Construction of EFIT File

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The EFIT equilibrium magnetic field is written

$$\mathbf{B} = T(\Psi) \, \nabla \phi + \nabla \phi \times \nabla \Psi, \tag{1}$$

whereas the TJ equilibrium field is written

$$\mathbf{B} = B_0 R_0 [g(r) \nabla \phi + R_0 f(r) \nabla \phi \times \nabla r]. \tag{2}$$

So,

$$T = B_0 R_0 g, \tag{3}$$

$$\frac{d\Psi}{dr} = B_0 R_0^2 f. (4)$$

Here, B_0 and R_0 are the toroidal magnetic field strength at the magnetic axis, and the major radius of the magnetic axis, respectively. Let $\Psi = 0$ at the plasma boundary, $r = \epsilon$. It follows that

$$\Psi(r) = B_0 R_0^2 \int_r^{\epsilon} f(r') dr'.$$
 (5)

Also,

$$q = \frac{r g}{f}. (6)$$

Let $\hat{r} = r/\epsilon$. It follows that

$$\frac{d\Psi}{d\hat{r}} = \epsilon B_0 R_0^2 f, \tag{7}$$

and

$$\Psi(\hat{r}) = \epsilon B_0 R_0^2 \int_{\hat{r}}^1 f(\hat{r}') d\hat{r}'.$$
 (8)

Furthermore,

$$T\frac{dT}{d\Psi} = B_0 \frac{g}{\epsilon f} \frac{dg}{d\hat{r}},\tag{9}$$

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and

$$\frac{dP}{d\Psi} = \frac{B_0}{\mu_0 R_0^2} \frac{1}{\epsilon f} \frac{dP}{d\hat{r}}.$$
 (10)

Let $\hat{r} f = \epsilon f_1 + \epsilon^3 f_3$. It follows that

$$\Psi(\hat{r}) = \epsilon^2 B_0 R_0^2 \int_{\hat{r}}^1 \frac{f_1(\hat{r}') + \epsilon^2 f_3(\hat{r}')}{\hat{r}'} d\hat{r}', \tag{11}$$

$$q(\Psi) = \frac{\hat{r}^2 (1 + \epsilon^2 g_2)}{f_1 + \epsilon^2 f_3},\tag{12}$$

$$T(\Psi) = B_0 R_0 \left(1 + \epsilon^2 g_2 \right), \tag{13}$$

$$p(\Psi) = \frac{B_0^2}{\mu_0} \epsilon^2 p_2(\hat{r}), \tag{14}$$

$$T\frac{dT}{d\Psi} = B_0 \frac{1}{f_1 + \epsilon^2 f_3} \hat{r} \frac{dg_2}{d\hat{r}} \left(1 + \epsilon^2 g_2 \right), \tag{15}$$

$$\frac{dP}{d\Psi} = \frac{B_0}{\mu_0 R_0^2} \frac{1}{f_1 + \epsilon^2 f_3} \hat{r} \frac{dp_2}{d\hat{r}}.$$
 (16)

At small \hat{r} , $g_2 = 0$, $f_1 = f_{1c}\hat{r}^2$, $f_3 = f_{3c}\hat{r}^2$, $\hat{r}p'_2 = p''_{2c}\hat{r}^2$, $\hat{r}g'_2 = g'_{2c}\hat{r}^2$, $\Psi = \Psi(0) + \epsilon^2 B_0 R_0^2 (1/2) (f_{1c} + \epsilon^2 f_{3c}) \hat{r}^2$, where $g'_{2c} = -2 (f_{1c}^2 + p''_{2c}/2)$, and $f_{3c} = -f_{1c} (\hat{H}_{2c}^2 + \hat{V}_{2c}^2)$. For $\hat{r} > 1$,

$$\Psi(\hat{r}) = \epsilon^2 B_0 R_0^2 \ln \hat{r} [f_1(1) + \epsilon^2 f_3(1)]. \tag{17}$$

TJ flux surfaces are parameterized via

$$R(\hat{r},\omega) = 1 - \epsilon \,\hat{r} \,\cos\omega + f_R(\hat{r},\omega),\tag{18}$$

$$Z(\hat{r},\omega) = \epsilon \,\hat{r} \,\sin\omega + f_Z(\hat{r},\omega),\tag{19}$$

where

$$f_R(\hat{r},\omega) = \epsilon^2 \sum_{j>0} H_j(\hat{r}) \cos[(j-1)\omega] + \epsilon^2 \sum_{j>1} V_j(\hat{r}) \sin[(j-1)\omega]$$
$$+ \epsilon^3 \left[L_3(\hat{r}) + \epsilon L_4(\hat{r}) \right] \cos \omega,$$
(20)

$$f_Z(\hat{r},\omega) = \epsilon^2 \sum_{j>1} H_j(\hat{r}) \sin[(j-1)\omega] - \epsilon^2 \sum_{j>1} V_j(\hat{r}) \cos[(j-1)\omega]$$

$$-\epsilon^{3} \left[L_{3}(\hat{r}) + \epsilon L_{4}(\hat{r}) \right] \sin \omega. \tag{21}$$

For $\hat{r} > 1$, let $L_3(\hat{r}) = L_3(1)$, $L_4(\hat{r}) = L_4(1)$, $H_j(\hat{r}) = H_j(1)$, and $V_j(\hat{r}) = V_j(1)$. To obtain the \hat{r} , ω coordinates that correspond to the point R, Z, our initial guess is

$$\hat{r} = \frac{[(R-1)^2 + Z^2]^{1/2}}{\epsilon},\tag{22}$$

$$\frac{\sin \omega}{\cos \omega} = \frac{Z}{1 - R}.\tag{23}$$

We can then iterate the following equations

$$\hat{r} = \frac{[(\hat{R} - 1)^2 + \hat{Z}^2]^{1/2}}{\epsilon},\tag{24}$$

$$\frac{\sin \omega}{\cos \omega} = \frac{\hat{Z}}{1 - \hat{R}},\tag{25}$$

where

$$\hat{R} = R - f_R(\hat{r}, \omega), \tag{26}$$

$$\hat{Z} = Z - f_Z(\hat{r}, \omega). \tag{27}$$

EFIT integer parameters are NRBOX, NZBOX, NPBOUND, NLIMITER. EFIT float parameters are RBOXLEN, ZBOXLEN, RBOXLFT, ZOFF. Also, R0EXP = R_0 , B0EXP = R_0 , RAXIS = R_0 , ZAXIS = 0, PSIAXIS = R_0 0, PSIBOUND = 0, CURRENT= R_0 1. The R_0 2 grid-points are

$$RR_0 = RBOXLFT + RBOXLEN * i/(NRBOX - 1),$$
 (28)

$$ZZ_0 = \text{ZOFF} - \text{ZBOXLEN}/2 + \text{ZBOXLEN} * j/(\text{NZBOX} - 1),$$
 (29)

for i = 0, NBBOX -1, j = 0, NZBOX -1. The profiles are evaluated on the grid PSI = PSIAXIS * (NRBOX -1 - i)/(NRBOX -1) for i = 0, NRBOX -1.