i}) + \left[F^{iI} \right] (u^{I}) \right\}, \tag{90}$$

$$(F^{I}) = \frac{n_i \, m_i}{\tau_{ii}} \left\{ \left[F^{Ii} \right] \, (u^i) - \left[F^{II} \right] \, (u^I) \right\}, \tag{91}$$

where

$$(F^a) = \begin{pmatrix} \vec{F}_{a0} \\ \vec{F}_{a1} \end{pmatrix}, \tag{92}$$

$$(u^a) = \begin{pmatrix} \vec{u}_0^a \\ \vec{u}_1^a \end{pmatrix}. \tag{93}$$

Here,

$$F_{00}^{ii} = \frac{\alpha \left(1 + m_i/m_I\right)}{\left(1 + x_{iI}\right)^{3/2}},\tag{94}$$

$$F_{01}^{ii} = \frac{3}{2} \frac{\alpha \left(1 + m_i/m_I\right)}{\left(1 + x_{iI}^2\right)^{5/2}},\tag{95}$$

$$F_{10}^{ii} = \frac{3}{2} \frac{\alpha \left(1 + m_i/m_I\right)}{\left(1 + x_{iI}^2\right)^{5/2}} = F_{01}^{ii}, \tag{96}$$

$$F_{11}^{ii} = \sqrt{2} + \frac{\alpha \left[13/4 + 4x_{iI}^2 + (15/2)x_{iI}^4\right]}{(1 + x_{iI}^2)^{5/2}},\tag{97}$$

$$F_{00}^{iI} = \frac{\alpha \left(1 + m_i/m_I\right)}{(1 + x_{iI}^2)^{3/2}} = F_{00}^{ii}, \tag{98}$$

$$F_{01}^{iI} = \frac{3}{2} \frac{T_i}{T_I} \frac{\alpha \left(1 + m_I/m_i\right)}{x_{iI} \left(1 + x_{Ii}^2\right)^{5/2}},\tag{99}$$

$$F_{10}^{iI} = \frac{3}{2} \frac{\alpha \left(1 + m_i/m_I\right)}{\left(1 + x_{iI}^2\right)^{5/2}} = F_{01}^{ii}, \tag{100}$$

$$F_{11}^{iI} = \frac{27}{4} \frac{T_i}{T_I} \frac{\alpha x_{iI}^2}{(1 + x_{iI}^2)^{5/2}},\tag{101}$$

$$F_{00}^{Ii} = \frac{m_i}{m_I} x_{Ii}^3 \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{3/2}} = \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{3/2}} = F_{00}^{ii},$$
(102)

$$F_{01}^{Ii} = \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{T_I}{T_i} \frac{\alpha \left(1 + m_i/m_I\right)}{x_{Ii} \left(1 + x_{iI}^2\right)^{5/2}} = \frac{3}{2} \frac{\alpha \left(1 + m_i/m_I\right)}{\left(1 + x_{iI}^2\right)^{5/2}} = F_{01}^{ii}, \tag{103}$$

$$F_{10}^{Ii} = \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}} = \frac{3}{2} \frac{T_i}{T_I} \frac{\alpha (1 + m_I/m_i)}{x_{II} (1 + x_{Ii}^2)^{5/2}} = F_{01}^{iI},$$
(104)

$$F_{11}^{Ii} = \frac{m_i}{m_I} x_{Ii}^3 \frac{27}{4} \frac{T_I}{T_i} \frac{\alpha x_{Ii}^2}{(1 + x_{Ii}^2)^{5/2}} = \frac{27}{4} \frac{\alpha x_{iI}^2}{(1 + x_{iI}^2)^{5/2}},$$
(105)

$$F_{00}^{II} = \frac{m_i}{m_I} x_{Ii}^3 \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{3/2}} = \frac{\alpha (1 + m_i/m_I)}{(1 + x_{iI}^2)^{3/2}} = F_{00}^{ii},$$
(106)

$$F_{01}^{II} = \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}} = \frac{3}{2} \frac{T_i}{T_I} \frac{\alpha (1 + m_I/m_i)}{x_{iI} (1 + x_{Ii}^2)^{5/2}} = F_{01}^{iI},$$
(107)

$$F_{10}^{II} = \frac{m_i}{m_I} x_{Ii}^3 \frac{3}{2} \frac{\alpha (1 + m_I/m_i)}{(1 + x_{Ii}^2)^{5/2}} = \frac{3}{2} \frac{T_i}{T_I} \frac{\alpha (1 + m_I/m_i)}{x_{II} (1 + x_{Ii}^2)^{5/2}} = F_{01}^{iI},$$
(108)

$$F_{11}^{II} = \frac{m_i}{m_I} x_{Ii}^3 \left\{ \sqrt{2} \alpha^2 + \frac{\alpha \left[13/4 + 4 x_{Ii}^2 + (15/2) x_{Ii}^4 \right]}{(1 + x_{Ii}^2)^{5/2}} \right\}$$

$$= \frac{T_i}{T_I} \left\{ \sqrt{2} \alpha^2 x_{Ii} + \frac{\alpha \left[15/2 + 4 x_{iI}^2 + (13/4) x_{iI}^4 \right]}{(1 + x_{iI}^2)^{5/2}} \right\}.$$
(109)

