

FLUIDS, FLOWS AND COMPLEXITY

PROBLEM SET 1 AND SOLUTIONS

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Changes for 2016: rearranged to more closely follow the lectures.

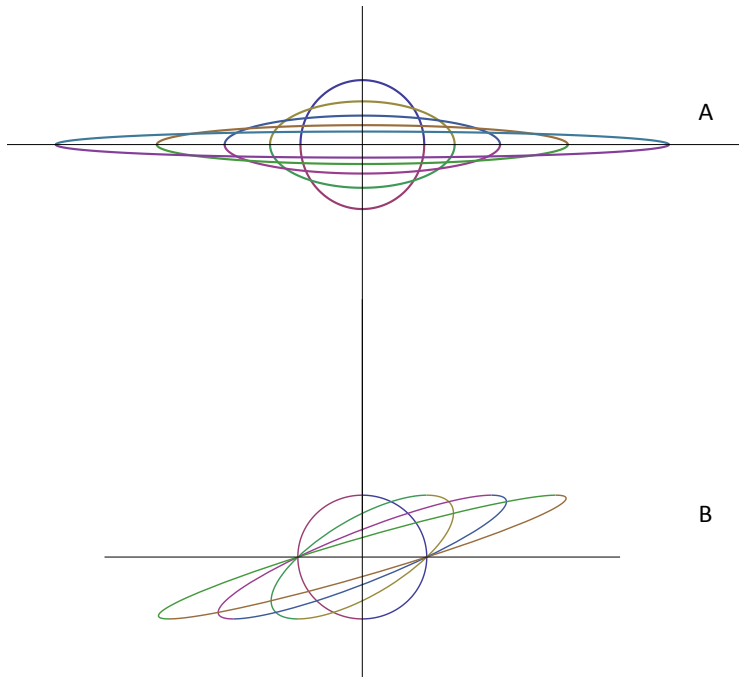
Thank you to Professor David Marshall who gave this course in previous years. I have used some of his problems and lecture notes. A few problems are from the web sites of Oxford Maths and Cambridge Maths.

1. Streamlines and flows

For (A) a 2D straining flow $\mathbf{u} = (\alpha x, -\alpha y)$ and (B) a simple shear flow $\mathbf{u} = (\gamma y, 0)$ where α and γ are constants:

- (a) Find the equation for a general streamline of the flow.
- (b) At $t = 0$ dye is introduced to mark the curve $x^2 + y^2 = a^2$. Find the equation for this material fluid curve for $t > 0$ and sketch how the curve evolves with time.
- (c) Does the area within the curve change in time, and why?
- (d) Which of the two flows stretches the curve faster at long times?

Solution



A:

- (a) Streamlines are tangent to the flow at every point. Therefore, for this 2D flow,

$$\frac{dx}{u_x} = \frac{dy}{u_y} \Rightarrow \frac{dx}{x} = -\frac{dy}{y} \Rightarrow \text{streamlines are } xy = \text{constant}.$$

(b)

$$u_x = \frac{dx}{dt} = \alpha x \text{ and } x = x(0) \text{ at } t = 0 \Rightarrow x(t) = x(0)e^{\alpha t}. \text{ Similarly } y(t) = y(0)e^{-\alpha t}.$$

Using $x^2(0) + y^2(0) = a^2$ this gives

$$x^2(t)e^{-2\alpha t} + y^2(t)e^{2\alpha t} = a^2$$

which is an ellipse with semi-axes of length $ae^{\alpha t}$ along x and $ae^{-\alpha t}$ along y . So the flow stretches along x and contracts along y but

(c) the area within the curve does not change with time because the flow is incompressible, $\text{div } \mathbf{u} = 0$.

B:

(a) for shear flow the streamlines are $y = \text{constant}$.

(b) $x(t) = x(0) + \gamma y(0)t, y(t) = y(0)$ so the circle of radius a evolves to an ellipse with major axis at an angle to the x -axis which decreases with time:

$$(x(t) - \gamma y(0)t)^2 + y^2(t) = a^2.$$

(c) the flow is again incompressible.

(d) A - the stretching is exponential in time.

2. Stream function and velocity potential

(a) Is the motion incompressible for the flows given by the following velocity potentials:

(i) $\phi = C(x^2 + y^2)$ (ii) $\phi = C(x^2 - y^2)$?

If so, determine the corresponding stream functions.

(b) Is the motion irrotational for the flows given by the following stream functions:

(iii) $\psi = C(x^2 + y^2)$ (iv) $\psi = C(x^2 - y^2)$?

If so, determine the corresponding velocity potentials.

(iii) Sketch the streamlines for all cases (a)–(d) and the lines of constant ϕ where possible.

