

(1D)

- Solución aproximada del modelo de almacenamiento térmico con intercambio de calor

El modelo de Schumann (1929) es una aproximación sencilla para el intercambio de calor entre un fluido que atraviesa un medio poroso a una determinada temperatura:

$$\left\{ \begin{array}{l} \epsilon \left(\frac{\partial T_f}{\partial t} + u_0 \frac{\partial T_f}{\partial x} \right) = \epsilon \alpha_f \frac{\partial^2 T_f}{\partial x^2} - \frac{h_v}{\rho_f C_f} (T_f - T_s), \\ (1-\epsilon) \frac{\partial T_s}{\partial t} = (1-\epsilon) \alpha_s \frac{\partial^2 T_s}{\partial x^2} + \frac{h_v}{\rho_s C_s} (T_f - T_s) \end{array} \right.$$

Con las condiciones de contorno:

$$T_f(x=0, t) = T_{in} \quad (\text{d.f.})$$

$$\frac{\partial T_f}{\partial x}(x=0, t) = 0$$

$$\frac{\partial T_s}{\partial x}(x=0, t) = \frac{\partial T_s}{\partial x}(x=1, t) = 0$$

y la condición inicial:

$$T_f(x, t=0) = T_s(x, t=0) = T_0$$

En el sistema. egs.,

α → difusividad térmica (m^2/s)

h_v → eff. de intercambio de calor ($\frac{W}{m^2 K}$)

ϵ → Porosidad

u_0 → Velocidad del flujo.

T_f, T_s → Temperaturas del fluido y del sólido.

Adimensionalización:

$$\eta = \frac{x}{L}, \quad \tau = \frac{t}{t_c}, \quad \theta_j = \frac{T_j - T_0}{T_0}$$

$$(1) \quad \epsilon \left(\frac{1}{t_c} \frac{\partial \theta_j}{\partial \tau} + \frac{u_0}{L} \frac{\partial \theta_j}{\partial \eta} \right) = \epsilon \frac{\alpha_j}{L^2} \frac{\partial^2 \theta_j}{\partial \eta^2} - \frac{h_v}{\rho_j C_p} (\theta_j - \theta_s)$$

$$(2) \quad (1-\epsilon) \frac{1}{t_c} \frac{\partial \theta_s}{\partial \tau} = (1-\epsilon) \frac{\alpha_s}{L^2} \frac{\partial^2 \theta_s}{\partial \eta^2} + \frac{h_v}{\rho_s C_p} (\theta_j - \theta_s)$$

$$\theta_j(0, \tau) = \frac{T_{in}}{T_0} - 1$$

$$\frac{\partial \theta_j}{\partial \eta}(1, \tau) = 0, \quad \frac{\partial \theta_s}{\partial \eta}(0, \tau) = \frac{\partial \theta_s}{\partial \eta}(1, \tau) = 0.$$

$$\theta_j(\eta, 0) = \theta_s(\eta, 0) = 0$$

$$\frac{h_v L}{\rho_j C_p u_0} \equiv \Lambda_j, \quad \frac{h_v L}{\rho_s C_p u_0} \equiv \Lambda_s.$$

El tiempo característico lo tomo para que sea del orden del término convectivo.

$$t_c = \frac{L}{u_0}$$

$$\epsilon \left(\frac{\partial \theta_j}{\partial \tau} + \frac{\partial \theta_s}{\partial \eta} \right) = \epsilon \left(\frac{\alpha_j L}{L u_0} \frac{\partial^2 \theta_j}{\partial \eta^2} - \Lambda_j (\theta_j - \theta_s) \right)$$

$$(1-\epsilon) \frac{\partial \theta_s}{\partial \tau} = (1-\epsilon) \left(\frac{\alpha_s L}{L u_0} \frac{\partial^2 \theta_s}{\partial \eta^2} + \Lambda_s (\theta_j - \theta_s) \right)$$

$$\frac{\alpha_j}{L u_0} \sim O(10^{-4}), \quad \frac{\alpha_s}{L u_0} \sim O(10^{-4})$$

$$\Lambda_j \sim O(10^2), \quad \Lambda_s \sim O(10^{-1})$$

En primera aprox. tomamos $\frac{\alpha_j}{L u_0} \sim \frac{\alpha_s}{L u_0} \sim \epsilon \ll 1$

sería un problema singular de perturbac. (a α_s límite)

Buscamos una solución para $\theta \rightarrow 0$,

$$(1) \frac{\partial \theta_1}{\partial t} + \frac{\partial \theta_1}{\partial z} = - \frac{1}{\lambda_s} (\theta_1 - \theta_s)$$

$$(2) \frac{\partial \theta_s}{\partial z} = \frac{1}{\lambda_s} (\theta_1 - \theta_s)$$

Renombrado: $\lambda_y \leftarrow \frac{1}{\lambda_s}$, $\lambda_s \leftarrow \frac{1}{1-\lambda_s}$.

$$\lambda_y \sim O(10^3), \lambda_s \sim O(10^{-1})$$

$$\text{De (1), } \theta_s = \frac{1}{\lambda_y} \left[\frac{\partial \theta_1}{\partial z} + \frac{\partial \theta_s}{\partial z} \right] + \theta_1.$$

$$\text{De (2), } \theta_1 = \frac{1}{\lambda_s} \frac{\partial \theta_s}{\partial z} + \theta_s$$

