

I. LOOP ALGORITHM WITH RATIO WEIGHTING

We present a modification to the recently proposed algorithm using loop updates with quantum Monte Carlo in the valence bond basis. This improvement aids in the measurement of the Renyi entanglement entropies.

A. Valence bond quantum Monte Carlo

Quantum Monte Carlo in the valence bond basis (VB QMC) is a ground state projection technique which, for the given Hamiltonian, yields a state proportional to the ground state wavefunction by applying a high power of that Hamiltonian to a trial state, $|\psi\rangle = \sum_i c_i |i\rangle$,

$$|0\rangle \propto \mathcal{H}^m |\psi\rangle = \sum_{i=0} E_i^m c_i |i\rangle \quad (1)$$

$$= E_0^m \left[c_0 |0\rangle + \sum_{i=1} \left(\frac{E_i}{E_0} \right)^m c_i |i\rangle \right], \quad (2)$$

where all but the first term of (2) vanishes as $m \rightarrow \infty$ when E_0 is the energy with the largest magnitude.

In the case of the Heisenberg model, the Hamiltonian is modified by a constant shift and rewritten in terms of *bond operators* H_{ij} which act on nearest-neighbour pairs of sites,

$$\mathcal{H} = \sum_{\langle i,j \rangle} (\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4}) = - \sum_{\langle i,j \rangle} H_{ij}. \quad (3)$$

This allows us to sample terms in the ground state wavefunction by applying strings of m bond operators to a trial state.

So far nothing we have done is basis specific, however, when these bond operators give a very simple results when applied to a valence bond state. A bond operator acting on two sites already joined by a valence bond will not change the state. In the other case, when the sites are not joined, the operator will join those sites and the resulting state will have a weight of $1/2$,

$$H_{ab} |(a, b)\rangle = |(a, b)\rangle \quad (4)$$

$$H_{ab} |(a, d)(c, b)\rangle = \frac{1}{2} |(a, b)(c, d)\rangle, \quad (5)$$

where (a, b) represents a valence bond between sites a and b . For non-bipartite lattices (5) could pick up a negative depending on the sublattice of the sites involved. In this algorithm we only use valence bonds including one site from each sublattice.

MEASURING OBSERVABLES!!!!

To simplify matters the list of m bond operators can be rewritten as just one operator P_i such that $|0\rangle \propto \sum_i P_i |\psi\rangle$.

From (4) and (5), we can generalize the action of P_i on some state as $P_i |\psi\rangle = W^i |\psi^i\rangle$, where ψ^i is the new state acquired from the application of the m operators in

P_i and $W^i = 2^{-m_i}$ contains all of the factors of $1/2$ from applying operators which change the valence bond state as in (5).

We can then measure observables...

$$\langle \mathcal{O} \rangle = \frac{\langle \psi_L | (-\mathcal{H})^m \mathcal{O} (-\mathcal{H})^m | \psi_R \rangle}{\langle \psi_L | (-\mathcal{H})^m (-\mathcal{H})^m | \psi_R \rangle} \quad (6)$$

$$= \frac{\sum_i (\langle \psi_L | P_i^\dagger \mathcal{O} \sum_j (P_j | \psi_R \rangle))}{\sum_i \sum_j \langle \psi_L | P_i^\dagger P_j | \psi_R \rangle} \quad (7)$$

$$= \frac{\sum_{ij} W_L^i W_R^j \langle \psi_L^i | \mathcal{O} | \psi_R^j \rangle}{\sum_{ij} W_L^i W_R^j \langle \psi_L^i | \psi_R^j \rangle} \quad (8)$$

We need to simulate two copies of the system to do these measurements. Come back to this section!!!! I think I forgot to explain a lot.