Note that $F_{0j}^{ii} = F_{0j}^{Ii}$ and $F_{0j}^{iI} = F_{0j}^{II}$, for j = 0, 1, which implies that $(F^I)_0 + (F^i)_0 = 0$. This is a statement of collisional momentum conservation. Furthermore, $F_{j0}^{ii} = F_{j0}^{iI}$ and $F_{j0}^{Ii} = F_{j0}^{II}.$ Now,

$$l_{jk}^{ee} = \frac{n_e \, m_e}{\tau_{ee}} \left[\left(\sum_a \frac{n_a \, e_a^2}{n_e \, e_e^2} \, M_{ea}^{jk} \right) + N_{ee}^{jk} \right] \tag{110}$$

$$= \frac{n_e \, m_e}{\tau_{ee}} \left(M_{ee}^{jk} + N_{ee}^{jk} + Z_{\text{eff}} \, M_{ei}^{jk} \right), \tag{111}$$

where a = e, i, I. It follows that

$$l_{00}^{ee} = -\frac{n_e \, m_e}{\tau_{ee}} \, Z_{\text{eff}},\tag{112}$$

$$l_{01}^{ee} = -\frac{n_e \, m_e}{\tau_{ee}} \, \frac{3}{2} \, Z_{\text{eff}},\tag{113}$$

$$l_{10}^{ee} = -\frac{n_e \, m_e}{\tau_{ee}} \, \frac{3}{2} \, Z_{\text{eff}},\tag{114}$$

$$l_{11}^{ee} = -\frac{n_e \, m_e}{\tau_{ee}} \left(\sqrt{2} + \frac{13}{4} \, Z_{\text{eff}} \right). \tag{115}$$

Likewise,

$$l_{jk}^{ea} = \frac{n_e \, m_e}{\tau_{ea}} \, Z_{\text{eff } a} \, N_{ea}^{jk}, \tag{116}$$

where

$$Z_{\text{eff }a} = \frac{n_a \, e_a^2}{n_e \, e_e^2},\tag{117}$$

for a = i, I. In particular,

$$Z_{\text{eff }i} = \frac{Z_I - Z_{\text{eff}}}{Z_I - 1},$$
 (118)

$$Z_{\text{eff }I} = \frac{Z_I (Z_{\text{eff}} - 1)}{Z_I - 1}.$$
 (119)

Hence,

$$l_{00}^{ei} = \frac{n_e \, m_e}{\tau_{ee}} \, Z_{\text{eff} \, i},\tag{120}$$

$$l_{01}^{ei} = 0, (121)$$

$$l_{10}^{ei} = \frac{n_e \, m_e}{\tau_{ee}} \, \frac{3}{2} \, Z_{\text{eff } i}, \tag{122}$$

$$l_{11}^{ei} = 0, (123)$$

and

$$l_{00}^{eI} = \frac{n_e \, m_e}{\tau_{ee}} \, Z_{\text{eff} \, I},\tag{124}$$

$$l_{01}^{eI} = 0, (125)$$

$$l_{10}^{eI} = \frac{n_e \, m_e}{\tau_{ee}} \, \frac{3}{2} \, Z_{\text{eff } I},\tag{126}$$

$$l_{11}^{eI} = 0. (127)$$

Thus,

$$(F^{e}) = \frac{n_{e} m_{e}}{\tau_{ee}} \left\{ - [F^{ee}] (u^{e}) + [F^{ei}] (u^{i}) + [F^{eI}] (u^{I}) \right\}, \tag{128}$$

where

$$F_{00}^{ee} = Z_{\text{eff}},\tag{129}$$

$$F_{01}^{ee} = \frac{3}{2} Z_{\text{eff}},\tag{130}$$

$$F_{10}^{ee} = \frac{3}{2} Z_{\text{eff}},$$
 (131)

$$F_{11}^{ee} = \sqrt{2} + \frac{13}{4} Z_{\text{eff}},$$
 (132)

$$F_{00}^{ei} = Z_{\text{eff }i},$$
 (133)

$$F_{01}^{ei} = 0, (134)$$

$$F_{10}^{ei} = \frac{3}{2} Z_{\text{eff} i}, \tag{135}$$

$$F_{11}^{ei} = 0, (136)$$

$$F_{00}^{eI} = Z_{\text{eff }I},$$
 (137)

$$F_{01}^{eI} = 0, (138)$$

$$F_{10}^{eI} = \frac{3}{2} Z_{\text{eff }I},\tag{139}$$

$$F_{11}^{eI} = 0. (140)$$

Note that $F_{j0}^{ee} = F_{j0}^{ei} + F_{j0}^{eI}$.

