



Inconsistency in the calculation of the time dependent climate feedback parameter.



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We introduce a **time-dependent climate feedback parameter** $\lambda(t)$ in Budyko's (1969) linear relationship between surface temperature and outgoing long-wave radiation. Using two different methods, we derive an energy balance model (EBM) with $\lambda(t)$ from energy conservation principles applied to the climate system.

We show the EBM with a time-variable λ : $N = \Delta F + \lambda(t)\Delta T_S$ (equation 1, N being the Earth energy imbalance, F the radiative forcing and T_S the surface temperature), is incomplete and should include another term $\Delta\lambda(t)\Delta T_{S0}$. The corrected EBM accurately reproduces the surface temperature response to abrupt CO₂ increase in multi-millennia experiments at all time-scales.

Our estimates of $\lambda(t)$ are consistent across simulations and time-scales, with much smaller departures from the mean value than previous estimates. We argue estimates of $\lambda(t)$ with equation 1 are erroneous and should be abandoned. To support this statement we show simulations where eq. 1 leads to absurd $\lambda(t)$ values. This has profound consequences that are developed in a companion poster by Guillaume-Castel and Meyssignac.

Theoretical development

The climate system is a **forced dynamical system**. The surface temperature follows

$$C \frac{dT_S}{dt} = F + R - H \quad (2)$$

Where:
H the heat exchange with the deep ocean
C is the heat capacity
T_S the surface temperature
F the radiative forcing
R the radiative response

On interannual and longer timescales, **R is linear with T_S** (Budyko, 1969): $R = \lambda T_S$ (3)

► λ is the climate feedback parameter.

If λ varies with time, the energy budget reads:

$$C \frac{dT_S}{dt} = F + \lambda(t)T_S - H \quad (4)$$

A given forcing F_0 is associated with steady state variables T_{S0} and λ_0 . On a steady state:

$$F_0 = \lambda_0 T_{S0} = R_0 \quad (5)$$

$$H = 0 \quad (6)$$

The preindustrial era/control experiments are considered to be on a **steady state**.

An increment forcing ΔF induces an increment radiative response ΔR . Equation 2 then becomes:

$$C \frac{dT_S}{dt} = F_0 + \Delta F + R_0 + \Delta R - H \quad (7)$$

Here, we derive the value of ΔR using three different methods.

Pattern effect hypothesis
 λ varies with the pattern of surface temperature warming, independently of the global average temperature:

$$R = \lambda(P)T_S = \lambda(t)T_S \quad (8)$$

Variations of λ are noted

$$\lambda(t) = \lambda_0 + \Delta\lambda(t) \quad (9)$$

Method I: Partial derivatives

We assume ΔR is small enough to be considered dR .

Following 8, we develop dR using the total derivative formula.

$$dR = \frac{\partial R}{\partial T_S} \Big|_{P=P_0} dT_S + \frac{\partial R}{\partial P} \Big|_{T_S=T_{S0}} dP$$

$$dR = \lambda_0 \frac{\partial R}{\partial T_S} \Big|_{P=P_0} dT_S + \frac{\partial R}{\partial P} \Big|_{T_S=T_{S0}} \cdot \frac{\partial \lambda}{\partial P} \Big|_{T_S=T_{S0}} dP$$

$$\Delta R = \lambda_0 \Delta T_S + T_{S0} \Delta \lambda$$

Method II: Perturbation theory

We assume a small deviation ΔF from F_0 induces $\Delta\lambda$. Ignoring H , the perturbed dynamical system therefore follows

$$C \frac{dT_S}{dt} = F_0 + \Delta F + (\lambda_0 + \Delta\lambda)T_S \quad (A)$$

Perturbation theory states that a solution to this system can be found close to a solution T_{S0} of the unperturbed system:

$$C \frac{dT_S}{dt} = F_0 + \lambda_0 T_S \quad (B)$$

We look for a 1st order solution to A with $T_S = T_{S0} + \Delta T_S$

Introducing $T_S = T_{S0} + \Delta T_S$ in A leads to:

$$C \frac{d(T_{S0} + \Delta T_S)}{dt} = F_0 + \Delta F + (\lambda_0 + \Delta\lambda)(T_{S0} + \Delta T_S)$$

$$= F_0 + \Delta F + \lambda_0 T_{S0} + T_{S0}\Delta\lambda + \lambda_0 \Delta T_S + \Delta\lambda \Delta T_S$$

$$C \frac{d\Delta T_S}{dt} = \Delta F + T_{S0}\Delta\lambda + \lambda_0 \Delta T_S$$

