Comparing Models of Delta-Notch Signalling

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Abstract

The Delta-Notch signalling pathway plays a crucial role in determining the fate of initially homogeneous cells in a wide array of biological systems. In this paper, we present a simple yet biologically realistic representation of Delta-Notch signalling and test it using three models: ODE, SDE, and agent-based. Using the Kolmogorov forward equations, we show that the agent-based model converges to the ODE model in expectation. We also demonstrate that the SDE and agent-based models can produce cell differentiation even from identical initial conditions. Finally, we study the effect of parameters such as stochastic noise, domain size, and boundary conditions on the behaviour of our models. Our findings provide insight into how modelling formalisms can affect simulations of the Delta-Notch signalling network with significant implications for future research.

Introduction

Cell differentiation is essential for the formation of structure in complex organisms. But surprisingly, only a small number of signalling pathways are responsible for this important function [3]. One of these pathways is the Delta-Notch signalling network, which has been shown to affect neurogenesis in vertebrates by determining which cells adopt neural or glial fates. Furthermore, mutations in Notch activity has been shown to be a contributing factor to many cancers [3]. The Delta-Notch signalling network is unique because its ligands Delta and Jagged are transmembrane proteins. This means that Delta-Notch signalling is largely restricted to neighbouring cells. However, new findings suggest that cellular protrusions can allow a small fraction of Delta-Notch signals to travel longer distances [1].

Binding of the Delta ligand from cell B to the Notch receptor on cell A results in the cleavage (removal) of both the ligand and receptor [3]. This interaction triggers the release of NICD (Notch Intracellular Domain) in cell A, which promotes Notch and inhibits Delta. It is also possible for a Delta ligand to bind to a Notch receptor on the same cell. This phenomenon, known as cis-inhibition, cleaves the Notch receptor but does not release NICD. Serrate (Jagged) is a ligand similar to Delta which can also bind to Notch. However, NICD promotes Serrate, which creates a double-positive feedaback loop.

Previous models of Delta-Notch signalling [4, 2] have made a number of assumptions which we also use for our models. First, cells interact through Delta-Notch signalling if and only if they are adjacent. Production of Notch is an increasing hill function of NICD. Production of Delta is a decreasing hill function of NICD. Both Notch and Delta decay at a rate proportional to their concentration. Low levels of Notch cause the primary cell fate, while high levels cause the secondary cell fate. The behaviour of cells with a specific fate depends on the organism and organ, which is outside the scope of these models.

The model of Delta-Notch signalling proposed by Collier et al. [4] makes other simplifications which violate the modern understanding of the signalling pathway. These

simplifications make it difficult to identify how individual chemical reactions are captured by the model. For instance, production of activated Notch (NICD) is exclusively a function of the Delta concentration in neighbouring cells, with no regard for the concentration of Notch receptors. This violates the mass-balance equations, since the quantity of Notch receptors changes dynamically in response to the concentration of Delta ligands in nearby cells. There are smaller issues, like the absence of cis-inhibition and the Serrate (Jagged) protein, but these omissions are not necessary to capture the behaviour of the system.

In this paper, we develop three models of the Delta-Notch signalling network based on a modern understanding of the biochemistry of the system. Our first model is an agent-based simulation of individual Delta, Notch, and NICD molecules. While this model precisely simulates the reaction kinetics under our assumptions, it is impossible to analyze analytically and slow to simulate computationally. Then, we define two models based on differential equations. The ordinary differential equation (ODE) model is equivalent the expected behaviour of the agent-based model (see Appendix C for a proof). The stochastic differential equation (SDE) model is an extension to the ODE model which contains additional noise terms. These terms are intended to simulate the stochasticity of the agent-based model. We compare these models with each other using a combination of analytical and computational methods in a variety of different domains. Furthermore, we show that our models qualitatively replicate the results of Collier et al. [4]. In particular we show that if a linear array of cells is given zero-flux boundary conditions, the pattern of alternating cell fates arises on the outside edges before gradually spreading inwards towards the center of the domain. Additionally, both our model and the Collier et al. model exhibit different ratios of primary to secondary fated cells depending on the parameters on a two-dimensional hexagonal lattice.

