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# *Preliminary identification and characterization of nonlinear wave-wave, wave-beam, and wave-particle interactions in beam-driven tokamak plasma*

Partial financial support from NNSA-JPHEDP grant DE-NA0003874 and DOE-FES grant DE-SC0021404 is gratefully acknowledged. GR and MK express deep gratitude to the DIII-D team for necessary support and impactful collaboration during this project.

# Abstract

We present the analysis of 18 DIII-D shots heated by neutral-beam injection (NBI). [Typical plasma parameters are  $n \sim 2 \times 10^{13} \text{ cm}^{-3}$ ,  $T \sim 1 \text{ keV}$ .] As energetic ions seeded by NBI resonate at the frequencies of various Alfvén eigenmodes (AEs), we observe rich toroidal Alfvén eigenmode (TAE) activity as these weakly-damped modes are driven. Both steady and modulated beam power have been investigated.

Statistical techniques were used to benignly filter out the polluting effect of edge-localized modes (ELMs) in the observed fluctuation spectra. Guided by recent simulations [Spong *et al* 2021; *Nucl. Fusion* 61, 116061] which identify nonlinear coupling between TAEs and zonal flows, we seek to correlate the observed nonlinear interaction between ensembles of AEs and lower-frequency MHD modes with coincident perturbations in the fast-ion distribution function.

Evidence for energy exchange due to this coupling is given by higher-order spectral techniques. In particular, we report on the evolution of consistent bicoherent features in magnetic fluctuation data.

# *Introduction to TAEs*

# Toroidal geometry facilitates eigenmodes<sup>1</sup>

Alfven eigenmodes (AEs) are an interaction of shear Alfven waves (SAW),

$$\omega_A(r) = k_{\parallel}(r)v_A(r)$$

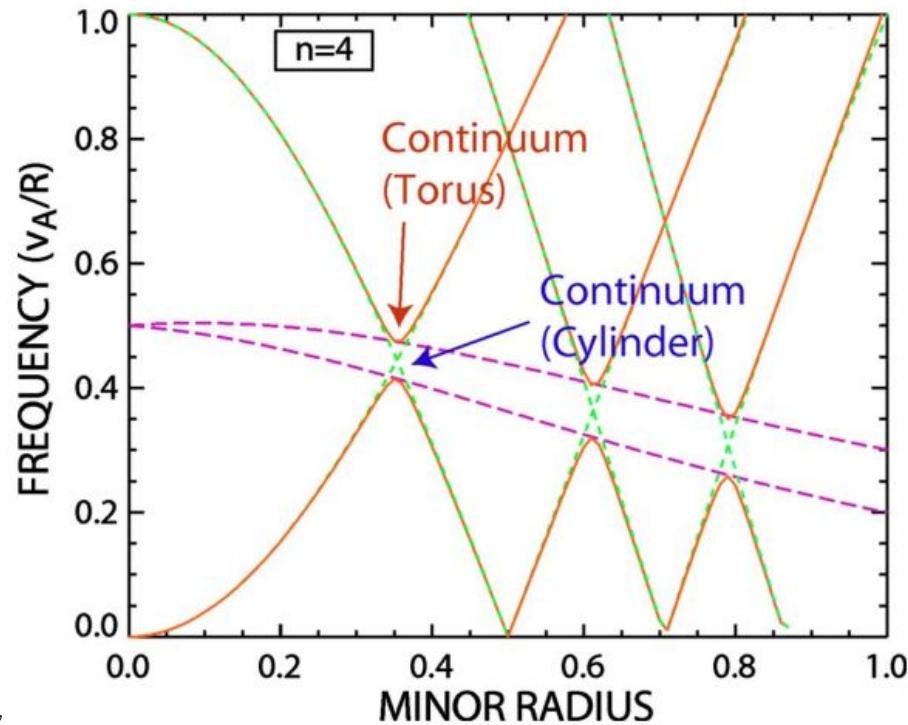
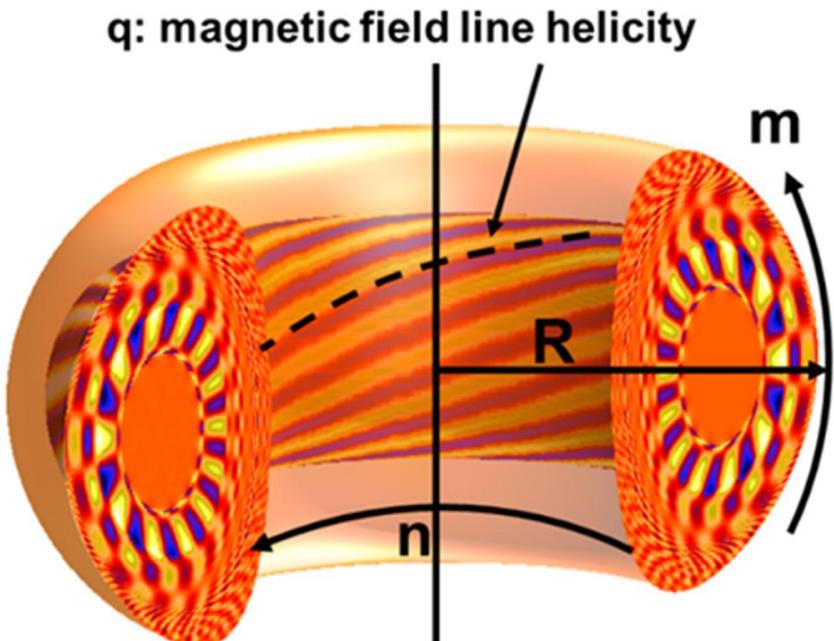
Parallel wave-vector is given approximately by

$$k_{\parallel m}(r) = \frac{1}{R} \left( n - \frac{m}{q(r)} \right)$$

Toroidal Alfven eigenmode (TAE) results from counter-propagation

$$k_{\parallel m}(r_0) = -k_{\parallel m+1}(r_0)$$

# Interaction of counter-propagating cylindrical modes provides frequency gap<sup>2</sup>



Garcia-Muñoz, M. et al. 2019 Active control of Alfvén eigenmodes in magnetically confined toroidal plasmas. Plasma Phys. Control. Fusion 61, 054007

[2] Heidbrink, W., Basic physics of Alfvén instabilities..., Physics of Plasmas 15, 055501 (2008); <https://doi.org/10.1063/1.2838239>

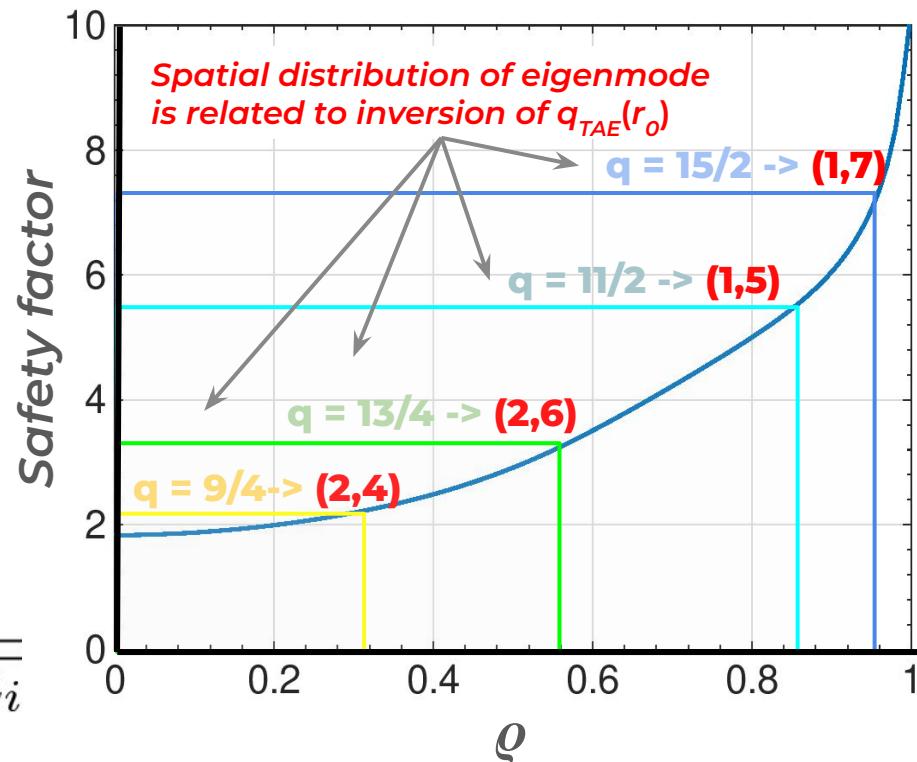
# Safety factor profile determines TAE localization<sup>1</sup>

Critical value of safety factor is

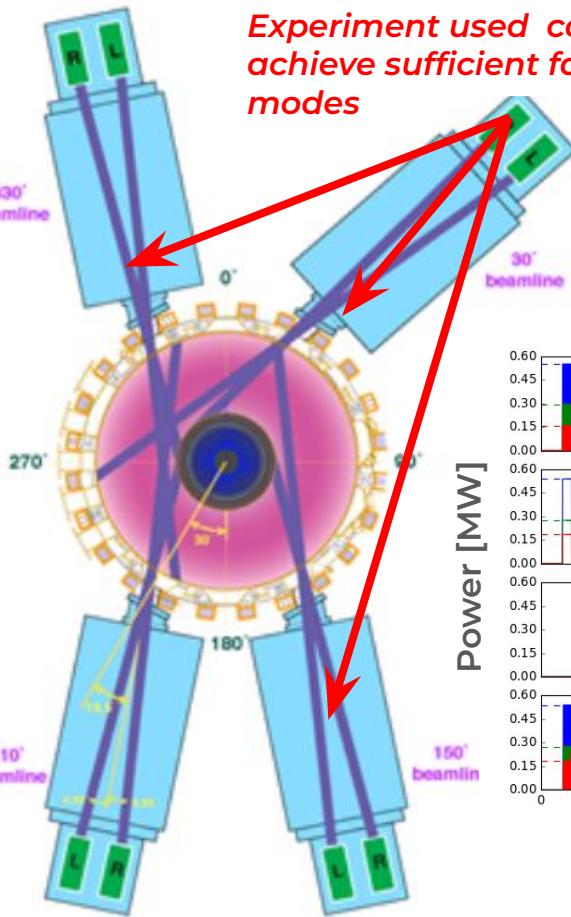
$$q_{\text{TAE}}(r_0) = \frac{m + 1/2}{n}$$

TAE frequency (in local frame) is

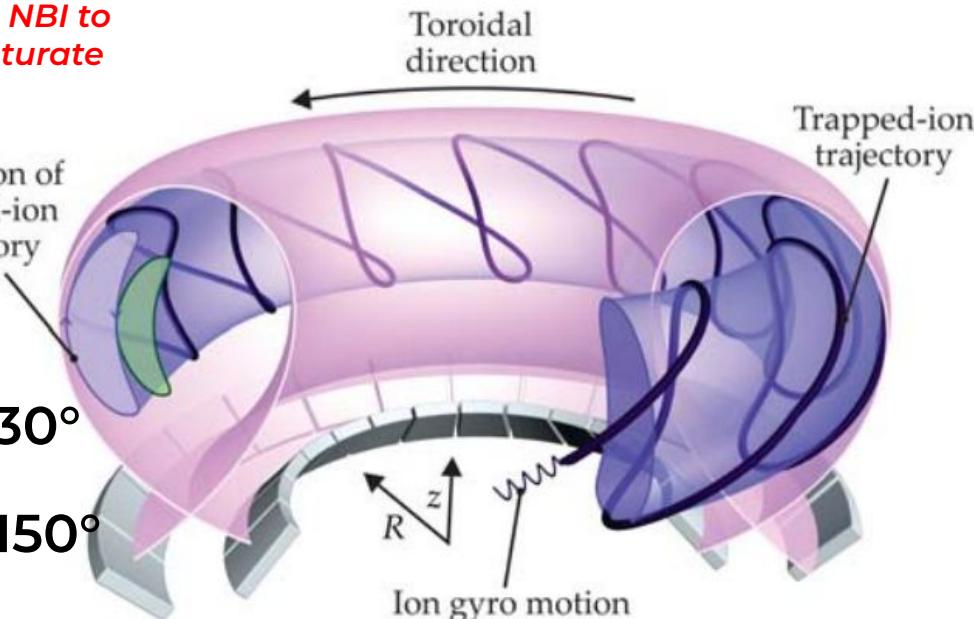
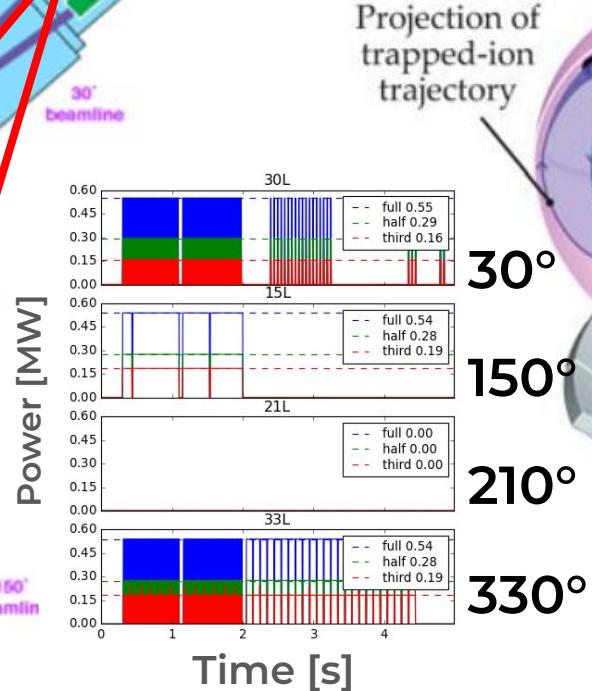
$$\begin{aligned} f_{\text{TAE}} &= \frac{v_A}{4\pi q_{\text{TAE}} R} \\ &\approx \frac{1}{(4\pi)^{3/2}} \frac{nB}{R(m + 1/2)\sqrt{m_i n_i}} \end{aligned}$$



# NBI seeds fast-ion population which drives TAEs<sup>2</sup>



Experiment used co-injected deuterium NBI to achieve sufficient fast-ion pressure to saturate modes



$$\omega + (m + l)\omega_\theta - n\omega_\zeta \approx 0$$

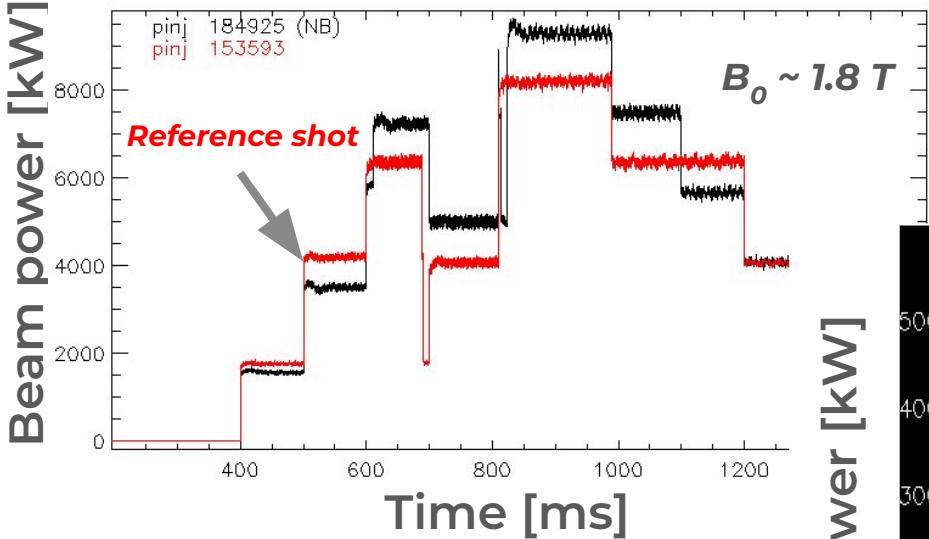
Pace, D., Heidbrink, W., and Van Zeeland, M 2015 Keeping fusion plasmas hot. Physics Today 68.10 <https://doi.org/10.1063/PT.3.2946>

Pawley, C.J. et al. 2017 Advanced control of neutral beam injected power in DIII-D. Fusion Engineering and Design. 123, 453-7.

