

# Chapter 2

## Properties of electromagnetic waves

- Basic information on electromagnetic waves in free space or other uniform media, but no rigorous development
- Maxwell's equations (brief reference to them)
- Plane wave properties
- Polarisation

# Maxwell's equations

An electric field is produced by a time-varying magnetic field

A magnetic field is produced by a time-varying electric field or a current

Electric field lines may either start and end on charges, or are continuous

Magnetic field lines are continuous

Will use del-operator next, definitions:

$$\nabla = \hat{\mathbf{x}} \frac{\partial}{\partial x} + \hat{\mathbf{y}} \frac{\partial}{\partial y} + \hat{\mathbf{z}} \frac{\partial}{\partial z}$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

# Maxwell's equations

$$\nabla \times \mathbf{E} = -\mu \frac{\partial \mathbf{H}}{\partial t}$$

Faraday's law. Time varying magnetic and electric fields.

$$\nabla \times \mathbf{H} = \mathbf{J} + \varepsilon \frac{\partial \mathbf{E}}{\partial t}$$

Ampère's law. Magnetic field and current.

$$\nabla \cdot (\varepsilon \mathbf{E}) = \rho$$

Gauss' law. Electric flux out of a closed surface surrounding the charge.

$$\nabla \cdot (\mu \mathbf{H}) = 0$$

No net magnetic flux out or in of a closed surface

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E: electric field (V/m)

$\mu$ : permeability (H/m)

H: magnetic field (A/m)

J: current density (A/m<sup>2</sup>)

$\varepsilon = \varepsilon_0 \varepsilon_r$ ,  $\varepsilon_0 = 8.854 \cdot 10^{-12}$  F/m

$\rho$ : charge density (C/m<sup>3</sup>)

$\mu = \mu_0 \mu_r$ ,  $\mu_0 = 4\pi \cdot 10^{-7}$  H/m

$\varepsilon$ : permittivity, dielectric constant (F/m)

$\varepsilon_r = \mu_r = 1$  in free space

# Wave equation in free space

In free space is  $\mathbf{J} = 0$  and  $\rho = 0$ .

Take the curl of the first equation

$$\nabla \times \left( \nabla \times \mathbf{E} + \mu_0 \frac{\partial \mathbf{H}}{\partial t} \right) = \nabla \times \nabla \times \mathbf{E} + \mu_0 \frac{\partial}{\partial t} \nabla \times \mathbf{H} = 0$$

$$\begin{aligned}\nabla \times \mathbf{E} &= -\mu \frac{\partial \mathbf{H}}{\partial t} \\ \nabla \times \mathbf{H} &= \mathbf{J} + \varepsilon \frac{\partial \mathbf{E}}{\partial t} \\ \nabla \cdot (\varepsilon \mathbf{E}) &= \rho \\ \nabla \cdot (\mu \mathbf{H}) &= 0\end{aligned}$$

Use  $\nabla \times \nabla \times \mathbf{E} = \nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}$

and second equation  $\nabla \times \mathbf{H} = \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t}$

Then  $-\nabla^2 \mathbf{E} + \mu_0 \varepsilon_0 \frac{\partial^2}{\partial t^2} \mathbf{E} = 0$  for  $\nabla \cdot \mathbf{E} = 0$  (in free space)

Wave equation:  $\nabla^2 \mathbf{E} - \mu_0 \varepsilon_0 \frac{\partial^2}{\partial t^2} \mathbf{E} = 0$

# Time-harmonic plane wave equation solution

With time dependency  $e^{j\omega t}$  a solution is

$$\mathbf{E} = E_0 \cos(\omega t - kz) \hat{\mathbf{x}}$$

for a wave propagating in the z-direction with the field varying in the x-direction.

The wave phase angle is  $\omega t - kz$ , where  $\omega$  is angle frequency,  $\omega = 2\pi f$  and  $k$  the wave number.  $E_0$  is the amplitude.

The wave number expresses the phase angle rate of change as a function of travelled distance  $z$ , i.e.

$$k = \frac{2\pi}{\lambda}$$

and



$$\cos(\omega t - kz) = \cos\left(\omega t - \frac{2\pi}{\lambda} z\right)$$

Convenient to use complex variables  $\mathbf{E} = E_0 e^{j(\omega t - kz)} \hat{\mathbf{x}}$

# Phase front velocity



Given a fixed point on the wave with an arbitrary phase,

$$\omega t - kz = \text{constant} \Rightarrow z = (\omega t - \text{constant})/k$$

then the phase front velocity  $v$

$$v = \frac{dz}{dt} = \frac{\omega}{k} = \lambda f$$

Given  $k^2 = \omega^2 \mu_0 \epsilon_0$  then  $\mathbf{E} = E_0 \cos(\omega t - \omega \sqrt{\mu_0 \epsilon_0} z) \hat{\mathbf{x}}$

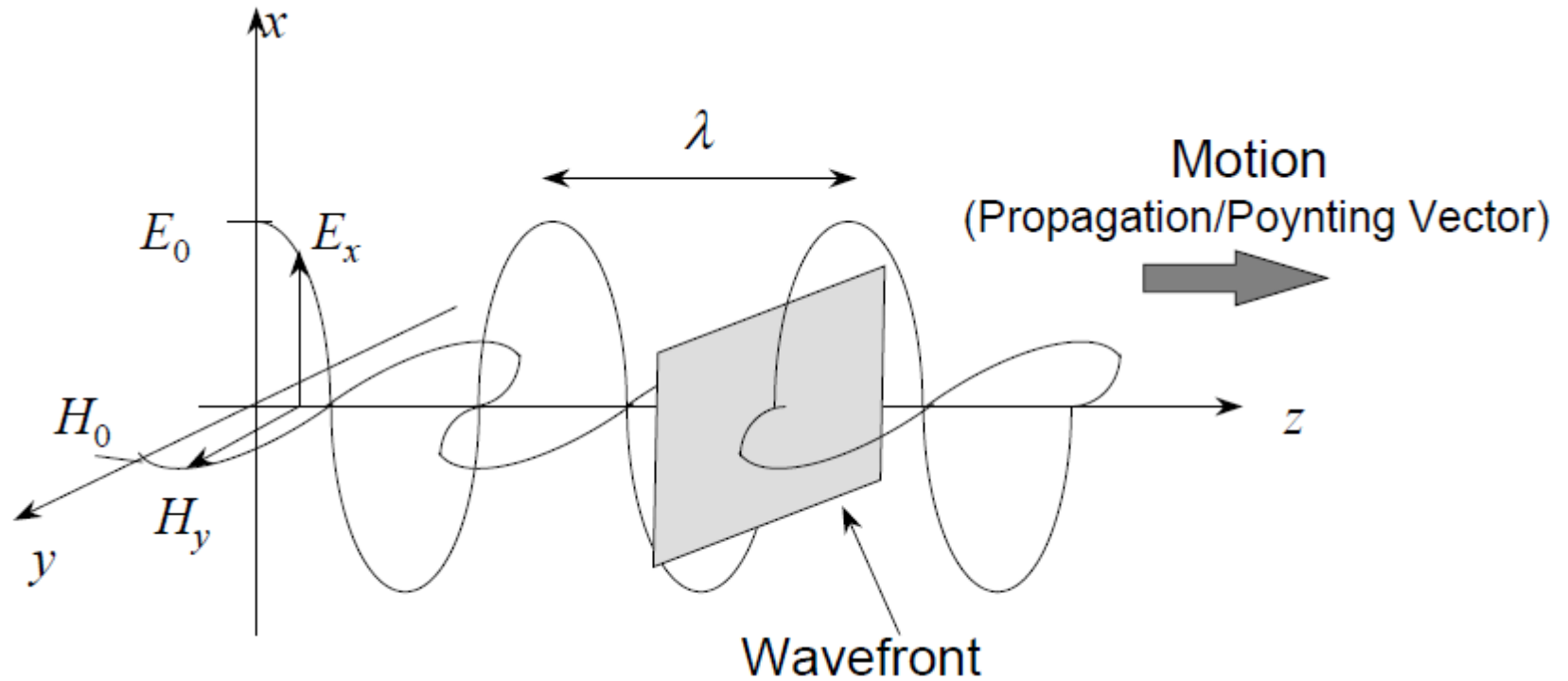
and

$$c = \frac{\omega}{k} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \approx 3 \cdot 10^8 \text{ m s}^{-1}$$

The ratio of velocity in free space to the velocity in the medium is the refractive index

$$n = \frac{c}{v}$$

# Plane wave propagation



# Wave impedance, $Z$

The solution for  $\mathbf{E}$  and  $\mathbf{H}$  satisfy Maxwell's equations provided

$$\frac{|\mathbf{E}|}{|\mathbf{H}|} = \frac{E_x}{H_y} = \frac{E_0}{H_0} = \sqrt{\frac{\mu}{\varepsilon}} = Z$$

for a given medium ( $\varepsilon, \mu$ ).