Methods

From the literature [4, 2], we assume Notch (N) receptor production has the form $H^+(I) = n_m I^2/(n_0^2 + I^2)$ where I is the NICD concentration and n_m, n_0 are hill function parameters. Similarly, Delta ligand (D) production is governed by a decreasing hill function of the for production has the form $H^-(I) = d_m d_0^2/(d_0^2 + I^2)$. Furthermore, we will assume that Notch and Delta both have decay rate γ . The variable k_T will denote the binding rate between Notch receptors and Delta ligands. γ_I is the decay rate of NICD. We will also assume that there is no cis-inhibition and no Serrate ligand.

Reaction	Rate	Description
		A Notch receptor from cell A binds to a delta
$N_A + D_B \rightarrow I_A$	$k_T N_A D_B$	ligand from cell B , triggering the release of
		NICD in cell A .
$N_A o \emptyset$	γN_A	A Notch receptor from cell A decays.
$D_A \to \emptyset$	γD_A	A Delta receptor from cell A decays.
$I_A o \emptyset$	$\gamma_I I_A$	A NICD molecule from cell A decays.
$\emptyset \to N_A$	$H^+(I_A)$	A Notch receptor in cell A is produced
$\emptyset \to D_A$	$H^-(I_A)$	A Delta ligand in cell A is produced

Table 1: A list of possible reactions and their rates in a cell A with neighbour B. Since there are 6 possible reactions per cell, a system with n cells will have 6n possible reactions.

We will assume that reaction times follow a Poisson point process. Suppose E is the set of all possible reaction events (the 6 reactions above for all cells). Furthermore, let |E| = N be the number of possible reactions. For each event, let $X_i \sim \text{Exp}(\lambda_i)$ denote the time until the next event of type i, where $1 \leq i \leq N$. Then, the waiting time T in between any two events is:

$$T = \min\{X_1, \dots, X_N\} \sim \operatorname{Exp}\left(\sum_{i=1}^N \lambda_i\right)$$

This results follows from the definition of an exponential distribution and the memoryless property. We can use Gillespie's algorithm to simulate the system. This is the *agent-based* model. Details of the computational implementation can be found in Appendix A.

In expectation, the agent-based model converges to the system of ordinary differential equations shown in Equation 1. A proof of this fact can be found in the appendix. This is the ODE model, which is described below for a single cell in a two-cell system. However, this model can be extended to arbitrary domains by defining equations dN_i/dt , dD_i/dt , and dI_i/dt for each $i \in \{1, ..., n\}$, where n is the number of cells. In a domain of arbitrary size, N_{ext} and D_{ext} denote the average Notch and Delta concentration over all neighbours.

$$\frac{dN}{dt} = \frac{n_m I^2}{n_0^2 + I^2} - k_T N D_{ext} - \gamma N$$

$$\frac{dD}{dt} = \frac{d_m d_0^2}{d_0^2 + I^2} - k_T D N_{ext} - \gamma D$$

$$\frac{dI}{dt} = k_T N D_{ext} - \gamma_I I$$
(1)

This model is a simplification of a model by Boareto et al. [2]. However, to more directly compare our results with those of Collier et al. [4], we chose to ignore Serrate and cis-inhibition. Adding NICD, which was absent from Collier et al., was necessary for the assignment of probabilities to each reaction event. An analysis of the system of ODEs in Equation 1 for fixed values of N_{ext} , D_{ext} reveals an N=0 steady state (the primary fate) for all parameter values. There is also a high Notch steady state (the secondary fate) which emerges from a fold bifurcation when D_{ext} is sufficiently large. See Appendix D for more information.