Then, using 7 finally leads to:

$$\Delta R = \lambda_0 \Delta T_S + T_{S0} \Delta \lambda$$

With $\lambda(t)$, the radiative response increment following ΔF is:

$$\Delta R(t) = \lambda_0 \Delta T_S(t) + T_{S0} \Delta \lambda(t) \neq \lambda(t) \Delta T_S(t)$$

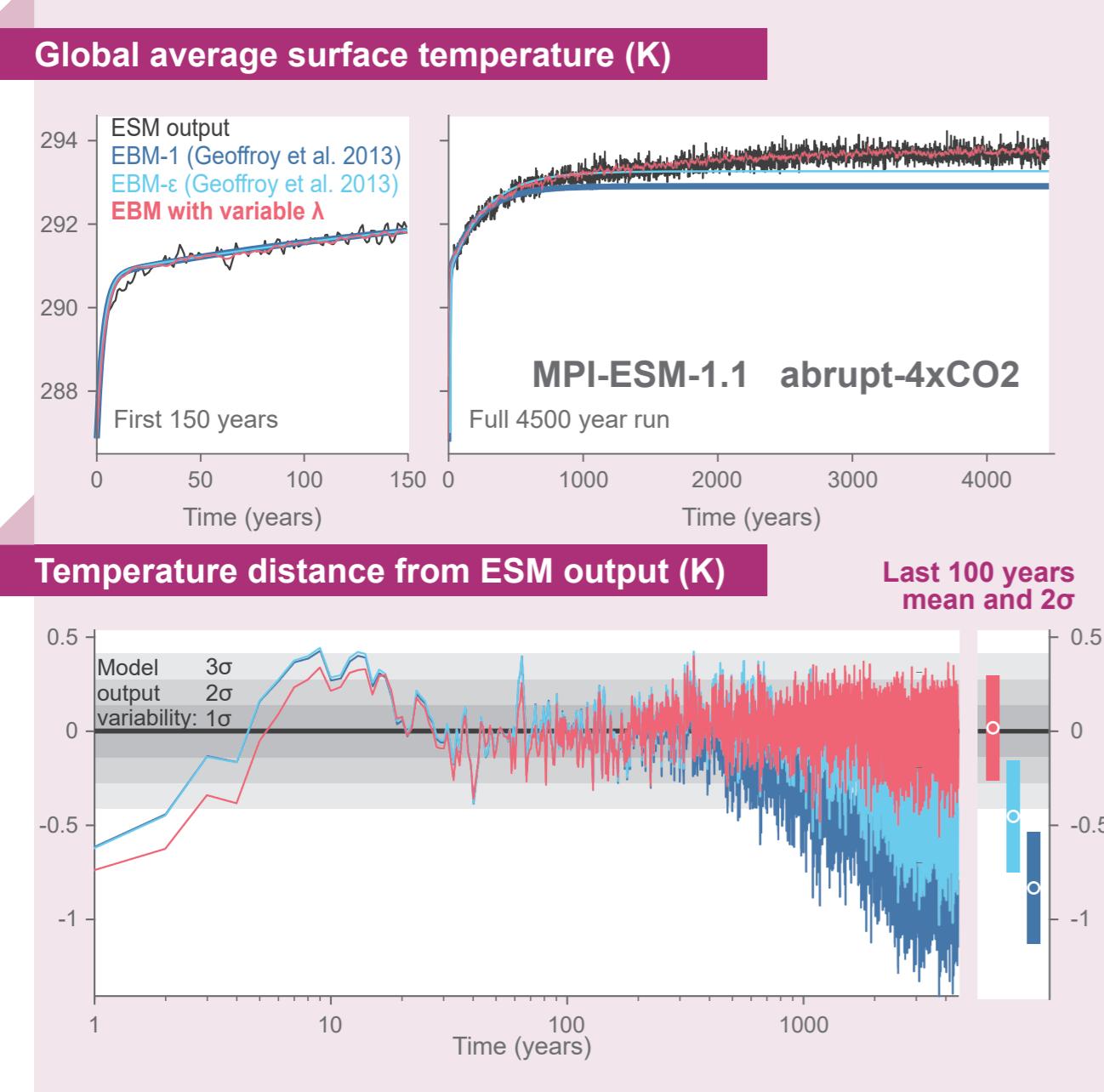
The surface temperature anomaly follows:

