Four-Field Resonant Layer Model

R. Fitzpatrick^a

Institute for Fusion Studies, Department of Physics, University of Texas at Austin, Austin TX 78712, USA

I. FOURIER-TRANSFORMED FOUR-FIELD EQUATIONS

The fundamental Fourier-transformed four-field equations take the form:

$$(g + i Q_e) \bar{\psi} = \frac{d(\bar{\phi} - \bar{N})}{dp} - p^2 \bar{\psi}, \tag{1}$$

$$g\,\bar{N} = -\mathrm{i}\,Q_e\,\bar{\phi} - D^2\,\frac{d(p^2\,\bar{\psi})}{dp} + c_\beta^2\,\frac{d\bar{V}}{dp} - P_\perp\,p^2\,\bar{N},$$
 (2)

$$(g + i Q_i) p^2 \bar{\phi} = \frac{d(p^2 \bar{\psi})}{dp} - P_{\varphi} p^4 \left(\bar{\phi} + \frac{\bar{N}}{\iota}\right), \tag{3}$$

$$g\,\bar{V} = \mathrm{i}\,Q_e\,\bar{\psi} + \frac{d\bar{N}}{dp} - P_\varphi\,p^2\,\bar{V}.\tag{4}$$

It follows that

$$\bar{\psi} = \frac{1}{g + i Q_e + p^2} \frac{d(\bar{\phi} - \bar{N})}{dp}.$$
 (5)

Now,

$$\frac{d(p^2 \,\bar{\psi})}{dp} = [(g + i \,Q_i) \,p^2 + P_{\varphi} \,p^4] \,\bar{\phi} + \frac{P_{\varphi}}{\iota} \,p^4 \,\bar{N},\tag{6}$$

so

$$c_{\beta}^{2} \frac{d\bar{V}}{dn} = (g + P_{\perp} p^{2} + \iota^{-1} D^{2} P_{\varphi} p^{4}) \bar{N} + [i Q_{e} + D^{2} (g + i Q_{i}) p^{2} + D^{2} P_{\varphi} p^{4}] \bar{\phi}.$$
 (7)

Let

$$\bar{J} = p^2 \, \bar{\psi},\tag{8}$$

$$\bar{Y} = \bar{\phi} - \bar{N}. \tag{9}$$

Hence, we obtain our final set of Fourier-transformed four-field equations:

$$\frac{d\bar{Y}}{dp} = \left(\frac{g + iQ_e + p^2}{p^2}\right)\bar{J},\tag{10}$$

a rfitzp@utexas.edu

$$\frac{d\bar{N}}{dp} = \left(\frac{-i\,Q_e}{p^2}\right)\bar{J} + \left(g + P_\varphi\,p^2\right)\bar{V},\tag{11}$$

$$\frac{d\bar{J}}{dp} = [(g + i Q_i) p^2 + P_{\varphi} p^4] \bar{Y} + [(g + i Q_i) p^2 + \iota_e^{-1} P_{\varphi} p^4] \bar{N},$$
(12)

$$c_{\beta}^{2} \frac{d\bar{V}}{dp} = [i Q_{e} + D^{2} (g + i Q_{i}) p^{2} + D^{2} P_{\varphi} p^{4}] \bar{Y}$$

+
$$\{g + i Q_e + [P_{\perp} + D^2 (g + i Q_i)] p^2 + \iota_e^{-1} D^2 P_{\varphi} p^4 \} \bar{N},$$
 (13)

where $\iota_e = \iota/(1+\iota)$.

II. SMALL-p BEHAVIOR OF FOURIER-TRANSFORMED FOUR-FIELD EQUATIONS

A. Introduction

Let us search for power-law solutions of Eqs. (10)–(13) at small-p. Given that we have four coupled first-order equations, we expect to find four independent power-law solutions.