To capture the inherent randomness in receptor-ligand interactions while maintaining computational efficiency, we implement a stochastic differential equation SDE model based on the Langevin equation formalism. This approach represents an intermediate level of detail between the ODE model and the fully discrete agent-based model. Starting from the ODE model described in Equation 1, we incorporate stochastic fluctuations by adding noise terms proportional to the square root of the reaction rates to get Equation 2.

$$dN = \left(\frac{n_m I^2}{n_0^2 + I^2} - k_T N D_{ext} - \gamma N\right) dt + \sigma \cdot N \cdot dW_N$$

$$dD = \left(\frac{d_m d_0^2}{d_0^2 + I^2} - k_T D N_{ext} - \gamma D\right) dt + \sigma \cdot D \cdot dW_D$$

$$dI = (k_T N D_{ext} - \gamma_I I) dt + \sigma \cdot I \cdot dW_I$$
(2)

In Equation 2, dW_N , dW_D , and dW_I are independent Wiener processes representing white noise. The structure of the noise terms follows from the assumption that noise in each reaction channel is proportional to the current concentration of those molecules.

To implement these models computationally, we employed the Euler-Maruyama method with time step dt. The deterministic increment calculated first, and then the Wiener increments are sampled from a normal distribution with mean 0 and standard deviation \sqrt{dt} . Then, each Wiener increment is scaled by its respective variable and added to the deterministic increment.

Results

Two-Cell Domains

Collier et al. [4] observed cell differentiation in their ODE model under small initial perturbations on a two-cell domain. In our first experiment, we attempted to replicate this result under all three modelling formalisms. Since the SDE and agent-based models are stochastic, we did not apply an initial perturbation to them. For the ODE model, which is deterministic, we perturbed the initial Notch concentration by a factor of X, where X is sampled from a $\mathcal{N}(1,0.1)$ random variable. The results of this experiment are shown in Figure 1.

Comparison of Models on Two-Cell Domain

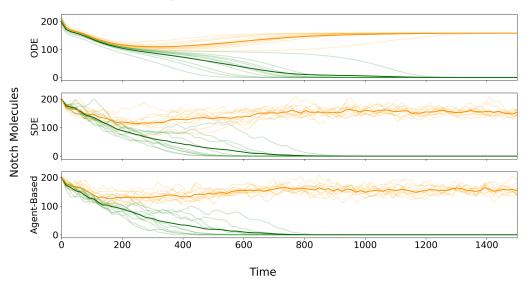


Figure 1: Results of ten simulations of each modelling formalism (ODE, SDE, agent-based) on a two-cell domain with periodic boundary conditions. An initial perturbation was applied to the ODE model to move the system off of an unstable steady state. In all simulations, the system eventually converged to a differentiated state. The number of Notch molecules in the primary-fated and secondary-fated cells are shown in green and yellow respectively. The bold line denotes the mean number of Notch molecules over all simulations.

Using stochastic formalisms such as the SDE and agent-based models, we can investigate more nuanced questions about the behaviour of the Delta-Notch signalling network. For our next experiment, we attempted to quantify how easy it would be force the two-cell system to adopt a certain outcome by perturbing the initial Notch concentration. We predicted that the cell with the lower initial number of Notch molecules would adopt the primary (low Notch) fate, while the other cell would adopt the secondary (high Notch) fate. While our prediction was generally true, it turns out the randomness of the agent-based model makes difficult to influence the outcome of a simulation using an initial perturbation. The results of this experiment are shown in Figure 2.

When we defined the SDE model in Equation 2, we introduced a parameter σ that represents the amount of noise. However, with the methods we know, it is difficult to estimate the value of this parameter analytically or numerically. Here, we present a computational method to get a qualitative estimate of σ which uses the differentiation time of the two-cell system. We define the differentiation time t^* to be the time at which one cell adopts the primary fate (N < 1). Through our experiments, we noticed that