In free space ( $\varepsilon_0, \mu_0$ ) the impedance becomes

$$Z_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}} \approx \sqrt{4\pi \cdot 10^{-7} \cdot \frac{36\pi}{10^{-9}}} = 120\pi \approx 377\Omega$$



# Power flow Poynting vector, $\mathbf{S}$

From laws of conservation of energy and momentum the instantaneous power flux per unit area  $\mathbf{S}$  is

$$\mathbf{S} = \mathbf{E} \times \mathbf{H}^*$$

for a given medium  $(\epsilon, \mu)$ .

More useful with time averaged over one period

$$S_{av} = \frac{1}{2} E_0 H_0 \hat{\mathbf{z}}$$

# Lossy media

Write  $n = n' - jn''$  in a lossy medium. Note  $k = \omega n/c$

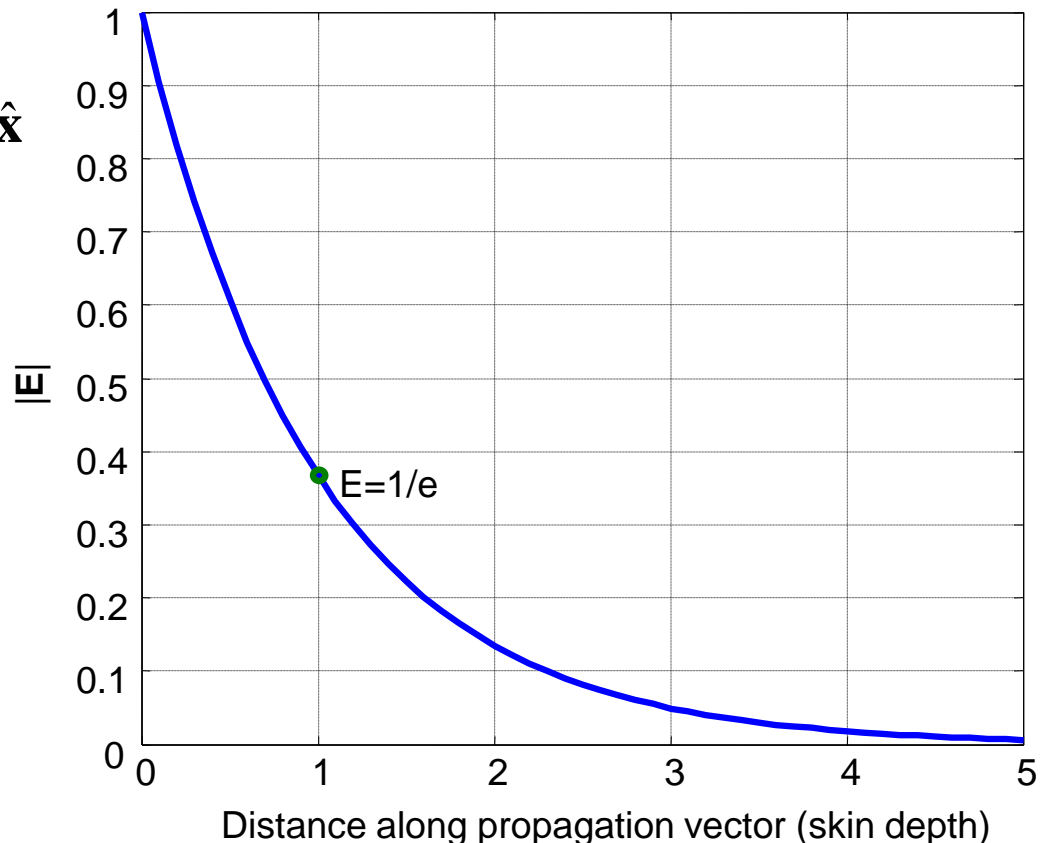
$$\mathbf{E} = E_0 \exp[j(\omega t - k' z) - \alpha z] \hat{\mathbf{x}}$$

Energy is removed from the wave and converted to heat.

The attenuation constant is  $\alpha$ .

Skin depth  $\delta$ : travelled distance where the field is reduced to  $1/e$  is

$$\delta = 1/\alpha$$



# Attenuation constant, wave number, wave impedance, wavelength and phase velocity (Table 2.1)

|                                                      | Exact expression                                                                                                 | Good dielectric<br>(insulator)<br>$(\sigma/\omega\epsilon)^2 \ll 1$ | Good conductor<br>$(\sigma/\omega\epsilon)^2 \gg 1$ |
|------------------------------------------------------|------------------------------------------------------------------------------------------------------------------|---------------------------------------------------------------------|-----------------------------------------------------|
| $n = ck/\omega$<br>in all cases                      |                                                                                                                  |                                                                     |                                                     |
| Attenuation constant<br>$\alpha$ [ $\text{m}^{-1}$ ] | $\omega \sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right]}$ | $\approx \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}}$              | $\approx \sqrt{\frac{\omega\mu\sigma}{2}}$          |
| Wavenumber<br>$k$ [ $\text{m}^{-1}$ ]                | $\omega \sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right]}$ | $\approx \omega \sqrt{\mu\epsilon}$                                 | $\approx \sqrt{\frac{\omega\mu\sigma}{2}}$          |
| Wave impedance<br>$Z$ [ $\Omega$ ]                   | $\sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}}$                                                             | $\approx \sqrt{\frac{\mu}{\epsilon}}$                               | $\approx \sqrt{\frac{\omega\mu}{2\sigma}} (1 + j)$  |
| Wavelength<br>$\lambda$ [m]                          | $\frac{2\pi}{k}$                                                                                                 | $\approx \frac{2\pi}{\omega \sqrt{\mu\epsilon}}$                    | $\approx 2\pi \sqrt{\frac{2}{\omega\mu\sigma}}$     |
| Phase velocity<br>$v$ [ $\text{m s}^{-1}$ ]          | $\frac{\omega}{k}$                                                                                               | $\approx \frac{1}{\sqrt{\mu\epsilon}}$                              | $\approx \sqrt{\frac{2\omega}{\mu\sigma}}$          |

## Example

A linear polarised 900 MHz wave travels in the  $z$ -direction in a medium with constitutive parameters  $\mu_r = 1$ ,  $\epsilon_r = 3$ , and  $\sigma = 0.01 \text{ Sm}^{-1}$ . The electric field at  $z = 0$  is  $1 \text{ Vm}^{-1}$ .

Calculate:

- a) the wave impedance
- b) the magnitude of the magnetic field at  $z = 0$
- c) the average power available in a  $0.5 \text{ m}^2$  area perpendicular to the direction of propagation at  $z = 0$
- d) the time taken for the wave to travel 10 cm
- e) the distance travelled by the wave before its field strength drops to one tenth its value at  $z = 0$

# Example solution

A linear polarised 900 MHz wave travels in the z-direction in a medium with constitutive parameters  $\mu_r = 1$ ,  $\epsilon_r = 3$ , and  $\sigma = 0.01 \text{ Sm}^{-1}$ .

The electric field at  $z = 0$  is  $1 \text{ Vm}^{-1}$ .

$$\frac{\sigma}{\omega\epsilon} = \frac{0.01}{2\pi \cdot 900 \cdot 10^6 \cdot 3 \cdot \frac{10^{-9}}{36\pi}} \approx 0.07 \ll 1 \quad \text{From Table 2.1: good insulator}$$

a) the wave impedance

$$Z \approx \sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_0}{\epsilon_r \epsilon_0}} = \frac{Z_0}{\sqrt{\epsilon_r}} \approx 218 \Omega$$

b) the magnitude of the magnetic field at  $z = 0$

$$H = \frac{E}{Z} = \frac{1}{218} \approx 4.6 \text{ Am}^{-1}$$

## Example solution continued

A linear polarised 900 MHz wave travels in the z-direction in a medium with constitutive parameters  $\mu_r = 1$ ,  $\epsilon_r = 3$ , and  $\sigma = 0.01 \text{ Sm}^{-1}$ .

The electric field at  $z = 0$  is  $1 \text{ Vm}^{-1}$ .