B. First Solution

Suppose that

$$\bar{Y}(p) = y_{-1} p^{-1} + y_1 p + \mathcal{O}(p^3),$$
 (14)

$$\bar{N}(p) = n_{-1} p^{-1} + n_1 p + \mathcal{O}(p^3), \tag{15}$$

$$\bar{J}(p) = j_0 + j_2 p^2 + \mathcal{O}(p^4), \tag{16}$$

$$\bar{V}(p) = v_2 p^2 + \mathcal{O}(p^4).$$
 (17)

Equations (10)–(13) yield

$$-y_{-1}p^{-2} + y_1 = (g + iQ_e)(j_0p^{-2} + j_2) + j_0 + \mathcal{O}(p^2),$$
(18)

$$-n_{-1}p^{-2} + n_1 = -iQ_e(j_0p^{-2} + j_2) + \mathcal{O}(p^2), \tag{19}$$

$$2 j_2 p = (g + i Q_i) (y_{-1} + n_{-1}) p + \mathcal{O}(p^3),$$
(20)

$$2 c_{\beta}^{2} v_{2} p = i Q_{e} (y_{-1} p^{-1} + y_{1} p) + (g + i Q_{e}) (n_{-1} p^{-1} + n_{1} p)$$

$$+ D^{2} (g + i Q_{i}) y_{-1} p + [P_{\perp} + D^{2} (g + i Q_{i})] n_{-1} p + \mathcal{O}(p^{3}).$$
(21)

It follows that

$$-y_{-1} = (g + i Q_e) j_0, \tag{22}$$

$$y_1 = (g + i Q_e) j_2 + j_0,$$
 (23)

$$-n_{-1} = -i Q_e j_0, (24)$$

$$n_1 = -\mathrm{i}\,Q_e\,j_2,\tag{25}$$

$$2j_2 = (g + iQ_i)(y_{-1} + n_{-1}), \tag{26}$$

$$0 = i Q_e y_{-1} + (g + i Q_e) n_{-1}, \tag{27}$$

$$2 c_{\beta}^{2} v_{2} = i Q_{e} y_{1} + (g + i Q_{e}) n_{1}$$

+
$$D^{2}(g + iQ_{i})y_{-1} + [P_{\perp} + D^{2}(g + iQ_{i})]n_{-1},$$
 (28)

which gives

$$y_{-1} = (g + i Q_e) a_{-1}, \tag{29}$$

$$y_1 = \left[\frac{1}{2} g \left(g + i Q_e \right) \left(g + i Q_i \right) - 1 \right] a_{-1}, \tag{30}$$

$$n_{-1} = -i Q_e a_{-1}, (31)$$

$$n_1 = -\frac{1}{2} g (i Q_e) (g + i Q_i) a_{-1},$$
(32)

$$j_0 = -a_{-1}, (33)$$

$$j_2 = \frac{1}{2} g (g + i Q_i) a_{-1}, \tag{34}$$

$$v_2 = \frac{\left[-i\,Q_e\,(1+P_\perp) + g\,(g+i\,Q_i)\,D^2\right]}{2\,c_\beta^2}\,a_{-1},\tag{35}$$

where a_{-1} is an arbitrary constant.

C. Second Solution

Suppose that

$$\bar{Y}(p) = y_0 + y_2 p^2 + \mathcal{O}(p^4),$$
 (36)

$$\bar{N}(p) = n_0 + n_2 p^2 + \mathcal{O}(p^4),$$
 (37)

$$\bar{J}(p) = j_3 p^3 + \mathcal{O}(p^5),$$
 (38)

$$\bar{V}(p) = v_3 \, p^3 + \mathcal{O}(p^5). \tag{39}$$

Equations (10)–(13) give

$$2 y_2 p = (g + i Q_e) j_3 p + \mathcal{O}(p^3), \tag{40}$$

$$2 n_2 p = -i Q_e j_3 p + \mathcal{O}(p^3), \tag{41}$$

$$3 j_3 p^2 = (g + i Q_i) (y_0 + n_0) p^2 + \mathcal{O}(p^4), \tag{42}$$

$$3 c_{\beta}^{2} v_{3} p^{2} = i Q_{e} (y_{0} + y_{2} p^{2}) + (g + i Q_{e}) (n_{0} + n_{2} p^{2})$$

