Appendix A

Composed Weibull distribution

In Chapter 3 the Composed Weibull distribution is described. In this appendix expressions are derived for the significant wave height $H_{1/3}$ and the ratio of a wave height with a certain probability of exceedance to the significant wave height.

A.I Mean of the highest I/N-part

As mentioned in Chapter 3, $H_{1/N}$ is defined by:

$$H_{1/N} = \frac{\int_{H_N}^{\infty} H f(H) dH}{\int_{H_N}^{\infty} f(H) dH} = \frac{\int_{H_N}^{\infty} H f(H) dH}{\frac{1}{N} \int_0^{\infty} f(H) dH} = N \int_{H_N}^{\infty} H f(H) dH$$
(A.1)

In this definition H_N is the wave height with an exceedance probability of 1/N (N>1). In coastal engineering practice the wave height with an exceedance probability of 1/N is often denoted by $H_{1/N}$ %. However, in this appendix H_N is used instead of $H_{1/N}$ %, in line with the definition of H_N . The determination of $H_{1/N}$ depends on the fact whether or not H_N exceeds H_{1r} .

Transitional wave height exceeds H_N

In Figure A.1 a wave height distribution with $H_{tr} > H_N$ is shown.

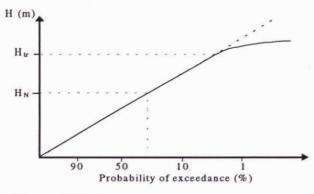


Figure A.1: Composed Weibull wave height distribution with H_{tr}>H_N.

In this case H_N is determined by:

$$Q_{\underline{H}_1}(H) \equiv \Pr_1\{\underline{H} \ge H\} = \frac{1}{N} = \exp\left[-\left(\frac{H_N}{H_1}\right)^{k_1}\right] \tag{A.2}$$

Yielding:

$$H_N = H_1[\ln(N)]^{\frac{1}{k_1}} \tag{A.3}$$

 $H_{1/N}$ is evaluated via

$$H_{1/N} = N \int_{H_N}^{H_U} H f_1(H) dH + N \int_{H_U}^{\infty} H f_2(H) dH$$
 (A.4)

The first integral can be rewritten into:

$$N \int_{H_N}^{H_{u}} H f_1(H) dH = N \int_{H_N}^{\infty} H f_1(H) dH - N \int_{H_{u}}^{\infty} H f_1(H) dH$$
 (A.5)

yielding:

$$H_{1/N} = N \int_{H_N}^{\infty} H f_1(H) dH - N \int_{H_{1r}}^{\infty} H f_1(H) dH + N \int_{H_{1r}}^{\infty} H f_2(H) dH$$
 (A.6)

Substitution of the probability density function (3.3) in (A.6) yields:

$$H_{UN} = N \int_{H_{N}}^{\infty} H \frac{k_{1}}{H_{1}^{k_{1}}} H^{k_{1}-1} \exp \left[-\left(\frac{H}{H_{1}} \right)^{k_{1}} \right] dH - N \int_{H_{N}}^{\infty} H \frac{k_{1}}{H_{1}^{k_{1}}} H^{k_{1}-1} \exp \left[-\left(\frac{H}{H_{1}} \right)^{k_{1}} \right] dH$$

$$+ N \int_{H_{N}}^{\infty} H \frac{k_{2}}{H_{2}^{k_{2}}} H^{k_{2}-1} \exp \left[-\left(\frac{H}{H_{2}} \right)^{k_{2}} \right] dH$$
(A.7)

Assume

$$t = \left(\frac{H}{H_i}\right)^{k_i} \qquad \Rightarrow \qquad H = H_i t^{\frac{1}{k_i}} \qquad \Rightarrow \qquad dH = \frac{H_i}{k_i} t^{\frac{1}{k_i} - 1} dt \qquad for \quad i = 1, 2. \quad (A.8)$$