4 Neoclassical Viscosities

$$\left\langle \vec{B} \cdot \vec{\nabla} \cdot (\vec{\Pi}^{a}) \right\rangle = \frac{n_a \, m_a}{\tau_{aa}} \left[\mu^{a} \right] (u_{\theta}^{a}), \tag{141}$$

where

$$\langle A \rangle = \oint \frac{A(\theta) \, d\theta}{\vec{B} \cdot \vec{\nabla} \theta} / \oint \frac{d\theta}{\vec{B} \cdot \vec{\nabla} \theta}, \tag{142}$$

$$(\vec{\Pi}^{a}) = \begin{pmatrix} \vec{\Pi}_{\parallel a} \\ \vec{\Theta}_{\parallel a} \end{pmatrix}, \tag{143}$$

$$(u_{\theta}^{a}) = \begin{pmatrix} u_{\theta 0}^{a} \langle B^{2} \rangle \\ u_{\theta 1}^{a} \langle B^{2} \rangle \end{pmatrix}, \tag{144}$$

$$[\mu^a] = \begin{pmatrix} \mu_{00}^a, & \mu_{01}^a \\ \mu_{10}^a, & \mu_{11}^a \end{pmatrix}. \tag{145}$$

Let

$$\mu_{00}^a = K_{00}^a, \tag{146}$$

$$\mu_{01}^a = \frac{5}{2} K_{00}^a - K_{01}^a = \mu_{10}^a, \tag{147}$$

$$\mu_{11}^a = K_{11}^a - 5K_{01}^a + \frac{25}{4}K_{00}^a, \tag{148}$$

and

$$\omega_{t\,a} = \frac{v_{T\,a}}{L_c},\tag{149}$$

$$\nu_{*a} = \frac{8}{3\pi} \frac{\langle B^2 \rangle}{\langle (\vec{b} \cdot \vec{\nabla} B)^2 \rangle} \frac{g \,\omega_{ta}}{v_{Ta}^2 \,\tau_{aa}},\tag{150}$$

$$g = \frac{1 - f_c}{f_c},\tag{151}$$

$$f_c = \frac{3}{4} \langle B^2 \rangle \int_0^{1/B_{\text{max}}} \frac{\lambda \, d\lambda}{\langle \sqrt{1 - \lambda \, B} \rangle}.$$
 (152)

We have

$$K_{ij}^{a} = g \frac{4}{3\sqrt{\pi}} \int_{0}^{\infty} \frac{e^{-x} x^{4+i+j} \hat{\nu}_{D}^{a}(x) dx}{\left[x^{2} + \nu_{*a} \hat{\nu}_{D}^{a}(x)\right] \left[x^{2} + (5\pi/8) \left(\omega_{ta} \tau_{aa}\right)^{-1} \hat{\nu}_{T}^{a}(x)\right]},$$
 (153)

and

$$\hat{\nu}_D^a(x) = \frac{3\sqrt{\pi}}{4} \sum_b \frac{n_b e_b^2}{n_a e_a^2} \left[\left(1 - \frac{x_{ab}}{2x} \right) \psi \left(\frac{x}{x_{ab}} \right) + \psi' \left(\frac{x}{x_{ab}} \right) \right] \frac{1}{x},\tag{154}$$

$$\hat{\nu}_{\epsilon}^{a}(x) = \frac{3\sqrt{\pi}}{2} \sum_{b} \frac{n_b e_b^2}{n_a e_a^2} \left[\frac{m_a}{m_b} \psi \left(\frac{x}{x_{ab}} \right) - \psi' \left(\frac{x}{x_{ab}} \right) \right] \frac{1}{x},\tag{155}$$

$$\hat{\nu}_T^a(x) = 3\,\hat{\nu}_D^a(x) + \hat{\nu}_{\epsilon}^a(x),\tag{156}$$

$$\psi(x) = \frac{2}{\sqrt{\pi}} \int_0^{x^2} \sqrt{t} \, e^{-t} \, dt = \operatorname{erf}(x) - \frac{2}{\sqrt{\pi}} x \, e^{-x^2}, \tag{157}$$

$$\psi'(x) = \frac{2}{\sqrt{\pi}} x e^{-x^2}.$$
 (158)

In particular,

$$K_{ij}^{i} = g \frac{4}{3\sqrt{\pi}} \int_{0}^{\infty} \frac{e^{-x} x^{2+i+j} \hat{\nu}_{D}^{i}(x) dx}{\left[x + \nu_{*i} \hat{\nu}_{D}^{i}(x)\right] \left[x + (5\pi/8) \left(\omega_{ti} \tau_{ii}\right)^{-1} \hat{\nu}_{T}^{i}(x)\right]},$$
(159)

$$\hat{\nu}_D^i = \frac{3\sqrt{\pi}}{4} \left[\left(1 - \frac{1}{2x} \right) \psi(x) + \psi'(x) \right] \frac{1}{x}$$

$$+\frac{3\sqrt{\pi}}{4}\alpha\left[\left(1-\frac{x_{iI}}{2x}\right)\psi\left(\frac{x}{x_{iI}}\right)+\psi'\left(\frac{x}{x_{iI}}\right)\right]\frac{1}{x},\tag{160}$$

$$\hat{\nu}_{\epsilon}^{i} = \frac{3\sqrt{\pi}}{2} \left[\psi(x) - \psi'(x) \right] \frac{1}{x}$$

$$+\frac{3\sqrt{\pi}}{2}\alpha \left[\frac{m_i}{m_I}\psi\left(\frac{x}{x_{iI}}\right) - \psi'\left(\frac{x}{x_{iI}}\right)\right]\frac{1}{x},\tag{161}$$