# *Experimental data*

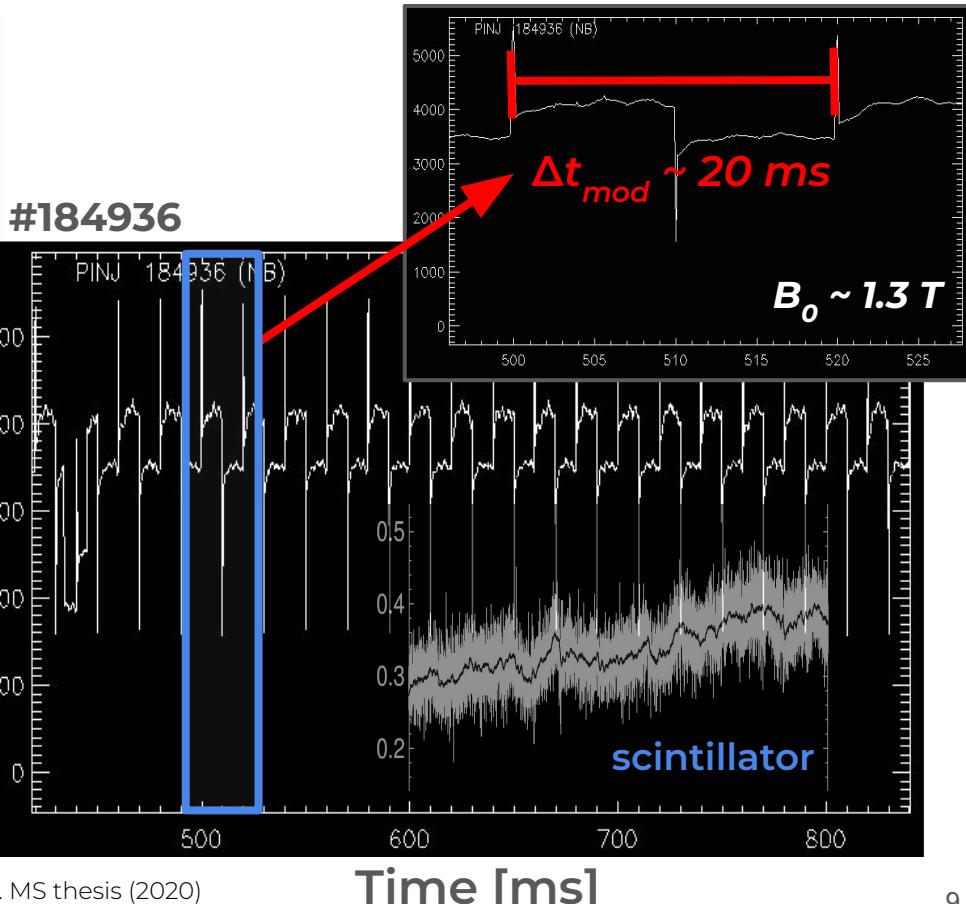
# Both steady and modulated beam power investigated

#184925



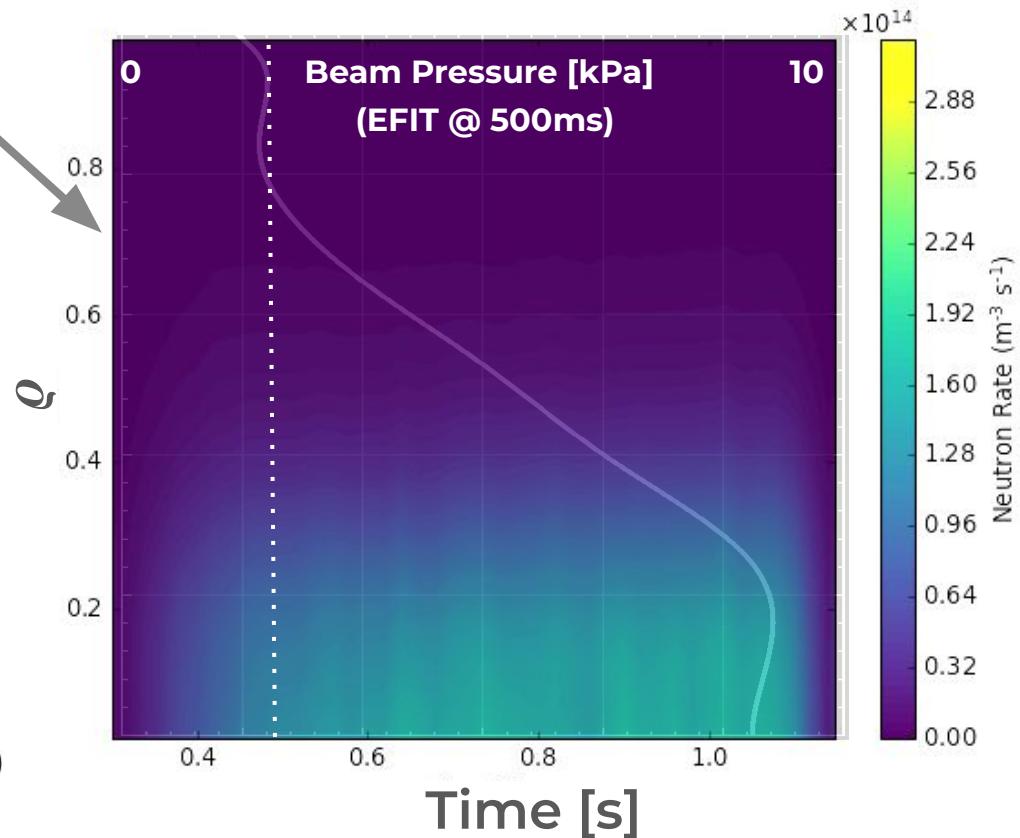
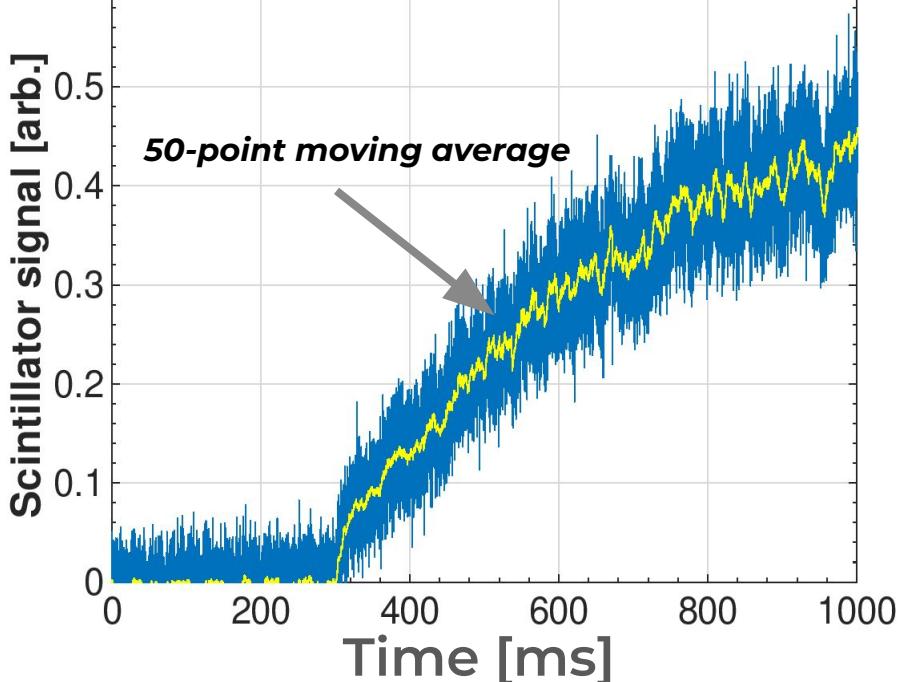
Favorable discharge parameters informed via investigation of archived data-set<sup>3</sup>

DIII-D platform delivers high reproducibility of plasma conditions and NBI programming



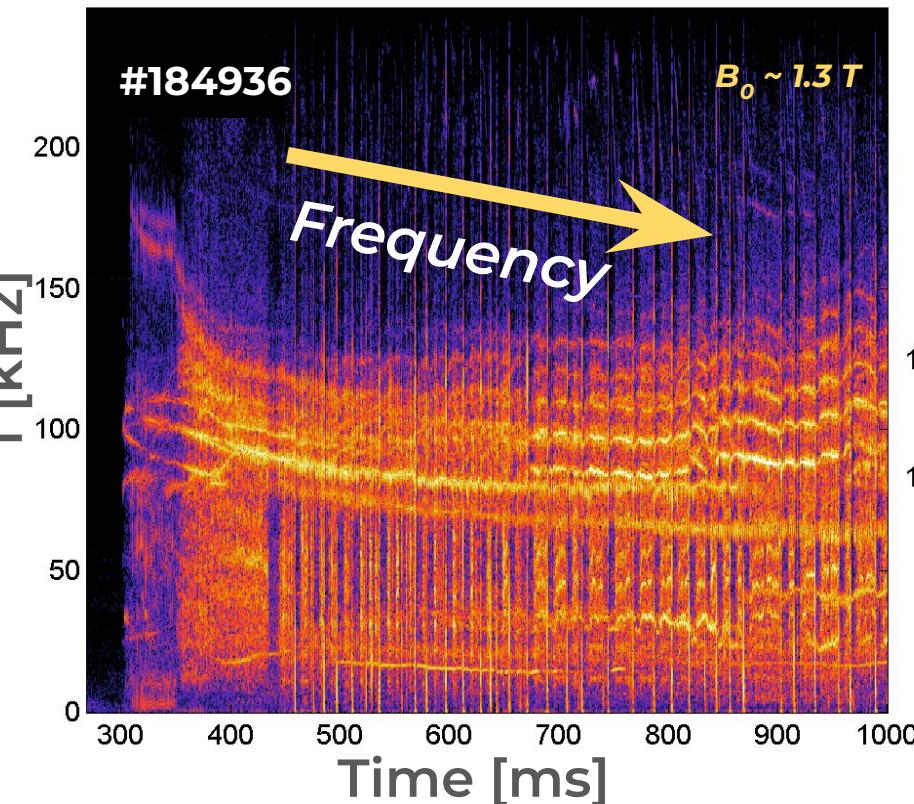
# Deposition of beam energy facilitates fusion reactions, enhancing neutron flux