$$+ D^{2} (g + i Q_{i}) y_{0} p^{2} + [P_{\perp} + D^{2} (g + i Q_{i})] n_{0} p^{2} + \mathcal{O}(p^{4}).$$

$$(43)$$

It follows that

$$2y_2 = (g + iQ_e)j_3, \tag{44}$$

$$2 n_2 = -i Q_e j_3,$$
 (45)

$$3j_3 = (g + iQ_i)y_0 + (g + iQ_i)n_0, \tag{46}$$

$$0 = i Q_e y_0 + (g + i Q_e) n_0, \tag{47}$$

$$3 c_{\beta}^{2} v_{3} = i Q_{e} y_{2} + (g + i Q_{e}) n_{2}$$

$$+ D^{2} (g + i Q_{e}) n_{0} + [P_{\perp} + D^{2} (g + i Q_{i})] n_{0}.$$
(48)

We get

$$y_0 = (g + i Q_e) a_0, \tag{49}$$

$$y_2 = \frac{1}{6} g (g + i Q_e) (g + i Q_i) a_0,$$
 (50)

$$n_0 = -\mathrm{i}\,Q_e\,a_0,\tag{51}$$

$$n_2 = -\frac{1}{6} g (i Q_e) (g + i Q_i) a_0,$$
 (52)

$$j_3 = \frac{1}{3} g (g + i Q_i) a_0, \tag{53}$$

$$v_3 = \frac{1}{3} \frac{\left[-i \, Q_e \, P_\perp + g \, (g + i \, Q_i) \, D^2 \right]}{c_\beta^2} \, a_0, \tag{54}$$

where a_0 is an arbitrary constant.

D. Third Solution

Suppose that

$$\bar{Y}(p) = y_2 p^2 + \mathcal{O}(p^4),$$
 (55)

$$\bar{N}(p) = n_0 + n_2 p^2 + \mathcal{O}(p^4),$$
 (56)

$$\bar{J}(p) = j_3 p^3 + \mathcal{O}(p^5),$$
 (57)

$$\bar{V}(p) = v_1 p + \mathcal{O}(p^3). \tag{58}$$

Equations (10)–(13) give

$$2y_2 p = (g + i Q_e) j_3 p + \mathcal{O}(p^3), \tag{59}$$

$$2 n_2 p = -i Q_e j_3 p + g v_1 p + \mathcal{O}(p^3), \tag{60}$$

$$3 j_3 p^2 = (g + i Q_i) n_0 p^2 + \mathcal{O}(p^4), \tag{61}$$

$$c_{\beta}^{2} v_{1} = (g + i Q_{e}) n_{0} + \mathcal{O}(p^{2}).$$
 (62)

It follows that

$$y_2 = \frac{1}{6} (g + i Q_e) (g + i Q_i) a_2, \tag{63}$$

$$n_0 = a_2, (64)$$

$$n_2 = \frac{1}{2} (g + i Q_i) \left(-\frac{1}{3} i Q_e + \frac{g}{c_\beta^2} \right) a_2,$$
 (65)

$$j_3 = \frac{1}{3} (g + i Q_i) a_2, \tag{66}$$

$$v_1 = \frac{(g + i Q_e)}{c_\beta^2} a_2,$$
 (67)

where a_2 is an arbitrary constant.