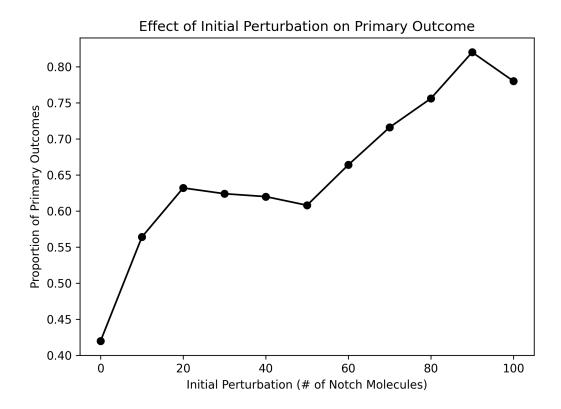


Figure 2: A Primary Outcome occurs when the cell with the lower initial Notch concentration ends up adopting the primary (low Notch) fate. This plot shows the relationship between initial perturbation and the sample proportion of primary outcomes for the agent-based model on a two-cell domain. Prior to the perturbation, each cell has 200 Notch molecules. A perturbation of x molecules removes x/2 Notch molecules from the first cell and adds x/2 Notch molecules to the second. Unsurprisingly, larger initial perturbations result in a higher proportion of primary outcomes, although this effect is not as extreme as we expected. Even in the largest perturbation we tested ($\pm 25\%$), only 80% of the simulations had the primary outcome. This suggests that it is difficult to reliably control cell fate by perturbation alone.

higher levels of noise lower the average differentiation time by rapidly perturbing the system into more "extreme" states. Using this relationship, we estimated that $\sigma=0.02$ by comparing simulations of the SDE model with the agent-based model. See Figure 3 for a summary of this result.

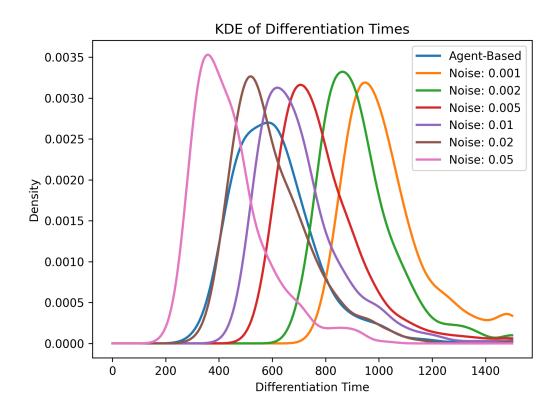


Figure 3: A plot of the KDE approximation of the distribution of differentiation times for the agent-based model and the stochastic ODE model with different levels of noise. All simulations were done on a two-cell domain with periodic boundary conditions. These simulations show that higher values of the noise parameter significantly decrease the average time it takes for the two cells to adopt different fates. Notably, a noise parameter of 2% causes the stochastic ODE model to align most closely with the results of agent-based simulations.

All of the previous experiments were run using a fixed set of parameter values. However, in order to evaluate the quality of our model, we wanted to see if we could replicate the key result of cell differentiation under different parameter regimes. We identified five main parameters which effect whether the two-cell system differentiates. These parameters are the maximum rates of Notch and Delta production n_m and d_m , the binding rate k_T , the Notch and Delta decay rate γ , and the NICD decay rate γ_I . Due to computational limitations, we could not perform a complete grid search of the 5-dimensional parameter space. Instead, we opted to explore all 10 of the 2-dimensional subspaces of our parameter space. The results of our stability analysis are shown in Figure 4.

Linear Domains

A linear domain is an array of cells in which each cell receives and emits Delta-Notch signals to 2 neighbours (1 if the cell is on the boundary of domain). Collier et al. [4] tested their model on linear domains of up to 40 cells with zero-flux boundary conditions. They