When this transformation is applied to (A.7) the following equation is obtained:

$$H_{1/N} = NH_{1} \int_{t_{1}}^{\infty} t^{\frac{1}{k_{1}}} \exp[-t]dt - NH_{1} \int_{t_{1}}^{\infty} t^{\frac{1}{k_{1}}} \exp[-t]dt$$

$$\left(\frac{H_{N}}{H_{1}}\right)^{k_{1}}$$

$$\left(\frac{H_{N}}{H_{1}}\right)^{k_{1}}$$

$$+ NH_{2} \int_{t_{1}}^{\infty} t^{\frac{1}{k_{2}}} \exp[-t]dt$$

$$\left(\frac{H_{N}}{H_{2}}\right)^{k_{2}}$$
(A.9)

With the incomplete gamma functions described in Section A.4 equation (A.9) is rewritten into:

$$H_{1/N} = NH_1 \left[\Gamma \left[\frac{1}{k_1} + 1, \ln(N) \right] - \Gamma \left[\frac{1}{k_1} + 1, \left(\frac{H_{\nu}}{H_1} \right)^{k_1} \right] \right] + NH_2 \Gamma \left[\frac{1}{k_2} + 1, \left(\frac{H_{\nu}}{H_2} \right)^{k_2} \right]$$
(A.10)

Transitional wave height exceeded by H_N

Figure A.2 shows a wave height distribution on a shallow foreshore for the case $H_{tr} < H_{N}$, which means that the transitional wave height is exceeded by the wave height with a certain probability of exceedance, H_{N} .

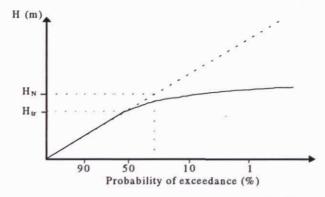


Figure A.2: Composed Weibull wave height distribution with H_{tr}<H_N.

When H_{tr} is exceeded by H_N , the determination of the $H_{1/N}$ is less complicated. H_N is determined by:

$$Q_{\underline{H}_2} \equiv \Pr_2\{\underline{H} \ge H\} = \frac{1}{N} = \exp\left[-\left(\frac{H_N}{H_2}\right)^{k_2}\right] \tag{A.11}$$

yielding:

$$H_N = H_2[\ln(N)]^{\frac{1}{k_2}}$$
 (A.12)

 $H_{1/N}$ is evaluated via

$$H_{1/N} = N \int_{H_N}^{\infty} H f_2(H) dH$$
 (A.13)

When the probability density function $f_2(H)$ from the Composed Weibull distribution (3.3) is substituted, the following equation is obtained:

$$H_{1/N} = N \int_{H_N}^{\infty} H \frac{k_2}{H_2^{k_2}} H^{k_2 - 1} \exp\left[-\left(\frac{H}{H_2}\right)^{k_2} \right] dH$$
 (A.14)

which can be transformed with (A.8) into:

$$H_{1/N} = NH_2 \int_{\left(\frac{H_N}{H_2}\right)^{k_2}}^{\infty} t^{\frac{1}{k_2}} \exp\left[-t\right] dt$$
 (A.15)

With the incomplete gamma functions (Section A.4) equation (A.15) is rewritten into:

$$H_{1/N} = N H_2 \Gamma \left[\frac{1}{k_2} + 1, \left(\frac{H_N}{H_2} \right)^{k_2} \right]$$
 (A.16)

With (A.12) the following equation is obtained:

$$H_{1/N} = N H_2 \Gamma \left[\frac{1}{k_2} + 1, \ln(N) \right]$$
 (A.17)

In coastal engineering design practice one often is interested in the ratio of an extreme wave height to the significant wave height, $H_{2\%}/H_{1/3}$, $H_{1\%}/H_{1/3}$ or $H_{0.1\%}/H_{1/3}$. The H_N to $H_{1/3}$ ratio for the Composed Weibull distribution is derived in the following section.

A.2 Extreme wave height to significant wave height ratio

In the evaluation of the H_N to $H_{1/3}$ ratio three different situations must be distinguished:

- 1. $H_{\rm N} < H_{\rm tr}$
- 2. $H_3 < H_{tr} < H_N$
- 3. $H_{tr} < H_3$

with $H_3 < H_N$.

H_N<H.,

In relatively deep water the transitional wave height is exceeding both H_3 and H_N , as is shown in Figure A.3. The H_N to $H_{1/3}$ ratio is obtained by evaluating the first part of the composed Weibull distribution $F_1(H)$. The wave height exceeded by 1/N-part of the waves

heights is determined by (A.3) and the significant wave height is determined by equation (A.10).