For seriousness add a definition of the innerproduct, but don't explain it. $\langle \psi_L | \psi_R \rangle = 2^{m_{\text{loops}} - N/2}$ where n_{loops} is the number of loops and N is the number of sites (or use n_{max} the max number of loops)

Check notation consistency. can use psi for a general state and V for a VB state.. Z for a spin state... just like sandvik

B. The loop algorithm for valence bond quantum Monte Carlo

Introduction to loop algorithm: Improvement on standard VB QMC. Similar to SSE. Each MC step gives a state that is hardly correlated with the previous state, unlike standard VB QMC. Also there's no rejection step. (Note: I didn't talk about acceptance/rejection. Add a sentence in about sampling operators according to some weight)

1. Bond Operators

In this method the bond operators of the previous section are split into two types of operators, diagonal operators $H_{ab}(1)$ and off-diagonal operators $H_{ab}(2)$, which act on interacting pairs of sites (labelled a and b here)

$$H_{ab}(1) = (\frac{1}{4} - S_a^z S_b^z) \quad (9)$$

$$H_{ab}(2) = -(S_a^x S_b^x + S_a^y S_b^y). \quad (10)$$

Bond operators are chosen such that they only act on sites with antiparallel spins, in which case the diagonal operator does not change the spin state, while the off-diagonal operator flips both spins.

More on bondops? How do they compare to bondops from the last section?

2. Initialization

We begin with two copies of the system, as in the last section, however, in this case we need to specify both the initial valence bond state and a compatible spin state. Compatible spin states have antiparallel spins on each site of a valence bond, as shown in Fig (**vb-spin-loop diagram**). There are, of course, many compatible spin states for a given valence bond configuration (two per bond) but different spin configurations are sampled over the simulation.

After the two initial trial valence bond states and spin states are chosen, $2m$ diagonal operators are randomly chosen such that they are always acting on pairs of antiparallel spins, and on nearest-neighbour sites in the case of the Heisenberg model. It is simpler to begin with the same initial spin states on the left and right, otherwise some off-diagonal operators need to be placed initially. The trial states can, however, be modified later in the simulation.

3. the next thing... build/flip loops .. update operators

The next step, after choosing the initial states and operators is building the loops, as shown in figure (**vb-spin-loop diagram**). These loops connect sites and operators, similar to the SSE algorithm except, instead of having loops which are periodic in imaginary time, the loops connect along the edges according to the valence bonds in the chosen trial states.

Once the loops are built each loop is given a chance to “flip” with probability $1/2$. That is, every spin in the loop will flip, causing some diagonal operators to become off-diagonal operators and vice versa.

Next, all the remaining diagonal operators are changed to act on randomly selected pairs of compatible sites. The off-diagonal operators remain unchanged. Through this process of loop flipping and choosing new operators we sample the necessary valence bond and spin states.

4. Measurements

To perform measurements of observables, as in (6), we can extract the two projected valence bond states (or even the spin states) from our simulation. This is accomplished by using at the states in the “middle” of the vb-loop-operator diagram (m operators from the edge states), denoted by the dotted line in figure (**vb-spin-loop diagram**). Follow the loops on the left side of the dotted line will give the valence bond state $|V_L\rangle$ and the right side will give $|V_R\rangle$. Without extracting the states first, the number of loops crossing the dotted line gives the overlap $\langle V_L | V_R \rangle$.

C. Measuring the Renyi entropies

- Intro to Renyi entropies *if not here it'll be *some-where* at least*

The Renyi entropies are defined as

$$S_\alpha(A) = \frac{1}{1-\alpha} \ln [\text{Tr}(\rho_A^\alpha)], \quad (11)$$

where $\rho_A = \text{Tr}_B(\rho_{AB})$ is the reduced density matrix of the total system traced out over region B .

- Swap/Permutation operator

- define Π_A

We define a permutation operator Π_α (as in Cardy's quench paper) acting on α copies of the system to cyclically permute the copies of region A between the systems. For $\alpha = 2$ this reduces to the *Swap* operator of (our renyi paper).

- * generally... like for a general state
 - * with a diagram
 - * for VBs
 - Proof that $\langle \Pi_A \rangle = \text{Tr}(\rho_A^\alpha)$ (prolly in an appendix? short version here anyway)
-

Using a Schmidt decomposition we can write the ground state (or any pure state) as

$$|0\rangle = \sum_i \lambda_i |i\rangle_A |i\rangle_B, \quad (12)$$

where i runs from 1 to the dimension of the smaller of the two Hilbert spaces of region A or B , and λ_i^2 are the eigenvalues of the reduced density matrices ρ_A and ρ_B , which have the same eigenvalues.

Using the permutation operator on α copies of the system we can obtain the value $\text{Tr}(\rho_A^\alpha)$, necessary to measure the Renyi entropies.

