TOY MODELS IN THEORETICAL PHYSICS

Raghav G. Jha, 2025, version 0.2

LAST EDITED: JUNE 20, 2025 https://rgjha.github.io

ISING MODEL WITH MIXED FIELDS

$$\mathcal{H} = -J\sum_{i} \sigma_{i}^{z} \sigma_{i+1}^{z} - h\sum_{i} \sigma_{i}^{x} - g\sum_{i} \sigma_{i}^{z}$$

$$\tag{1}$$

q-state Potts model (classical)

$$\mathcal{H} = -J \sum_{\langle i,j \rangle} \delta_{\sigma_i,\sigma_j}, \quad \sigma_i \in 1, \cdots, q$$
 (2)

SHERRINGTON-KIRPATRICK MODEL (SK)

$$\mathcal{H} = -\sum_{i < j < N} J_{ij} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x, \quad J_{ij} \in \mathcal{N}(0, J^2)$$
(3)

KOGUT-SUSSKIND HAMILTONIAN (PURE GAUGE)

$$\mathcal{H} = \frac{g^2}{2} \sum_{\ell} E_{\ell}^2 - \frac{1}{2g^2} \sum_{\square} \text{Tr}(U_{\square} + U_{\square}^{\dagger}) \tag{4}$$

AKLT (AFFLECK-KENNEDY-LIEB-TASAKI) HAMILTONIAN

$$\mathcal{H} = \sum_{\langle i,j \rangle} \left(\vec{S}_i \cdot \vec{S}_j + \frac{1}{3} \left(\vec{S}_i \cdot \vec{S}_j \right)^2 \right) \tag{5}$$

where $\vec{S} = (S_1, S_2, S_3)$

$$S_{1} = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad S_{2} = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad S_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$
 (6)

Notes: It has the famous Haldane gap.

 $\bigstar \mathcal{N} = 4 \text{ SYM} - (3+1)\text{-dimensions}$

$$\mathcal{L} = \text{Tr}\left[F^{\mu\nu}F_{\mu\nu} + (D_{\mu}X_{i})^{2} - \frac{1}{2}[X_{i}, X_{j}]^{2} + \Psi^{T}D\!\!\!/\Psi + \Psi^{T}\gamma_{i}[X, \Psi]\right]$$
(7)

BFSS (0+1)-dimensions

$$\mathcal{L} = \text{Tr}\left((D_t X_i)^2 - [X_i, X_j]^2 \right) + \Psi^T D \Psi + \Psi^T \gamma_i [X, \Psi]$$
(8)

 $\overline{\mathrm{BMN/PWMM}}$ - (0+1)-dimensions with i,j,k $=1\cdots 3$ and $M=4\cdots 9$

$$\mathcal{L} = \mathcal{L}_{BFSS} + Tr \left[\left(\frac{\mu}{3} X_I \right)^2 + \left(\frac{\mu}{6} X_M \right)^2 + \frac{\mu}{4} \Psi_{\alpha}^T \gamma_{\alpha\beta}^{123} \Psi_{\beta} + \frac{\sqrt{2}\mu}{3} \epsilon_{IJK} X_I X_J X_K \right]. \tag{9}$$

SYK in (0+1)

$$\mathcal{H} = \frac{1}{4!} \sum_{i,j,k,l=1}^{N} J_{ijkl} \chi_i \chi_j \chi_k \chi_l$$
 (10)

IKKT (0+0)-dimensions, with $i,j = 1 \cdots 10$

$$\mathcal{L} = \text{Tr}([X_I, X_J]^2) + \Psi^T D \Psi + \Psi^T \gamma_I [X, \Psi]$$
(11)

EINSTEIN-HILBERT ACTION

$$S = \frac{c^4}{16\pi G} \int d^4x \left(R - 2\Lambda\right) \tag{12}$$

NAMBU-GOTO (NG) ACTION

$$S = -T \int d^2 \sigma \sqrt{-(\dot{X}^2)(X')^2 + (\dot{X} \cdot X')^2}$$
 (13)

 $\dot{X}^{\mu} = \partial X^{\mu}/\partial \tau, (X')^{\mu} = \partial X^{\mu}/\partial \sigma$

POLYAKOV ACTION

$$S = -\frac{1}{4\pi\alpha'} \int d^2\sigma \sqrt{g} g^{\alpha\beta} \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu} \eta_{\mu\nu}$$
 (14)

CHERN-SIMONS ACTION

$$S = \frac{k}{4\pi} \int d^3x \, \epsilon^{\mu\nu\rho} \, \text{Tr} \left(A_{\mu} \partial_{\nu} A_{\rho} - \frac{2i}{3} A_{\mu} A_{\nu} A_{\rho} \right) \tag{15}$$

PRINCIPAL CHIRAL FIELD

$$\mathcal{L} = \frac{\beta}{2} \operatorname{Tr} \left(\partial_{\mu} g^{-1} \partial_{\mu} g \right) \text{ where } g \in SU(N)$$
 (16)