**Scintillator data compares favorably with output of RABBIT<sup>4</sup> simulation**

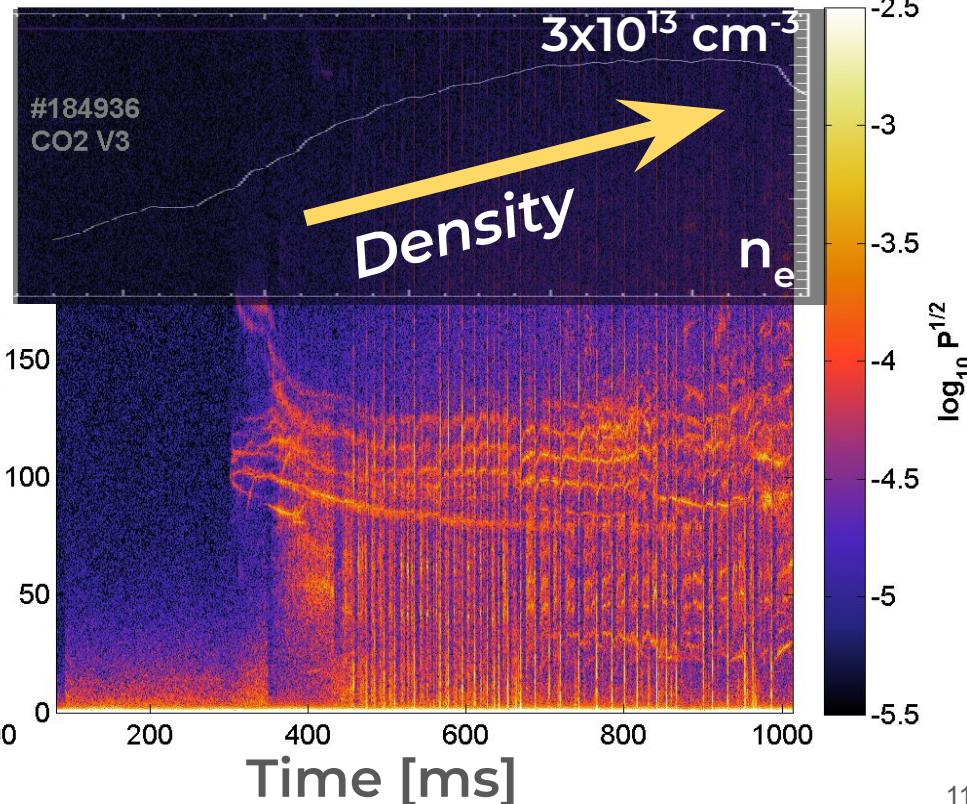


# Density ramp interrogates the mechanism of TAE coupling via scan of frequency & amplitude

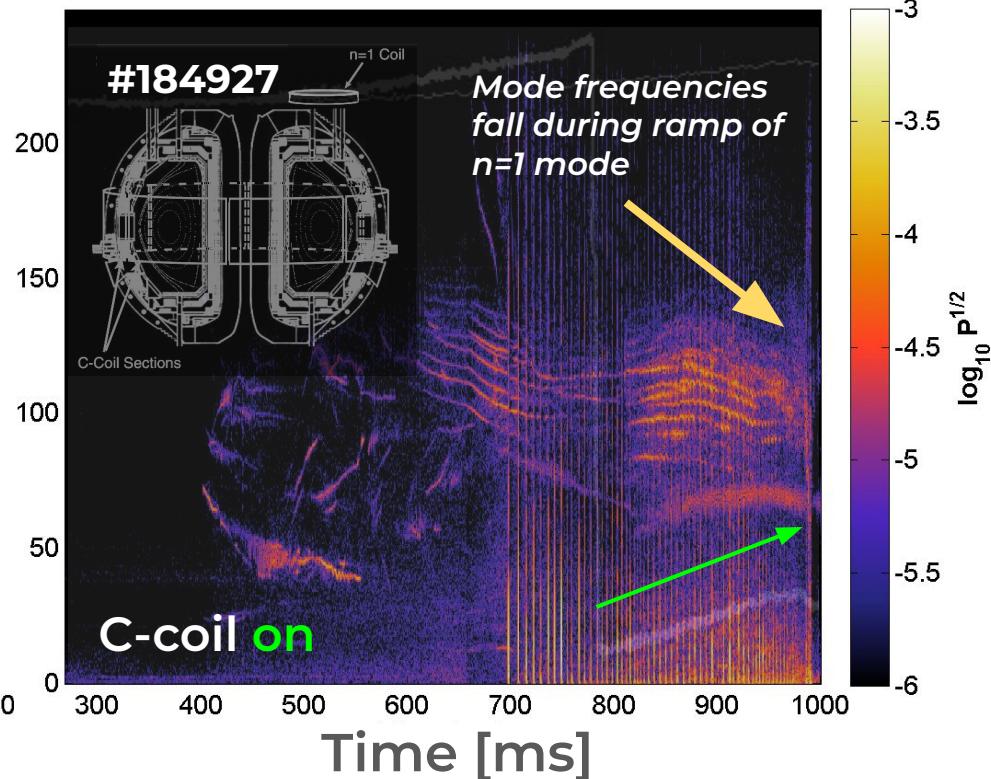
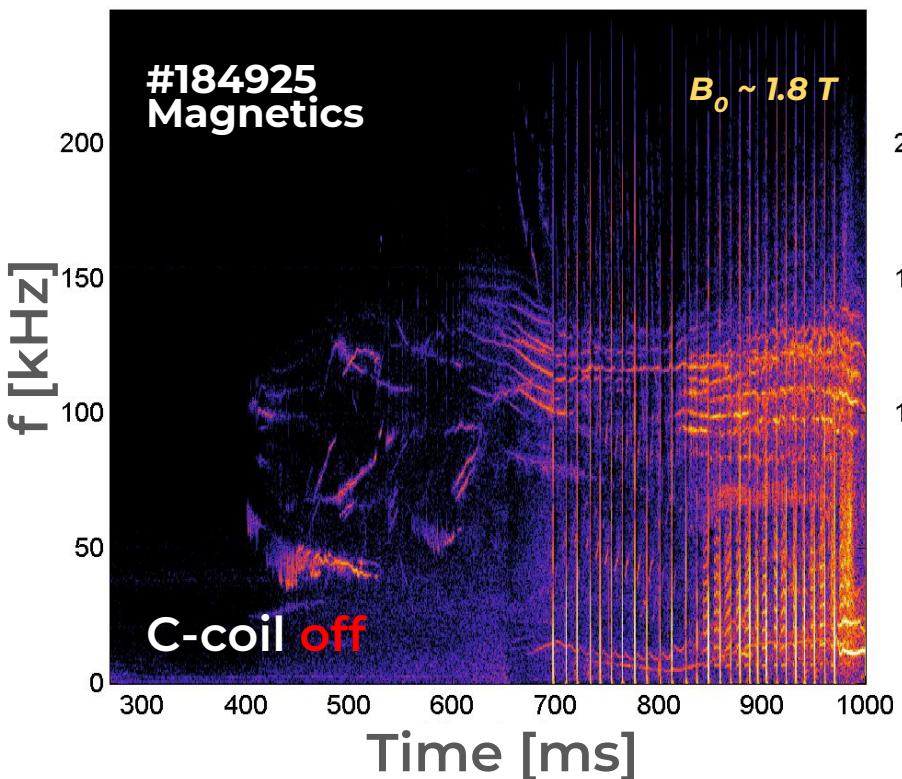
Mirnov coil



Interferometer

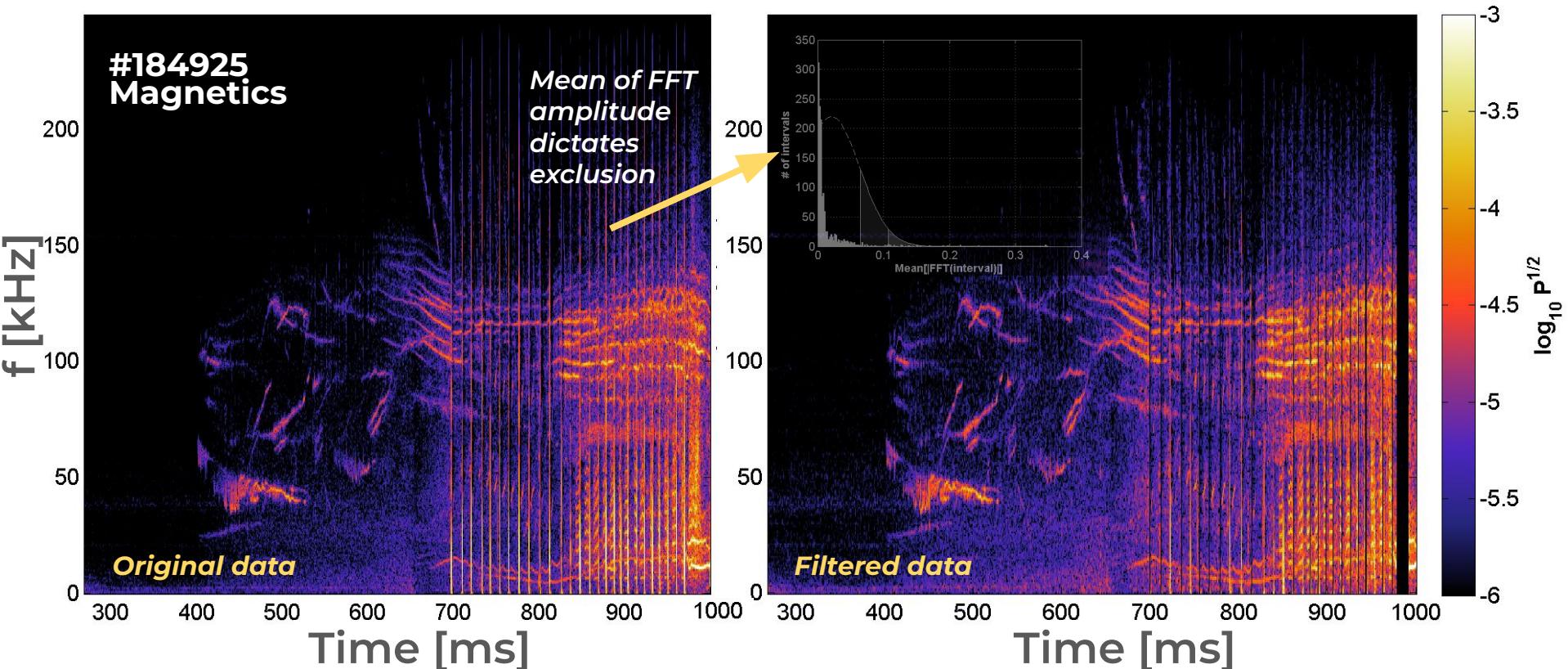


# Eigenmodes augmented by C-coil perturbation

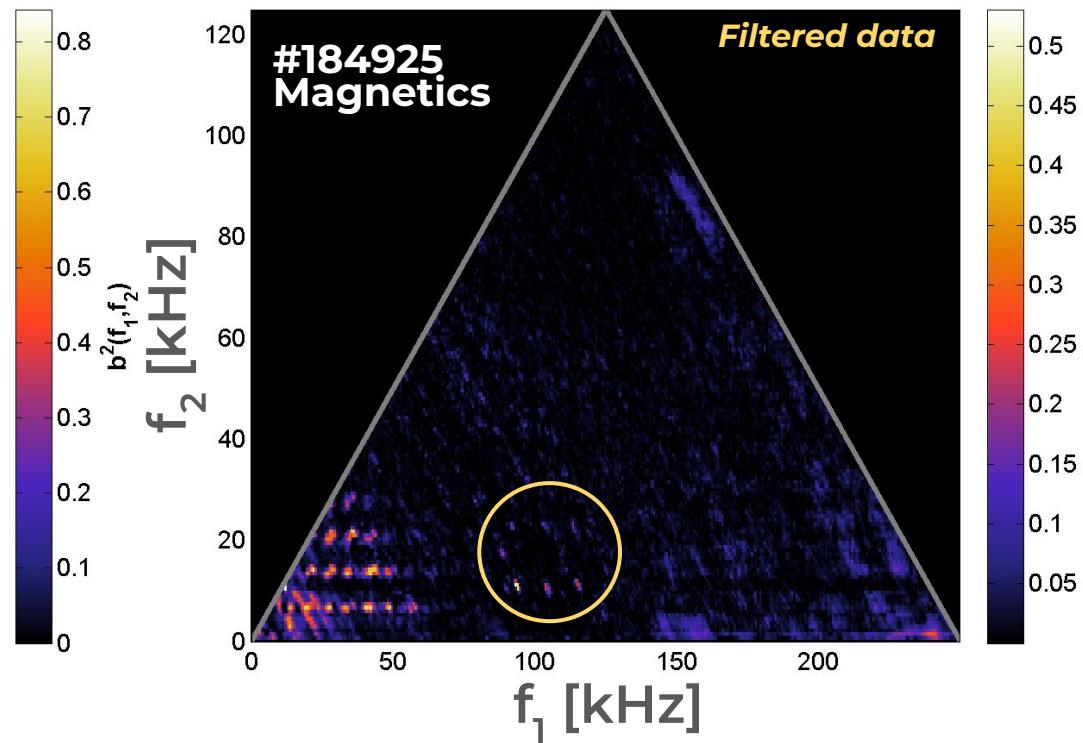
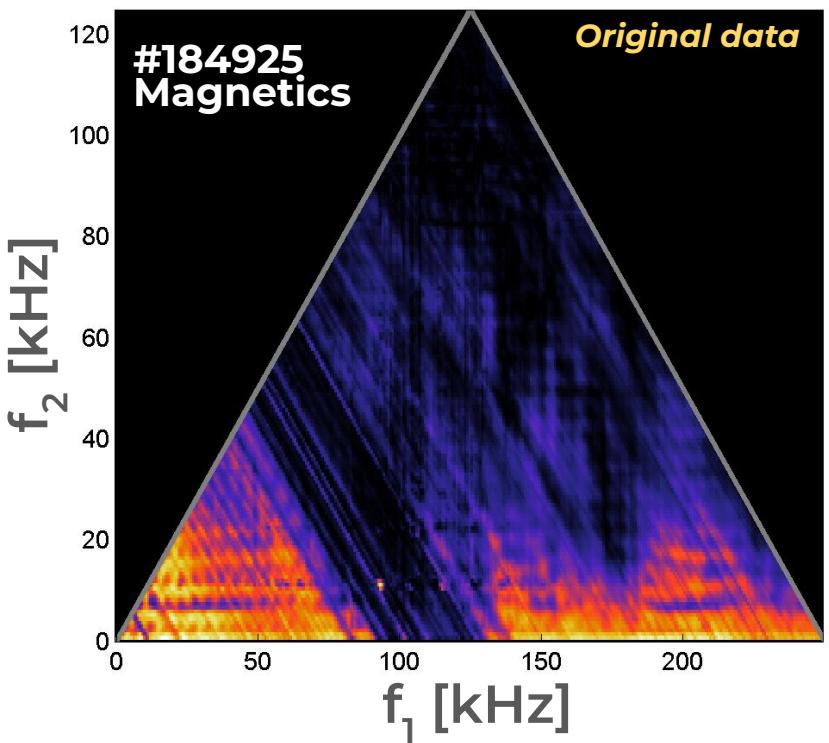