- Measuring in VB QMC???
-
-
-

- double (or α -tuple) the system
- use the swap (or permutation) operator
- other measurements will be screwed up

In addition to the standard measurements (*examples*) we can use VB QMC to measure the Renyi entanglement entropies (Eq. (11)) for integer values of $\alpha \geq 2$. The method for measuring S_2 using the double projector algorithm is described in Ref[?].

To measure S_α we need to α -tuple the system that is simulated, so for the double projector algorithm 2α non-interacting copies of that system are required.

We are then able to take the expectation value of a *Swap* operator in the case of $\alpha = 2$ or a cyclical (*cyclic or cyclical?*) permutation operator for $\alpha > 2$. (*Permutation operator can also be a combination of $(\alpha - 1)$ Swap operators.*)

- Define the swap operator in: product basis, vb basis, etc?
- Explain how it gives $\text{Tr}(\rho_A^\alpha)$

1. Swap in double projector VBQMC

A pictorial example of what the *Swap* operator for $\alpha = 4$ does to regions *A* and *B* of the system is shown in figure ??*. Omg explain this so much better. Might need a different diagram for the double projector algorithm.*

2. Swap in loop VB QMC

Figures 2 and 3.

Same as double projector... just extract states to swap.

D. Ratio weighting

1. In double projector VBQMC

2. in the loop VBQMC

II. ——— O - L - D - S - T - U - F - F ———

A. Pre-initialization? Set-up?

1. Lattice

The standard loop algorithm simulation contains two independent copies of the system to be studied, amounting to one VB-spin-operator list. To measure S_2 an additional VB-spin-operator list is required, amounting to four copies of the original system. In general we will have 2α copies of the system, resulting in α VB-spin-operator structures.

We must choose an initial VB state for each copy of the system (not necessarily all different) with bonds only between different sublattices. For each of these initial VB states a compatible initial spin state must also be chosen

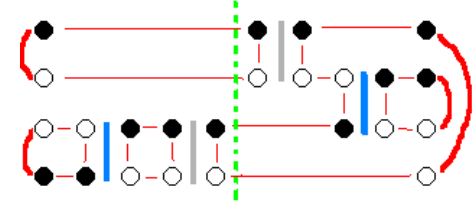


FIG. 1:

such that each valence bond has one site with spin up \uparrow and the other with spin down \downarrow .

It is simplest to begin with the same spin states for both initial states within a single VB-spin-operator list, though the VB states on either side may be different.

2. Bond Operators

Valence bond quantum Monte Carlo simulations are a type of ground state projection technique, and as such, we must decide how many times to apply the Hamiltonian to our trial state(s) in each step. This value m is a convergence parameter for the simulation. **((Something about the balance between large and small m .. need convergence.. don't want to waste simulations time. Doesn't scale properly per site... (((seems like you need less than a number of operators per site... like the number of ops/site that works for small systems is more than enough for large systems.. PLOTS))))**

- talk about convergence in the number of bond operators (plot proly, eh?)

B. Main Program

1. Building the loops

2. Flip...?

- keep the loops the same, flip all the spins.
- causes some diagonal operators to become off-diagonal and vice versa

3. $\langle \text{Swap}_A \rangle$ measurement

- renyi measurement
- need to multiply by weight factor to get the actual measurement values
- energy (and other measurements) could be doubled (or alpha-ed) because we have multiple copies of the system

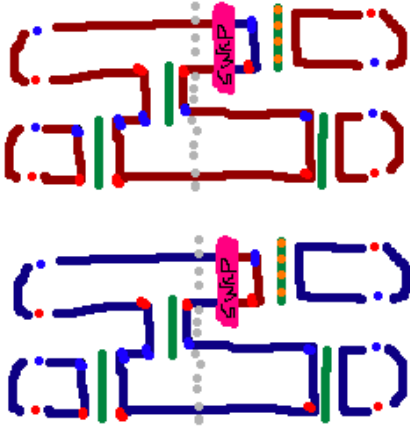


FIG. 2: A possible VB-spin-operator diagram for a 4-site system with $m=2$ operators per copy of the system, (**something** about ratio and used to calculate the 2nd renyi entropy) (also, make sure the two copies have **different** operators and loops, but still compatible spin configurations)



FIG. 3: The cyclical permutation of states in region A for four copies of a system.

4. Updating the operator list

- keep off-diagonal bond operators, move diagonal bond operators to a random, compatible spot.

C. Swap operator for $\alpha > 2$