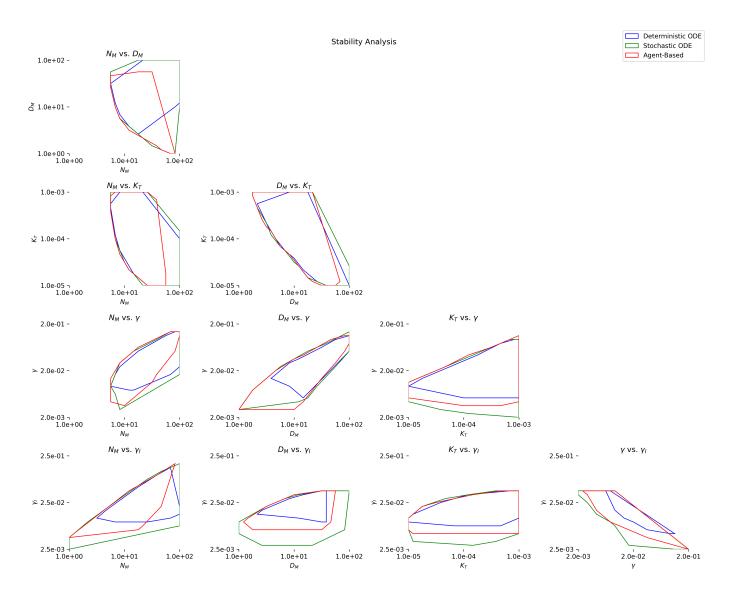


Figure 4: Results of two-at-a-time stability analysis of the deterministic, stochastic, and agent-based modelling formalisms on a two-cell domain with periodic boundary conditions. The regions enclosed by the red, blue, and green lines are regions in which cell differentiation occurs for each of the three models. Regions were obtained by simulating a 25×25 grid of test points and taking the convex hull of the points at which cell differentiation occurred. This analysis suggests that the deterministic and stochastic ODE models are on average less sensitive to changes in parameter values than the agent-based model. In particular, the agent-based model does not achieve cell differentiation for large values of the Notch and Delta production parameters n_m and d_m .

observed that cells on the boundary of the domain adopted the primary fate first. Then, the alternating pattern of primary and secondary fate cells gradually spread inwards towards the center of the domain. In our first experiment on a linear domain, we replicated this result as shown in Figure.

We also attempted to broadly classify the types of patterns we observed on linear domains. In order to do this, we first need to define the notion of a pattern. Since the "default" or "regular" outcome of a simulation on a linear domain is an alternating array of cells with each fate, we will call this pattern 0:[], 1:[]. Then, each contiguous group of cells with the same fate (where 0 is the primary fate and 1 is the secondary fate) will add entries to the pattern arrays. For example, if there are 2 contiguous pairs of cells with the primary fate and 1 contiguous triplet of cells with the secondary fate, the simulation will have the pattern 0:[2,2], 1:[3]. There is a lot more nuance related to pattern formation that we didn't have time to explore, but one interesting result is shown in Table 2.

Pattern	Frequency (Odd Domain)	Frequency (Even Domain)
0:[], 1:[]	87	403
0:[], 1:[2]	801	123
0:[], 1:[2,2]	27	424
0:[], 1:[2,2,2]	48	1

Table 2: Frequency of selected cell patterns in 1000 simulations of the agent-based model on 9-cell (even) and 10-cell (odd) domains with zero-flux boundary conditions. We observe significant variations in patterns across even and odd domains. Similarly drastic variation is observed across different modelling formalisms and boundary conditions (not shown).

Hexagonal Domains

Discussion

References

- [1] Francisco Berkemeier and Karen M. Page. "Coupling dynamics of 2D Notch-Delta signalling". In: *Mathematical Biosciences* 360 (June 1, 2023), p. 109012. ISSN: 0025-5564. DOI: 10.1016/j.mbs.2023.109012. URL: https://www.sciencedirect.com/science/article/pii/S0025556423000536 (visited on 03/27/2025).
- [2] Marcelo Boareto et al. "Jagged-Delta asymmetry in Notch signaling can give rise to a Sender/Receiver hybrid phenotype". In: *Proceedings of the National Academy of Sciences* 112.5 (Feb. 3, 2015). Publisher: Proceedings of the National Academy of Sciences, E402-E409. DOI: 10.1073/pnas.1416287112. URL: https://www.pnas.org/doi/full/10.1073/pnas.1416287112 (visited on 03/14/2025).
- [3] Sarah J. Bray. "Notch signalling: a simple pathway becomes complex". In: Nature Reviews Molecular Cell Biology 7.9 (Sept. 2006). Publisher: Nature Publishing Group, pp. 678-689. ISSN: 1471-0080. DOI: 10.1038/nrm2009. URL: https://www.nature.com/articles/nrm2009 (visited on 03/14/2025).
- [4] Joanne R. Collier et al. "Pattern Formation by Lateral Inhibition with Feedback: a Mathematical Model of Delta-Notch Intercellular Signalling". In: Journal of Theoretical Biology 183.4 (Dec. 21, 1996), pp. 429-446. ISSN: 0022-5193. DOI: 10.1006/jtbi.1996.0233. URL: https://www.sciencedirect.com/science/article/pii/S0022519396902337 (visited on 03/14/2025).

