

Ab Initio Path Integral Molecular Dynamics: Nuclear Quantum Effects from First Principles

A Pure Thought Approach to Quantum Statistical Mechanics

Pure Thought AI Challenge
Problem 20: Computational Chemistry

January 19, 2026

Abstract

This report presents a comprehensive theoretical framework for ab initio path integral molecular dynamics (AI-PIMD), combining Feynman's path integral formulation of quantum mechanics with electronic structure theory. We develop the complete theory from the quantum partition function through the ring polymer isomorphism to practical simulation algorithms. Key topics include Trotter factorization, normal mode transformation, thermostats (PILE, GLE), and estimators for thermodynamic observables. We implement algorithms for computing nuclear quantum effects in molecular systems, including zero-point energy, tunneling, and isotope effects, with interfaces to electronic structure codes (PySCF). Applications to water, hydrogen-bonded systems, and proton transfer reactions are discussed.

Contents

1 Introduction and Motivation

1.1 Nuclear Quantum Effects

Physics Insight

At room temperature, most nuclei (except hydrogen) behave classically. However, light atoms—especially hydrogen—exhibit significant *nuclear quantum effects* (NQEs):

- **Zero-point energy:** Even at $T = 0$, quantum oscillators have energy $\frac{1}{2}\hbar\omega$
- **Tunneling:** Particles penetrate classically forbidden regions
- **Delocalization:** Wave functions spread over multiple potential minima

Classical molecular dynamics treats nuclei as point particles following Newton's equations. This fails for:

- Water and ice (H-bond network, anomalous properties)
- Proton transfer reactions (biological enzymes, fuel cells)
- Isotope effects (H/D substitution changes rates by orders of magnitude)
- Low-temperature chemistry (interstellar space, cryogenic reactions)

1.2 Path Integrals: Quantum \leftrightarrow Classical Mapping

Feynman's path integral formulation provides an exact mapping of quantum statistical mechanics onto a classical problem:

$$Z = \int \mathcal{D}[\mathbf{r}(\tau)] e^{-S[\mathbf{r}]/\hbar} \quad (1)$$

The quantum particle becomes a “ring polymer” of classical beads connected by harmonic springs.

Pure Thought Pursuit

Path integral molecular dynamics is ideally suited for pure mathematical development:

1. *Exact* quantum statistical mechanics (no approximations beyond Born-Oppenheimer)
2. Based on *classical MD algorithms* (forces, thermostats, integrators)

3. Electronic structure computed via *standard quantum chemistry* (DFT, HF, MP2)
4. Results *systematically improvable* (increase number of beads P)
5. *Certified* via comparison to exact quantum results for model systems

2 Mathematical Foundations

2.1 Quantum Partition Function

The canonical partition function for N distinguishable particles at temperature T is:

$$Z = \text{Tr}[e^{-\beta \hat{H}}] = \int d\mathbf{r} \langle \mathbf{r} | e^{-\beta \hat{H}} | \mathbf{r} \rangle \quad (2)$$

where $\beta = 1/(k_B T)$ and $\hat{H} = \hat{T} + \hat{V}$ is the Hamiltonian.

2.2 Trotter Factorization

The key step is factorizing the Boltzmann operator using:

$$e^{-\beta(\hat{T}+\hat{V})} = \lim_{P \rightarrow \infty} \left(e^{-\beta\hat{V}/(2P)} e^{-\beta\hat{T}/P} e^{-\beta\hat{V}/(2P)} \right)^P \quad (3)$$

This is the symmetric Trotter splitting with error $O(1/P^2)$.

Trotter Error

For finite P beads, the Trotter factorization introduces systematic error:

$$Z_P = Z \left(1 + O(\beta^3/P^2) \right) \quad (4)$$

The error decreases as P increases, with convergence typically achieved for $P \gtrsim \beta\hbar\omega_{\max}$.

