

Ryan George July 14, 2010

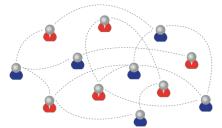
The defining feature of neural systems is plasticity.



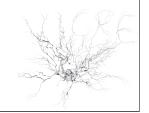
The defining feature of neural systems is plasticity.

Network

Neurotransmitters sent at synapses



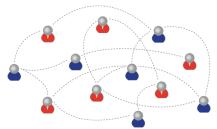
synaptic plasticity



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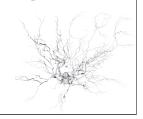
synaptic plasticity

Individual

Current passed along membrane



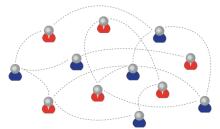
excitability



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Neurotransmitters sent at synapses



synaptic plasticity

Internal

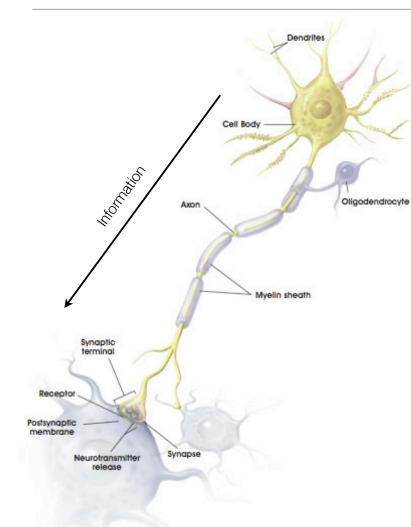
Current passed along membrane



excitability



Neuronal Structure



Our interest: dendrite ion channels, "conductance"

Conductance affects excitability

Non-uniform distribution

Voltage Clamp (Hoffman et al.)

Electron Microscopy (Lorincz et al.)

We want a non-invasive method for recovering the distribution of ion channels along a dendrite.



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"non-invasive": only stimulating and recording voltage from the soma

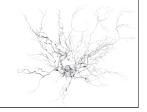


Overview

1. The Cable Equation

2. Dirichlet Series and Conductance Recovery

3. Expansion to trees



Modeling the Dendrite







Modeling the Dendrite

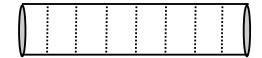






Modeling the Dendrite

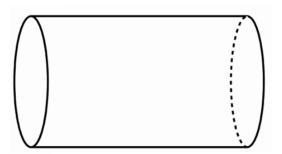




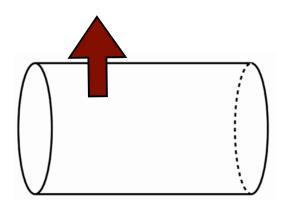
Transmembrane potential v(x,t)



Modeling the Dendrite: Current Balance

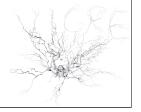


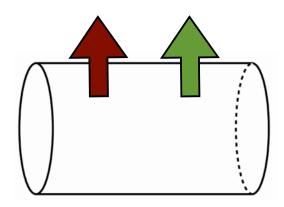




$$C_m \frac{\partial v}{\partial t}(x,t)$$

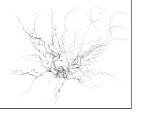
capacitive

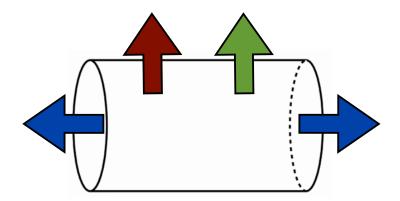




$$C_m \frac{\partial v}{\partial t}(x,t) + g_L(x)v(x,t)$$

capacitive leakage current current



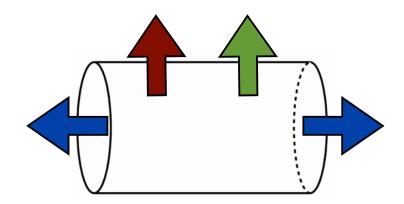


$$C_m \frac{\partial v}{\partial t}(x,t) + g_L(x)v(x,t) - \frac{a}{2R_a} \frac{\partial^2 v}{\partial x^2}(x,t) = 0$$

capacitive

leakage current axial current





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The Cable Equation

$$C_m \frac{\partial v}{\partial t}(x,t) + g_L(x)v(x,t) - \frac{a}{2R_a} \frac{\partial^2 v}{\partial x^2}(x,t) = 0$$