Massless Schwinger (1+1)

$$\mathcal{L} = \frac{1}{2} (\epsilon^{\mu\nu} \partial_{\nu} A_{\mu})^2 - ej^{\mu} A_{\mu} + \overline{\Psi} \partial \Psi$$
 (17)

▲ Massive Thirring model

$$\mathcal{L} = \overline{\Psi}i\partial \Psi - m_F \overline{\Psi}\Psi - \frac{g}{2}(\overline{\Psi}\gamma^{\mu}\Psi)^2; \quad \partial = \gamma^{\mu}\partial_{\mu}$$
 (18)

CLASSICAL XY MODEL

$$\mathcal{H} = -J\sum_{\langle ij\rangle}\cos(\theta_i - \theta_j) - h\sum_i\cos\theta_i, \quad \theta \in [0, 2\pi)$$
(19)

GENERALIZED XY MODEL (FRACTIONAL VORTICES)

$$\mathcal{H} = -J \sum_{\langle ij \rangle} \cos(q(\theta_i - \theta_j)), \quad \theta \in [0, 2\pi), \quad q > 2$$
(20)

O(N) non-linear σ in 1+1

$$\mathcal{L} = \frac{1}{2g} \sum_{i=1}^{N} (\partial^{\mu} \hat{n}_i)^2 \tag{21}$$

▲ SINE-GORDON

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - \frac{m^4}{\lambda} \left(1 - \cos \left(\frac{\sqrt{\lambda} \phi}{m} \right) \right)$$
 (22)

★ Heisenberg model [1928, Solved by Bethe (1931)]

$$\mathcal{H}_{XXX} = \frac{J}{2} \sum_{i=1}^{N-1} \left(\sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y + \sigma_i^z \sigma_{i+1}^z \right)$$
 (23)

Majumdar-Ghosh model

$$\mathcal{H} = J_1 \sum_{i=1}^{N-1} \vec{S}_i \cdot \vec{S}_{i+1} + J_2 \sum_{i=1}^{N-2} \vec{S}_i \cdot \vec{S}_{i+2}$$
(24)

FERMI-HUBBARD MODEL

$$\mathcal{H} = -t \sum_{\langle i,j \rangle, \sigma} \left(c_{i\sigma}^{\dagger} c_{j\sigma} + c_{j\sigma}^{\dagger} c_{i\sigma} \right) + U \sum_{i} n_{i\uparrow} n_{i\downarrow}, \quad ; \quad \{ c_{i\alpha}, c_{j\beta}^{\dagger} \} = \delta_{ij} \delta_{\alpha\beta}, \quad ; \alpha, \beta = \uparrow, \downarrow$$
 (25)

Bose-Hubbard Model

$$\mathcal{H} = -t \sum_{\langle i,j \rangle} \left(b_i^{\dagger} b_j + b_j^{\dagger} b_i \right) + \frac{U}{2} \sum_i n_i (n_i - 1) - \mu \sum_i n_i \quad ; \quad [b_i, b_j^{\dagger}] = \delta_{ij}$$
 (26)

Blume-Capel model

$$\mathcal{H} = -J\sum_{\langle i,j\rangle} Z_i Z_j + \Delta \sum_i Z_i^2 - h \sum_i Z_i \quad ; \quad Z = \text{Pauli} Z = \text{diag}(1, -1)$$
 (27)

O'Brien-Fendley model (tri-critical Ising)

$$\mathcal{H} = -\sum_{j=1}^{N-1} Z_j Z_{j+1} - g \sum_{j=1}^{N} X_j - h \sum_{j=1}^{N} Z_j + \underbrace{\lambda \sum_{j=1}^{N-2} (X_j Z_{j+1} Z_{j+2} + Z_j Z_{j+1} X_{j+2})}_{\text{Staggered term}}$$
(28)

Nambu-Jona-Lasinio model

$$\mathcal{L} = \bar{\Psi} i \partial \!\!\!/ \Psi + \frac{\lambda}{4} \left[(\bar{\Psi} \Psi)(\bar{\Psi} \Psi) - (\bar{\Psi} \gamma^5 \Psi)(\bar{\Psi} \gamma^5 \Psi) \right] \quad ; \quad \bar{\Psi} = \Psi^{\dagger} \gamma^0$$
 (29)

GROSS-NEVEU MODEL

$$\mathcal{L} = \bar{\Psi}_k (i\partial \!\!\!/ + m) \Psi_k - \frac{g}{2} (\bar{\Psi}_k \Psi_k)^2 \quad ; \quad k = 1, \cdots, N(\text{flavors})$$
 (30)

SU-SCHRIEFFER-HEEGER (SSH) MODEL

$$\mathcal{H} = \sum_{i} \left[t_1 (c_{Ai}^{\dagger} c_{Bi} + \text{H.c}) - t_2 (c_{Bi}^{\dagger} c_{A,i+1} + \text{H.c}) \right] \quad ; A, B \equiv \text{flavors}, \tag{31}$$