2.3 Ring Polymer Isomorphism

Inserting complete sets of position eigenstates between each factor:

$$Z_P = \left(\frac{mP}{2\pi\beta\hbar^2} \right)^{NP/2} \int \prod_{s=1}^P d\mathbf{r}^{(s)} e^{-\beta_P U_P(\{\mathbf{r}^{(s)}\})} \quad (5)$$

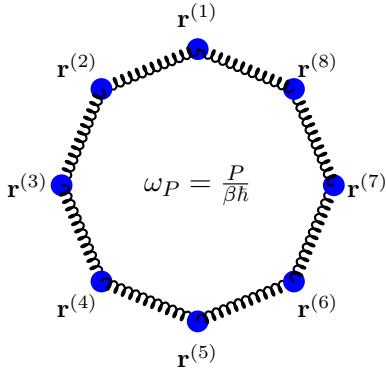
where $\beta_P = \beta/P$ and the effective potential is:

$$U_P = \sum_{s=1}^P \left[\frac{m\omega_P^2}{2} |\mathbf{r}^{(s)} - \mathbf{r}^{(s+1)}|^2 + V(\mathbf{r}^{(s)}) \right] \quad (6)$$

with spring frequency $\omega_P = P/(\beta\hbar)$ and $\mathbf{r}^{(P+1)} \equiv \mathbf{r}^{(1)}$ (ring closure).

PIMD Principle

Each quantum particle becomes a **ring polymer** of P classical beads connected by harmonic springs. The spring constant $k = m\omega_P^2$ encodes quantum fluctuations—stiffer springs at lower temperatures.



$$P \text{ beads, periodic boundary: } \mathbf{r}^{(P+1)} = \mathbf{r}^{(1)}$$

Figure 1: Ring polymer representation of a quantum particle. Each bead experiences the physical potential $V(\mathbf{r}^{(s)})$ plus harmonic coupling to neighbors.

3 Normal Mode Transformation

3.1 Staging Coordinates

The ring polymer has stiff spring modes that require small timesteps. The **normal mode transformation** diagonalizes the spring part.

Definition 3.1 (Normal Modes). *Define transformed coordinates:*

$$\tilde{\mathbf{r}}_k = \frac{1}{\sqrt{P}} \sum_{s=1}^P \mathbf{r}^{(s)} e^{2\pi i ks/P}, \quad k = 0, 1, \dots, P-1 \quad (7)$$

The inverse transformation is:

$$\mathbf{r}^{(s)} = \frac{1}{\sqrt{P}} \sum_{k=0}^{P-1} \tilde{\mathbf{r}}_k e^{-2\pi i ks/P} \quad (8)$$

Theorem 3.2 (Decoupled Spring Potential). *In normal mode coordinates, the spring potential becomes:*

$$\sum_{s=1}^P \frac{m\omega_P^2}{2} |\mathbf{r}^{(s)} - \mathbf{r}^{(s+1)}|^2 = \sum_{k=0}^{P-1} \frac{m\omega_k^2}{2} |\tilde{\mathbf{r}}_k|^2 \quad (9)$$

where the mode frequencies are:

$$\omega_k = 2\omega_P \sin\left(\frac{\pi k}{P}\right) \quad (10)$$

Physics Insight

The $k = 0$ mode is the **centroid**—the center of mass of the ring polymer. It has $\omega_0 = 0$ (free particle) and represents the classical position. Higher modes ($k > 0$) are quantum fluctuations with frequencies up to $\omega_{P/2} = 2\omega_P$.