A Gillespie's Algorithm

Gillespie's algorithm allows us to simulate many types reactions simultaneously using random time steps. Recall that T_i is a random variable that represents the distribution of the next reaction of any type. The random variables X_1, \ldots, X_N are the waiting time distributions for each reaction $1, \ldots, N$, where $N = (6 \times \# \text{ Cells})$.

Step 1: Sample a value t_i from T using inverse transform sampling.

Step 2: Increment the current time t by t_i .

Step 3: Determine which event took place. To do this, recall that:

$$P(X_k = \min\{X_1, \dots, X_N\}) = \frac{\lambda_k}{\lambda_1 + \dots + \lambda_N}$$

Use this fact to partition the interval [0,1]. Then sample y_i from a Unif(0,1) distribution, identify its partition, and find the corresponding event.

Step 4: Update the state based on the event determined in the previous step.

Remark: The expected time for reaction i to occur is given by $\mathbb{E}[X_i] = 1/\lambda_i$. Therefore, the rate at which event i occurs is λ_i , which is equal to the rate at which the event occurs in the original ODE. This means that the trajectories generated from the Gillespie's algorithm implementation of the agent-based model can be compared directly to the trajectories generated from the ODE itself, without any need for scaling.

B Kolmogorov Forward Equations

Derivation of Kolmogorov Forward Equation

Let us define P(n, d, i, t) to be the probability of having n Notch receptors, d Delta ligands, and i NICD molecules at time t. The Kolmogorov forward equation describes how P(n, d, i, t) changes over a small time step Δt :

$$\begin{split} P(n,d,i,t+\Delta t) &= P(n,d,i,t) \\ &+ P(n-1,d,i,t) \cdot H^+(i) \Delta t \quad \text{(Notch production)} \\ &+ P(n,d-1,i,t) \cdot H^-(i) \Delta t \quad \text{(Delta production)} \\ &+ P(n+1,d,i,t) \cdot \gamma(n+1) \Delta t \quad \text{(Notch decay)} \\ &+ P(n,d+1,i,t) \cdot \gamma(d+1) \Delta t \quad \text{(Delta decay)} \\ &+ P(n,d,i+1,t) \cdot \gamma_I(i+1) \Delta t \quad \text{(NICD decay)} \\ &+ P(n+1,d,i-1,t) \cdot k_T(n+1) D_{ext} \Delta t \quad \text{(Trans-activation)} \\ &- P(n,d,i,t) \cdot [H^+(i) + H^-(i) + \gamma n + \gamma d + \gamma_I i + k_T n D_{ext}] \Delta t \end{split}$$

Rearranging to find the time derivative:

$$\frac{P(n,d,i,t+\Delta t) - P(n,d,i,t)}{\Delta t} = P(n-1,d,i,t) \cdot H^{+}(i)$$

$$+ P(n,d-1,i,t) \cdot H^{-}(i)$$

$$+ P(n+1,d,i,t) \cdot \gamma(n+1)$$

$$+ P(n,d+1,i,t) \cdot \gamma(d+1)$$

$$+ P(n,d,i+1,t) \cdot \gamma_{I}(i+1)$$

$$+ P(n+1,d,i-1,t) \cdot k_{T}(n+1)D_{ext}$$

$$- P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext}]$$

Taking the limit as $\Delta t \to 0$:

$$\begin{split} \frac{\partial P(n,d,i,t)}{\partial t} &= P(n-1,d,i,t) \cdot H^{+}(i) \\ &+ P(n,d-1,i,t) \cdot H^{-}(i) \\ &+ P(n+1,d,i,t) \cdot \gamma(n+1) \\ &+ P(n,d+1,i,t) \cdot \gamma(d+1) \\ &+ P(n,d,i+1,t) \cdot \gamma_{I}(i+1) \\ &+ P(n+1,d,i-1,t) \cdot k_{T}(n+1) D_{ext} \\ &- P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I} i + k_{T} n D_{ext}] \end{split}$$

Matrix Formulation

The Kolmogorov forward equation can be expressed in matrix form as:

$$\frac{d\vec{P}(t)}{dt} = \mathbf{A}\vec{P}(t)$$

where $\vec{P}(t)$ is a column vector containing the probabilities for all possible states, and **A** is the transition rate matrix.