- c) the average power available in a  $0.5 \text{ m}^2$  area perpendicular to the direction of propagation at  $z = 0$

$$P = SA = \frac{EH}{2} A = \frac{1 \cdot 0.005}{2} = 1.25 \text{ mW}$$

- d) the time taken for the wave to travel 10 cm

$$t = \frac{d}{v} = d \sqrt{\mu \epsilon} = d \sqrt{\mu_0 \epsilon_r \epsilon_0} = \frac{d}{c} \sqrt{\epsilon_r} \approx \frac{0.1}{3 \cdot 10^8} \sqrt{3} \approx 0.6 \text{ ns}$$

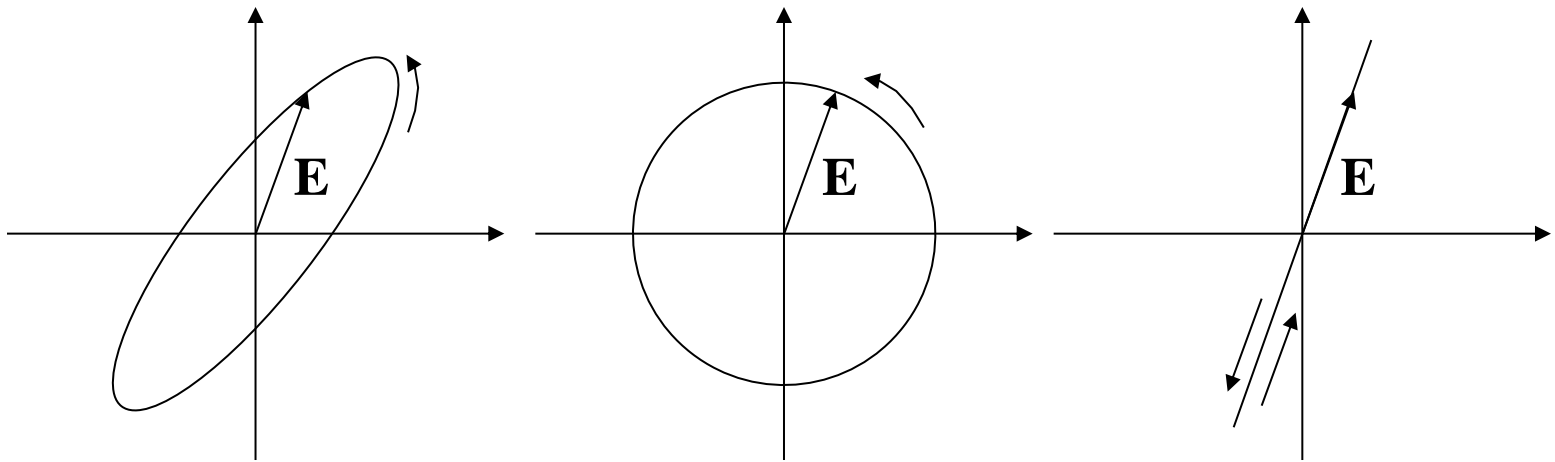
- e) the distance travelled by the wave before its field strength drops to one tenth its value at  $z = 0$

$$E(z) = E(0)e^{-z/\delta} \Rightarrow z = -\delta \ln \frac{E(z)}{E(0)} \quad \text{where} \quad \delta = \frac{1}{\alpha} \approx \frac{2}{\sigma} \sqrt{\frac{\epsilon}{\mu}} = \frac{2}{\sigma Z}$$

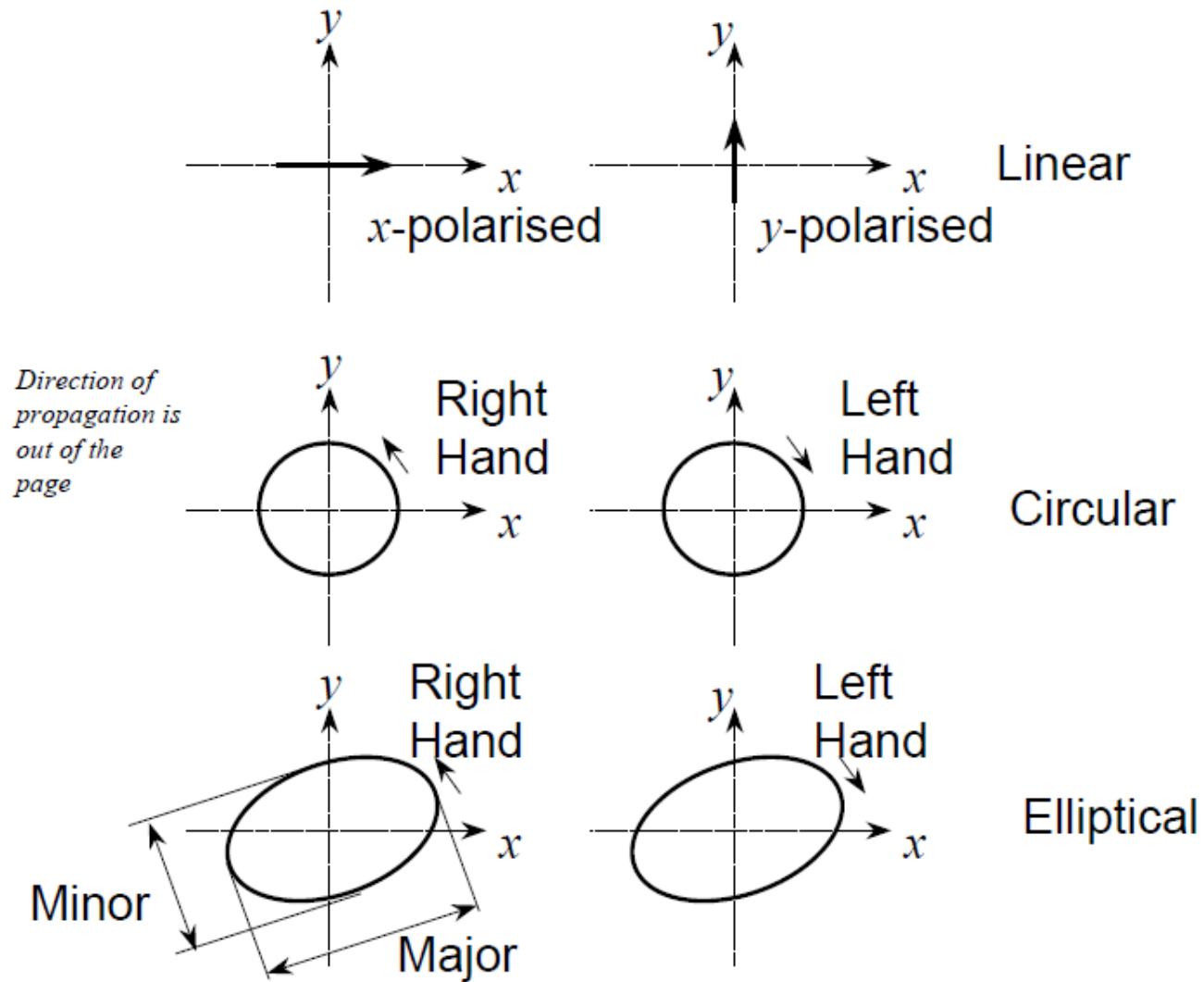
$$z = -\frac{2}{\sigma Z} \ln \frac{E(z)}{E(0)} = -\frac{2}{0.01 \cdot 218} \ln \frac{1}{10} = 2.11 \text{ m}$$

# Polarisation

Generally, along the propagation direction the tip of the time-varying electric field vector will form an ellipse, where circle or line is degenerated special cases