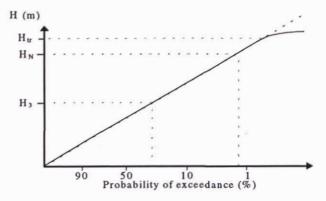


Figure A.3: Composed Weibull wave height distribution with H_N<H_{1r}.

The H_N to $H_{1/3}$ ratio becomes

$$\frac{H_{N}}{H_{1/3}} = \frac{H_{1} \left[\ln(N)\right]_{k_{1}}^{\frac{1}{k_{1}}}}{3H_{1} \left[\Gamma\left[\frac{1}{k_{1}}+1,\ln(N)\right]-\Gamma\left[\frac{1}{k_{1}}+1,\left(\frac{H_{\nu}}{H_{1}}\right)^{k_{1}}\right]\right] + 3H_{2} \Gamma\left[\frac{1}{k_{2}}+1,\left(\frac{H_{\nu}}{H_{2}}\right)^{k_{2}}\right]}$$
(A.18)

$H_3 < H_{tr} < H_N$

When the waves propagate into shallower water, the wave height at which the waves deviate from the Rayleigh distribution decreases.

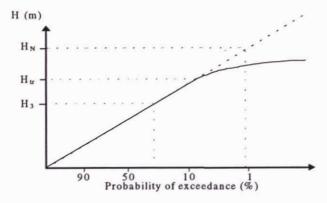


Figure A.4: Composed Weibull wave height distribution with H3<H1r<HN.

As shown in Figure A.4, the transitional wave height is exceeded by the wave height with a probability of exceedance of 1/N. Therefore H_N is determined by the second part of the composed wave height distribution via equation (A.12). The wave height with an exceedance probability of 1/3, H_3 , is now obtained via

$$H_3 = H_1 \left[\ln(3) \right]^{\frac{1}{k_1}} \tag{A.19}$$

when H_3 is smaller than the transitional wave height. Equation (A.10) provides the mean of the highest 1/3-part of the wave height distribution, $H_{1/3}$. Thus, for $H_3 < H_{tr} < H_N$ the following equation determines the H_N to $H_{1/3}$ ratio:

$$\frac{H_N}{H_{1/3}} = \frac{H_2[\ln(N)]^{\frac{1}{k_2}}}{3H_1\left[\Gamma\left[\frac{1}{k_1}+1, \ln(3)\right] - \Gamma\left[\frac{1}{k_1}+1, \left(\frac{H_{\nu}}{H_1}\right)^{k_1}\right]\right] + 3H_2\Gamma\left[\frac{1}{k_2}+1, \left(\frac{H_{\nu}}{H_2}\right)^{k_2}\right]}$$
(A.20)

H,, < H3

In Figure A.5 the wave height, with an exceedance probability of 1/3, H_3 , exceeds the transitional wave height.

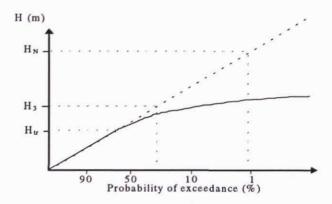


Figure 3.5: Composed Weibull height distribution with H1r<H3.

All wave heights exceeding H_{tr} obey the second part of the wave height distribution, $F_2(H)$, H_N is determined by (A.12) and the significant wave height by

$$H_{1/3} = 3 H_2 \Gamma \left[\frac{1}{k_2} + 1, \ln(3) \right]$$
 (A.21)

Therefore the following H_N to $H_{1/3}$ ratio is found:

$$\frac{H_N}{H_{1/3}} = \frac{\left[\ln(N)\right]^{\frac{1}{k_2}}}{3\Gamma\left[\frac{1}{k_2} + 1, \ln(3)\right]}$$
(A.22)

A.3 Evaluation of the $H_{1\%}$ to $H_{1/3}$ ratio

In this section a specific ratio of a wave height with a certain probability of exceedance to the significant wave height is evaluated as an example. Since in coastal engineering practice $H_{1N\%}$ is often used to represent a wave height with a certain probability of exceedance, in this section the same notation is used.