To construct the matrix **A**, we first establish a one-to-one mapping between the three-dimensional state space (n, d, i) and a one-dimensional index k. For a system with maximum values n_{max} , d_{max} , and i_{max} , we define:

$$k(n, d, i) = n \cdot (d_{max} + 1) \cdot (i_{max} + 1) + d \cdot (i_{max} + 1) + i$$

for $0 \le n \le n_{max}$, $0 \le d \le d_{max}$, $0 \le i \le i_{max}$.

The inverse mapping gives:

$$n(k) = \left\lfloor \frac{k}{(d_{max} + 1) \cdot (i_{max} + 1)} \right\rfloor$$
$$d(k) = \left\lfloor \frac{k \mod ((d_{max} + 1) \cdot (i_{max} + 1))}{i_{max} + 1} \right\rfloor$$
$$i(k) = k \mod (i_{max} + 1)$$

With this mapping, we can now define the elements of matrix **A**. Let k and k' be the indices corresponding to states (n, d, i) and (n', d', i') respectively:

$$A_{k',k} = \begin{cases} H^{+}(i) & \text{if } (n',d',i') = (n+1,d,i) & \text{(Notch production)} \\ H^{-}(i) & \text{if } (n',d',i') = (n,d+1,i) & \text{(Delta production)} \\ \gamma(n) & \text{if } (n',d',i') = (n-1,d,i) & \text{(Notch decay)} \\ \gamma(d) & \text{if } (n',d',i') = (n,d-1,i) & \text{(Delta decay)} \\ \gamma_{I}(i) & \text{if } (n',d',i') = (n,d,i-1) & \text{(NICD decay)} \\ k_{T}(n)D_{ext} & \text{if } (n',d',i') = (n-1,d,i+1) & \text{(Trans-activation)} \\ -\sum_{k''\neq k} A_{k'',k} & \text{if } k' = k & \text{(Diagonal terms)} \end{cases}$$

Note that $A_{k',k}$ represents the transition rate from state k to state k'. The diagonal elements $A_{k,k}$ are the negative sum of all outgoing rates from state k:

$$A_{k,k} = -[H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext} + k_{T}dN_{ext}]$$

where (n, d, i) is the state corresponding to index k.

The resulting matrix \mathbf{A} is a sparse matrix with the following properties:

- Dimension: $(n_{max} + 1) \cdot (d_{max} + 1) \cdot (i_{max} + 1) \times (n_{max} + 1) \cdot (d_{max} + 1) \cdot (i_{max} + 1)$
- Number of non-zero elements: $\approx 7 \cdot (n_{max} + 1) \cdot (d_{max} + 1) \cdot (i_{max} + 1)$

To illustrate, let's consider a minimal system with $n_{max} = d_{max} = i_{max} = 1$. This gives us 8 possible states: (0,0,0), (0,0,1), (0,1,0), (0,1,1), (1,0,0), (1,0,1), (1,1,0), (1,1,1) The corresponding transition matrix **A** would be:

$$\mathbf{A} = \begin{pmatrix} -\alpha_{000} & \gamma_I & \gamma & 0 & \gamma & 0 & 0 & 0 \\ H^+(0) & -\alpha_{001} & 0 & \gamma & 0 & \gamma & 0 & 0 \\ H^-(0) & 0 & -\alpha_{010} & \gamma_I & 0 & 0 & \gamma & 0 \\ 0 & H^-(0) & H^+(0) & -\alpha_{011} & 0 & 0 & 0 & \gamma \\ H^+(0) & 0 & 0 & 0 & -\alpha_{100} & \gamma_I & \gamma & 0 \\ 0 & H^+(0) & 0 & 0 & H^+(1) & -\alpha_{101} & 0 & \gamma \\ 0 & 0 & H^+(0) & 0 & H^-(1) & 0 & -\alpha_{110} & \gamma_I \\ 0 & 0 & 0 & H^+(0) & 0 & H^-(1) & H^+(1) & -\alpha_{111} \end{pmatrix}$$

where:

$$\alpha_{ndi} = H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext} + k_{T}dN_{ext}$$

Additionally, we need to add the terms for trans-activation, which would connect states like (1,0,0) to (0,0,1) with rate $k_T \cdot 1 \cdot D_{ext}$.