Solution

(a) (i) $\mathbf{u} = \nabla\phi = (2Cx, 2Cy)$, $\nabla \cdot \mathbf{u} \neq 0$ so compressible

(ii) $\mathbf{u} = \nabla\phi = (2Cx, -2Cy)$, $\nabla \cdot \mathbf{u} = 0$, incompressible, $\psi = 2Cxy$ (cf Q1A)

(b) (iii) $\mathbf{u} = \left(\frac{d\psi}{dy}, -\frac{d\psi}{dx}\right) = (2Cy, -2Cx)$, $\nabla \wedge \mathbf{u} \neq 0$ so not irrotational

(iv) $\mathbf{u} = \left(\frac{d\psi}{dy}, -\frac{d\psi}{dx}\right) = (-2Cy, -2Cx)$, $\nabla \wedge \mathbf{u} = 0$, irrotational, $\phi = -2Cxy$

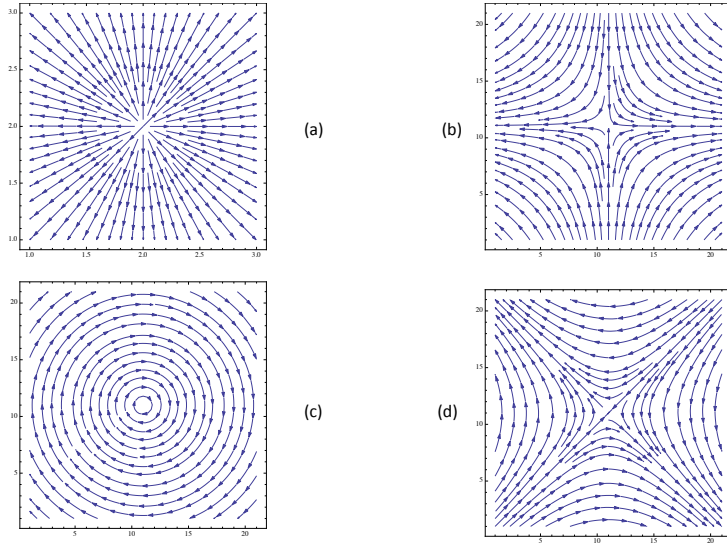


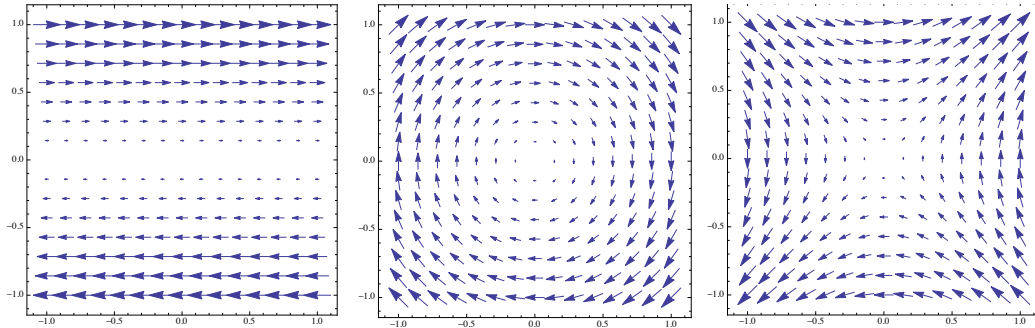
Figure 1: Lines of constant ϕ are perpendicular to the streamlines for (a), (b), (d).

3. Velocity gradient tensor

Show that a simple shear flow $\mathbf{u} = (\alpha y, 0, 0)$ can be decomposed into a sum of a rotation and a straining flow (i) pictorially (ii) in terms of the velocity gradient tensor.

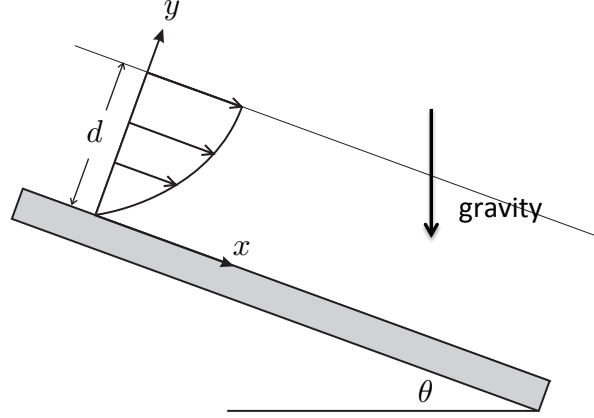
Solution

$$\begin{aligned} \text{shear} &= \text{rotation} + \text{straining} \\ \mathbf{u} &= (\alpha y, 0, 0) = \left(\frac{\alpha y}{2}, -\frac{\alpha x}{2}, 0\right) + \left(\frac{\alpha y}{2}, \frac{\alpha x}{2}, 0\right) \\ \text{velocity gradient tensor} &= \begin{pmatrix} 0 & \alpha \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & \alpha/2 \\ -\alpha/2 & 0 \end{pmatrix} + \begin{pmatrix} 0 & \alpha/2 \\ \alpha/2 & 0 \end{pmatrix} \end{aligned}$$



4. Solving Navier-Stokes: flow down an inclined plane

Consider a steady, two-dimensional, incompressible, viscous flow down an inclined plane under the influence of gravity. Define the axes as shown in the diagram, and assume that the velocity \mathbf{u} depends only on y .