# Polarisation states





# Conclusion

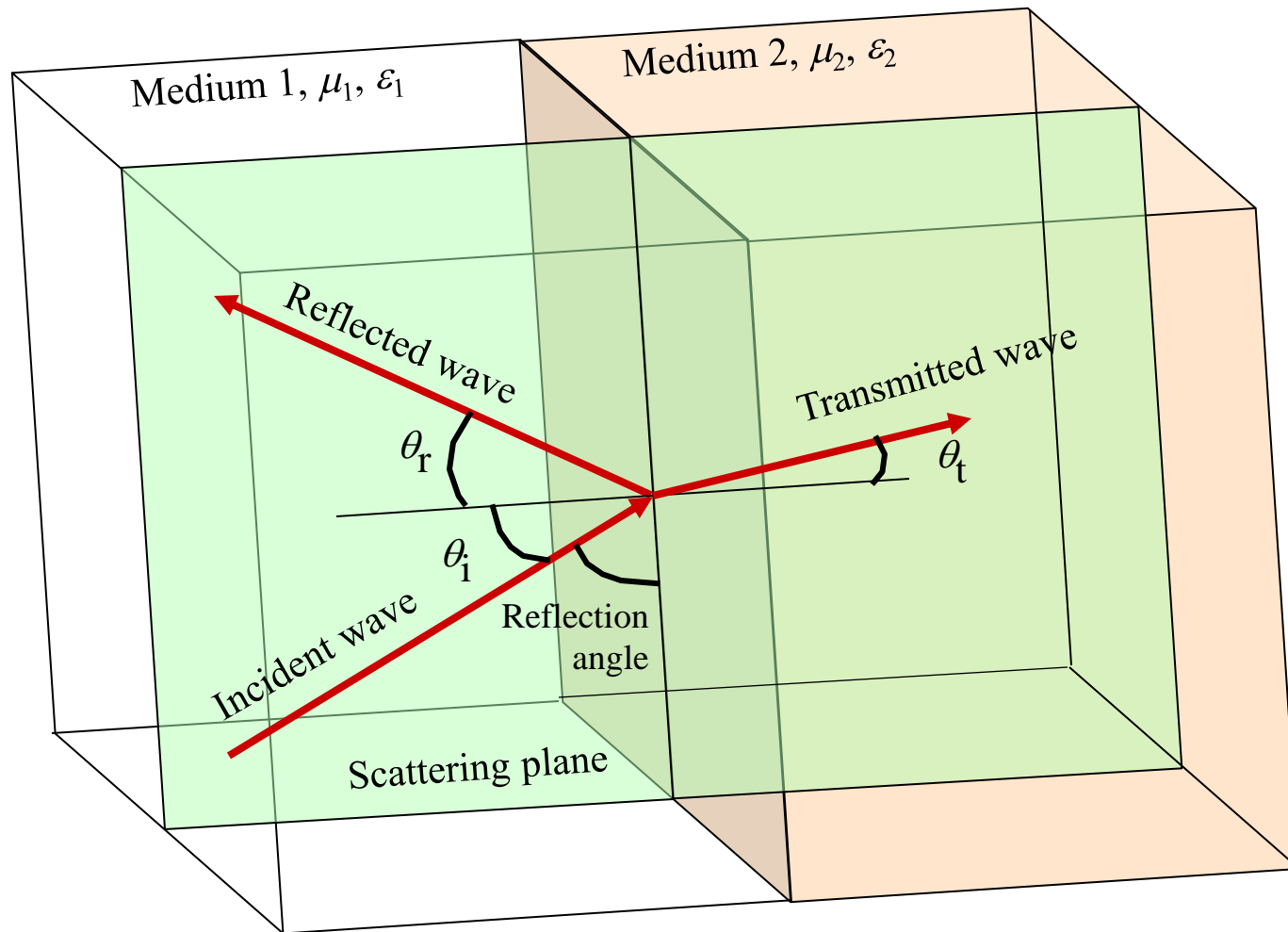
- Plane wave in uniform media
- Plane wave properties and interaction with the medium
- Poynting vector
- Polarisation

# Chapter 3

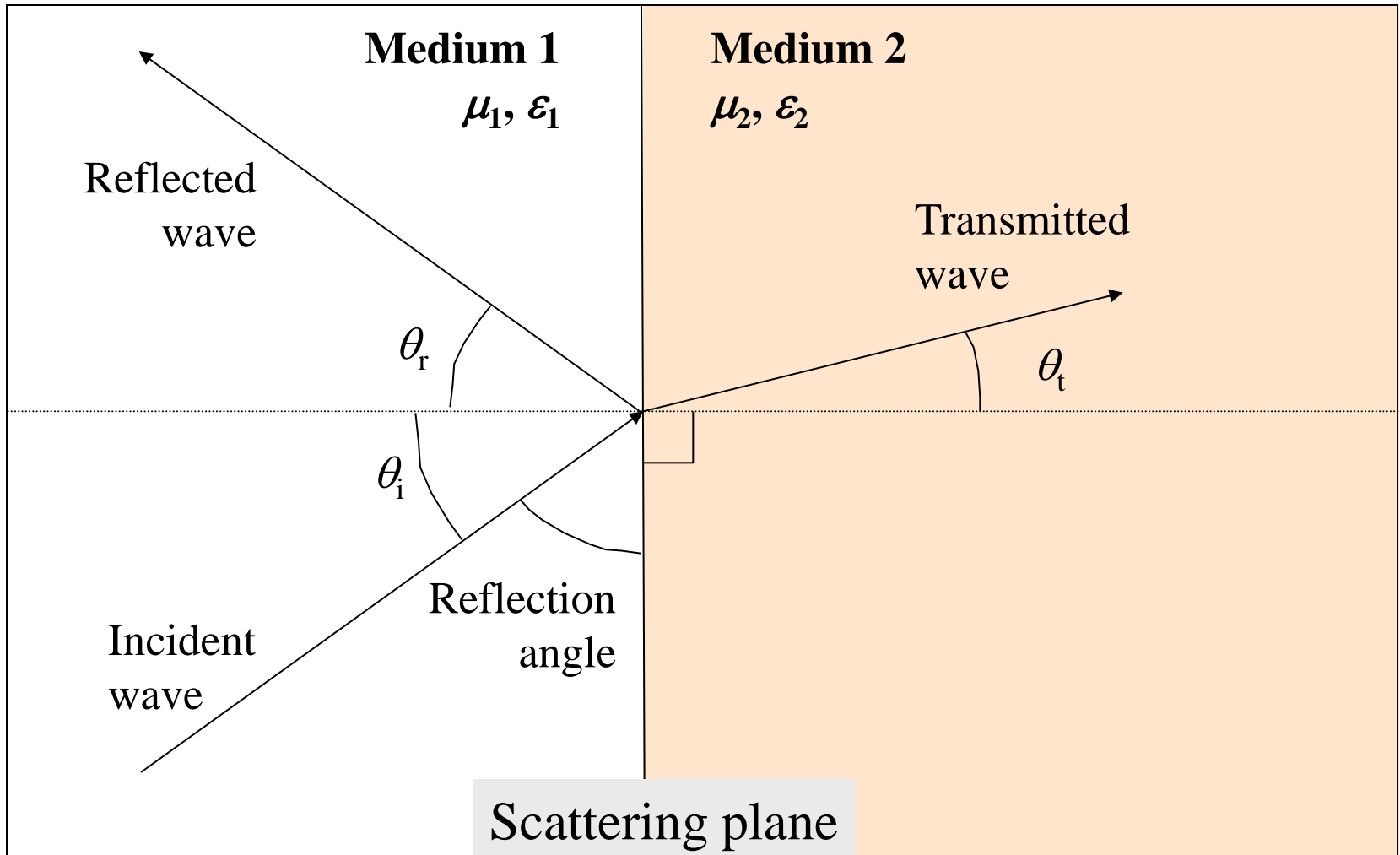
## Propagation mechanisms

- Reflection and refraction: reflected and transmitted wave
- Rough surfaces
- Diffraction

# Reflection and refraction



# Reflection and refraction



# Snell's law of refraction

Wave number  $k$  and Fermat's principle minimising the electrical path length

$$\frac{\sin \theta_i}{\sin \theta_t} = \frac{k_2}{k_1}$$

Refractive index  $n$ , the ratio of free space phase velocity  $c$  to the phase velocity in the medium

$$n = \frac{c}{v} = \frac{ck}{\omega}$$

$$\frac{\sin \theta_i}{\sin \theta_t} = \frac{n_2}{n_1}$$

# Fresnel reflection and transmission coefficients

$$R_{\parallel} = \frac{E_{r\parallel}}{E_{i\parallel}} = \frac{Z_1 \cos \theta_i - Z_2 \cos \theta_t}{Z_2 \cos \theta_t + Z_1 \cos \theta_i} \quad R_{\perp} = \frac{E_{r\perp}}{E_{i\perp}} = \frac{Z_2 \cos \theta_i - Z_1 \cos \theta_t}{Z_2 \cos \theta_i + Z_1 \cos \theta_t}$$
$$T_{\parallel} = \frac{E_{t\parallel}}{E_{i\parallel}} = \frac{2Z_2 \cos \theta_i}{Z_2 \cos \theta_t + Z_1 \cos \theta_i} \quad T_{\perp} = \frac{E_{t\perp}}{E_{i\perp}} = \frac{2Z_2 \cos \theta_i}{Z_2 \cos \theta_i + Z_1 \cos \theta_t}$$

Total reflected field:

$$\mathbf{E}_r = E_{r\parallel} \mathbf{a}_{\parallel} + E_{r\perp} \mathbf{a}_{\perp} = E_{i\parallel} R_{\parallel} \mathbf{a}_{\parallel} + E_{i\perp} R_{\perp} \mathbf{a}_{\perp}$$

$$\mathbf{E}_i = E_{i\parallel} \mathbf{a}_{\parallel} + E_{i\perp} \mathbf{a}_{\perp}$$

where  $\mathbf{a}_{\parallel}$  and  $\mathbf{a}_{\perp}$  are unit vectors parallel and normal to the scattering plane, respectively

# Reflected and transmitted field

$$\mathbf{E}_r = \mathbf{R}\mathbf{E}_i$$

where

$$\mathbf{E}_r = \begin{bmatrix} E_{r\parallel} \\ E_{r\perp} \end{bmatrix}, \mathbf{R} = \begin{bmatrix} R_{\parallel} & 0 \\ 0 & R_{\perp} \end{bmatrix}, \mathbf{E}_i = \begin{bmatrix} E_{i\parallel} \\ E_{i\perp} \end{bmatrix}$$

similarly

$$\mathbf{E}_t = \mathbf{T}\mathbf{E}_i$$

where

$$\mathbf{E}_t = \begin{bmatrix} E_{t\parallel} \\ E_{t\perp} \end{bmatrix}, \mathbf{T} = \begin{bmatrix} T_{\parallel} & 0 \\ 0 & T_{\perp} \end{bmatrix}, \mathbf{E}_i = \begin{bmatrix} E_{i\parallel} \\ E_{i\perp} \end{bmatrix}$$

# Loss media

Fresnel reflection and transmission coefficients still hold given

$$Z = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}}$$



## Example

A linear polarised plane wave in free space with an electric field amplitude 30 V/m is incident with angle  $30^\circ$  onto a plane boundary of a lossless non-magnetic medium with  $\epsilon_r=2$ . The electric field is parallel to the plane of incidence.