Combinando para aislar θ_1 ,

$$\left(\frac{\partial \theta_s}{\partial z} = \frac{1}{\lambda_y} \left[\frac{\partial^2 \theta_1}{\partial z^2} + \frac{\partial \theta_1}{\partial z \partial h} \right] + \frac{\partial \theta_1}{\partial z} \right)$$

$$\cancel{\theta_1} = \frac{1}{\lambda_s \lambda_y} \left[\frac{\partial^2 \theta_1}{\partial z^2} + \frac{\partial^2 \theta_1}{\partial z \partial h} \right] + \frac{\lambda_s}{\lambda_s} \frac{\partial \theta_1}{\partial z} +$$

$$+ \frac{\lambda_s}{\lambda_y} \left[\frac{\partial \theta_1}{\partial z} + \frac{\partial \theta_s}{\partial z} \right] + \cancel{\theta_1}$$

$$0 = \frac{\partial^2 \theta_1}{\partial z^2} + \frac{\partial^2 \theta_1}{\partial z \partial h} + (\lambda_y + \lambda_s) \frac{\partial \theta_1}{\partial z} +$$

$$+ \frac{\partial \theta_s}{\partial z}$$

Resolver para θ_1 y luego obtener θ_s

→ Sep. Variables -

$$\Theta_1 = X(\eta) T(\varepsilon)$$

$$0 = X T'' + X' T' + (\lambda_2 + \lambda_s) \tilde{X} T' + X' T$$

$$0 = \frac{T''}{T} + \frac{X' T'}{X T} + \lambda_2 \frac{T'}{T} + \frac{X'}{X}$$

$$\frac{T''}{T} + \lambda_2 \frac{T'}{T} + \frac{X' T'}{X T} = -\frac{X'}{X}$$

$$\frac{T''}{T} + \frac{T'}{T} \left(\lambda_2 + \frac{X'}{X} \right) = -\frac{X'}{X}$$

$$\frac{d}{d\varepsilon} \left(\frac{T''}{T} \right) + \left(\lambda_2 + \frac{X'}{X} \right) \frac{d}{d\varepsilon} \left(\frac{T'}{T} \right) = 0.$$

$$\cancel{\frac{T''' T - T'' T'}{T^2}} + \left(\lambda_2 + \frac{X'}{X} \right) \cancel{\frac{T'' T - T'^2}{T^2}} = 0.$$

$$\frac{T''' T - T'' T'}{T^2} = - \left(\lambda_2 + \frac{X'}{X} \right) (T'' T - T'^2)$$

$$\frac{T''' T - T'' T'}{T'' T - T'^2} = - \left(\lambda_2 + \frac{X'}{X} \right) = -\lambda.$$

$$\frac{X'}{X} = \lambda - \lambda_2 \equiv \mu \Rightarrow X = C_1 e^{\mu h}$$

$$X' = \mu C_1 e^{\mu h}$$

$$\frac{X'}{X} = \mu$$

$$\frac{T''}{T} + \lambda_2 \frac{T'}{T} + \mu \frac{T'}{T} = -\mu$$

$$T' + T' \left(\lambda_2 + \mu \right) + T \mu = 0.$$

$$T = C_2 e^{-\frac{1}{2} \sqrt{\lambda^2 - 4\mu} h} + C_2 e^{\frac{1}{2} \sqrt{\lambda^2 - 4\mu} h}$$

La solución se recupera de

$$\theta_T = X_T = e^{\mu T} \left(C_1 e^{-\frac{1}{2}\zeta(\sqrt{\lambda^2 - 4\mu} - \lambda)} + C_2 e^{\frac{1}{2}\zeta(\sqrt{\lambda^2 - 4\mu} + \lambda)} \right)$$

$$\theta_T(\eta, 0) = 0 = e^{\mu T} (C_1 + C_2) \Rightarrow \boxed{C_1 = -C_2}$$

$$\theta_T(0, \zeta) = \frac{T\eta}{T_0} - 1 = g \left(e^{-\frac{1}{2}\zeta(\sqrt{\lambda^2 - 4\mu} - \lambda)} - e^{\frac{1}{2}\zeta(\sqrt{\lambda^2 - 4\mu} + \lambda)} \right)$$

$$\boxed{C_1 = \left(1 - \frac{T\eta}{T_0} \right) \frac{e^{\frac{1}{2}\zeta(-\lambda + \sqrt{\lambda^2 - 4\mu})}}{e^{\zeta\sqrt{\lambda^2 - 4\mu}} - 1}}$$

$$\frac{\partial g}{\partial \zeta}(1, \zeta) = 0 \Rightarrow \mu e^{\mu} \boxed{\mu = 0 \Rightarrow \lambda = \lambda_0}$$

$$C_1 = \left(1 - \frac{T\eta}{T_0} \right) \frac{1}{e^{\zeta\lambda_0} - 1}$$

$$\theta_T = \left(1 - \frac{T\eta}{T_0} \right) \frac{1}{e^{\zeta\lambda_0} - 1} \left(1 - e^{\lambda_0 \zeta} \right)$$