Boundary Conditions:

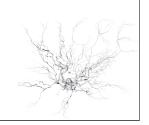
$$\frac{\partial v}{\partial x}(0,t) = -\frac{R_a}{\pi a^2} I_{\text{stim}}(t)$$

$$\frac{\partial v}{\partial x}(\ell, t) = 0$$

$$v(0,t) = v_{\rm obs}(t)$$

Assume C_m , R_a , a known.

Seek $g_L(x)$.



The Eigenproblem

Separation of Variables: $v(x,t) = \sum_{n=0}^{\infty} q_n(x)p_n(t)$

The space-dependent part must satisfy:

$$\frac{a}{2R_a}q_n''(x) = q_n(x) \left(\vartheta_n + g_L(x)\right)$$
$$q_n'(0) = q_n'(\ell) = 0$$

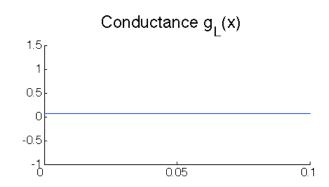
These are orthogonal in the sense:

$$\int_0^\ell q_m(x)q_n(x)\mathrm{d}x = \delta_{mn}$$



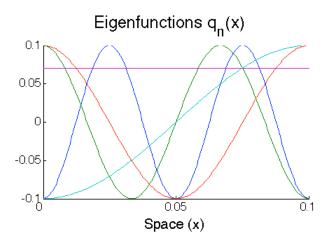
The Eigenfunctions

When
$$g_L(x) = g_0$$
,



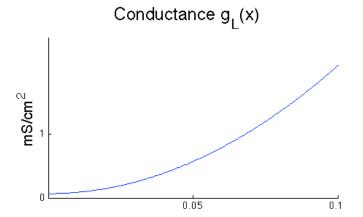
$$\vartheta_n = -\frac{an^2\pi^2}{2R_a\ell^2} - g_0$$

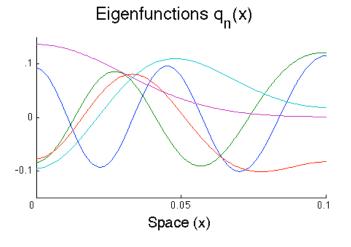
$$q_n(x) = \cos(n\pi x/\ell)$$

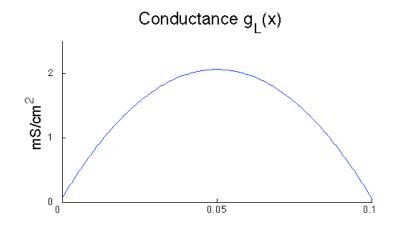


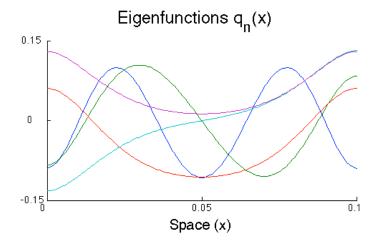


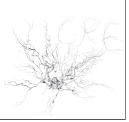
Examples of Eigenfunctions











Recall:
$$v(x,t) = \sum_{n=0}^{\infty} q_n(x)p_n(t)$$



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Use orthogonality of q:

$$\int_0^l v(x,t)q_n(x)dx = p_n(t)\left(\int_0^l q_n(x)q_n(x)dx\right) = p_n(t)$$



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Differentiate and Substitute:

$$p'_n(t) = \int_0^l \frac{\partial v}{\partial t}(x, t)q_n(x)dx$$

$$= \int_0^l \frac{1}{C_m} \left(\frac{a}{2R_a} \frac{\partial^2 v}{\partial x^2} - g_L(x)v \right) q_n(x)dx$$