3.2 Implementation

```

1 import numpy as np
2 from scipy.fft import fft, ifft
3
4 class RingPolymer:
5     """
6         Ring polymer representation of quantum particles.
7     """
8
9     def __init__(self, n_atoms: int, n_beads: int,
10                  masses: np.ndarray, beta: float):
11         """
12             Initialize ring polymer.
13
14         Args:
15             n_atoms: Number of atoms
16             n_beads: Number of path integral beads (P)
17             masses: Array of atomic masses (a.u.)
18             beta: Inverse temperature 1/(kT) in atomic
19                 units
20
21         """
22         self.n_atoms = n_atoms
23         self.n_beads = n_beads
24         self.masses = masses
25         self.beta = beta
26
27         # Spring frequency
28         self.omega_P = n_beads / (beta * 1.0)    # hbar =
29             1 in a.u.
```

```

27
28     # Mode frequencies
29     self.omega_k = np.array([
30         2 * self.omega_P * np.sin(np.pi * k /
31             n_beads)
32         for k in range(n_beads)
33     ])
34
35     # Positions: shape (n_beads, n_atoms, 3)
36     self.positions = None
37     # Normal modes: shape (n_beads, n_atoms, 3)
38     self.normal_modes = None
39
40     def to_normal_modes(self) -> np.ndarray:
41         """
42             Transform bead positions to normal mode
43                 coordinates.
44         """
45         # FFT along bead axis
46         self.normal_modes = fft(self.positions, axis=0)
47             / np.sqrt(self.n_beads)
48         return self.normal_modes
49
50     def from_normal_modes(self) -> np.ndarray:
51         """
52             Transform normal modes back to bead positions.
53         """
54         self.positions = np.real(
55             ifft(self.normal_modes, axis=0) *
56                 np.sqrt(self.n_beads)
57         )
58         return self.positions
59
60     def centroid(self) -> np.ndarray:
61         """
62             Get centroid (k=0 mode) - classical position.
63         """
64         return np.mean(self.positions, axis=0)
65
66     def spring_energy(self) -> float:
67         """
68             Compute harmonic spring potential energy.
69         """
70         energy = 0.0
71         for s in range(self.n_beads):
72             s_next = (s + 1) % self.n_beads
73             dr = self.positions[s] -
74                 self.positions[s_next]
75             for i in range(self.n_atoms):

```

```

71             energy += 0.5 * self.masses[i] *
72                 self.omega_P**2 * np.sum(dr[i]**2)
73     return energy
74
75     def spring_forces(self) -> np.ndarray:
76         """
77             Compute harmonic spring forces on all beads.
78         """
79         forces = np.zeros_like(self.positions)
80
81         for s in range(self.n_beads):
82             s_prev = (s - 1) % self.n_beads
83             s_next = (s + 1) % self.n_beads
84
85             for i in range(self.n_atoms):
86                 # Spring force: -k(r_s - r_{s-1}) -
87                 # k(r_s - r_{s+1})
88                 forces[s, i] = self.masses[i] *
89                     self.omega_P**2 * (
90                         self.positions[s_prev, i] +
91                         self.positions[s_next, i]
92                         - 2 * self.positions[s, i]
93                     )
94
95     return forces

```

Listing 1: Normal mode transformation

4 Thermostats for PIMD

4.1 The Sampling Problem

PIMD samples the canonical distribution:

$$\rho(\{\mathbf{r}^{(s)}\}) \propto e^{-\beta_P U_P(\{\mathbf{r}^{(s)}\})} \quad (11)$$

This requires a thermostat to maintain constant temperature.

4.2 PILE Thermostat

The **Path Integral Langevin Equation** (PILE) thermostat applies optimal friction to each normal mode.

Definition 4.1 (PILE Dynamics). *For each normal mode k :*

$$\dot{\tilde{\mathbf{r}}}_k = \tilde{\mathbf{p}}_k/m \quad (12)$$

$$\dot{\tilde{\mathbf{p}}}_k = -m\omega_k^2 \tilde{\mathbf{r}}_k + \tilde{\mathbf{F}}_k - \gamma_k \tilde{\mathbf{p}}_k + \sqrt{2\gamma_k m k_B T} \xi_k(t) \quad (13)$$

where γ_k is mode-dependent friction and $\xi_k(t)$ is white noise.