H_{1%}to H_{1/3} ratio in deep water

In this section the $H_{1\%}$ to $H_{1/3}$ ratio is given for deep and extremely shallow water. In deep water, where all the wave heights obey the Rayleigh distribution $H_{tr} \rightarrow \infty$. This means that (A.18) reduces to

$$\lim_{H_{y}\to\infty} \left(\frac{H_{N}}{H_{1/3}}\right) = \frac{\left[\ln(N)\right]^{\frac{1}{k_{1}}}}{3\Gamma\left[\frac{1}{k_{1}}+1, \ln(3)\right]}$$
(A.23)

In Section 3.3.2 assumptions are proposed for the exponents k_1 and k_2 , i.e. k_1 =2 and k_2 =3.5. If k_1 =2 and N=100 are substituted in equation (A.23), the following $H_{1\%}$ to $H_{1/3}$ ratio is obtained:

$$\frac{H_{1\%}}{H_{1/3}} = \frac{\left[\ln(100)\right]^{\frac{1}{2}}}{3\Gamma\left[\frac{1}{2}+1, \ln(3)\right]} \approx 1.52 \tag{A.24}$$

This is in line with the $H_{1\%}$ to $H_{1/3}$ ratio of the Rayleigh distribution, since from (2.4):

$$Q_{\underline{H}}(H) \equiv \Pr\{\underline{H} > H\} = \exp\left[-2\left(\frac{H}{H_{1/3}}\right)^2\right]$$
(A.25)

the same $H_{1\%}$ to $H_{1/3}$ ratio is obtained:

$$\frac{H_{1\%}}{H_{1/3}} = \sqrt{\frac{1}{2} \ln(Q_H^{-1})} = \sqrt{\frac{1}{2} \ln(100)} \approx 1.52$$
 (A.26)

H_{1%}to H_{1/3} ratio in extremely shallow water

In extremely shallow water, where $H_{1\%}$ and H_{3} exceed the transitional wave height, equation (A.22) was found to represent the $H_{1\%}$ to $H_{1/3}$ ratio. When k_{2} = 3.5 is substituted in (A.22) the following constant $H_{1\%}$ to $H_{1/3}$ ratio is obtained:

$$\frac{H_{1\%}}{H_{1/3}} = \frac{\left[\ln(100)\right]^{\frac{1}{3.5}}}{3\Gamma\left[\frac{1}{3.5} + 1, \ln(3)\right]} \approx 1.28$$
(A.27)

This means that, when $H_{tr} < H_3$, the H_N to $H_{1/3}$ ratio fully depends on the shape of the second part of the composed wave height distribution.

A.4 Incomplete gamma functions

The gamma function is defined by:

$$\Gamma(a) = \int_{0}^{\infty} t^{a-1} \exp[-t] dt \qquad 0 < t < \infty$$
(A.28)

This gamma function is a generalization of the factorial function. The gamma function has the following properties.

$$\Gamma(a+1) = a\Gamma(a) \tag{A.29}$$

$$\Gamma(a+1) = a!$$
 for $a = 1, 2, ..., n$. (A.30)

The incomplete gamma function is defined by (Abramowitz and stegun 1965):

$$P(a,x) = \frac{\gamma(a,x)}{\Gamma(a)} = \frac{1}{\Gamma(a)} \int_{0}^{x} t^{a-1} \exp[-t] dt \qquad (a > 0)$$
 (A.31)

in which $\Gamma(a)$ is the complete gamma function. Therefore $\gamma(a,x)$ is defined by:

$$\gamma(a,x) = \int_{0}^{x} t^{a-1} \exp[-t] dt$$
 (A.32)

The complement of P(a,x) is defined by:

$$Q(a,x) = 1 - P(a,x) = \frac{\Gamma(a,x)}{\Gamma(a)} = \frac{1}{\Gamma(a)} \int_{x}^{\infty} t^{a-1} \exp[-t] dt$$
 (A.33)

yielding

$$\Gamma(a,x) = \int_{x}^{\infty} t^{a-1} \exp[-t] dt$$
 (A.34)

Appendix C

Table 7.1: Characteristic nondimensional wave heights \tilde{H}_{x} as function of $\tilde{H}_{\rm tr}$.