and

$$K_{ij}^{I} = g \frac{4}{3\sqrt{\pi}} \int_{0}^{\infty} \frac{e^{-x} x^{2+i+j} \hat{\nu}_{D}^{I}(x) dx}{\left[x + \nu_{*I} \hat{\nu}_{D}^{I}(x)\right] \left[x + (5\pi/8) \left(\omega_{tI} \tau_{II}\right)^{-1} \hat{\nu}_{T}^{I}(x)\right]},$$
 (162)

$$\hat{\nu}_D^I = \frac{3\sqrt{\pi}}{4} \left[\left(1 - \frac{1}{2x} \right) \psi(x) + \psi'(x) \right] \frac{1}{x} + \frac{3\sqrt{\pi}}{4} \alpha^{-1} \left[\left(1 - \frac{x_{Ii}}{2x} \right) \psi\left(\frac{x}{x_{Ii}}\right) + \psi'\left(\frac{x}{x_{Ii}}\right) \right] \frac{1}{x}, \tag{163}$$

$$\hat{\nu}_{\epsilon}^{I} = \frac{3\sqrt{\pi}}{2} \left[\psi(x) - \psi'(x) \right] \frac{1}{x} + \frac{3\sqrt{\pi}}{2} \alpha^{-1} \left[\frac{m_{I}}{m_{i}} \psi\left(\frac{x}{x_{Ii}}\right) - \psi'\left(\frac{x}{x_{Ii}}\right) \right] \frac{1}{x}, \tag{164}$$

and

$$K_{ij}^{e} = g \frac{4}{3\sqrt{\pi}} \int_{0}^{\infty} \frac{e^{-x} x^{4+i+j} \hat{\nu}_{D}^{e}(x) dx}{[x^{2} + \nu_{*e} \hat{\nu}_{D}^{e}(x)] [x^{2} + (5\pi/8) (\omega_{te} \tau_{ee})^{-1} \hat{\nu}_{T}^{e}(x)]},$$
(165)

$$\hat{\nu}_D^e = \frac{3\sqrt{\pi}}{4} \left[\left(1 - \frac{1}{2x} \right) \psi(x) + \psi'(x) \right] + \frac{3\sqrt{\pi}}{4} Z_{\text{eff}}, \tag{166}$$

$$\hat{\nu}_{\epsilon}^{e} = \frac{3\sqrt{\pi}}{2} \left[\psi(x) - \psi'(x) \right]. \tag{167}$$

5 Parallel Force and Heat Balance Equations

$$\langle \vec{B} \cdot \vec{\nabla} \cdot \vec{\Pi}_a \rangle = \langle \vec{B} \cdot \vec{F}_{a0} \rangle + \langle \vec{B} \cdot \vec{F}_{a0}^{cx} \rangle + n_a \, e_a \, \langle \vec{E} \cdot \vec{B} \rangle, \tag{168}$$

$$\langle \vec{B} \cdot \vec{\nabla} \cdot \vec{\Theta}_a \rangle = \langle \vec{B} \cdot \vec{F}_{a1} \rangle + \langle \vec{B} \cdot \vec{F}_{a1}^{cx} \rangle. \tag{169}$$

Now,

$$\langle \vec{B} \cdot \vec{F}^{i} \rangle = \frac{n_{i} m_{i}}{\tau_{ii}} \left\{ - \left[F^{ii} \right] \left(u_{\parallel}^{i} \right) + \left[F^{iI} \right] \left(u_{\parallel}^{I} \right) \right\}, \tag{170}$$

$$\langle \vec{B} \cdot \vec{F}^{I} \rangle = \frac{n_i \, m_i}{\tau_{ii}} \left\{ \left[F^{Ii} \right] \left(u_{\parallel}^i \right) - \left[F^{II} \right] \left(u_{\parallel}^I \right) \right\}, \tag{171}$$

$$\langle \vec{B} \cdot \vec{F}^{e} \rangle = \frac{n_e \, m_e}{\tau_{ee}} \left\{ - \left[F^{ee} \right] \left(u_{\parallel}^e \right) + \left[F^{ei} \right] \left(u_{\parallel}^i \right) + \left[F^{eI} \right] \left(u_{\parallel}^I \right) \right\}, \tag{172}$$

$$\langle \vec{B} \cdot \vec{F}^{i cx} \rangle = -n_i \, m_i \, \langle \sigma \, v \rangle_i^{cx} \, [X^i](\tilde{u}_{\parallel}^i), \tag{173}$$

$$\langle \vec{B} \cdot \vec{F}^{I cx} \rangle = 0, \tag{174}$$

$$\langle \vec{B} \cdot \vec{F}^{e c x} \rangle = 0, \tag{175}$$

where

$$(u_{\parallel}^{a}) = \begin{pmatrix} \langle u_{\parallel 0}^{a} B \rangle \\ \langle u_{\parallel 1}^{a} B \rangle \end{pmatrix}, \tag{176}$$

$$(\tilde{u}_{\parallel}^{a}) = \begin{pmatrix} \langle n_{n} u_{\parallel 0}^{a} B \rangle \\ \langle n_{n} u_{\parallel 1}^{a} B \rangle \end{pmatrix}, \tag{177}$$

$$[X^i] = \begin{pmatrix} 1, & 0\\ 0, & E_n/T_i \end{pmatrix}. \tag{178}$$

Here, E_n is the neutral kinetic energy, n_n the neutral number density, and $\langle \sigma v \rangle_i^{cx}$ the majority ion charge exchange rate coefficient.