$$C \frac{d\Delta T_S}{dt} = \Delta F + T_{S0} \Delta \lambda + \lambda_0 \Delta T_S$$

Validation and Consequences

Reproducing the surface temperature dynamics

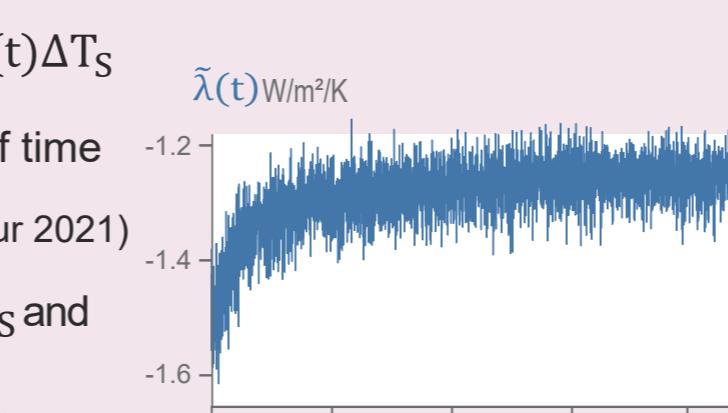
Numerical integration with 3-layer ocean



Comparing time varying λ in abrupt CO₂ increase experiments

In the literature

- $C \frac{d\Delta T_S}{dt} = \Delta F + \lambda(t)\Delta T_S$
 - Many definitions of time varying λ exist (Rugenstein & Armour 2021)
 - λ is defined using T_S and F anomalies, e.g.:
- $$\tilde{\lambda}(t) = \frac{N - \Delta F}{\Delta T_S}$$



With our formalism

- $C \frac{d\Delta T_S}{dt} = \Delta F + T_{S0}\Delta\lambda + \lambda_0\Delta T_S$
 - λ is given using total variables by :
- $$\lambda(t) = \frac{N - F_0 - \Delta F}{T_{S0} + \Delta T_S}$$

- Similar patterns
- λ deviates less from the mean than $\tilde{\lambda}$: note the different scales!

The constant temperature experiment

Test if $\Delta R = \lambda(t)\Delta T_S$ can be true using a constant temperature experiment, with varying λ : needs to include the pattern effect.

Experiment setup

- There is no CMIP experiment with pattern but fixed T_S
- Workaround with 2 experiments: amip-future4K (amip with 4K patterned increase) and amip-p4K (amip with 4K uniform increase)
- Compute $\tilde{\lambda}$, λ_0 and $\Delta\lambda$ from amip-future4K minus amip-p4K

Results

	$\langle \tilde{\lambda} \rangle$	$\lambda_0 (\Delta\lambda) \times 10^{-3}$
HadGEM3	-16.9	-2.3 1.5
MRI-ESM2	-7.4	-1.1 4.1
CESM2	-6.8	-1.9 4.3
MIROC6	-4.3	-2.0 1.6
IPSL-CM6A	-1.6	-1.6 0
CanESM5	-1.3	-1.4 0

The usual ΔR leads to absurd values of λ
 ΔR requires the supplementary term $T_{S0}\Delta\lambda(t)$

The true surface energy budget equation is:

$$C \frac{d\Delta T_S}{dt} = \Delta F + \Delta\lambda T_{S0} + \lambda_0 \Delta T_S - \Delta H$$

Budyko (1969). The effect of solar radiation variations on the climate of the Earth. tellus, 21(5), 611-619.

Rugenstein & Armour (2021). Three flavors of radiative feedbacks and their implications for estimating Equilibrium Climate Sensitivity. Geophysical Research Letters, 48(15), e2021GL092983.