Calculate

- a) angle of reflection
- b) angle of refraction
- c) amplitude of transmitted electric and magnetic fields
- d) amplitude of electric field

## Example solution

- a) Wave and boundary are plane. Snell's law gives reflection angle of  $30^\circ$ .  
 b) Use Snell's law

$$\theta_t = \sin^{-1}\left(\frac{n_1}{n_2} \sin \theta_i\right) = \sin^{-1}\left(\frac{\sqrt{\mu_{r1}\epsilon_{r1}}}{\sqrt{\mu_{r2}\epsilon_{r2}}} \sin \theta_i\right) = \sin^{-1}\left(\frac{\sqrt{1}}{\sqrt{2}} \sin 30\right) = 20.7^\circ$$

- c) Find  $T_{\parallel}$  with  $Z_1 = Z_0 = 377\Omega$  and  $Z_2 = Z_0/\sqrt{\epsilon_{r2}} = 377/\sqrt{2}\Omega$

$$T_{\parallel} = \frac{2 \cos \theta_i}{\cos \theta_t + \frac{Z_1}{Z_2} \cos \theta_i} = \frac{2 \cos 30^\circ}{\cos 20.7^\circ + \sqrt{2} \cos 30^\circ} \approx 0.8$$

$$E_t = E_i T_{\parallel} = 0.8 \cdot 30 = 24 \text{Vm}^{-1} \quad H_t = \frac{E_t}{Z_2} = \frac{24\sqrt{2}}{377} \approx 0.09 \text{Am}^{-1}$$

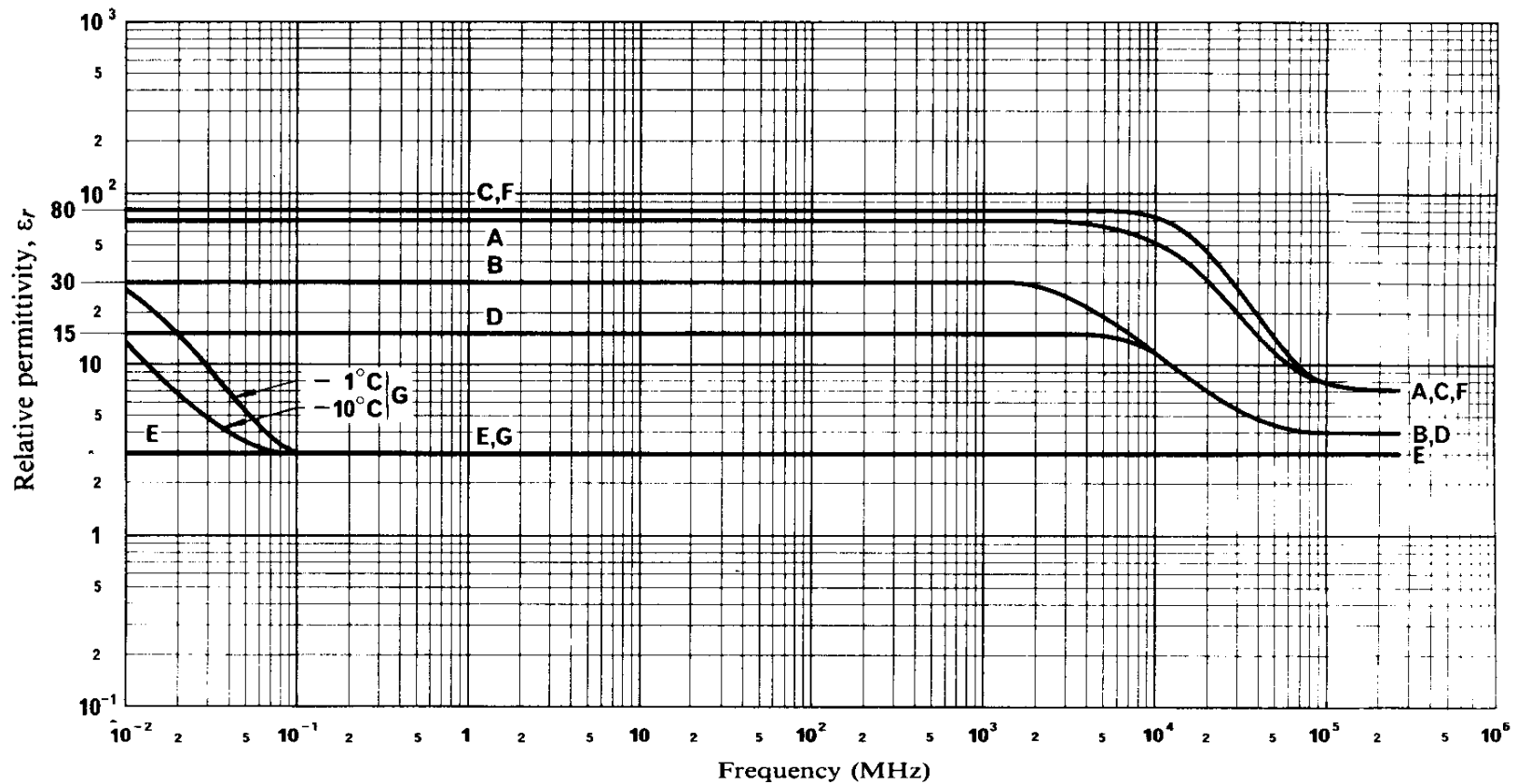
- d) Reflected electric field

$$E_r = R_{\parallel} E_i = \frac{\frac{Z_1}{Z_2} \cos \theta_i - \cos \theta_t}{\cos \theta_t + \frac{Z_1}{Z_2} \cos \theta_i} E_i \approx -8.54 \cdot 30 = -256.2 \text{Vm}^{-1}$$

# Typical transmission and reflection coefficients

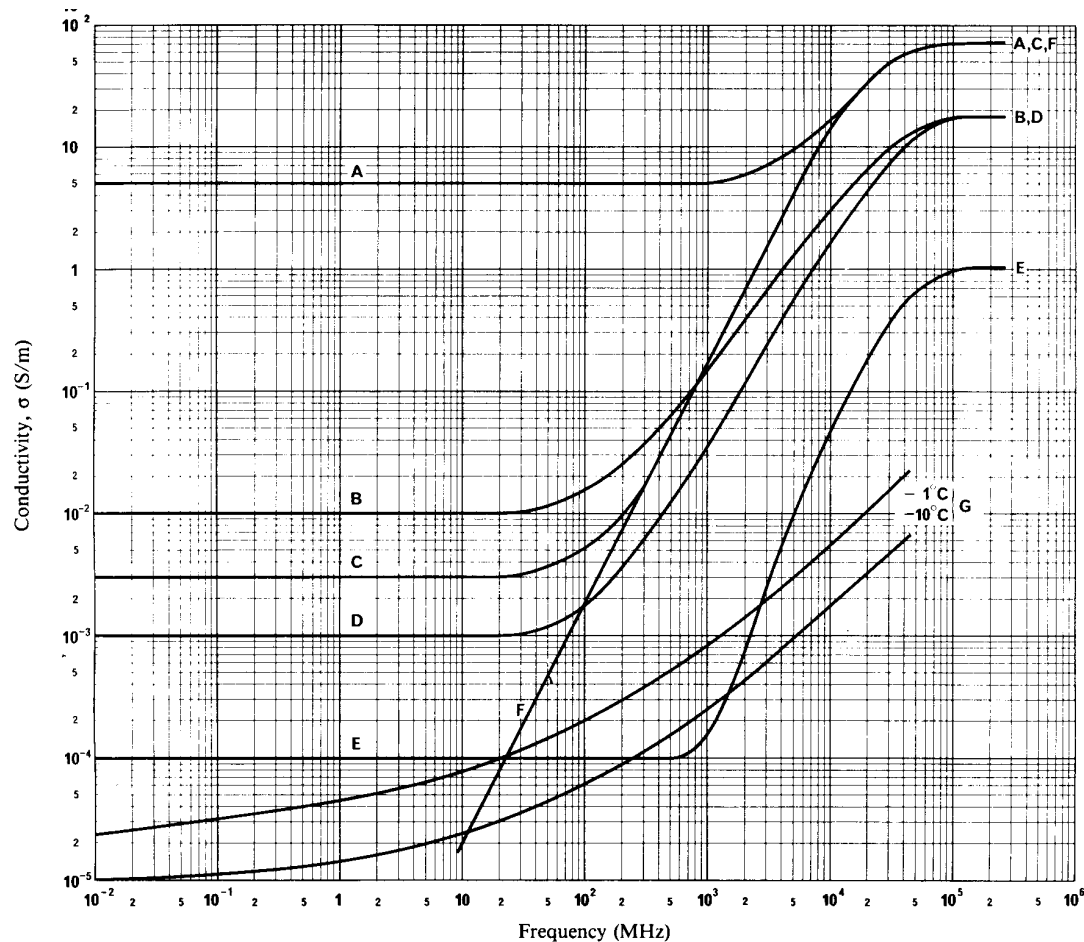
| Surface        | Conductivity, $\sigma$ [S m <sup>-1</sup> ] | Relative dielectric constant, $\epsilon_r$ |
|----------------|---------------------------------------------|--------------------------------------------|
| Dry ground     | 0.001                                       | 4-7                                        |
| Average ground | 0.005                                       | 15                                         |
| Wet ground     | 0.02                                        | 25-30                                      |
| Sea water      | 5                                           | 81                                         |
| Fresh water    | 0.01                                        | 81                                         |