| $\widetilde{H}_{\scriptscriptstyle tr}$ | $	ilde{H}_1$ | \widetilde{H}_2 | $	ilde{H}_{	ext{\tiny 1/3}}$ | $	ilde{H}_{\scriptscriptstyle 1/10}$ | $\widetilde{H}_{2\%}$ | $\widetilde{H}_{1\%}$ | $\widetilde{H}_{\scriptscriptstyle 0.1\%}$ |
|---|----------------|-------------------|------------------------------|--------------------------------------|-----------------------|-----------------------|--|
| 0.05 | 9.949 | 1.029 | 1.249 | 1.438 | 1.520 | 1.592 | 1.788 |
| 0.10 | 5.916 | 1.029 | 1.249 | 1.438 | 1.520 | 1.592 | 1.788 |
| 0.15 | 4.365 | 1.029 | 1.249 | 1.438 | 1.520 | 1.592 | 1.788 |
| 0.20 | 3.518 | 1.029 | 1.249 | 1.438 | 1.520 | 1.592 | 1.788 |
| 0.25 | 2.976 | 1.029 | 1.249 | 1.438 | 1.520 | 1.593 | 1.788 |
| 0.30 | 2.595 | 1.029 | 1.249 | 1.438 | 1.520 | 1.593 | 1.788 |
| 0.35 | 2.312 | 1.030 | 1.249 | 1.438 | 1.520 | 1.593 | 1.788 |
| 0.40 | 2.092 | 1.030 | 1.250 | 1.438 | 1.520 | 1.593 | 1.789 |
| 0.45 | 1.916 | 1.030 | 1.250 | 1.438 | 1.521 | 1.593 | 1.789 |
| 0.50 | 1.772 | 1.030 | 1.250 | 1.439 | 1.521 | 1.594 | 1.790 |
| 0.55 | 1.651 | 1.031 | 1.251 | 1.440 | 1.522 | 1.595 | 1.791 |
| 0.60 | 1.549 | 1.032 | 1.252 | 1.441 | 1.523 | 1.596 | 1.792 |
| 0.65 | 1.462 | 1.033 | 1.253 | 1.442 | 1.525 | 1.598 | 1.794 |
| 0.70 | 1.387 | 1.035 | 1.255 | 1.445 | 1.528 | 1.600 | 1.797 |
| 0.75 | 1.322 | 1.037 | 1.258 | 1.448 | 1.531 | 1.604 | 1.801 |
| 0.80 | 1.265 | 1.040 | 1.262 | 1.452 | 1.535 | 1.608 | 1.806 |
| 0.85 | 1.216 | 1.043 | 1.266 | 1.457 | 1.540 | 1.614 | 1.812 |
| 0.90 | 1.174 | 1.048 | 1.271 | 1.463 | 1.547 | 1.621 | 1.820 |
| 0.95 | 1.137 | 1.053 | 1.278 | 1.470 | 1.555 | 1.629 | 1.829 |
| 1.00 | 1.106 | 1.059 | 1.285 | 1.479 | 1.564 | 1.638 | 1.840 |
| 1.05 | 1.079 | 1.066 | 1.294 | 1.489 | 1.575 | 1.650 | 1.852 |
| 1.10 | 1.056 | 1.075 | 1.304 | 1.501 | 1.587 | 1.663 | 1.867 |
| 1.15 | 1.037 | 1.084 | 1.314 | 1.514 | 1.601 | 1.677 | 1.883 |
| 1.20 | 1.021 | 1.094 | 1.321 | 1.528 | 1.616 | 1.693 | 1.901 |
| 1.25 | 1.008 | 1.105 | 1.327 | 1.544 | 1.632 | 1.710 | 1.920 |
| 1.30 | 0.998 | 1.117 | 1.331 | 1.561 | 1.650 | 1.729 | 1.941 |
| 1.35 | 0.989 | 1.130 | 1.335 | 1.578 | 1.669 | 1.748 | 1.963 |
| 1.40 | 0.983 | 1.144 | 1.338 | 1.597 | 1.689 | 1.769 | 1.987 |
| 1.45 | 0.978 | 1.158 | 1.341 | 1.617 | 1.709 | 1.791 | 2.011 |
| 1.50 | 0.974 | 1.172 | 1.344 | 1.637 | 1.731 | 1.813 | 2.036 |
| 1.55 | 0.972 | 1.187 | 1.347 | 1.654 | 1.753 | 1.836 | 2.062 |
| 1.60 | 0.970 | 1.202 | 1.350 | 1.669 | 1.775 | 1.860 | 2.088 |
| 1.65 | 0.970 | 1.218 | 1.354 | 1.682 | 1.798 | 1.884 | 2.115 |
| 1.70 | 0.969 | 1.233 | 1.357 | 1.694 | 1.821 | 1.908 | 2.142 |
| 1.75 | 0.970 0.971 | 1.249 | 1.361 | 1.704 | 1.844 | 1.932 | 2.170 |
| 1.80 | 0.971 | 1.265 | 1.364 | 1.714 | 1.868 | 1.957 | 2.197 |
| 1.85 | 0.972 | 1.281 | 1.368 | 1.722 1.730 | 1.891 | 1.981 2.006 | 2.225 |
| 1.90 1.95 | 0.974 | 1.297 | 1.372 | 1.730 | 1.915 1.926 | 2.006 | 2.252 2.280 |
| 2.00 | 0.973 | 1.312 1.328 | 1.375 1.379 | 1.737 | 1.926 | 2.055 | 2.307 |
| 2.00 | 0.977 | 1.328 | 1.379 | 1.744 | 1.929 | 2.033 | 2.334 |
| 2.05 | 0.979 | 1.344 | 1.382 | 1.756 | 1.933 | 2.103 | 2.361 |
| 2.10 | 0.983 | 1.339 | 1.389 | 1.761 | 1.936 | 2.105 | 2.388 |
| 2.20 | 0.985 | 1.390 | 1.392 | 1.766 | 1.944 | 2.109 | 2.414 |
| 2.25 | 0.986 | 1.404 | 1.395 | 1.770 | 1.947 | 2.113 | 2.440 |