KITAEV CHAIN MODEL

$$\mathcal{H} = \sum_{j} \left[-t(c_{j}^{\dagger}c_{j+1} + c_{j+1}^{\dagger}c_{j}) - \mu c_{j}^{\dagger}c_{j} + \underbrace{\Delta}_{\text{p-wave SC pairing}} (c_{j}c_{j+1} + c_{j+1}^{\dagger}c_{j}^{\dagger}) \right]$$
(32)

JAYNES-CUMMINGS-HUBBARD MODEL

$$\mathcal{H} = \sum_{i} \left[\omega_{c} a_{i}^{\dagger} a_{i} + \omega_{a} \sigma_{i}^{+} \sigma_{i}^{-} + g(a_{i}^{\dagger} \sigma_{i}^{-} + a_{i} \sigma_{i}^{+}) \right] - J \sum_{\langle i,j \rangle} \left(a_{i}^{\dagger} a_{j} + a_{j}^{\dagger} a_{i} \right)$$
(33)

HUBBARD-HOLSTEIN MODEL

$$\mathcal{H} = -t \sum_{\langle i,j \rangle, \sigma} \left(c_{i\sigma}^{\dagger} c_{j\sigma} + \text{h.c.} \right) + U \sum_{i} n_{i\uparrow} n_{i\downarrow} + \omega_0 \sum_{i} b_i^{\dagger} b_i + g \sum_{i,\sigma} (b_i^{\dagger} + b_i) n_{i\sigma}$$
 (34)

LIPKIN-MESHKOV-GLICK (LMG) MODEL

$$\mathcal{H} = -\frac{\epsilon}{2} \sum_{i=1}^{N} \underbrace{\left(\hat{a}_{i,\uparrow}^{\dagger} \hat{a}_{i,\uparrow} + \hat{a}_{i,\downarrow}^{\dagger} \hat{a}_{i,\downarrow}\right)}_{2\hat{J}_{z}} + \frac{V}{2} \sum_{i,j} \left(\hat{a}_{i,\uparrow}^{\dagger} \hat{a}_{j,\uparrow}^{\dagger} \hat{a}_{j,\downarrow} \hat{a}_{i,\downarrow} + \hat{a}_{i,\downarrow}^{\dagger} \hat{a}_{j,\downarrow}^{\dagger} \hat{a}_{j,\uparrow} \hat{a}_{j,\uparrow} \hat{a}_{i,\uparrow}\right)$$
(35)

and $J_{+} = \sum_{i} \hat{a}_{i,\uparrow}^{\dagger} \hat{a}_{i,\downarrow}$ and $J_{-} = J_{+}^{\dagger}$

KITAEV HONEYCOMB MODEL

$$\mathcal{H} = -J_x \sum_{\langle ij \rangle_x} \sigma_i^x \sigma_j^x - J_y \sum_{\langle ij \rangle_y} \sigma_j^y \sigma_j^y - J_z \sum_{\langle ij \rangle_z} \sigma_i^z \sigma_j^z$$
(36)

DICKE MODEL

$$\mathcal{H} = \omega_c a^{\dagger} a + \omega \sum_{j=1}^{N} \sigma_j^z + \frac{\lambda}{\sqrt{N}} \sum_{j=1}^{N} \underbrace{\left(a^{\dagger} \sigma_j^- + a \sigma_j^+\right)}_{\text{co-rotating}} + \frac{\lambda}{\sqrt{N}} \sum_{j=1}^{N} \underbrace{\left(a \sigma_j^- + a^{\dagger} \sigma_j^+\right)}_{\text{counter-rotating}}$$
(37)

TAVIS-CUMMINS MODEL

$$\mathcal{H} = \omega_c \, a^{\dagger} a + \frac{\omega_a}{2} \sum_{j=1}^{N} \sigma_j^z + \sqrt{N} \sum_{j=1}^{N} \left(\sigma_j^+ a + \sigma_j^- a^{\dagger} \right) \tag{38}$$

Aubry-André model

$$\mathcal{H} = -J \sum_{j=1}^{N} \left(c_j^{\dagger} c_{j+1} + \text{h.c.} \right) + (2J + \lambda) \sum_{j=1}^{N} \underbrace{\cos[2\pi(\gamma j + \phi)]}_{\text{quasi-periodic potential}} c_j^{\dagger} c_j$$
 (39)

KITAEV TORIC CODE

$$\mathcal{H} = -\sum_{s} A_s - \sum_{\square} B_{\square} = -\prod_{i \in +} X_i - \prod_{i \in \square} Z_i \tag{40}$$

PXP MODEL (FOR QUANTUM SCARS)

$$\mathcal{H} = X_1 P_2 + \sum_{j=2}^{N-1} P_{j-1} X_j P_{j+1} + P_{N-1} X_N$$
(41)

where $P_n = (\mathbb{1} - Z_n)/2$