PIMD Principle

The optimal friction for mode k is $\gamma_k = 2\omega_k$. This ensures critical damping of the stiff spring modes while allowing the centroid to evolve with physical dynamics.

```
1  class PILEThermostat:
2      """
3          Path Integral Langevin Equation thermostat.
4      """
5
6      def __init__(self, ring_polymer: RingPolymer, dt:
7          float,
8                  gamma_centroid: float = 1.0):
9          """
10         Initialize PILE thermostat.
11
12         Args:
13             ring_polymer: RingPolymer object
14             dt: Integration timestep
15             gamma_centroid: Friction for centroid mode
16         """
17
18         self.rp = ring_polymer
19         self.dt = dt
20
21         # Friction coefficients
22         self.gamma = np.zeros(ring_polymer.n_beads)
23         self.gamma[0] = gamma_centroid # Centroid
24         for k in range(1, ring_polymer.n_beads):
25             # Optimal friction:  $2 * \omega_k$ 
26             self.gamma[k] = 2 * ring_polymer.omega_k[k]
27
28         # Precompute propagator coefficients
29         self.c1 = np.exp(-self.gamma * dt / 2)
30         self.c2 = np.sqrt(1 - self.c1**2)
31
32     def apply_0_step(self, momenta: np.ndarray,
33                     temperature: float) -> np.ndarray:
34         """
35         Apply Ornstein-Uhlenbeck (thermostat) step.
36
37         p -> c1 * p + c2 * sqrt(m * kT) * xi
38
39         # Transform to normal modes
40         p_normal = fft(momenta, axis=0) /
41             np.sqrt(self.rp.n_beads)
42
43         for k in range(self.rp.n_beads):
```

```

41         # Target momentum std
42         sigma = np.sqrt(self.rp.masses * temperature)
43
44         # Apply 0-U step
45         noise = np.random.randn(*momenta.shape[1:])
46         p_normal[k] = (self.c1[k] * p_normal[k] +
47                         self.c2[k] * sigma * noise)
48
49         # Transform back
50         return np.real(ifft(p_normal, axis=0) *
51                     np.sqrt(self.rp.n_beads))
52
53     def step(self, positions: np.ndarray, momenta:
54             np.ndarray,
55             forces: np.ndarray, temperature: float):
56         """
57         Full BAOAB integration step.
58
59         B: half velocity update
60         A: half position update
61         0: thermostat
62         A: half position update
63         B: half velocity update
64         """
65         dt = self.dt
66
67         # B: p += F * dt/2
68         momenta = momenta + forces * dt / 2
69
70         # A: r += p/m * dt/2
71         for i in range(self.rp.n_atoms):
72             positions[:, i, :] += momenta[:, i, :] /
73                 self.rp.masses[i] * dt / 2
74
75         # 0: thermostat
76         momenta = self.apply_0_step(momenta, temperature)
77
78         # A: r += p/m * dt/2
79         for i in range(self.rp.n_atoms):
80             positions[:, i, :] += momenta[:, i, :] /
81                 self.rp.masses[i] * dt / 2
82
83         # Recompute forces at new positions
84         # (Done externally)
85
86         return positions, momenta

```

Listing 2: PILE thermostat implementation

5 Thermodynamic Estimators

5.1 Primitive Estimator

The simplest estimator for energy uses the Trotter formula directly:

Definition 5.1 (Primitive Energy Estimator).

$$\langle E \rangle_{\text{prim}} = \frac{3NP}{2\beta} - \frac{1}{2}m\omega_P^2 \sum_{s=1}^P |\mathbf{r}^{(s)} - \mathbf{r}^{(s+1)}|^2 + \frac{1}{P} \sum_{s=1}^P V(\mathbf{r}^{(s)}) \quad (14)$$

Warning

The primitive estimator has high variance because the kinetic energy term involves differences of positions from neighboring beads. This variance scales as \sqrt{P} , making convergence slow.