| \widetilde{H}_{tr} | ${	ilde H}_1$ | ${	ilde H}_2$ | ${	ilde H}_{1/3}$ | $	ilde{H}_{{\scriptscriptstyle 1/10}}$ | ${	ilde H}_{2\%}$ | ${	ilde H}_{1\%}$ | ${\widetilde H}_{0.1\%}$ |
|----------------------|---------------|---------------|-------------------|--|-------------------|-------------------|--------------------------|
| 2.30 | 0.988 | 1.419 | 1.397 | 1.774 | 1.951 | 2.116 | 2.465 |
| 2.35 | 0.989 | 1.433 | 1.400 | 1.778 | 1.954 | 2.120 | 2.490 |
| 2.40 | 0.991 | 1.448 | 1.402 | 1.781 | 1.957 | 2.123 | 2.515 |
| 2.45 | 0.992 | 1.462 | 1.404 | 1.784 | 1.960 | 2.126 | 2.539 |
| 2.50 | 0.993 | 1.475 | 1.406 | 1.786 | 1.962 | 2.129 | 2.563 |
| 2.55 | 0.994 | 1.489 | 1.407 | 1.789 | 1.965 | 2.132 | 2.586 |
| 2.60 | 0.995 | 1.502 | 1.409 | 1.791 | 1.967 | 2.134 | 2.609 |
| 2.65 | 0.996 | 1.515 | 1.410 | 1.792 | 1.969 | 2.136 | 2.616 |
| 2.70 | 0.997 | 1.528 | 1.411 | 1.794 | 1.970 | 2.137 | 2.618 |
| 2.75 | 0.997 | 1.540 | 1.412 | 1.795 | 1.971 | 2.139 | 2.620 |
| 2.80 | 0.998 | 1.553 | 1.413 | 1.796 | 1.973 | 2.140 | 2.621 |
| 2.85 | 0.998 | 1.565 | 1.413 | 1.797 | 1.974 | 2.141 | 2.623 |
| 2.90 | 0.999 | 1.577 | 1.414 | 1.797 | 1.974 | 2.142 | 2.624 |
| 2.95 | 0.999 | 1.589 | 1.414 | 1.798 | 1.975 | 2.143 | 2.625 |
| 3.00 | 0.999 | 1.601 | 1.414 | 1.798 | 1.976 | 2.144 | 2.625 |
| 3.05 | 0.999 | 1.612 | 1.415 | 1.799 | 1.976 | 2.144 | 2.626 |
| 3.10 | 0.999 | 1.623 | 1.415 | 1.799 | 1.977 | 2.144 | 2.626 |
| 3.15 | 1.000 | 1.635 | 1.415 | 1.799 | 1.977 | 2.145 | 2.627 |
| 3.20 | 1.000 | 1.646 | 1.415 | 1.799 | 1.977 | 2.145 | 2.627 |
| 3.25 | 1.000 | 1.657 | 1.415 | 1.799 | 1.977 | 2.145 | 2.627 |
| 3.30 | 1.000 | 1.668 | 1.415 | 1.800 | 1.977 | 2.145 | 2.628 |
| 3.35 | 1.000 | 1.679 | 1.416 | 1.800 | 1.978 | 2.146 | 2.628 |
| 3.40 | 1.000 | 1.689 | 1.416 | 1.800 | 1.978 | 2.146 | 2.628 |
| 3.45 | 1.000 | 1.700 | 1.416 | 1.800 | 1.978 | 2.146 | 2.628 |
| 3.50 | 1.000 | 1.711 | 1.416 | 1.800 | 1.978 | 2.146 | 2.628 |