Let

$$(E_{\parallel}) = \begin{pmatrix} \langle E_{\parallel} B \rangle / \langle B \rangle \\ 0 \end{pmatrix}, \tag{179}$$

$$\Omega_a = \frac{e_a \langle B \rangle}{m_a},\tag{180}$$

$$[\tilde{\mu}^{I}] = \alpha^{2} \frac{m_{i}}{m_{I}} x_{Ii}^{3} [\mu^{I}] = \alpha^{2} \frac{T_{i}}{T_{I}} x_{Ii} [\mu^{I}].$$
(181)

Follows that

$$\begin{bmatrix}
[F^{ii}], & -[F^{iI}] \\
-[F^{Ii}], & [F^{II}]
\end{bmatrix}
\begin{bmatrix}
(u_{\parallel}^{i}) \\
(u_{\parallel}^{I})
\end{bmatrix}
\simeq -
\begin{bmatrix}
[\mu^{i}](u_{\theta}^{i}) \\
[\tilde{\mu}^{I}](u_{\theta}^{I})
\end{bmatrix} -
\begin{bmatrix}
\tau_{ii} \langle \sigma v \rangle_{i}^{cx} [X^{i}](\tilde{u}_{\parallel i}) \\
0
\end{bmatrix}, (182)$$

and

$$[\mu^{e}](u_{\theta}^{e}) = -[F^{ee}](u_{\parallel}^{e}) + [F^{ei}](u_{\parallel}^{i}) + [F^{eI}](u_{\parallel}^{I}) + \Omega_{e} \tau_{ee}(E_{\parallel}). \tag{183}$$

Follows from Eq. (12) that

$$(u_{\parallel}^{a}) = (u_{\theta}^{a}) + (V_{\psi}^{a}),$$
 (184)

and

$$\left(\tilde{u}_{\parallel}^{a}\right) = \frac{\langle n_{n} \rangle}{y} \left(u_{\theta}^{a}\right) + \langle n_{n} \rangle \left(V_{\psi}^{a}\right), \tag{185}$$

where

$$V_{\psi}^{a} = \begin{pmatrix} V_{\psi 0}^{a} \\ V_{\psi 1}^{a} \end{pmatrix}, \tag{186}$$

and

$$y = \frac{\langle n_n \rangle \langle B^2 \rangle}{\langle n_n B^2 \rangle}. (187)$$

Thus,

$$\begin{bmatrix} [F^{ii} + \mu^{i} + Y^{i}/y], & -[F^{iI}] \\ -[F^{Ii}], & [F^{II} + \tilde{\mu}^{I}] \end{bmatrix} \begin{bmatrix} (u_{\theta}^{i}) \\ (u_{\theta}^{I}) \end{bmatrix} = -\begin{bmatrix} [F^{ii} + Y^{i}], & -[F^{iI}] \\ -[F^{Ii}], & [F^{II}] \end{bmatrix} \begin{bmatrix} (V_{\psi}^{i}) \\ (V_{\psi}^{I}) \end{bmatrix},$$
(188)

where

$$[Y^{i}] = \tau_{ii} \langle \sigma v \rangle_{i}^{cx} \langle n_{n} \rangle [X^{i}]. \tag{189}$$

Can write

$$(V_{\psi}^{a}) = (V_{*}^{a}) + (V_{E}), \tag{190}$$

where

$$V_{*0}^{a}(\psi) = -\frac{I T_{a}}{e_{a}} \frac{d \ln p_{a}}{d \psi}, \tag{191}$$

$$V_{*1}^{a}(\psi) = \frac{I T_{a}}{e_{a}} \frac{d \ln T_{a}}{d \psi}, \tag{192}$$

$$V_{E0}(\psi) = -I \frac{d\Phi}{d\psi},\tag{193}$$

$$V_{E1}(\psi) = 0. (194)$$

Fact that $F_{j0}^{ii} = F_{j0}^{iI}$ and $F_{j0}^{Ii} = F_{j0}^{II}$ means that we can rewrite Eq. (188) as

$$\begin{bmatrix}
[F^{ii} + \mu^{i} + Y^{i}/y], & -[F^{iI}] \\
-[F^{Ii}], & [F^{II} + \tilde{\mu}^{I}]
\end{bmatrix}
\begin{bmatrix}
(u_{\theta}^{i}) \\
(u_{\theta}^{I})
\end{bmatrix} = -\begin{bmatrix}
[F^{ii} + Y^{i}], & -[F^{iI}] \\
-[F^{Ii}], & [F^{II}]
\end{bmatrix}
\begin{bmatrix}
(V_{*}^{i}) \\
(V_{*}^{I})
\end{bmatrix}$$

$$-\tau_{ii} \langle \sigma \, v \rangle_{i}^{cx} \langle n_{n} \rangle \begin{bmatrix}
(V_{E}) \\
(0)
\end{bmatrix}. \tag{195}$$