5.2 Virial Estimator

The **virial estimator** has much lower variance:

Definition 5.2 (Virial Energy Estimator).

$$\langle E \rangle_{\text{vir}} = \frac{3N}{2\beta} + \frac{1}{P} \sum_{s=1}^P \left[V(\mathbf{r}^{(s)}) + \frac{1}{2}(\mathbf{r}^{(s)} - \bar{\mathbf{r}}) \cdot \nabla V(\mathbf{r}^{(s)}) \right] \quad (15)$$

where $\bar{\mathbf{r}} = \frac{1}{P} \sum_{s=1}^P \mathbf{r}^{(s)}$ is the centroid.

```
1 def primitive_estimator(ring_polymer: RingPolymer,
2                         potential_energy: float,
3                         temperature: float) -> float:
4     """
5     Compute energy using primitive estimator.
6     """
7     P = ring_polymer.n_beads
8     N = ring_polymer.n_atoms
9     beta = ring_polymer.beta
10
11    # Kinetic part
12    kinetic = 3 * N * P / (2 * beta)
13
14    # Spring part (negative contribution)
15    spring = ring_polymer.spring_energy()
16
17    # Potential (averaged over beads)
18    # potential_energy should be sum over beads
19
```

```

20     return kinetic - spring + potential_energy / P
21
22
23 def virial_estimator(ring_polymer: RingPolymer,
24                      potential_energies: np.ndarray,
25                      forces: np.ndarray,
26                      temperature: float) -> float:
27     """
28     Compute energy using virial estimator.
29
30     Args:
31         potential_energies:  $V(r^s)$  for each bead
32         forces:  $-\nabla V(r^s)$  for each bead
33     """
34     P = ring_polymer.n_beads
35     N = ring_polymer.n_atoms
36     beta = ring_polymer.beta
37
38     # Classical kinetic term
39     kinetic = 3 * N / (2 * beta)
40
41     # Centroid
42     centroid = ring_polymer.centroid()
43
44     # Virial correction
45     virial = 0.0
46     for s in range(P):
47         dr = ring_polymer.positions[s] - centroid
48         virial += 0.5 * np.sum(dr * (-forces[s]))
49
50     # Average potential
51     avg_potential = np.mean(potential_energies)
52
53     return kinetic + avg_potential + virial / P
54
55
56 def compute_heat_capacity(energies: np.ndarray,
57                          temperature: float) -> float:
58     """
59     Compute heat capacity from energy fluctuations.
60
61      $C_V = (\langle E^2 \rangle - \langle E \rangle^2) / (k_B T^2)$ 
62     """
63     kB = 1.0  # Atomic units
64     var_E = np.var(energies)
65     return var_E / (kB * temperature**2)