- What are the boundary conditions for \mathbf{u} at $y = 0$? Using incompressibility show that the y -component of the velocity is zero throughout the flow.
- Write down the x - and y -components of the Navier-Stokes equation.
- From the y -component show that the pressure

$$p = p_0 + \rho g(d - y) \cos \theta$$

where p_0 is the pressure at the free surface $y = d$.

- From the x -component, using the appropriate boundary conditions at $y = 0$ and the zero tangential stress condition $\nu du_x/dy = 0$ at the free surface $y = d$ show that

$$u_x = \frac{g}{2\nu} y(2d - y) \sin \theta.$$

- Show that the volume flux per unit distance along z is $gd^3 \sin \theta / (3\nu)$.

Solution

- At $y = 0$ the boundary condition is no-slip, $\mathbf{u} = \mathbf{0}$. By symmetry only y -derivatives are non-zero. Therefore, from the continuity equation, $\partial u_y / \partial y = 0$ so $u_y = \text{constant} = 0$ to match the boundary condition.

- The Navier Stokes equation is

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{1}{\rho} \nabla p = \nu \nabla^2 \mathbf{u} + \mathbf{g}.$$

The flow is time-independent. So the x and y components of the Navier-Stokes equation are

$$0 = \nu \frac{d^2 u_x}{dy^2} + g \sin \theta, \quad (1)$$

$$0 = -\frac{dp}{dy} - \rho g \cos \theta. \quad (2)$$

- (c) Integrate equation (2) with $p = p_0$ at $y = d$.
- (d) Integrate equation (1) with $u_x = 0$ at $y = 0$ and $\nu du_x/dy = 0$ at $y = d$.
- (e) The volume flux per unit distance along z is

$$\int_0^d u_x dy.$$

5. Reynolds number

Estimate the magnitude of the Reynolds number for:

- (a) flow past the wing of a jumbo jet,
- (b) a human swimmer,
- (c) a thick layer of treacle draining off a spoon,
- (d) a bacterium swimming in water.

Take the kinematic viscosity ν to be $10^{-6}\text{m}^2\text{s}^{-1}$ for water, $1.5 \times 10^{-5}\text{m}^2\text{s}^{-1}$ for air and $10^{-1}\text{m}^2\text{s}^{-1}$ for treacle.

Solution

$$\text{Re} = UL/\nu$$

Order of magnitude estimates:

plane	$\text{Re} \sim \frac{150 \times 10}{1.5 \times 10^{-5}} \sim 10^8$
human swimming	$\text{Re} \sim \frac{1 \times 1}{10^{-6}} \sim 10^6$
treacle	$\text{Re} \sim \frac{10^{-3} \times 10^{-2}}{10^{-1}} \sim 10^{-4}$
bacterium swimming	$\text{Re} \sim \frac{10^{-5} \times 10^{-5}}{10^{-6}} \sim 10^{-4}$

6. Dynamical similarity and dimensionless variables

Determine the conditions for the dynamical similarity of steady incompressible flow of an electrically conducting fluid in a magnetic field, governed by the equations

$$\nabla \cdot \mathbf{u} = 0, \tag{3}$$

$$\nabla \cdot \mathbf{B} = 0, \tag{4}$$

$$\mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla p + \frac{1}{\rho \mu} (\nabla \wedge \mathbf{B}) \wedge \mathbf{B} + \nu \nabla^2 \mathbf{u}, \tag{5}$$

$$\mathbf{u} \cdot \nabla \mathbf{B} = \mathbf{B} \cdot \nabla \mathbf{u} + \frac{1}{\sigma \mu} \nabla^2 \mathbf{B}. \tag{6}$$

You will need to define a length scale L , a velocity scale U and a magnetic field scale B_0 . Notation: \mathbf{u} =velocity, \mathbf{B} =magnetic field, p =pressure, ρ =density, ν =kinematic viscosity, μ =magnetic permeability, σ =electrical conductivity.

Comment on the physical meaning of the dimensionless control parameters.