WL | delft hydraulics

Notation

Roman letters:

A : parameter based on the exponent in the Glukhovskiy distribution (-)

d : water depth (m).

 d_t : water depth at the toe of the sea-defence work (m). \tilde{d} : ratio of mean wave height to water depth, H_m/d (-).

 d^* : ratio of root mean square wave height to water depth, H_{rms}/d (-).

H : wave height (m).

 H_i : scale wave height of the Composed Weibull distribution (m).

 $H_{1/3}$: significant wave height defined as the mean of the highest 1/3-part of the wave

heights in a wave field (m).

 $H_{1/3,0}$: incident significant wave height (m).

 $H_{1/3,t}$: significant wave height at the toe of the sea-defence work (m).

 H_m : mean wave height (m).

 H_{m0} : spectral significant wave height, $H_{m0} = 4\sqrt{m_0}$ (m).

 H_N : the wave height with an exceedance probability of 1/N (m).

 $H_{1/N}$: mean of the highest 1/N-part of the wave heights in a wave field (m).

 H_{rms} : root mean square wave height (m).

H_t: transition wave height of the split Weibull distribution (Van Vledder 1993) (m).

 H_{tr} : transitional wave height of the Composed Weibull distribution (m). \tilde{H}_x : nondimensionalized characteristic wave height, $\tilde{H}_x = H_x/H_{rms}$ (-).

 k_i : exponent of the Composed Weibull distribution (-).

 L_f : foreshore length (m).

 L_{op} : deep water wavelength based on the peak period T_{p} (m).

 $L_{0.2}$: local wave length based on the average zero-crossing period $T_{0.2}$ (m).

 m_n : n^{th} moment of the frequency spectrum (m²).

 m_0 : variance of the water surface elevation, i.e. the total wave energy (m²).

 $T_{0,2}$: average zero-crossing period, $T_{0,2} = \sqrt{m_0/m_2}$ (s).

Greek letters:

 α : slope of the foreshore (-).

 α_{H1} : empirical coefficient of the scale wave height forecasting function (-).

 α_{tr} : empirical coefficient of the transitional wave height forecasting function (-). β : coefficient of the Modified Glukhovskiy distribution (Klopman 1996) (-).

 β_{tr} : empirical coefficient of the transitional wave height forecasting function (-).

 ε_{rms} : root mean square error (-)

 γ_{tr} : empirical breaker coefficient of the transitional wave height (-).

 κ : exponent of the Glukhovskiy distribution (-).

 κ^* : exponent of the Modified Glukhovskiy distribution (Klopman 1996) (-). σ : scale wave height of the split Weibull distribution (Van Vledder 1993) (m).

 Ψ : degree of saturation, $\Psi = \sqrt{m_0} / d$ (-).