E. Fourth Solution

Suppose that

$$\bar{Y}(p) = y_3 p^3 + \mathcal{O}(p^5),$$
 (68)

$$\bar{N}(p) = n_1 p + \mathcal{O}(p^3), \tag{69}$$

$$\bar{J}(p) = j_4 p^4 + \mathcal{O}(p^6),$$
 (70)

$$\bar{V}(p) = v_0 + v_2 p^2 + \mathcal{O}(p^4). \tag{71}$$

Equations (10)–(13) give

$$3y_3 p^2 = (g + iQ_e) j_4 p^2 + \mathcal{O}(p^4), \tag{72}$$

$$n_1 = q v_0 + \mathcal{O}(p^2),$$
 (73)

$$4 j_4 p^3 = (g + i Q_i) n_1 p^3 + \mathcal{O}(p^5), \tag{74}$$

$$2 c_{\beta}^{2} v_{2} p = (g + i Q_{e}) n_{1} p + \mathcal{O}(p^{3}).$$
(75)

It follows that

$$y_3 = \frac{1}{12} g (g + i Q_e) (g + i Q_i) a_3,$$
 (76)

$$n_1 = g a_3, \tag{77}$$

$$j_4 = \frac{1}{4} g (g + i Q_i) a_3, \tag{78}$$

$$v_2 = \frac{g(g + iQ_e)}{2c_\beta^2} a_3, \tag{79}$$

where a_3 is an arbitrary constant.

F. General Solution

We conclude that, at small-p, the most general solution for $\bar{Y}(p)$ and $\bar{N}(p)$ takes the form

$$\bar{Y}(p) = (g + i Q_e) p^{-1} a_{-1} + (g + i Q_e) a_0 + \mathcal{O}(p),$$
 (80)

$$\bar{N}(p) = (-i Q_e) p^{-1} a_{-1} + (-i Q_e) a_0 + a_2 + \mathcal{O}(p).$$
(81)

III. MATRIX DIFFERENTIAL EQUATION

Let

$$\underline{u} = \begin{pmatrix} \bar{Y} \\ \bar{N} \end{pmatrix}, \tag{82}$$

$$\underline{v} = \begin{pmatrix} \bar{J} \\ c_{\beta}^2 \bar{V} \end{pmatrix}. \tag{83}$$

Equations (10)–(13) can be written in the form

$$\frac{d\underline{u}}{dp} = \underline{\underline{A}}\,\underline{v},\tag{84}$$

$$\frac{d\underline{v}}{dp} = \underline{\underline{B}}\,\underline{u},\tag{85}$$

where

$$A_{11} = \frac{g + i Q_e + p^2}{p^2},\tag{86}$$

$$A_{21} = \frac{-i Q_e}{p^2},\tag{87}$$

$$A_{22} = \frac{g + P_{\varphi} p^2}{c_{\beta}^2},\tag{88}$$

$$B_{11} = (g + i Q_i) p^2 + P_{\varphi} p^4, \tag{89}$$

$$B_{12} = (g + i Q_i) p^2 + \iota_e^{-1} P_{\varphi} p^4, \tag{90}$$

$$B_{21} = i Q_e + D^2 (g + i Q_i) p^2 + D^2 P_{\varphi} p^4, \tag{91}$$

$$B_{22} = g + i Q_e + [P_{\perp} + D^2 (g + i Q_i)] p^2 + \iota_e^{-1} D^2 P_{\varphi} p^4.$$
(92)

Thus, we obtain the following matrix differential equation:

$$\frac{d}{dp}\left(\underline{\underline{A}}^{-1}\frac{d\underline{u}}{dp}\right) = \underline{\underline{B}}\underline{u}.\tag{93}$$

IV. RICCATI MATRIX DIFFERENTIAL EQUATION

Let

$$p\frac{d\underline{u}}{dp} = \underline{\underline{W}}\underline{u}.\tag{94}$$

The previous equation can be combined with Eq. (93) to give

$$\left(p\frac{d\underline{\underline{W}}}{dp} - \underline{\underline{W}} + \underline{\underline{W}}\underline{\underline{W}} + \underline{\underline{A}}p\frac{d\underline{\underline{A}}^{-1}}{dp}\underline{\underline{W}} - p^2\underline{\underline{A}}\underline{\underline{B}}\right)\underline{\underline{u}} = \underline{0},$$
(95)

which implies that

$$p\frac{d\underline{\underline{W}}}{dp} = \underline{\underline{W}} - \underline{\underline{W}}\underline{\underline{W}} - \underline{\underline{A}}p\frac{d\underline{\underline{A}}^{-1}}{dp}\underline{\underline{W}} + p^2\underline{\underline{A}}\underline{\underline{B}}.$$
 (96)