# **Wave-wave *interaction***

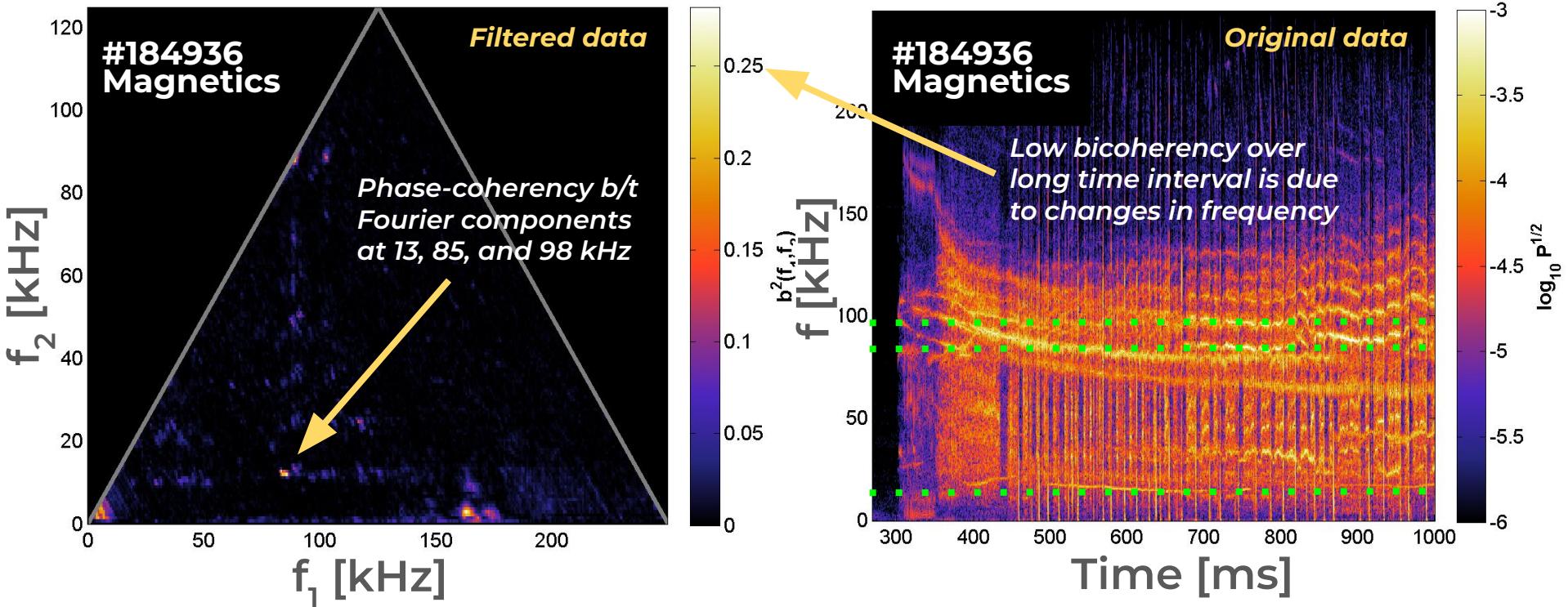
# Acquired spectrograms are ELM-filtered conveniently



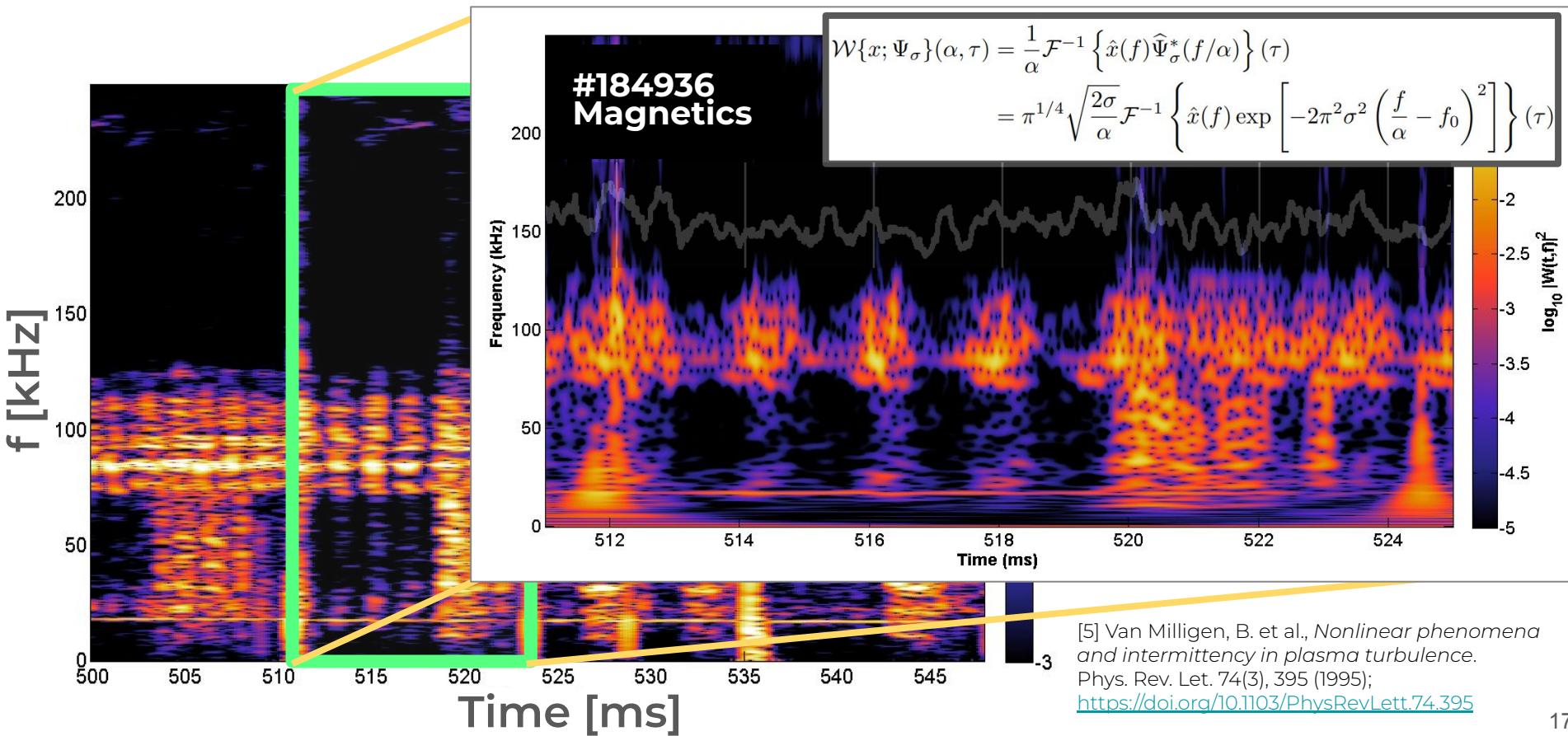
# ELM-filtered bispectra reveal signatures of 3-wave interaction



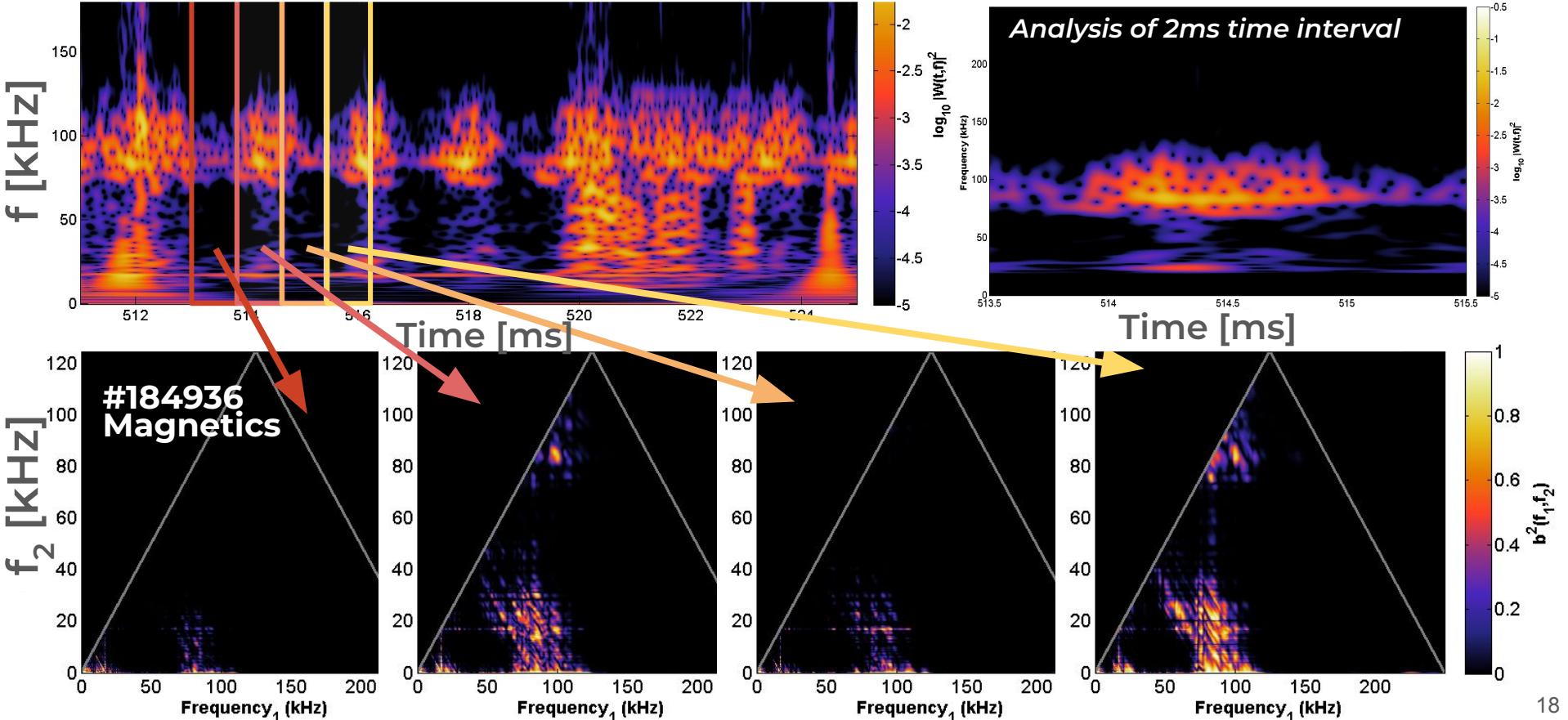
# Nonlinear wave-wave coupling inferred from magnetics data



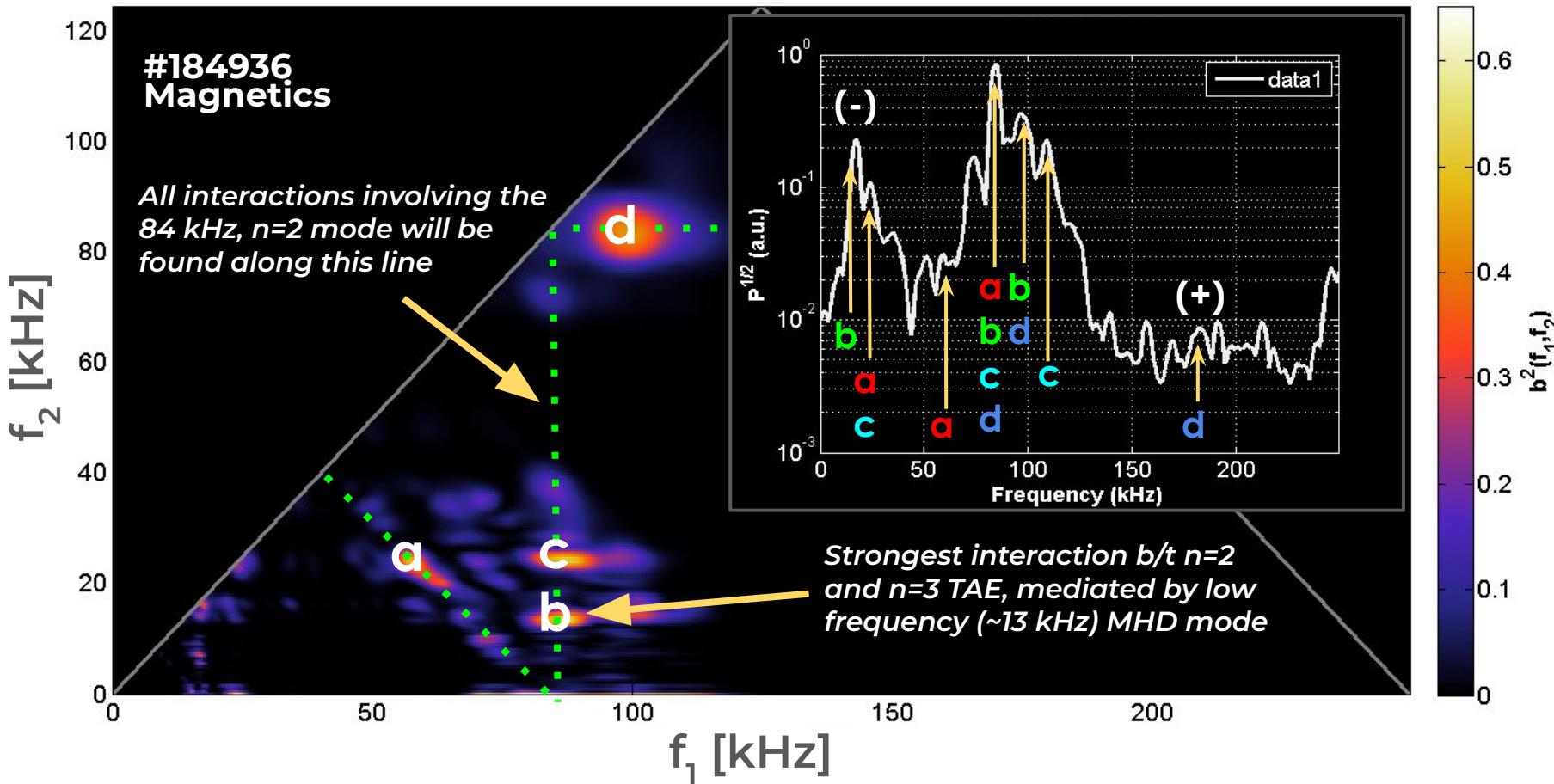
# Morlet wavelet optimizes simultaneous resolution in time and frequency<sup>5</sup>