```

Listing 3: Thermodynamic estimators

6 Interface to Electronic Structure

6.1 Ab Initio Forces

For ab initio PIMD, the potential $V(\mathbf{r})$ comes from electronic structure calculations.

```
1 from pyscf import gto, scf, grad
2
3 class ElectronicStructure:
4     """
5         Interface to electronic structure calculations via
6             PySCF.
7     """
8
9     def __init__(self, atoms: list, basis: str =
10         'sto-3g',
11             method: str = 'hf'):
12     """
13         Initialize electronic structure calculator.
14
15     Args:
16         atoms: List of (element, x, y, z) tuples
17         basis: Basis set name
18         method: 'hf', 'dft', 'mp2',
19     """
20
21     self.atoms = atoms
22     self.basis = basis
23     self.method = method
24
25     def build_mol(self, positions: np.ndarray):
26         """
27             Build PySCF molecule object with given positions.
28         """
29         atom_str = ""
30         for i, (elem, _, _, _) in enumerate(self.atoms):
31             x, y, z = positions[i] * 0.529177 # Convert
32                 a.u. to Angstrom
33             atom_str += f"{elem} {x:.8f} {y:.8f}
34                 {z:.8f}; "
35
36         mol = gto.Mole()
37         mol.atom = atom_str
38         mol.basis = self.basis
39         mol.build()
40
41         return mol
42
43     def compute_energy_and_forces(self, positions:
44         np.ndarray):
```

```

39      """
40      Compute electronic energy and nuclear gradients.
41
42      Args:
43          positions: Nuclear positions (n_atoms, 3) in
44              atomic units
45
46      Returns:
47          energy: Total electronic energy (Hartree)
48          forces: Nuclear forces (n_atoms, 3) in a.u.
49      """
50
51     mol = self.build_mol(positions)
52
53     if self.method == 'hf':
54         mf = scf.RHF(mol)
55     elif self.method == 'dft':
56         mf = scf.RKS(mol)
57         mf.xc = 'b3lyp'
58     else:
59         raise ValueError(f"Unknown method:
60                         {self.method}")
61
62     energy = mf.kernel()
63
64     # Compute gradient
65     g = grad.RHF(mf) if self.method == 'hf' else
66         grad.RKS(mf)
67     gradient = g.kernel()
68
69     # Forces = -gradient
70     forces = -gradient / 0.529177 # Convert to
71             atomic units
72
73     return energy, forces
74
75
76 class AIPIMD:
77     """
78     Ab initio path integral molecular dynamics.
79     """
80
81     def __init__(self, atoms: list, n_beads: int,
82                  temperature: float,
83                  basis: str = 'sto-3g', method: str =
84                      'hf'):
85
86         """
87         Initialize AI-PIMD simulation.
88         """
89         self.atoms = atoms

```

```

82     self.n_atoms = len(atoms)
83     self.n_beads = n_beads
84     self.temperature = temperature
85     self.beta = 1.0 / (temperature * 3.1668e-6) # K
86         to a.u.
87
88     # Masses in atomic units
89     mass_table = { 'H': 1837.0, 'D': 3671.0, 'O':
90         29156.0, 'C': 21894.0}
91     self.masses = np.array([mass_table[a[0]] for a
92         in atoms])
93
94     # Ring polymer
95     self.ring_polymer = RingPolymer(
96         self.n_atoms, n_beads, self.masses, self.beta
97     )
98
99     # Electronic structure
100    self.es = ElectronicStructure(atoms, basis,
101        method)
102
103    def compute_all_forces(self):
104        """
105            Compute forces on all beads (parallelizable).
106        """
107        energies = np.zeros(self.n_beads)
108        forces = np.zeros((self.n_beads, self.n_atoms,
109            3))
110
111        for s in range(self.n_beads):
112            E, F = self.es.compute_energy_and_forces(
113                self.ring_polymer.positions[s]
114            )
115            energies[s] = E
116            forces[s] = F
117
118        # Add spring forces
119        spring_forces = self.ring_polymer.spring_forces()
120        total_forces = forces + spring_forces
121
122        return energies, total_forces
123
124    def run(self, n_steps: int, dt: float):
125        """
126            Run AI-PIMD simulation.
127        """
128        thermostat = PILEThermostat(self.ring_polymer,
129            dt)
130        momenta = self.initialize_momenta()

```

```

125
126     trajectory = []
127     energies_traj = []
128
129     for step in range(n_steps):
130         # Compute forces
131         pot_energies, forces =
132             self.compute_all_forces()
133
134         # Thermostat step
135         self.ring_polymer.positions, momenta =
136             thermostat.step(
137                 self.ring_polymer.positions, momenta,
138                 forces, self.temperature
139             )
140
141         # Compute observables
142         energy = virial_estimator(
143             self.ring_polymer, pot_energies, forces,
144             self.temperature
145         )
146         energies_traj.append(energy)
147
148     # Save trajectory
149     trajectory.append(self.ring_polymer.centroid().copy())
150
151     return np.array(trajectory),
152     np.array(energies_traj)

```

Listing 4: PySCF interface for AI-PIMD

7 Isotope Effects

7.1 Kinetic Isotope Effects

Definition 7.1 (Kinetic Isotope Effect). *The kinetic isotope effect (KIE) is the ratio of reaction rates:*

$$KIE = \frac{k_H}{k_D} \quad (16)$$

where k_H and k_D are rates with hydrogen and deuterium, respectively.