Now,

$$\underline{\underline{A}}^{-1} = \begin{pmatrix} C_{11} & 0 \\ C_{21} & C_{22} \end{pmatrix}, \tag{97}$$

where

$$C_{11} = \frac{p^2}{g + i Q_e + p^2},\tag{98}$$

$$C_{21} = \frac{i c_{\beta}^{2} Q_{e}}{(g + i Q_{e} + p^{2}) (g + P_{\varphi} p^{2})},$$
(99)

$$C_{22} = \frac{c_{\beta}^2}{g + P_{\omega} \, p^2}.\tag{100}$$

So, if

$$p\frac{d\underline{\underline{A}}^{-1}}{dp} = \begin{pmatrix} D_{11} & 0\\ D_{21} & D_{22} \end{pmatrix} \tag{101}$$

then

$$D_{11} = \frac{2p^2(g + iQ_e)}{(g + iQ_e + p^2)^2},$$
(102)

$$D_{21} = -\frac{2 i c_{\beta}^{2} Q_{e} p^{2} [g + P_{\varphi} (g + i Q_{e}) + 2 P_{\varphi} p^{2}]}{(g + i Q_{e} + p^{2})^{2} (g + P_{\varphi} p^{2})^{2}},$$
(103)

$$D_{22} = -\frac{2 c_{\beta}^2 P_{\varphi} p^2}{(g + P_{\varphi} p^2)^2}.$$
 (104)

Furthermore, if

$$\underline{\underline{A}} p \frac{d\underline{\underline{A}}^{-1}}{dp} = \begin{pmatrix} E_{11} & 0 \\ E_{21} & E_{22} \end{pmatrix}$$
 (105)

then

$$E_{11} = \frac{2(g + iQ_e)}{g + iQ_e + p^2},$$
(106)

$$E_{21} = -\frac{2i Q_e (g + 2 P_{\varphi} p^2)}{(g + i Q_e + p^2) (g + P_{\varphi} p^2)},$$
(107)

$$E_{22} = -\frac{2P_{\varphi}p^2}{q + P_{\varphi}p^2}. (108)$$

Finally, if

$$p^2 \underline{\underline{A}} \underline{\underline{B}} = \begin{pmatrix} F_{11} & F_{12} \\ F_{21} & F_{22} \end{pmatrix} \tag{109}$$

then

$$F_{11} = p^2 (g + i Q_e + p^2) (g + i Q_i + P_{\varphi} p^2), \tag{110}$$

$$F_{12} = p^2 (g + i Q_e + p^2) (g + i Q_i + \iota_e^{-1} P_{\varphi} p^2), \tag{111}$$

$$F_{21} = -i Q_e p^2 (g + i Q_i + P_{\varphi} p^2)$$

$$+ c_{\beta}^{-2} p^{2} (g + P_{\varphi} p^{2}) [i Q_{e} + D^{2} (g + i Q_{i}) p^{2} + D^{2} P_{\varphi} p^{4}],$$
(112)

$$F_{22} = -i Q_e p^2 (g + i Q_i + \iota_e^{-1} P_{\varphi} p^2)$$

$$+ c_{\beta}^{-2} p^{2} (g + P_{\varphi} p^{2}) [g + i Q_{e} + [P_{\perp} + \iota_{e}^{-1} D^{2} (g + i Q_{i})] p^{2} + \iota_{e}^{-1} D^{2} P_{\varphi} p^{4}].$$
 (113)

Here, we have modified the expression for F_{22} in order to make the layer equations more stable.