Let

$$\begin{bmatrix} [L^{ii}], & [L^{iI}] \\ [L^{Ii}], & [L^{II}] \end{bmatrix} = \begin{bmatrix} [F^{ii} + \mu^i + Y^i/y], & -[F^{iI}] \\ -[F^{Ii}], & [F^{II} + \tilde{\mu}^I] \end{bmatrix}^{-1} \begin{bmatrix} [F^{ii} + Y^i], & -[F^{iI}] \\ -[F^{Ii}], & [F^{II}] \end{bmatrix}, (196)$$

and

$$\begin{bmatrix} [G^{ii}], & [G^{iI}] \\ [G^{Ii}], & [G^{II}] \end{bmatrix} = \tau_{ii} \langle \sigma v \rangle_i^{cx} \langle n_n \rangle \begin{bmatrix} [F^{ii} + \mu^i + Y^i/y], & -[F^{iI}] \\ -[F^{Ii}], & [F^{II} + \tilde{\mu}^I] \end{bmatrix}^{-1}.$$
(197)

Follows that

$$(u_{\theta}^{i}) = -[G^{ii}](V_{E}) - [L^{ii}](V_{*}^{i}) - [L^{iI}](V_{*}^{I}), \tag{198}$$

$$(u_{\theta}^{I}) = -[L^{II}](V_{*}^{I}) - [L^{Ii}](V_{*}^{i}), \tag{199}$$

$$(u_{\parallel}^{i}) = ([1] - [G^{ii}])(V_{E}) + (1 - [L^{ii}])(V_{*}^{i}) - [L^{iI}](V_{*}^{I}), \tag{200}$$

$$(u_{\parallel}^{I}) = (V_{E}) + (1 - [L^{II}])(V_{*}^{I}) - [L^{Ii}](V_{*}^{i}). \tag{201}$$

Equation (183) becomes

$$[F^{ee} + \mu^{e}] (u_{\theta}^{e}) = -([F^{ee}] - [F^{ei}] - [F^{eI}]) (V_{E}) - [F^{ei}] [G^{ii}] (V_{E}) - [F^{ee}] (V_{*}^{e}) + \{ [F^{ei}] (1 - [L^{ii}]) - [F^{eI}] [L^{Ii}] \} (V_{*}^{i}) + \{ [F^{eI}] (1 - [L^{II}]) - [F^{ei}] [L^{iI}] \} (V_{*}^{I}) + \Omega_{ee} \tau_{ee} (E_{\parallel}).$$
 (202)

However, because $F_{j0}^{ee} = F_{j0}^{ei} + F_{j0}^{eI}$, the previous equation becomes

$$[F^{ee} + \mu^e](u_\theta^e) = -[F^{ei}][G^{ii}](V_E) - [F^{ee}](V_*^e)$$

$$+ \left\{ [F^{eI}] \left(1 - [L^{ii}] \right) - [F^{II}] [L^{Ii}] \right\} (V_*^i)$$

$$+ \left\{ [F^{eI}] \left(1 - [L^{II}] \right) - [F^{ei}] [L^{iI}] \right\} (V_*^I) + \Omega_{ee} \tau_{ee} (E_{\parallel}).$$
(203)

Let

$$[Q^{ee}] = [F^{ee} + \mu^e]^{-1}, \tag{204}$$

$$[P^{ei}] = [Q^{ee}] [F^{ei}] [G^{ii}], \tag{205}$$

$$[L^{ee}] = [Q^{ee}][F^{ee}],$$
 (206)

$$[L^{ei}] = -[Q^{ee}] \{ [F^{ei}] (1 - [L^{ii}]) - [F^{eI}] [L^{Ii}] \},$$
(207)

$$[L^{eI}] = -[Q^{ee}] \{ [F^{eI}] (1 - [L^{II}]) - [F^{ei}] [L^{iI}] \}.$$
(208)

Follows that

$$(u_{\theta}^{e}) = -[P^{ei}](V_{E}) - [L^{ee}](V_{*}^{e}) - [L^{ei}](V_{*}^{i}) - [L^{eI}](V_{*}^{I}) + \Omega_{e} \tau_{ee} [Q^{ee}](E_{\parallel}),$$
(209)

$$(u_{\parallel}^{e}) = ([1] - [P^{ei}]) (V_{E}) + (1 - [L^{ee}]) (V_{*}^{e}) - [L^{ei}] (V_{*}^{i}) - [L^{eI}] (V_{*}^{I}) + \Omega_{e} \tau_{ee} [Q^{ee}] (E_{\parallel}),$$
(210)

Finally,

$$(j_{\parallel}) \equiv \sum_{a} n_{a} e_{a} (u_{\parallel}^{a})$$

$$= -\left(n_{i} e_{i} [G^{ii}] + n_{e} e_{e} [P^{ei}]\right) (V_{E})$$

$$+ \left\{n_{i} e_{i} (1 - [L^{ii}]) - n_{I} e_{I} [L^{Ii}] - n_{e} e_{e} [L^{ei}]\right\} (V_{*}^{i})$$

$$+ \left\{n_{I} e_{I} (1 - [L^{II}]) - n_{i} e_{i} [L^{iI}] - n_{e} e_{e} [L^{eI}]\right\} (V_{*}^{I})$$

$$+ n_{e} e_{e} (1 - [L^{ee}]) (V_{*}^{e}) + n_{e} e_{e} \Omega_{e} \tau_{ee} [Q^{ee}] (E_{\parallel}), \tag{211}$$

where we have made use of the fact that $n_i e_i + n_I e_I + n_e e_e = 0$.