Integrate by parts (not shown):

$$\int_0^{l_i} \frac{\partial^2 v}{\partial x^2}(x,t) q_n(x) dx = q_n(0) R_2 I_{stim}(t) / (\pi a^2) + \int_0^l q_n''(x) v(x,t) dx$$

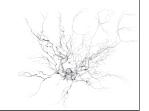


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Use the identity:

$$\frac{a_i}{2R_a} \int_0^l q_n(x) q_m''(x) dx - \int_0^l q_n(x) q_m(x) g_L(x) dx - \theta_n \int_0^l q_n(x) q_m(x) dx = 0$$



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Arrive at:

$$C_m p_t'(t) - \vartheta_n p(t) = q_n(0) I_{stim}(t) / (2\pi a)$$



General Solution

Solution to the time-dependent portion:

$$p_n(t) = \frac{q_n(0)}{2\pi a C_m} \int_0^t I_{stim}(s) e^{(t-s)\vartheta_n} ds$$



General Solution

Solution to the time-dependent portion:

$$p_n(t) = \frac{q_n(0)}{2\pi a C_m} \int_0^t I_{stim}(s) e^{(t-s)\vartheta_n} ds$$

Solution of cable equation:

$$v(x,t) = \sum_{n=0}^{\infty} q_n(x)p_n(t) = \sum_{n=0}^{\infty} \frac{q_n(0)q_n(x)}{2C_m a\pi} \int_0^t I_{stim}(s)e^{(t-s)\theta_n} ds$$

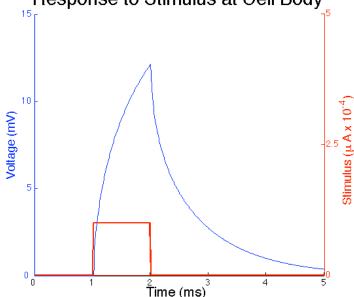


Dirichlet Series

Set x = 0. This gives us voltage at the cell body:

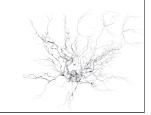
$$v(0,t) = \sum_{n=0}^{\infty} \frac{q_n(0)^2}{2C_m a\pi} \int_0^t I_{stim}(s) e^{(t-s)\theta_n} ds$$

Response to Stimulus at Cell Body



This solution is a Dirichlet series dependent on $\vartheta_n, q_n(0)^2$

(eigenvalues, norming constants)



Dirichlet Series

A function has a unique Dirichlet Series representation.

If two representations:
$$\sum_{n=1}^{\infty}a_ne^{\vartheta_nt}=v(0,t)=\sum_{n=1}^{\infty}b_ne^{\vartheta_nt}$$

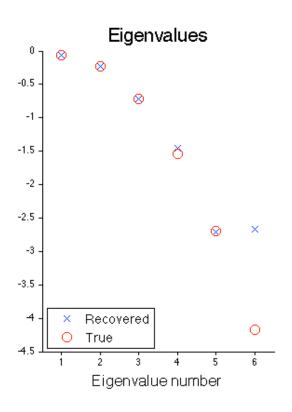
Then
$$\sum_{n=1}^{\infty} (b_n - a_n)e^{\vartheta_n t} = 0 \quad \text{for all } t$$

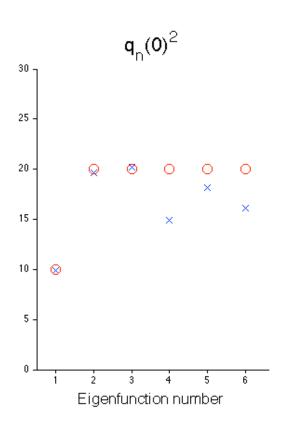
Thus for given v(0,t) there is exactly one set $\{\vartheta_n,q_n(0)^2\}$.



Numerical Recovery of Eigenvalues

'Peeling' - Fit $\{\vartheta_n, q_n(0)^2\}$ to numerical solution.







Eigenvalues to Coefficient

Inverse Sturm-Liouville Problem

Borg and Levinson's Theorem:

$$\{\vartheta_n, q_n(0)^2\} \iff g_L(x)$$

We recover conductance from eigenvalues and norming constants.