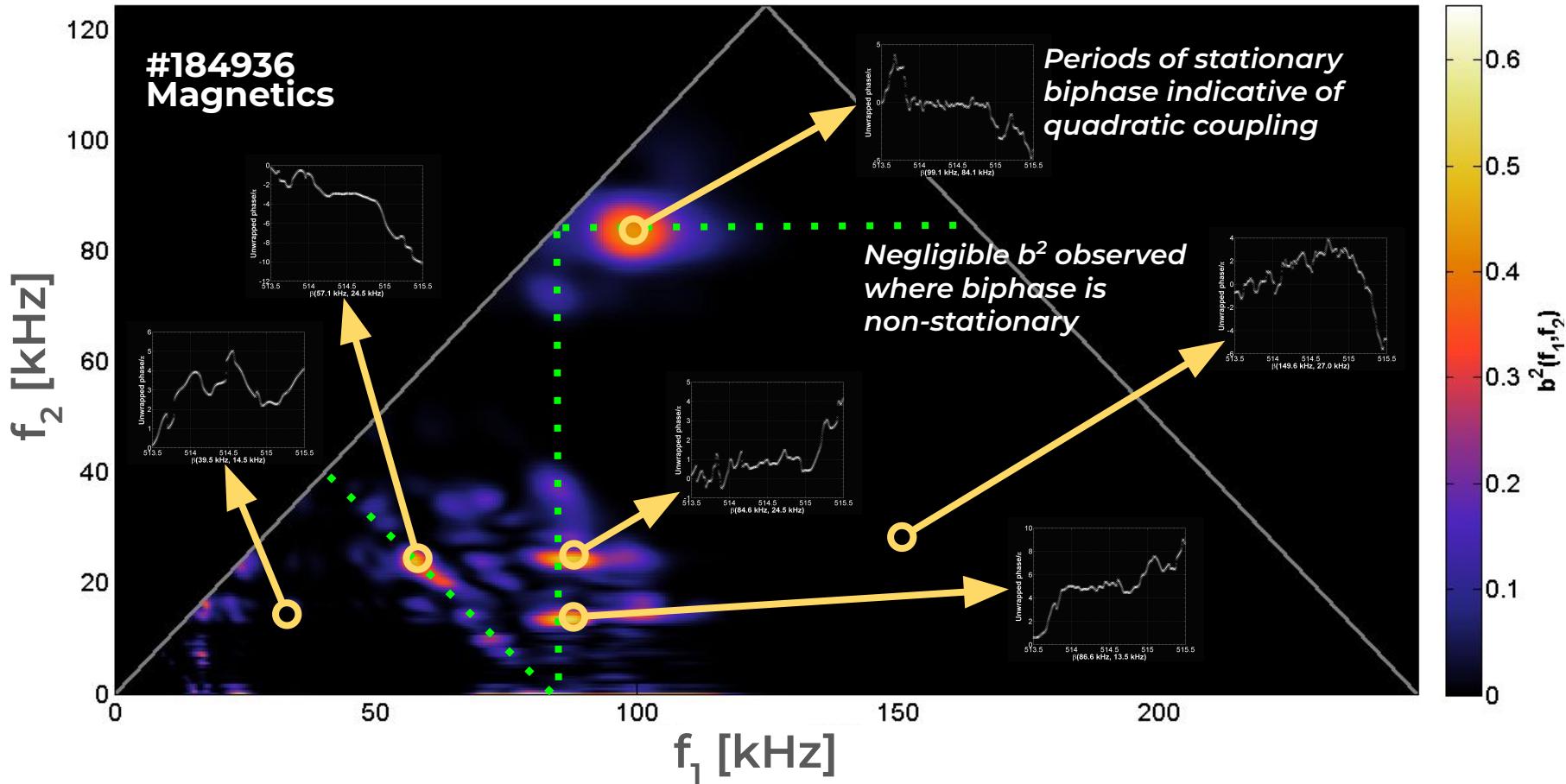
# Wavelet-based bicoherence enables highly time- and frequency-resolved assessment of phase-coherency



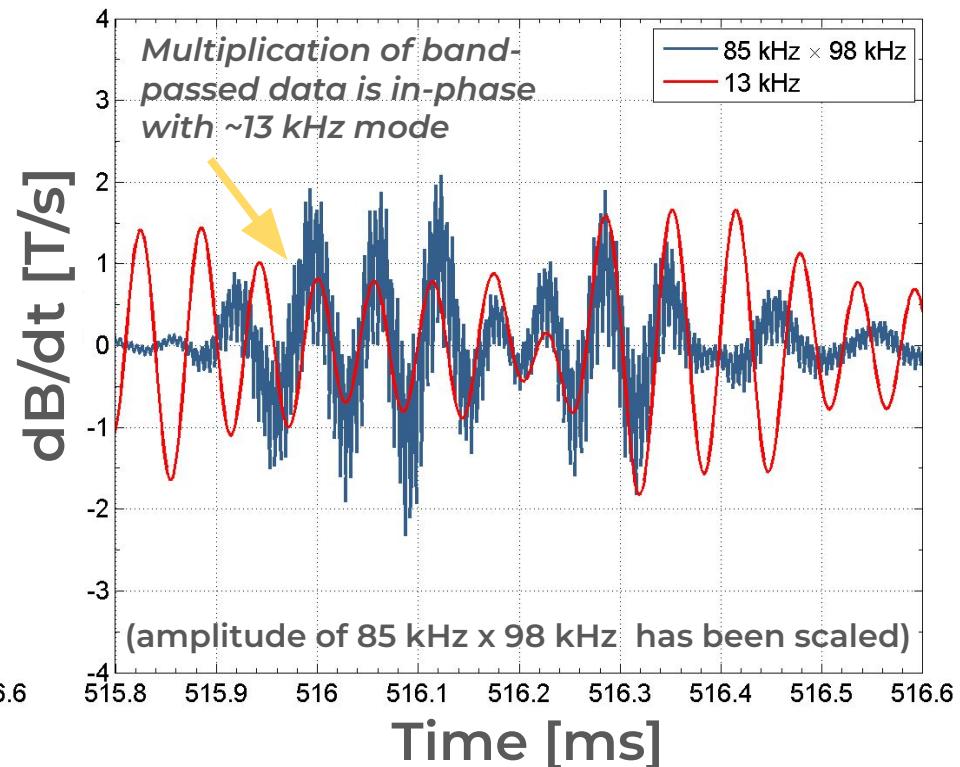
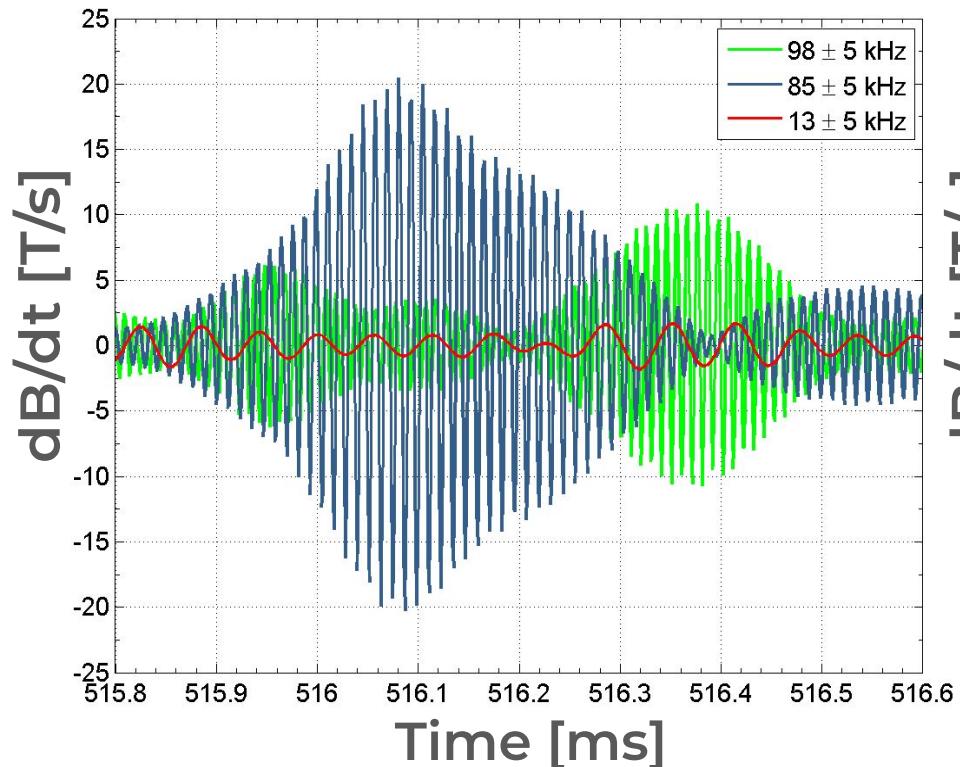
# Mode-mode interaction evinced by coherent sum and difference frequencies



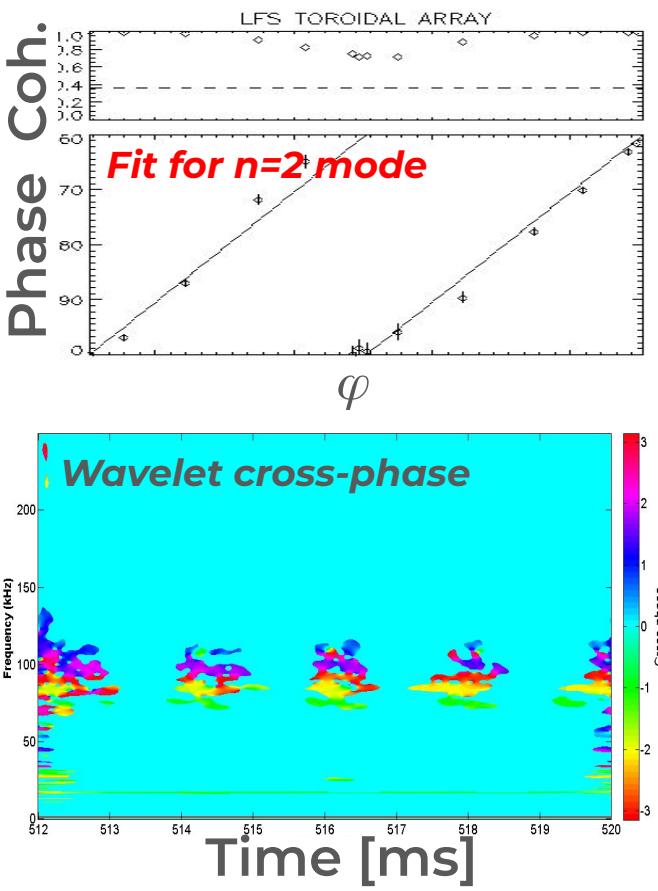
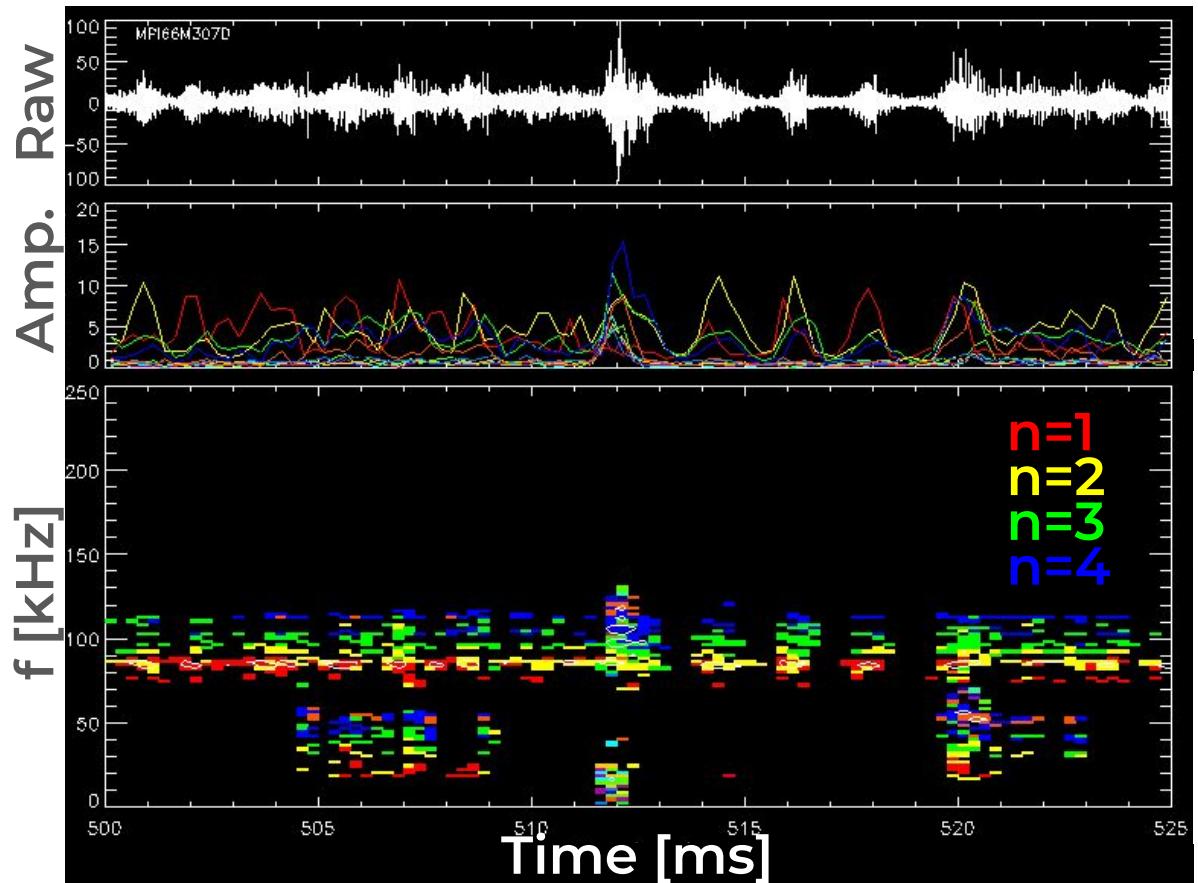
# Biphase evolution is consistent with nonlinear 3-wave coupling