Hence, we obtain the following Riccati matrix differential equation:

$$p\frac{d\underline{\underline{W}}}{dp} = \underline{\underline{W}} - \underline{\underline{W}}\underline{\underline{W}} - \underline{\underline{E}}\underline{\underline{W}} + \underline{\underline{F}}.$$
 (114)

Furthermore, if

$$\underline{\underline{W}} = \begin{pmatrix} W_{11} & W_{12} \\ W_{21} & W_{22} \end{pmatrix} \tag{115}$$

then

$$p\frac{dW_{11}}{dp} = W_{11} - W_{11}W_{11} - W_{12}W_{21} - E_{11}W_{11} + F_{11}, \tag{116}$$

$$p\frac{dW_{12}}{dp} = W_{12} - W_{11}W_{12} - W_{12}W_{22} - E_{11}W_{12} + F_{12}, \tag{117}$$

$$p\frac{dW_{21}}{dp} = W_{21} - W_{21}W_{11} - W_{22}W_{21} - E_{21}W_{11} - E_{22}W_{21} + F_{21},$$
(118)

$$p\frac{dW_{22}}{dp} = W_{22} - W_{21}W_{12} - W_{22}W_{22} - E_{21}W_{12} - E_{22}W_{22} + F_{22}.$$
 (119)

V. SMALL-p BEHAVIOR OF RICCATI MATRIX DIFFERENTIAL EQUATION

At p=0, let $\underline{\underline{E}}=\underline{\underline{E}}^{(0)}$, where the elements of $\underline{\underline{E}}^{(0)}$ are constants. It follows from Eqs. (105)–(108) that

$$E_{11}^{(0)} = 2, (120)$$

$$E_{12}^{(0)} = 0, (121)$$

$$E_{21}^{(0)} = -\frac{2i\,Q_e}{g + i\,Q_e},\tag{122}$$

$$E_{22}^{(0)} = 0. (123)$$

Likewise, at small p, we can write $\underline{\underline{F}} = p^2 \underline{\underline{F}}^{(2)}$, where the elements of $\underline{\underline{F}}^{(2)}$ are constants, where use has been made of Eqs. (110)–(113).

Suppose that $\underline{\underline{W}} = \underline{\underline{W}}^{(0)}$ at p = 0, where all of the elements of $\underline{\underline{W}}^{(0)}$ are constants. To lowest-order in p, Eq. (114) gives

$$\underline{\underline{0}} = \underline{\underline{W}}^{(0)} - \underline{\underline{W}}^{(0)} \underline{\underline{W}}^{(0)} - \underline{\underline{E}}^{(0)} \underline{\underline{W}}^{(0)}, \tag{124}$$

or

$$(\underline{\underline{1}} - \underline{\underline{W}}^{(0)} - \underline{\underline{E}}^{(0)}) \underline{\underline{W}}^{(0)} = \underline{\underline{0}}. \tag{125}$$

Hence, we deduce that

$$\underline{\underline{W}}^{(0)} = \underline{\underline{1}} - \underline{\underline{E}}^{(0)} = \begin{pmatrix} -1 & 0 \\ -E_{21}^{(0)} & 1 \end{pmatrix}. \tag{126}$$

At small-p, let

$$\underline{u}(p) = \underline{u}_{-1} \, p^{-1} + \underline{u}_0, \tag{127}$$

$$\underline{\underline{W}}(p) = \underline{\underline{W}}^{(0)} + \underline{\underline{W}}^{(1)} p, \tag{128}$$

where the elements of \underline{u}_{-1} , \underline{u}_{0} , and $\underline{\underline{W}}^{(1)}$ are constants. Equation (94) gives

$$\underline{W}^{(0)}\,\underline{u}_{-1} = -\underline{u}_{-1},\tag{129}$$

$$\underline{\underline{W}}^{(0)}\underline{u}_0 + \underline{\underline{W}}^{(1)}\underline{u}_{-1} = \underline{0}. \tag{130}$$