6 Fluid Velocities

Let

$$\omega_E(\psi) = -\frac{d\Phi}{d\psi},\tag{212}$$

$$\omega_{*a}(\psi) = -\frac{T_a}{e_a} \frac{d \ln p_a}{d\psi},\tag{213}$$

$$\eta_a(\psi) = \frac{d \ln T_a}{d \ln n_a}.$$
 (214)

It follows that

$$V_{*0}^{a}(\psi) = I \,\omega_{*a},\tag{215}$$

$$V_{*1}^{a}(\psi) = -I\left(\frac{\eta_a}{1+\eta_a}\right)\omega_{*a},\tag{216}$$

$$V_{E0}(\psi) = I\,\omega_E,\tag{217}$$

$$V_{E1}(\psi) = 0. (218)$$

We have $e_e = -e$, $e_i = e$, $e_I = Z_I e$. Thus,

$$\frac{V_{\theta}^{i}}{B_{\theta}} \frac{\langle B^{2} \rangle}{I} = -G_{00}^{ii} \omega_{E} - L_{00}^{ii} \omega_{*i} - L_{00}^{iI} \omega_{*I} + L_{01}^{ii} \left(\frac{\eta_{i}}{1 + \eta_{i}}\right) \omega_{*i} + L_{01}^{iI} \left(\frac{\eta_{I}}{1 + \eta_{I}}\right) \omega_{*I},$$
(219)

$$\frac{V_{\theta}^{I}}{B_{\theta}} \frac{\langle B^{2} \rangle}{I} = -L_{00}^{II} \omega_{*I} - L_{00}^{Ii} \omega_{*i} + L_{01}^{II} \left(\frac{\eta_{I}}{1 + \eta_{I}} \right) \omega_{*I} + L_{01}^{Ii} \left(\frac{\eta_{i}}{1 + \eta_{i}} \right) \omega_{*i},$$
(220)

$$\frac{\langle V_{\parallel}^{i} B \rangle}{I} = (1 - G_{00}^{ii}) \,\omega_{E} + \omega_{*i} - L_{00}^{ii} \,\omega_{*i} - L_{00}^{iI} \,\omega_{*I} + L_{01}^{ii} \left(\frac{\eta_{i}}{1 + \eta_{i}}\right) \omega_{*i} + L_{01}^{iI} \left(\frac{\eta_{I}}{1 + \eta_{I}}\right) \omega_{*I}, \tag{221}$$

$$\frac{\langle V_{\parallel}^{I}B\rangle}{I} = \omega_{E} + \omega_{*I} - L_{00}^{II} \,\omega_{*I} - L_{00}^{Ii} \,\omega_{*i} + L_{01}^{II} \left(\frac{\eta_{I}}{1 + \eta_{I}}\right) \omega_{*I} + L_{01}^{Ii} \left(\frac{\eta_{i}}{1 + \eta_{i}}\right) \omega_{*i}.$$
(222)

Also,

$$\frac{V_{\phi}^{i}}{R} = \omega_{E} + \omega_{*i} + \frac{I^{2}}{R^{2} \langle B^{2} \rangle} \left[-G_{00}^{ii} \omega_{E} - L_{00}^{ii} \omega_{*i} - L_{00}^{iI} \omega_{*I} + L_{01}^{ii} \left(\frac{\eta_{i}}{1 + \eta_{i}} \right) \omega_{*i} \right] + L_{01}^{iI} \left(\frac{\eta_{I}}{1 + \eta_{I}} \right) \omega_{*I} \right],$$
(223)

$$\frac{V_{\phi}^{I}}{R} = \omega_{E} + \omega_{*I} + \frac{I^{2}}{R^{2} \langle B^{2} \rangle} \left[-L_{00}^{II} \omega_{*I} - L_{00}^{Ii} \omega_{*i} + L_{01}^{II} \left(\frac{\eta_{I}}{1 + \eta_{I}} \right) \omega_{*I} + L_{01}^{Ii} \left(\frac{\eta_{i}}{1 + \eta_{i}} \right) \omega_{*i} \right].$$
(224)

7 Fluid Angular Rotation Velocities

Let

$$\omega_{\theta i}(\psi) = \frac{V_{\theta}^{i}}{B_{\theta}} \frac{I}{R^{2}},\tag{225}$$

$$\omega_{\theta I}(\psi) = \frac{V_{\theta}^{I}}{B_{\theta}} \frac{I}{R^{2}},\tag{226}$$

$$\omega_{\phi i}(\psi) = \frac{V_{\phi}^{i}}{R},\tag{227}$$

$$\omega_{\phi I}(\psi) = \frac{V_{\phi}^{I}}{R},\tag{228}$$

where the right-hand sides are evaluated on the outboard mid-plane. It follows that