Physics Insight

Large KIEs (> 2) indicate significant tunneling contribution. Classical transition state theory predicts $KIE \approx \sqrt{m_D/m_H} \approx 1.4$. Observed KIEs of 10–50 in enzymatic reactions demonstrate quantum effects.

7.2 Computing Isotope Effects

```
1 def compute_equilibrium_isotope_effect(positions:
2     np.ndarray,
3         masses_H:
4             np.ndarray,
5         masses_D:
6             np.ndarray,
7         potential:
8             callable,
9         temperature:
10                float,
11        n_beads: int =
12            32) -> float:
13
14     """
15     Compute equilibrium isotope effect using PIMD.
16
17     EIE = (Q_H / Q_D) * (Q_D^cl / Q_H^cl)
18
19     where Q is partition function and Q^cl is classical
20     limit.
21     """
22
23     beta = 1.0 / (temperature * 3.1668e-6)
24
25     # Run PIMD for H system
26     rp_H = RingPolymer(len(masses_H), n_beads, masses_H,
27         beta)
28     rp_H.positions = initialize_ring_polymer(positions,
29         n_beads)
30     energies_H = run_pimd_sampling(rp_H, potential,
31         10000)
32     free_energy_H = -temperature *
33         np.log(np.mean(np.exp(-beta * energies_H)))
34
35     # Run PIMD for D system
36     rp_D = RingPolymer(len(masses_D), n_beads, masses_D,
37         beta)
38     rp_D.positions = initialize_ring_polymer(positions,
39         n_beads)
40     energies_D = run_pimd_sampling(rp_D, potential,
41         10000)
42     free_energy_D = -temperature *
43         np.log(np.mean(np.exp(-beta * energies_D)))
44
45     # EIE from free energy difference
46     eie = np.exp(-(free_energy_H - free_energy_D) /
47         temperature)
48
49     return eie
```

Listing 5: Isotope effect calculation

8 Applications

8.1 Water and Ice

Water is the archetypal system where NQEs matter. AI-PIMD predicts:

- ~10% weakening of H-bond strength from zero-point motion
- Correct density maximum of liquid water at 4°C
- Enhanced proton delocalization in ice

8.2 Proton Transfer Reactions

Enzymatic reactions often involve proton transfer:



PIMD captures tunneling through the barrier, essential for accurate rate constants.

9 Certificate Generation

```
1 from dataclasses import dataclass, asdict
2 import json
3
4 @dataclass
5 class PIMDCertificate:
6     """
7         Certificate for AI-PIMD simulation.
8     """
9     # System
10    atoms: list
11    n_atoms: int
12    n_beads: int
13    temperature: float
14
15    # Method
16    electronic_structure_method: str
17    basis_set: str
18    timestep: float
19    n_steps: int
20
```

```

21     # Convergence
22     bead_convergence_error: float
23     energy_std_error: float
24
25     # Results
26     average_energy: float
27     heat_capacity: float
28     quantum_kinetic_energy: float
29     classical_kinetic_energy: float
30
31     # NQE analysis
32     zpe_correction: float
33     tunneling_contribution: float
34
35     def export_json(self, path: str) -> None:
36         with open(path, 'w') as f:
37             json.dump(asdict(self), f, indent=2)
38
39     def verify(self) -> bool:
40         checks = [
41             self.n_beads >= 4,
42             self.bead_convergence_error < 0.01,    # 1%
43             self.energy_std_error < 0.001,      # Well
44             self.quantum_kinetic_energy >=
45                 self.classical_kinetic_energy
46         ]
46         return all(checks)