# Relative permittivity, $\epsilon_r$ (ITU-R Rec.P.527)



- A: sea water (average salinity), 20° C
- B: wet ground
- C: fresh water, 20° C
- D: medium dry ground
- E: very dry ground
- F: pure water, 20° C
- G: ice (fresh water)

# Conductivity, $\sigma$ (ITU-R Rec.P.527)

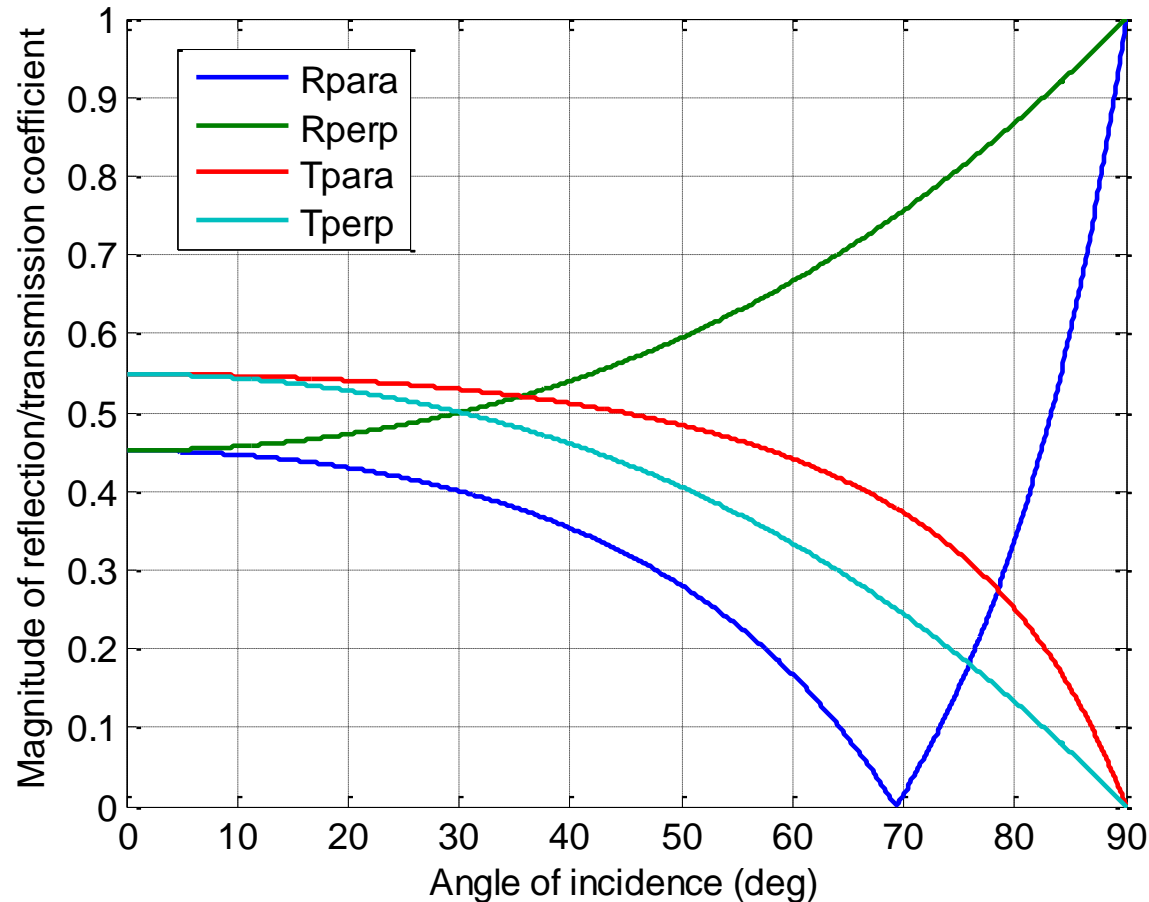


- A: sea water (average salinity), 20° C
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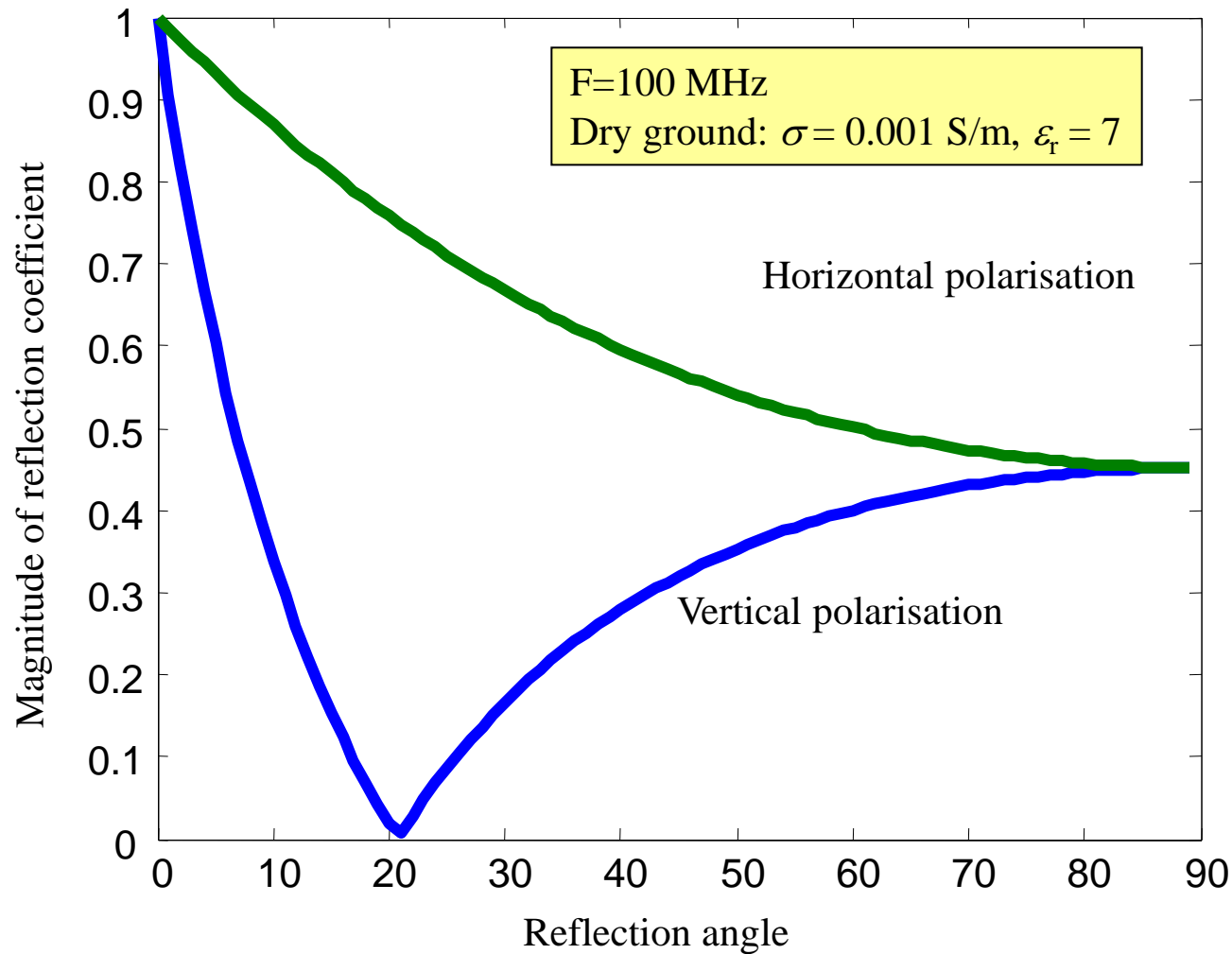
# Reflection and transmission coefficients