Thus, making use of Eq. (126), we get

$$\begin{pmatrix} -1 & 0 \\ -E_{21}^{(0)} & 1 \end{pmatrix} \begin{pmatrix} y_{-1} \\ n_{-1} \end{pmatrix} = -\begin{pmatrix} y_{-1} \\ n_{-1} \end{pmatrix}, \tag{131}$$

which implies that

$$E_{21}^{(0)} y_{-1} = -\frac{2 i Q_e}{g + i Q_e} y_{-1} = 2 n_{-1}, \tag{132}$$

in accordance with Eqs. (29) and (31), where use has been made of Eq. (122). Thus, if we write

$$y_{-1} = (g + i Q_e) a_{-1}, \tag{133}$$

$$n_{-1} = -i Q_e a_{-1}, \tag{134}$$

$$y_0 = (g + i Q_e) a_0, (135)$$

$$n_0 = -i Q_e a_0 + a_2, \tag{136}$$

in accordance with Eqs. (80) and (81), then we deduce from Eqs. (126) and (130) that

$$\frac{\pi}{\hat{\Delta}_s} \equiv \frac{a_0}{a_{-1}} = W_{11}^{(1)} - W_{12}^{(1)} \frac{(i Q_e)}{g + i Q_e}.$$
 (137)

VI. LARGE-p BEHAVIOR OF RICCATI MATRIX DIFFERENTIAL EQUATION

In the large-p limit, it is clear from Eqs. (110)–(113) that $\underline{\underline{F}}(p) = p^6 \underline{\underline{F}}^{(6)} + p^8 \underline{\underline{F}}^{(8)}$, where the elements of $\underline{\underline{F}}^{(6)}$ and $\underline{\underline{F}}^{(8)}$ are constants. On the other hand, Eqs. (105)–(108) imply that $\underline{\underline{E}}(p) = \underline{\underline{F}}^{(0)}$, where the elements of $\underline{\underline{F}}^{(0)}$ are constants. Thus, if we write $\underline{\underline{W}}(p) = p^2 \underline{\underline{W}}^{(2)} + p^4 \underline{\underline{W}}^{(4)}$, where the elements of $\underline{\underline{W}}^{(2)}$ and $\underline{\underline{W}}^{(4)}$ are constants, then Eq. (114) gives

$$\underline{\underline{W}}^{(4)}\underline{\underline{W}}^{(4)} = \underline{\underline{F}}^{(8)},\tag{138}$$

$$\underline{\underline{W}}^{(2)}\underline{\underline{W}}^{(4)} + \underline{\underline{W}}^{(4)}\underline{\underline{W}}^{(2)} = \underline{\underline{F}}^{(6)}.$$
(139)

Now, according to Eqs. (110)–(113),

$$F_{11}^{(8)} = 0, (140)$$

$$F_{12}^{(8)} = 0, (141)$$

$$F_{21}^{(8)} = c_{\beta}^{-2} D^2 P_{\varphi}^2, \tag{142}$$

$$F_{22}^{(8)} = \iota_e^{-1} \, c_\beta^{-2} \, D^2 \, P_\varphi^2, \tag{143}$$

so Eq. (138) yields

$$W_{11}^{(4)} = 0, (144)$$

$$W_{12}^{(4)} = 0, (145)$$

$$W_{21}^{(4)} = -\iota_e^{1/2} \, c_\beta^{-1} \, D \, P_\varphi, \tag{146}$$

$$W_{22}^{(4)} = -\iota_e^{-1/2} c_\beta^{-1} D P_\varphi, \tag{147}$$

where we have chosen the sign of the square root that is associated with well-behaved solutions at large-p. We are also assuming at $\iota_e > 0$. Equations (110)–(113) also give

$$F_{11}^{(6)} = P_{\varphi},$$
 (148)

$$F_{12}^{(6)} = \iota_e^{-1} P_{\varphi}, \tag{149}$$

$$F_{21}^{(6)} = c_{\beta}^{-2} D^2 g P_{\varphi} + c_{\beta}^{-2} D^2 (g + i Q_i) P_{\varphi},$$
(150)

$$F_{22}^{(6)} = c_{\beta}^{-2} \iota_e^{-1} D^2 g P_{\varphi} + c_{\beta}^{-2} [P_{\perp} + \iota_e^{-1} D^2 (g + i Q_i)] P_{\varphi}.$$
 (151)