# Interpretation supported by correlation between beat dynamics and low-frequency fluctuations

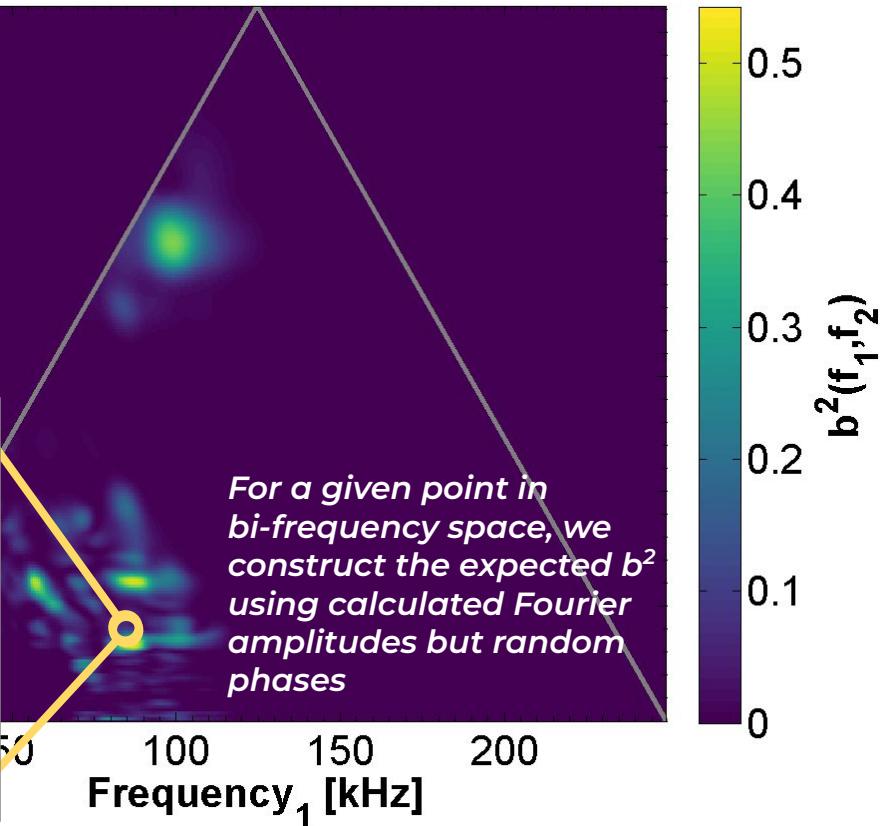
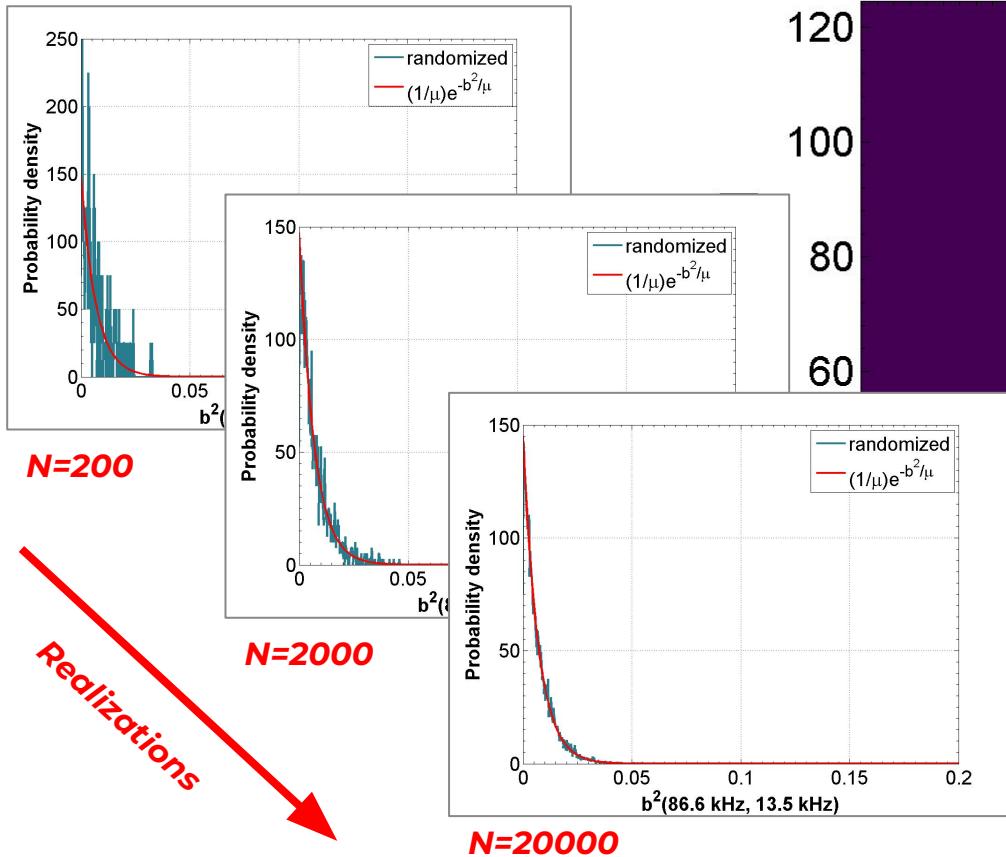


## Mode identification benchmarked by cross-phase analysis<sup>6</sup>

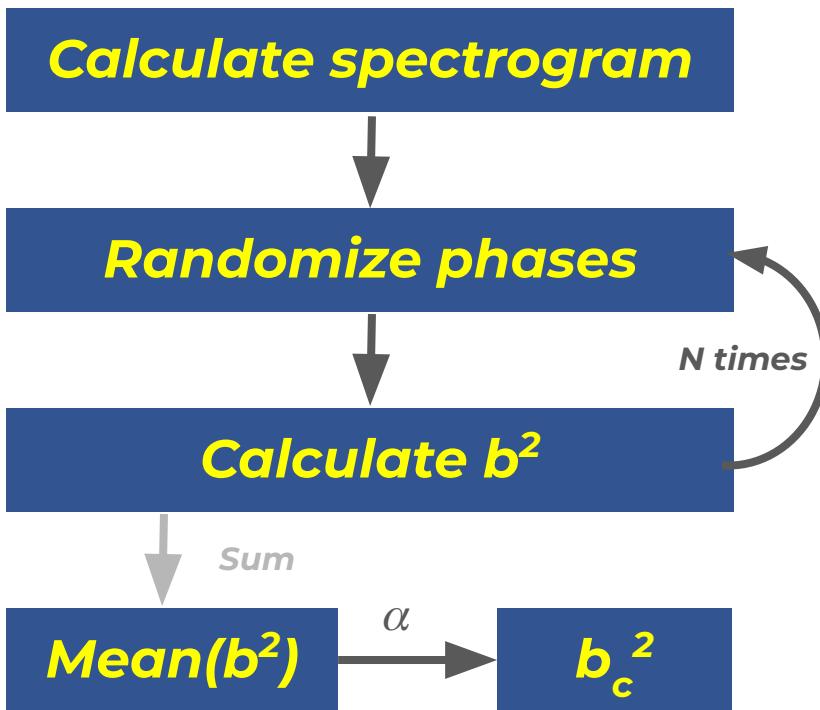


# ***Assessment of noise floor***

# Expected value of bicoherence is typically well-approximated by exponential distribution



# Confidence interval may be derived using quantile<sup>7</sup>



Modelled distribution given by:

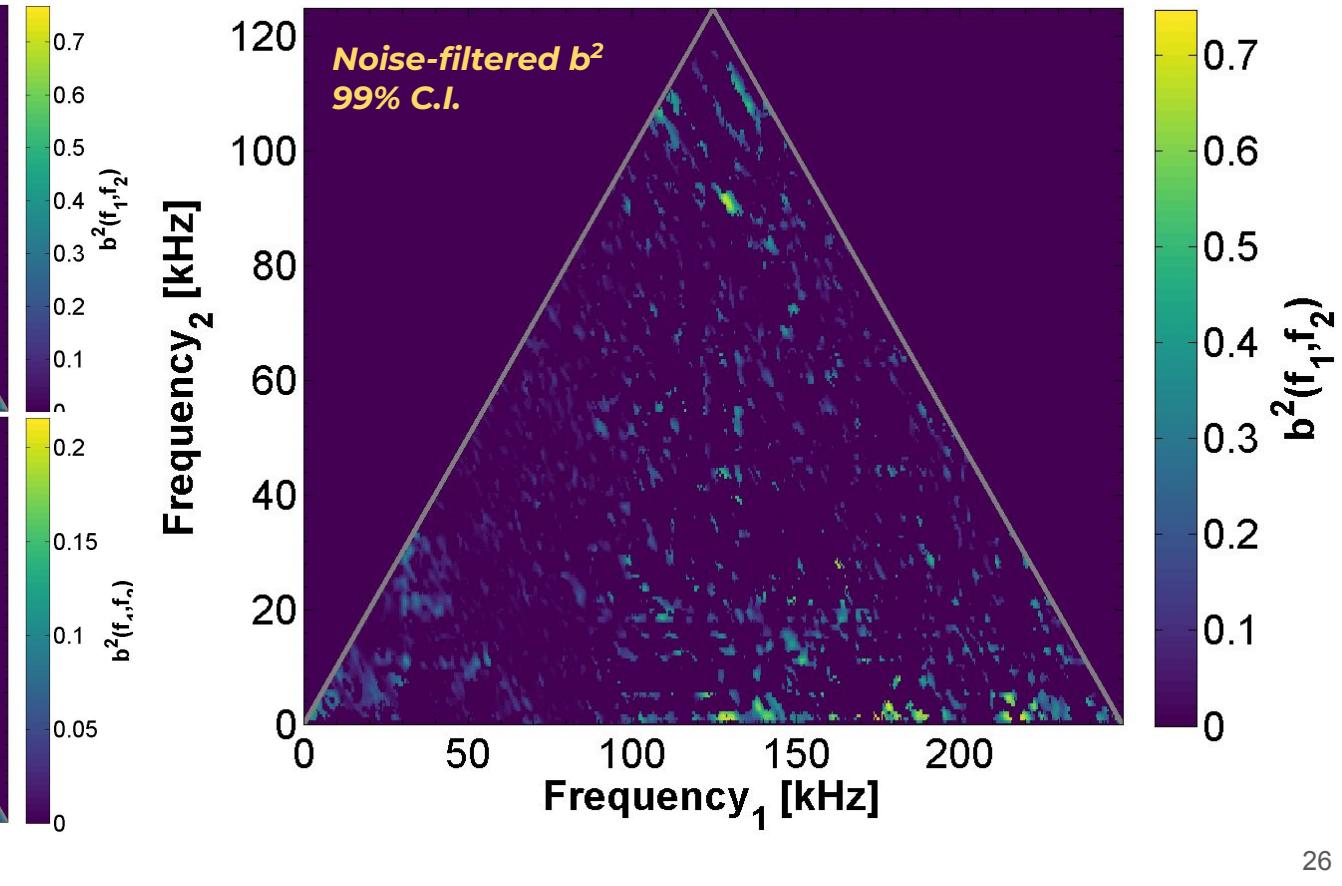
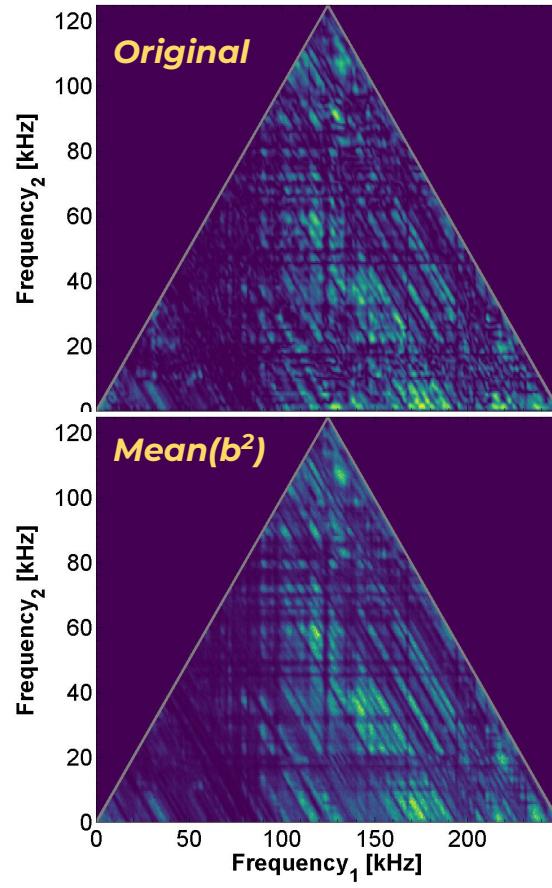
$$\text{PDF}(b^2) = \frac{e^{-b^2/\mu}}{\mu}$$