F=100 MHz

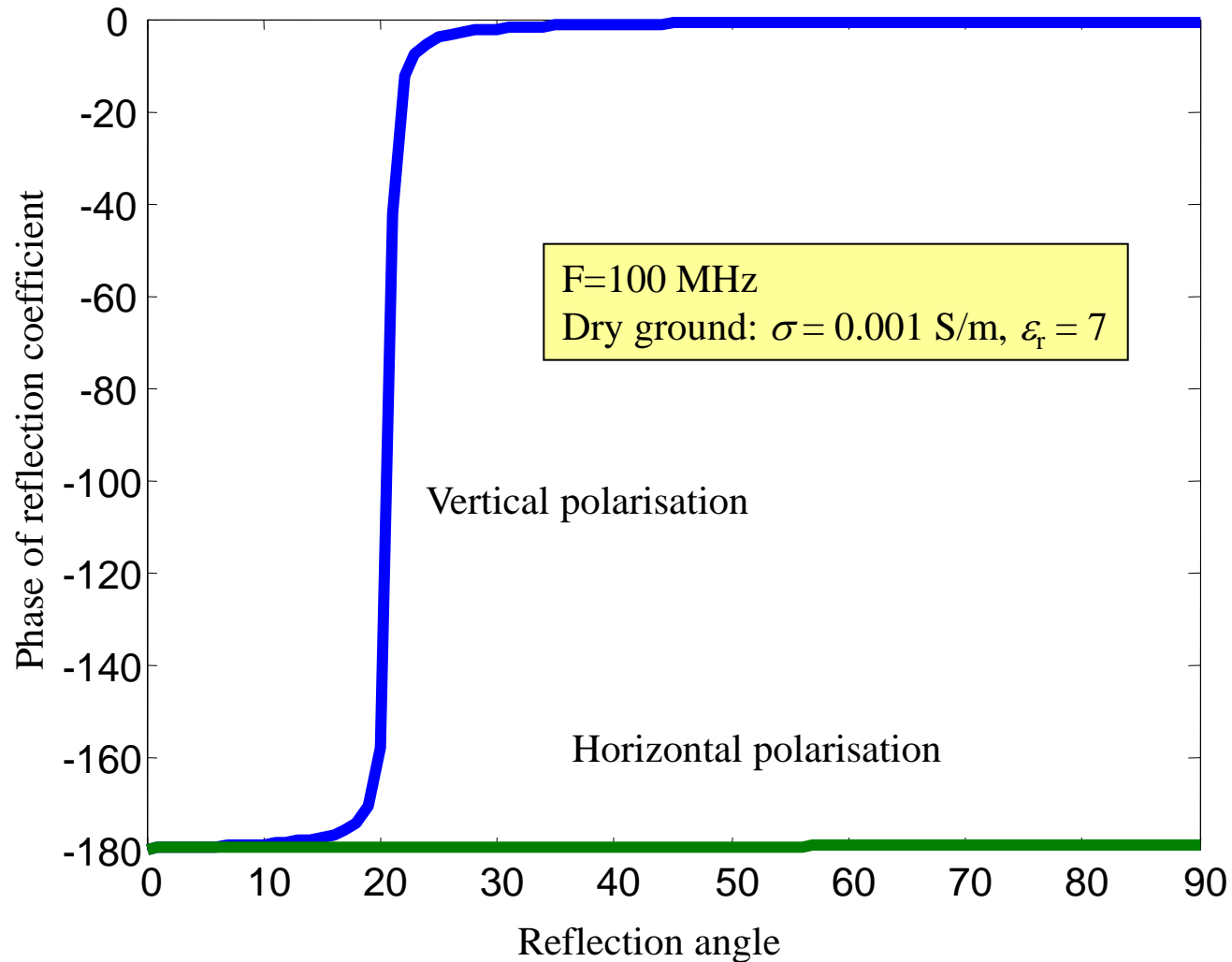
Dry ground:  $\sigma = 0.001$  S/m,  $\epsilon_r = 7$



# Example absolute value of reflection coefficient



# Example phase of reflection coefficient





# Brewster's angle

- The reflection example:
  - $\parallel$  represents the vertical polarisation
  - $\perp$  represents the horizontal polarisation
- Vertical polarisation reflected amplitude goes to zero at angle  $\theta_B$ , Brewster's angle

$$\theta_B = \tan^{-1} \frac{n_2}{n_1}$$

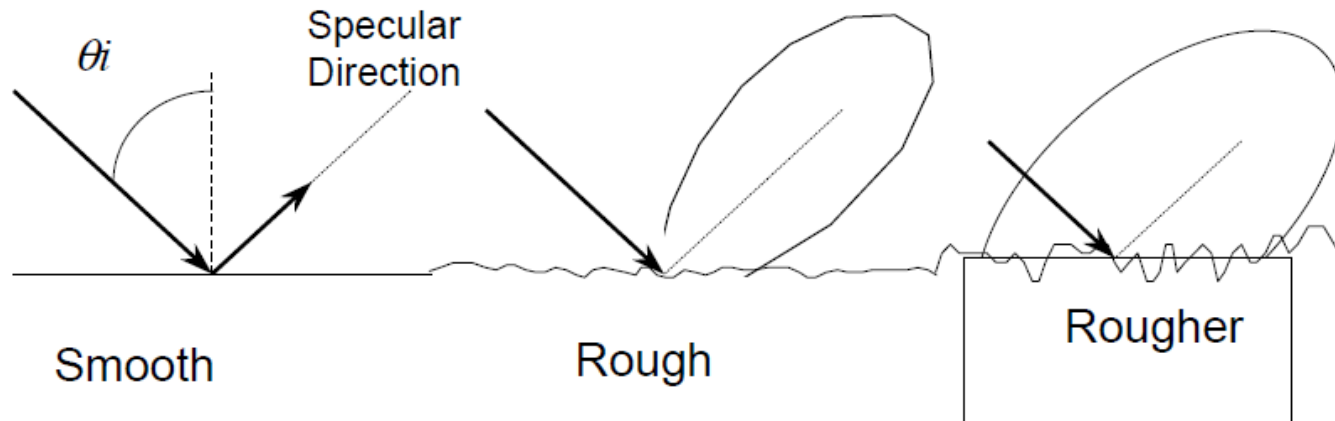
- At reflection angle of 0, i.e., incidence angle of 90, the reflected amplitude is -1. This is the grazing angle
- Note that vertical polarisation undergoes a phase shift from 0 to -180

# Change of polarisation state

$$R_{co} = \frac{1}{2}(R_{\parallel} + R_{\perp}) \quad \text{and} \quad R_{cx} = \frac{1}{2}(R_{\parallel} - R_{\perp})$$

|                        | $\theta_i < \theta_B$                          | $\theta_i > \theta_B$                          |
|------------------------|------------------------------------------------|------------------------------------------------|
| Right-hand circular    | Left-hand elliptical                           | Right-hand elliptical                          |
| Left-hand circular     | Right-hand elliptical                          | Left-hand elliptical                           |
| Right -hand elliptical | Left-hand elliptical<br>(axial ratio changed)  | Right-hand elliptical<br>(axial ratio changed) |
| Left-hand elliptical   | Right-hand elliptical<br>(axial ratio changed) | Left-hand elliptical<br>(axial ratio changed)  |