Method

1. Record (simulate) from soma

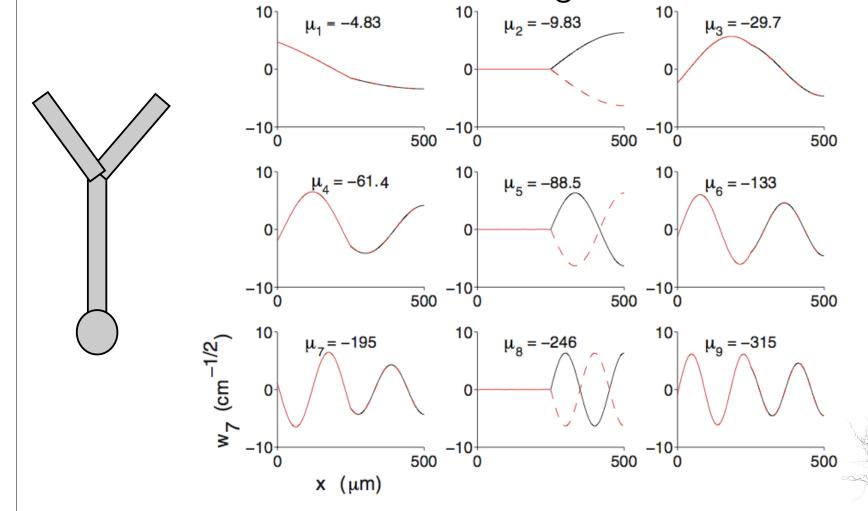
2. Find eigenvalues, norming constants

3. Recover Conductance

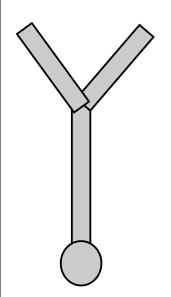


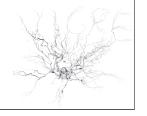
Expansion to Trees

Branched Eigenfunctions

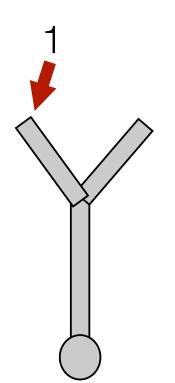


Expansion to Trees





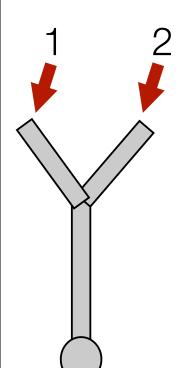
Expansion to Trees



$$v_1(0,t) = \sum_{n=0}^{\infty} \frac{q_{n,1}(0)q_{n,3}(\ell_3)}{A_s h} \int_0^t I_{\text{stim}}(s) \exp((t-s)(\vartheta_n)) ds$$

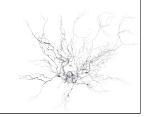


Expansion to Trees

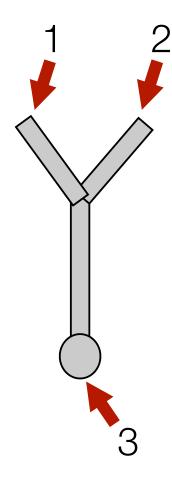


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$$v_2(0,t) = \sum_{n=0}^{\infty} \frac{q_{n,2}(0)q_{n,3}(\ell_3)}{A_s h} \int_0^t I_{\text{stim}}(s) \exp((t-s)(\vartheta_n)) ds$$



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$$v_3(0,t) = \sum_{n=0}^{\infty} \frac{q_{n,3}(\ell_3)^2}{A_s h} \int_0^t I_{\text{stim}}(s) \exp((t-s)(\vartheta_n)) ds$$



Conclusion

Cable:

Tree:

1. Record (simulate) from soma and branches

2. Find eigenvalues,

eigenfunctions at 0



$$\{\vartheta_n, q_n(0)^2\}$$



3. Recover Conductance

 $g_L(x)$

$$v_i(0,t), \ i=1,...,N$$

$$\{\theta_n, q_{n,1}, ..., q_{n,N}\}$$

$$g_L^i(x), i = 1, ..., N$$



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Thank you for your time.