where  $\mu$  is mean of random-phase realizations

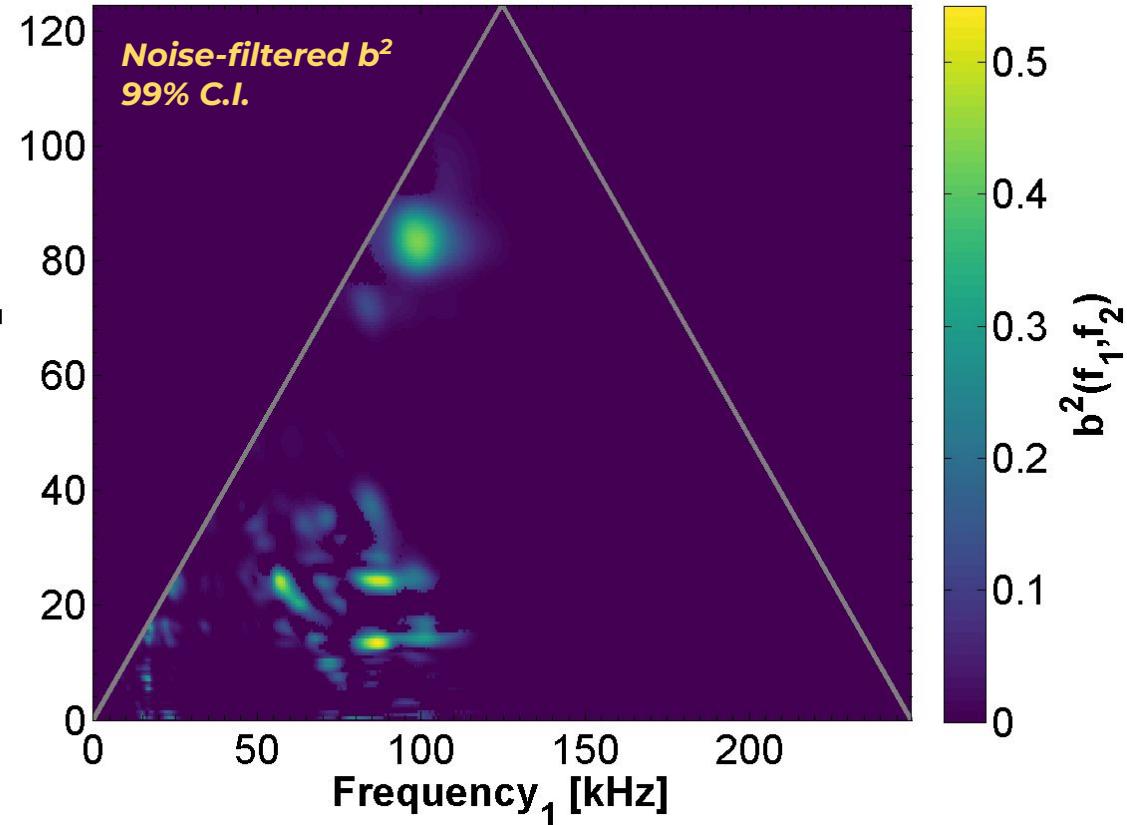
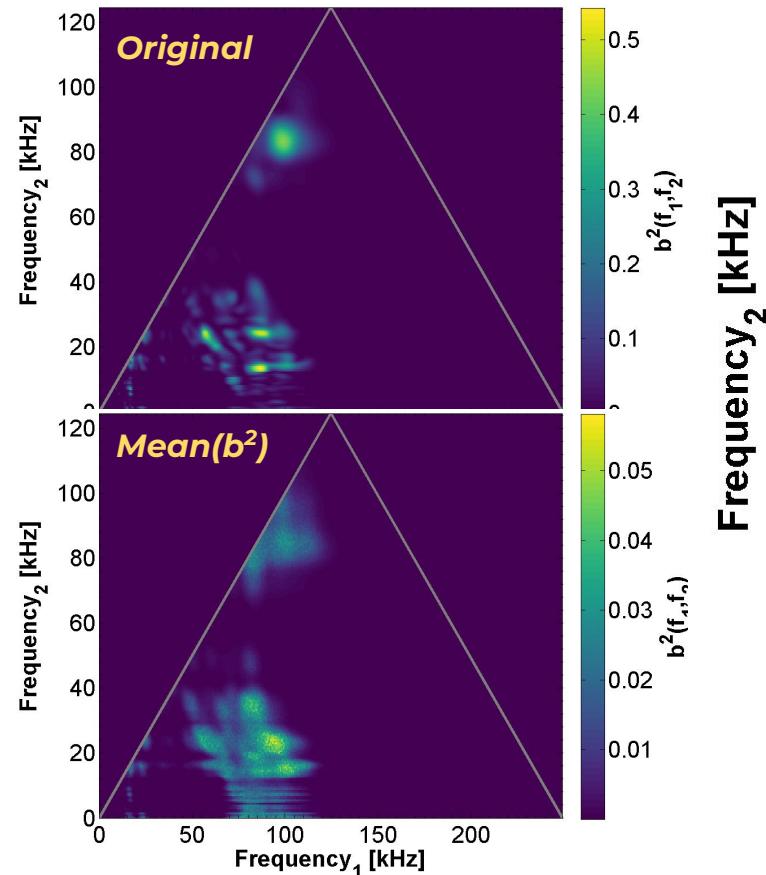
Critical value of bicoherence for confidence level  $0 < \alpha < 1$ :

$$b_c^2 = -\mu \log(1 - \alpha)$$

# Technique provides robust filtering of spurious bicoherence for non-stationary signals



# Consistent phase-coherency is unlikely to be discarded



***What's next?***

# Priorities for 2023

Use FAR3d and TRANSP simulations to provide insight into wave-wave and wave-particle interactions

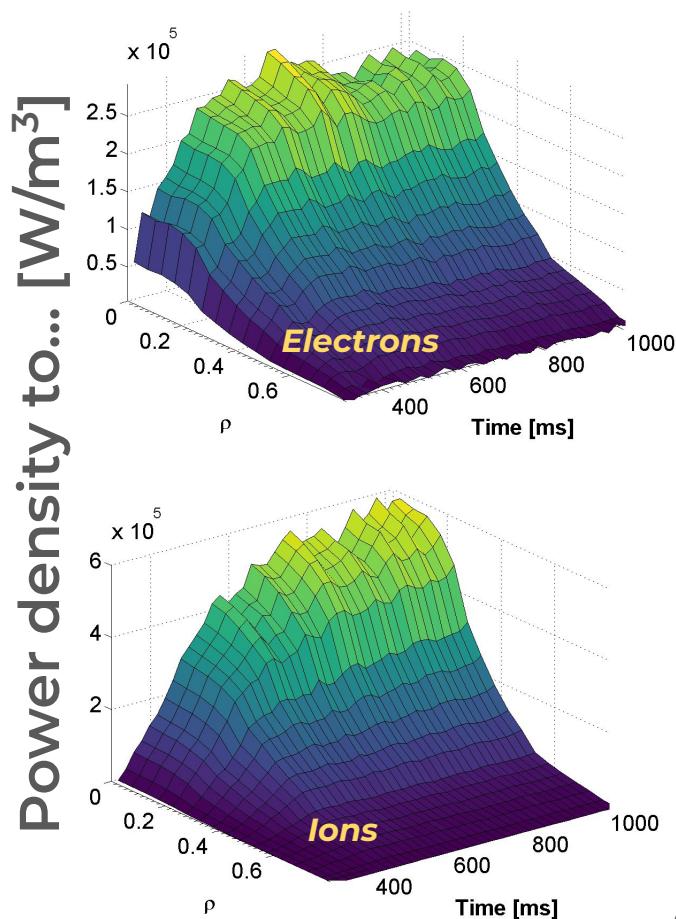
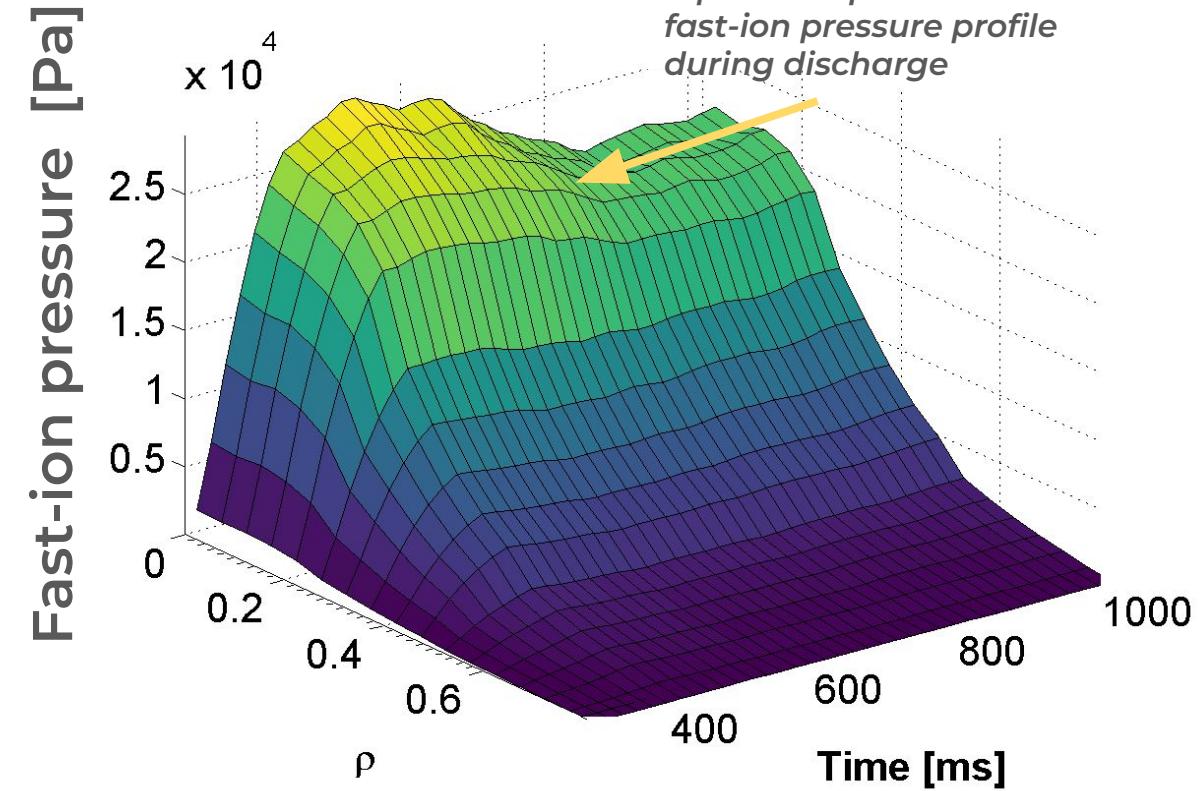
Correlate fluctuations in density and magnetic field with perturbations in fast-ion distribution function (e.g., FILD fluctuation analysis)

Quantify role of nonlinear coupling and energy transfer in mediating saturated amplitude of TAEs

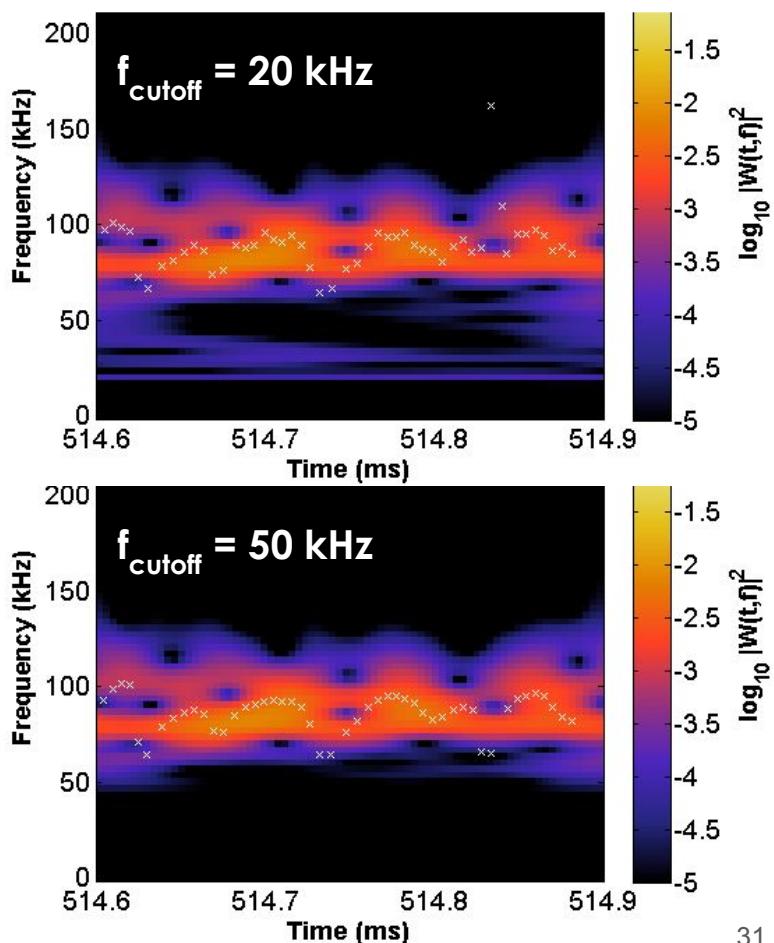
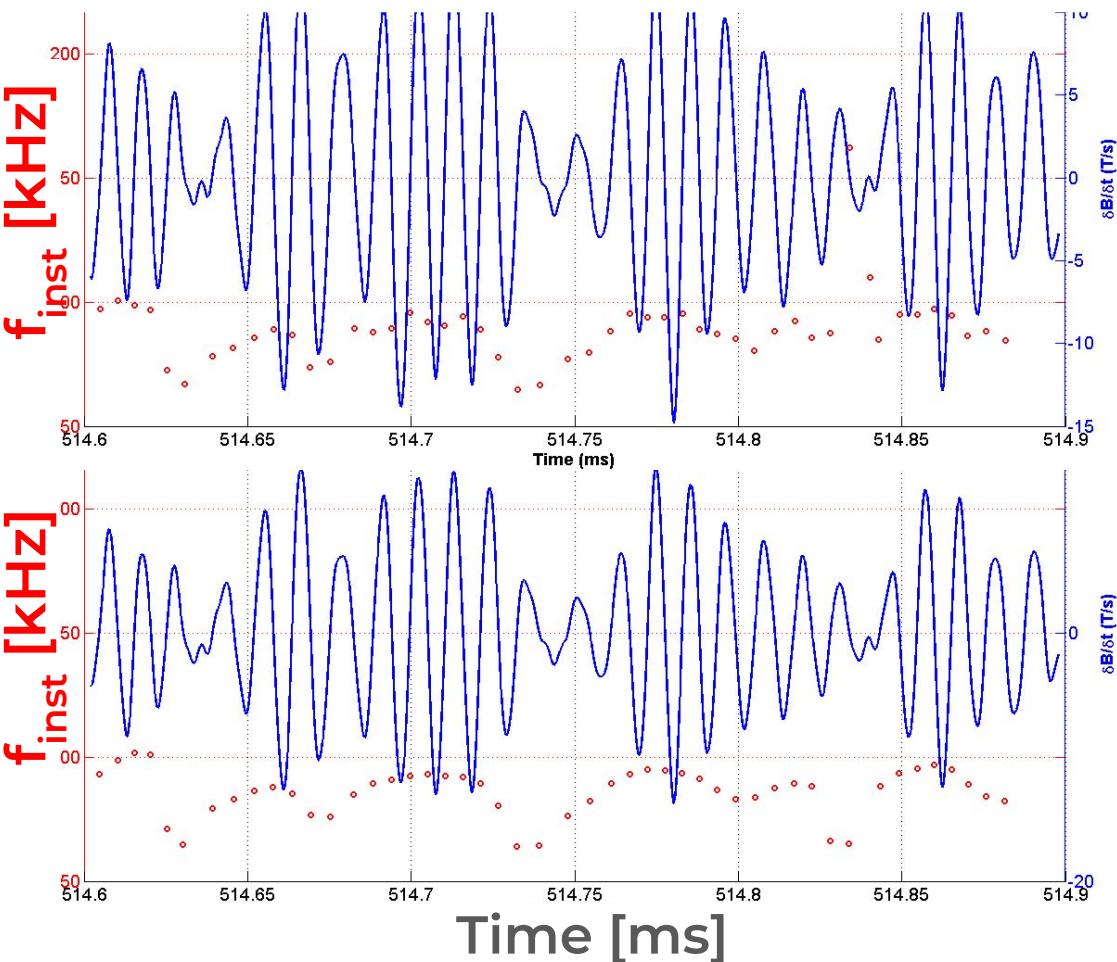
Assess efficacy of Berk-Breizman model to explain observed AM/FM, or other nonlinear effects