# Rough surface

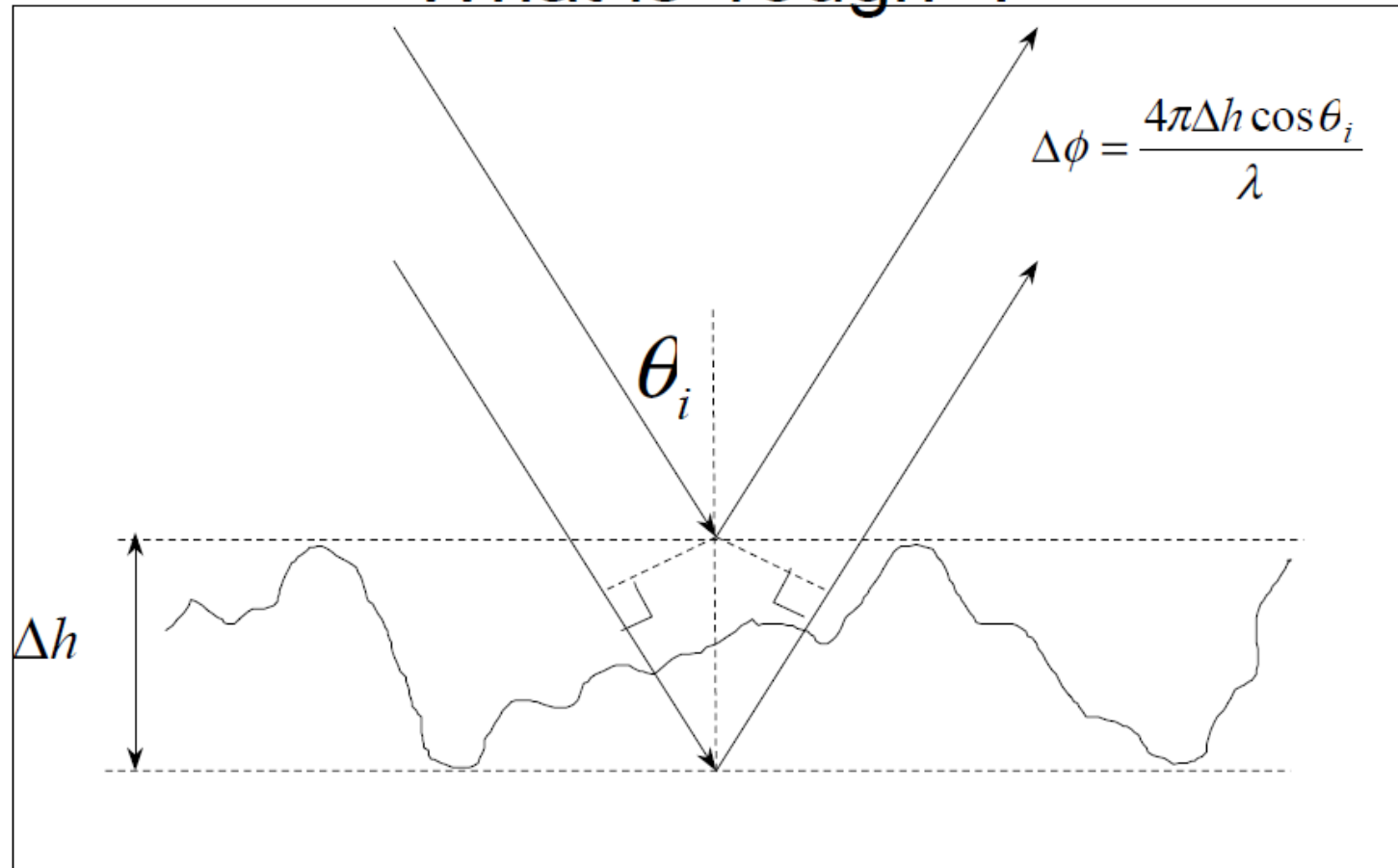


Roughness depends on :

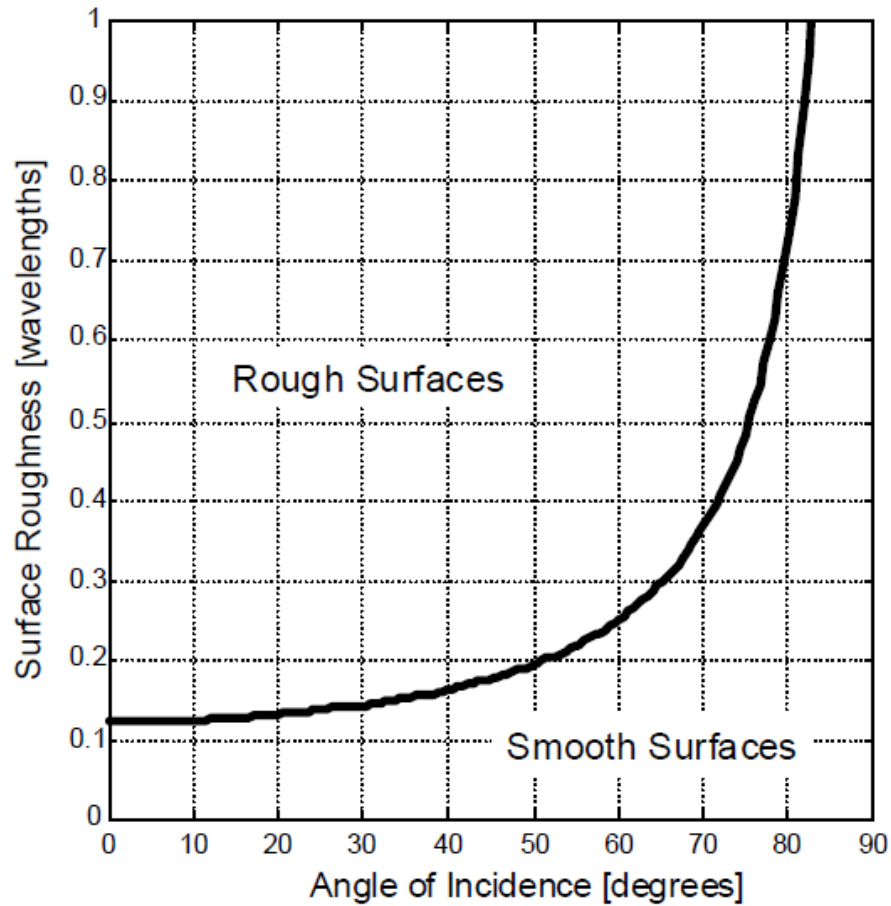
- Surface height range
- Angle of incidence
- Wavelength

# Definition of rough surface

## What is 'rough' ?



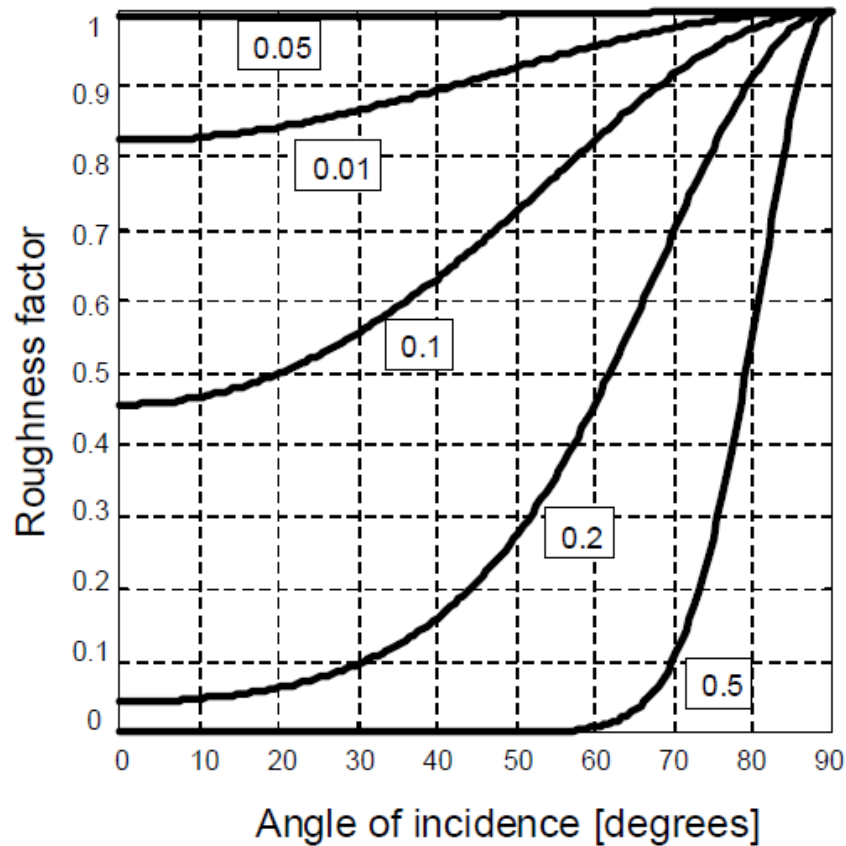
# Rayleigh criterion



$$\Delta\phi < \frac{\pi}{2}$$

$$\Delta h < \frac{\lambda}{8 \cos \theta_i}$$

# Roughness factor $f$



$$R_{eff} = Rf(\sigma_s)$$

$$f(\sigma_s) = \exp\left[-\frac{1}{2}\left(\frac{4\pi\sigma_s \cos\theta}{\lambda}\right)^2\right]$$

# Conclusion

- Reflection and refraction
- Transmission
- Rough surface scattering