Solution

Define dimensionless variables

$$\tilde{\mathbf{u}} = \frac{\mathbf{u}}{U}, \quad \tilde{\mathbf{B}} = \frac{\mathbf{B}}{B_0}, \quad \tilde{\mathbf{x}} = \frac{\mathbf{x}}{L}.$$

Substitute into equations (3)–(6) to give

$$\begin{aligned} \nabla \cdot \tilde{\mathbf{u}} &= 0, \\ \nabla \cdot \tilde{\mathbf{B}} &= 0, \\ \tilde{\mathbf{u}} \cdot \tilde{\nabla} \tilde{\mathbf{u}} &= -\frac{1}{U^2 \rho} \tilde{\nabla} p + \frac{B_0^2}{\rho \mu U^2} (\tilde{\nabla} \wedge \tilde{\mathbf{B}}) \wedge \tilde{\mathbf{B}} + \frac{\nu}{LU} \tilde{\nabla}^2 \tilde{\mathbf{u}}, \\ \tilde{\mathbf{u}} \cdot \tilde{\nabla} \tilde{\mathbf{B}} &= \tilde{\mathbf{B}} \cdot \tilde{\nabla} \tilde{\mathbf{u}} + \frac{1}{\sigma \mu LU} \tilde{\nabla}^2 \tilde{\mathbf{B}}. \end{aligned}$$

Dimensionless ratios are:

UL/ν Reynolds number; ratio of inertial to viscous terms in N-S;

$B_0^2/\mu\rho U^2$ ratio of magnetic energy per unit volume to kinetic energy per unit volume; measures the importance of magnetic effects in controlling the flow;

$\sigma\mu UL$ magnetic Reynolds number.

7. More dynamical similarity

(a) A flat plate of width L is placed at a right angle to the flow in a wind tunnel, in which the upstream wind speed is U .

Show that the expected scaling of the pressure variations is

$$(i) \Delta p \sim \frac{\rho \nu U}{L} \text{ in the limit } Re \ll 1,$$

$$(ii) \Delta p \sim \rho U^2 \text{ in the limit } Re \gg 1.$$

(b) In the wind tunnel vortices are shed behind the plate at a frequency of 0.5 s^{-1} . The same plate is now placed into a water channel. Calculate the flow rate required, as a multiple of that in the wind tunnel, to produce dynamically similar behaviour, and calculate the frequency of the vortex shedding.

Solution

(a) The Navier Stokes equation is

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{1}{\rho} \nabla p = \nu \nabla^2 \mathbf{u}.$$

For $Re \ll 1$ the pressure term is balanced by the viscous term so

$$\frac{\Delta p}{\rho L} \sim \frac{\nu U}{L^2} \quad \rightarrow \quad \Delta p \sim \frac{\rho \nu U}{L}.$$

For $Re \gg 1$ the pressure term is balanced by the inertia so

$$\frac{\Delta p}{\rho L} \sim \frac{U^2}{L} \rightarrow \Delta p \sim \rho U^2.$$

(b) The Reynolds number is UL/ν . L is unchanged so for similar flows we require

$$\frac{U_{water}}{U_{wind}} = \frac{\nu_{water}}{\nu_{wind}} = 0.067.$$

frequency $\sim U/L$ so

$$\text{freq}_{water} = \frac{U_{water}}{U_{wind}} \text{freq}_{air} = 0.034.$$

8. Vorticity

- (a) What is meant by the vorticity of a fluid flow? Illustrate your answer by discussing:
 (i) a rectilinear flow that has vorticity eg simple shear (again) $\mathbf{u} = (\alpha y, 0, 0)$.
 (ii) a rotating flow that does not have vorticity eg $u_r = 0, u_\theta = A/r$ (in plane polar co-ordinates).
 (b) What is the vorticity of a flow in rigid body motion with angular velocity $\boldsymbol{\Omega}$?

Solution

The vorticity, $\boldsymbol{\omega} = \nabla \wedge \mathbf{v}$.

- (a)(i) $\boldsymbol{\omega} = -\alpha \hat{\mathbf{z}}$
 (ii) $\boldsymbol{\omega} = 0$ ($r \neq 0$)

$$(b) \mathbf{u} = \boldsymbol{\Omega} \wedge \mathbf{r} \Rightarrow \boldsymbol{\omega} = \nabla \wedge (\boldsymbol{\Omega} \wedge \mathbf{r}) = (\mathbf{r} \cdot \nabla) \boldsymbol{\Omega} - (\boldsymbol{\Omega} \cdot \nabla) \mathbf{r} + \boldsymbol{\Omega} (\nabla \cdot \mathbf{r}) - \mathbf{r} (\nabla \cdot \boldsymbol{\Omega}) = 2\boldsymbol{\Omega}$$

using $\boldsymbol{\Omega}$ is constant, $(\nabla \cdot \mathbf{r}) = 3$ and $(\boldsymbol{\Omega} \cdot \nabla) \mathbf{r} = \boldsymbol{\Omega}$.