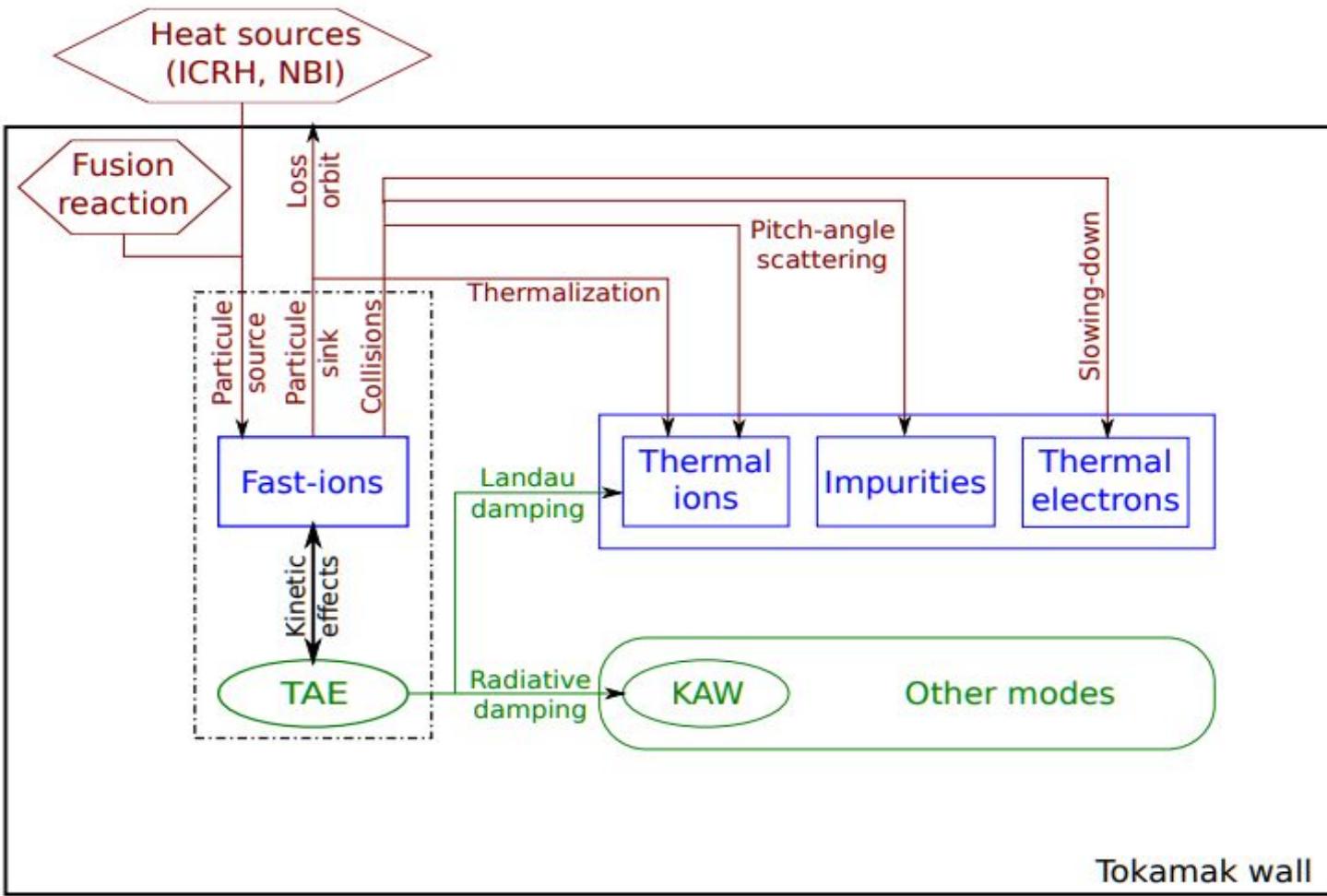
Develop “flowchart”

# Mechanism of saturation is under investigation



# Frequency modulation likely in HP-filtered data





# **Appendix: Bicoherence primer**

# Bispectrum is generalization of power spectrum<sup>9</sup>

Fourier transform of auto/cross-correlation = power/cross-spectrum

Bispectrum = higher order transform; identifies nonlinear interactions via

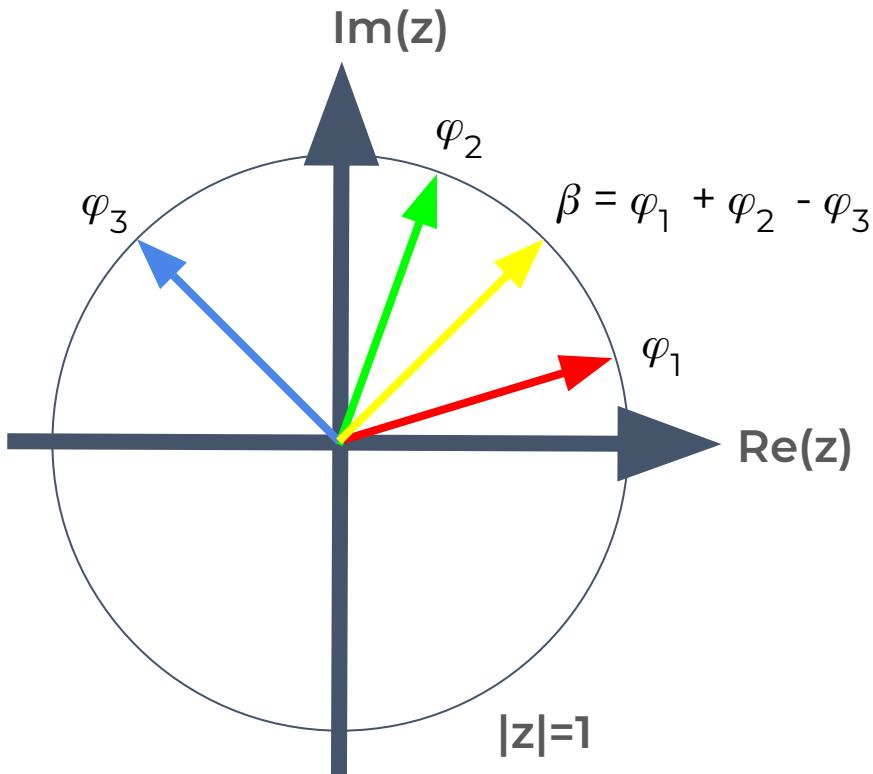
$$\mathcal{B}_{x_1 x_2 x_3}(f_1, f_2) = \left\langle \hat{x}_1(f_1) \hat{x}_2(f_2) \overline{\hat{x}_3(f_1 + f_2)} \right\rangle$$

Auto/cross-bicoherence = normalized auto/cross-bispectrum

$$b_{x_1 x_2 x_3}^2(f_1, f_2) = \frac{|\mathcal{B}_{x_1 x_2 x_3}(f_1, f_2)|^2}{\left\langle |\hat{x}_1(f_1) \hat{x}_2(f_2)|^2 \right\rangle \left\langle |\hat{x}_3(f_1 + f_2)|^2 \right\rangle}$$

Detects complementary phase relationships between frequency triples

# Bicoherence determined by phase-coherency

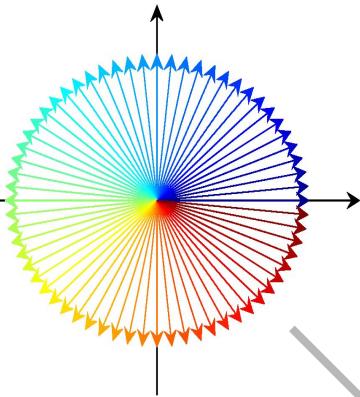


$$\mathcal{B}_{x_1 x_2 x_3}(f_1, f_2) \propto \left\langle e^{i\beta(f_1, f_2, \tau)} \right\rangle$$

- When Fourier amplitudes are slowly varying in time,  $B(f_1, f_2)$  depends entirely on the dynamics of biphase  $\beta$
- Crucially, the bispectrum will tend to null when  $\beta$  is random or linear in time
- A static biphase thus corresponds to nonzero values of bicoherence
- Oscillatory biphase does not generally lead to vanishing bicoherence

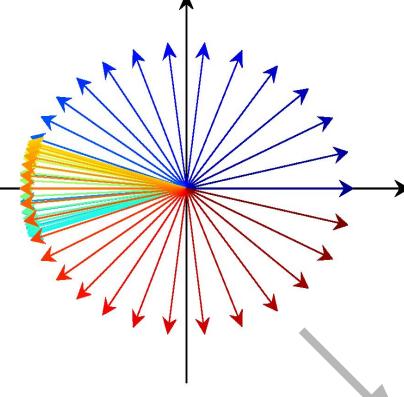
# Phase and amplitude modulation require careful interpretation

$$\beta(t) = 2\pi t$$

 $\Sigma$ 

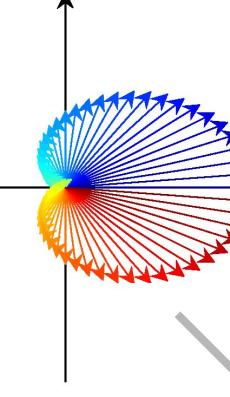
$$\langle \exp(i\beta(t)) \rangle = 0$$

$$\beta(t) = \pi[2t + (1/2)\sin(2\pi t)]$$

 $\Sigma$ 

$$\langle \exp(i\beta(t)) \rangle \neq 0$$

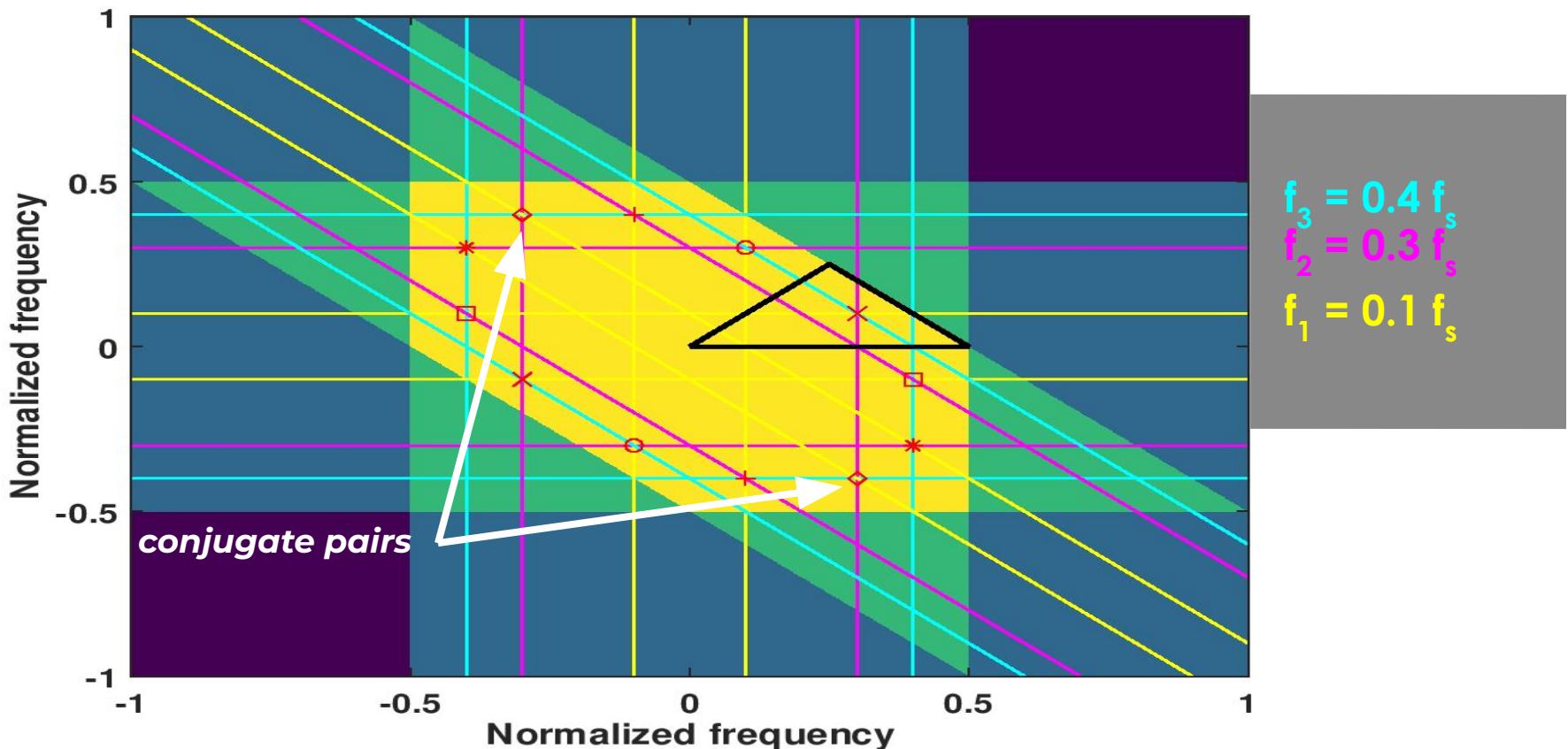
$$\beta(t) = 2\pi t$$

 $\Sigma$ 

$$\langle B(t) \exp(i\beta(t)) \rangle \neq 0$$

 $dB/dt \neq 0$

# Bi-frequency space is 12-fold degenerate for auto-bicoherence analyses



# **Thanks for your time!**

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