$$\omega_{\theta i} = K_{\theta} \, \omega_{\text{nc} \, i}, \tag{229}$$

$$\omega_{\theta I} = K_{\theta} \, \omega_{\text{nc} I},\tag{230}$$

$$\omega_{\phi i} = \omega_E + \omega_{*i} + K_\theta \,\omega_{\text{nc}\,i},\tag{231}$$

$$\omega_{\phi I} = \omega_E + \omega_{*I} + K_\theta \,\omega_{\text{nc}\,I},\tag{232}$$

where

$$K_{\theta} = \frac{I^2}{R^2 \langle B^2 \rangle},\tag{233}$$

and

$$\omega_{\text{nc}\,i} = -G_{00}^{\,ii}\,\omega_E - L_{00}^{\,ii}\,\omega_{*\,i} - L_{00}^{\,iI}\,\omega_{*\,I} + L_{01}^{\,iI}\left(\frac{\eta_i}{1+\eta_i}\right)\omega_{*\,i} + L_{01}^{\,iI}\left(\frac{\eta_I}{1+\eta_I}\right)\omega_{*\,I},\tag{234}$$

$$\omega_{\text{nc}\,I} = -L_{00}^{II}\,\omega_{*\,I} - L_{00}^{Ii}\,\omega_{*\,i} + L_{01}^{II}\left(\frac{\eta_I}{1+\eta_I}\right)\omega_{*\,I} + L_{01}^{Ii}\left(\frac{\eta_i}{1+\eta_i}\right)\omega_{*\,i}. \tag{235}$$

8 Natural Frequencies

Let

$$\omega_{0\ln} = \omega_E + \omega_{*e},\tag{236}$$

$$\omega_{0 \, \text{eb}} = \omega_E, \tag{237}$$

$$\omega_{0\,\text{nl}} = \omega_E + \omega_{*i} + \omega_{\text{nc}\,i}. \tag{238}$$

The corresponding natural frequencies are

$$\varpi_{0\ln} = -n\,\omega_{0\ln},\tag{239}$$

$$\varpi_{0\,\mathrm{eb}} = -n\,\omega_{0\,\mathrm{eb}},\tag{240}$$

$$\varpi_{0\,\mathrm{nl}} = -n\,\omega_{0\,\mathrm{nl}}.\tag{241}$$

9 Current Densities

Can write

$$j_{\parallel} = j_{\text{bootstrap}} + j_{\text{ohmic}},$$
 (242)

where

$$\frac{\langle j_{\text{bootstrap}} B \rangle}{I} = -\left\{ P_{00}^{ei} - \left(\frac{Z_{I} - Z_{\text{eff}}}{Z_{\text{eff}} - 1} \right) G_{00}^{ii} \right\} n_{e} e \frac{d\Phi}{d\psi}
- \left\{ 1 - L_{00}^{ii} - \left(\frac{Z_{\text{eff}} - 1}{Z_{I} - Z_{\text{eff}}} \right) L_{00}^{Ii} + \left(\frac{Z_{I} - 1}{Z_{I} - Z_{\text{eff}}} \right) L_{00}^{ei} \right\} \frac{dp_{i}}{d\psi}
- \left\{ 1 - L_{00}^{II} - \left(\frac{Z_{I} - Z_{\text{eff}}}{Z_{\text{eff}} - 1} \right) L_{00}^{iI} + \left(\frac{Z_{I} - 1}{Z_{\text{eff}} - 1} \right) L_{00}^{eI} \right\} \frac{dp_{I}}{d\psi}
- \left(1 - L_{00}^{ee} \right) \frac{dp_{e}}{d\psi}
- \left\{ L_{01}^{ii} + \left(\frac{Z_{\text{eff}} - 1}{Z_{I} - Z_{\text{eff}}} \right) L_{01}^{Ii} - \left(\frac{Z_{I} - 1}{Z_{I} - Z_{\text{eff}}} \right) L_{01}^{ei} \right\} n_{i} \frac{dT_{i}}{d\psi}
- \left\{ L_{01}^{II} + \left(\frac{Z_{I} - Z_{\text{eff}}}{Z_{\text{eff}} - 1} \right) L_{01}^{iI} - \left(\frac{Z_{I} - 1}{Z_{\text{eff}} - 1} \right) L_{01}^{eI} \right\} n_{I} \frac{dT_{I}}{d\psi}
- L_{01}^{ee} n_{e} \frac{dT_{e}}{d\psi} \tag{243}$$

and

$$\langle j_{\text{ohmic}} B \rangle = Q_{00}^{ee} \, \sigma_{ee} \, \langle E_{\parallel} B \rangle,$$
 (244)

where

$$\sigma_{ee} = \frac{n_e \, e^2 \, \tau_{ee}}{m_e}.\tag{245}$$

10 Model Poloidal Flow Damping Term

Our model poloidal flow damping force density in the plasma equation of motion takes the form

$$\frac{m_i n_i}{\tau_{\theta}} \left(V_{\theta} - V_{\theta}^{nc} \right) \mathbf{e}_{\theta}, \tag{246}$$

where

$$\frac{1}{\tau_{\theta}} = \left(\frac{I^2}{\langle |\nabla \psi|^2 \rangle}\right) \frac{\mu_{00}^i}{\tau_{ii}},\tag{247}$$

and V_{θ}^{nc} is specified in Eq. (219).