Thus, Eq. (139) gives

$$W_{12}^{(2)} W_{21}^{(4)} = F_{11}^{(6)}, (152)$$

$$W_{12}^{(2)} W_{22}^{(4)} = F_{12}^{(6)}. (153)$$

Thus, we obtain

$$W_{12}^{(2)} = -\iota_e^{-1/2} c_\beta D^{-1}. {154}$$

Now, if

$$\underline{W}\,\underline{u} = \lambda\,\underline{u} \tag{155}$$

then Eq. (94) gives

$$p\frac{d\underline{u}}{dp} = \lambda \,\underline{u}.\tag{156}$$

Of course, for a well behaved solution, λ must be negative at large-p. As we have seen, the large-p limit of Eq. (114) is

$$\underline{\underline{W}}\,\underline{\underline{W}} = \underline{\underline{F}}.\tag{157}$$

Hence, if

$$\underline{\underline{F}}\,\underline{u} = \Lambda\,\underline{u} \tag{158}$$

then

$$\lambda^2 = \Lambda. \tag{159}$$

The eigenvalue problem for the F-matrix reduces to

$$\Lambda^2 - (F_{11} + F_{22}) \Lambda + F_{11} F_{22} - F_{12} F_{21} = 0.$$
(160)

Now,

$$F_{11} + F_{22} \simeq F_{22}^{(8)} p^8 = \iota_e^{-1} c_\beta^{-2} D^2 P_\omega^2 p^8,$$
 (161)

$$F_{11} F_{22} - F_{12} F_{21} \simeq \left[F_{11}^{(6)} F_{22}^{(6)} - F_{12}^{(6)} F_{21}^{(6)} \right] p^{12} = c_{\beta}^{-2} P_{\perp} P_{\varphi}^{2} p^{12},$$
 (162)

SO

$$\Lambda_1 = \iota_e^{-1} \, c_\beta^{-2} \, D^2 \, P_\varphi^2 \, p^8, \tag{163}$$

$$\Lambda_2 = \iota_e \, P_\perp \, D^{-2} \, p^4. \tag{164}$$

Assuming that $\iota_e > 0$, we deduce that

$$\lambda_1 = -\iota_e^{-1} \, c_\beta^{-1} \, D \, P_\varphi \, p^4, \tag{165}$$

$$\lambda_2 = -\iota_e^{1/2} P_{\perp}^{1/2} D^{-1} p^2. \tag{166}$$

Now, the eigenvalue problem for the W-matrix reduces to

$$\lambda^{2} - W_{22}^{(4)} p^{4} \lambda + (W_{11}^{(2)} W_{22}^{(4)} - W_{12}^{(2)} W_{21}^{(4)}) p^{6} = 0.$$
 (167)

So, we need

$$\lambda_1 = W_{22}^{(4)} \, p^4, \tag{168}$$

which is satisfied, and

$$\lambda_2 = \left[W_{11}^{(2)} - \frac{W_{12}^{(2)} W_{21}^{(4)}}{W_{22}^{(4)}} \right] p^2, \tag{169}$$

which implies that

$$W_{11}^{(2)} = -\iota_e^{1/2} P_\perp^{1/2} D^{-1} - \iota_e^{-1/2} c_\beta D^{-1}.$$
(170)

Hence, the large-p boundary condition for the W-matrix is

$$\underline{\underline{W}}(p) = \begin{pmatrix}
-\iota_e^{1/2} P_{\perp}^{1/2} D^{-1} p^2 - \iota_e^{-1/2} c_{\beta} D^{-1} p^2, & -\iota_e^{-1/2} c_{\beta} D^{-1} p^2 \\
-\iota_e^{1/2} c_{\beta}^{-1} D P_{\varphi} p^4, & -\iota_e^{-1/2} c_{\beta}^{-1} D P_{\varphi} p^4
\end{pmatrix}.$$
(171)