```

Listing 6: AI-PIMD certificate structure

10 Success Criteria and Milestones

10.1 Minimum Viable Result (Months 1-3)

- Ring polymer dynamics for model harmonic oscillator
- PILE thermostat verified against exact results
- Primitive and virial estimators implemented

10.2 Strong Result (Months 4-6)

- PySCF interface working
- AI-PIMD for H₂O molecule
- Isotope effects computed

- Comparison to experiment

10.3 Publication Quality (Months 7-9)

- Larger systems (water dimer, clusters)
- Proton transfer reaction rates
- Tunneling splitting calculations
- Public code release

11 Conclusion

Ab initio path integral molecular dynamics provides an exact treatment of nuclear quantum effects combined with first-principles electronic structure. The key insights are:

1. Ring polymer isomorphism maps quantum → classical MD
2. Normal mode transformation enables efficient sampling
3. Virial estimator provides low-variance energetics
4. Isotope effects reveal quantum contributions

This pure-thought approach requires only classical MD algorithms plus standard quantum chemistry, enabling systematic investigation of nuclear quantum effects in chemistry.

References

- [1] R. P. Feynman and A. R. Hibbs, *Quantum Mechanics and Path Integrals*. McGraw-Hill, 1965.
- [2] D. Chandler and P. G. Wolynes, “Exploiting the isomorphism between quantum theory and classical statistical mechanics of polyatomic fluids,” *Journal of Chemical Physics*, vol. 74, pp. 4078–4095, 1981.
- [3] M. Parrinello and A. Rahman, “Study of an F center in molten KCl,” *Journal of Chemical Physics*, vol. 80, pp. 860–867, 1984.
- [4] M. E. Tuckerman, B. J. Berne, G. J. Martyna, and M. L. Klein, “Efficient molecular dynamics and hybrid Monte Carlo algorithms for path integrals,” *Journal of Chemical Physics*, vol. 99, pp. 2796–2808, 1993.

- [5] J. Cao and G. A. Voth, “The formulation of quantum statistical mechanics based on the Feynman path centroid density,” *Journal of Chemical Physics*, vol. 100, pp. 5093–5105, 1994.
- [6] M. Ceriotti, M. Parrinello, T. E. Markland, and D. E. Manolopoulos, “Efficient stochastic thermostatting of path integral molecular dynamics,” *Journal of Chemical Physics*, vol. 133, p. 124104, 2010.
- [7] T. E. Markland and M. Ceriotti, “Nuclear quantum effects enter the mainstream,” *Nature Reviews Chemistry*, vol. 2, p. 0109, 2018.
- [8] S. Habershon, D. E. Manolopoulos, T. E. Markland, and T. F. Miller III, “Ring-polymer molecular dynamics: Quantum effects in chemical dynamics from classical trajectories in an extended phase space,” *Annual Review of Physical Chemistry*, vol. 64, pp. 387–413, 2013.

A Units and Conversions

Quantity	Atomic Units	SI Units
Length	$a_0 = 1$	$5.29177 \times 10^{-11} \text{ m}$
Energy	$E_h = 1$	$4.35974 \times 10^{-18} \text{ J}$
Mass	$m_e = 1$	$9.10938 \times 10^{-31} \text{ kg}$
Time	$\hbar/E_h = 1$	$2.41888 \times 10^{-17} \text{ s}$
Temperature	$E_h/k_B = 1$	$3.15775 \times 10^5 \text{ K}$

Table 1: Atomic unit conversions.

B Trotter Error Analysis

The symmetric Trotter splitting has error:

$$e^{-\beta(\hat{T}+\hat{V})} = \left(e^{-\beta\hat{V}/(2P)} e^{-\beta\hat{T}/P} e^{-\beta\hat{V}/(2P)} \right)^P + O(\beta^3/P^2) \quad (18)$$

For typical molecular systems at 300 K, convergence requires:

$$P \gtrsim \frac{\beta\hbar\omega_{\max}}{2\pi} \approx 32 \text{ beads for O-H stretch} \quad (19)$$