For a fully labelled matrix with all transitions explicitly shown, the matrix for $n_{max} = d_{max} = i_{max} = 1$ would be:

$$\mathbf{A} = \begin{bmatrix} A_{11} & A_{12} & A_{13} & A_{14} & A_{15} & A_{16} & A_{17} & A_{18} \\ B_{11} & B_{12} & B_{13} & B_{14} & B_{15} & B_{16} & B_{17} & B_{18} \\ C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} & C_{17} & C_{18} \\ D_{11} & D_{12} & D_{13} & D_{14} & D_{15} & D_{16} & D_{17} & D_{18} \\ E_{11} & E_{12} & E_{13} & E_{14} & E_{15} & E_{16} & E_{17} & E_{18} \\ F_{11} & F_{12} & F_{13} & F_{14} & F_{15} & F_{16} & F_{17} & F_{18} \\ G_{11} & G_{12} & G_{13} & G_{14} & G_{15} & G_{16} & G_{17} & G_{18} \\ H_{11} & H_{12} & H_{13} & H_{14} & H_{15} & H_{16} & H_{17} & H_{18} \end{bmatrix}$$

where

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A_{11} = -(H^{+}(0) + H^{-}(0)),
                                                                                  A_{12}=\gamma_I,
A_{13} = \gamma, A_{15} = \gamma,
                                                                                  A_{14}=0,
                                                                                  A_{16} = 0,
A_{17} = 0,
                                                                                  A_{18} = 0,
B_{11} = H^{+}(0),
                                                                                  B_{12} = -(H^{+}(1) + H^{-}(1) + \gamma_{I}),
                                                                                  B_{14}=\gamma,
B_{13}=0,
                                                                                  B_{16} = \gamma,
B_{18} = 0,
B_{15} = k_T D_{ext},
B_{17} = 0,
C_{11} = H^{-}(0),
                                                                                  C_{12} = 0,
C_{13} = -(H^{+}(0) + H^{-}(0) + \gamma + k_T N_{ext}),
                                                                                  C_{14} = \gamma_I
C_{15} = 0,
                                                                                  C_{16} = 0,
C_{17} = \gamma
                                                                                  C_{18} = 0,
D_{11} = 0,
                                                                                  D_{12} = H^{-}(1),
                                                                                  D_{14} = -(H^{+}(1) + H^{-}(1) + \gamma + \gamma_{I} + k_{T}N_{ext}),
D_{13} = H^{+}(0),
D_{15} = 0,
                                                                                  D_{16} = k_T D_{ext},
D_{17} = 0,
                                                                                  D_{18}=\gamma,
E_{11} = H^{+}(0),
                                                                                  E_{12} = 0,
E_{13} = 0,
                                                                                  E_{14}=0,
E_{15} = -(H^{+}(0) + H^{-}(0) + \gamma + k_T D_{ext}),
                                                                                  E_{16} = \gamma_{I},
                                                                                  E_{18} = 0,
F_{11} = 0,
                                                                                  F_{12} = H^+(1),
F_{13} = 0,
                                                                                   F_{14} = 0,
F_{15} = H^{+}(0),
                                                                                   F_{16} = -(H^{+}(1) + H^{-}(1) + \gamma + \gamma_{I} + k_{T}D_{ext}),
F_{17} = 0,
G_{11} = 0,
                                                                                  G_{12} = 0,
G_{13} = H^{+}(0),
                                                                                  G_{14} = 0,
G_{15} = H^{-}(0),
                                                                                  G_{16} = 0,
G_{17} = -(H^{+}(0) + H^{-}(0) + \gamma + \gamma + k_T D_{ext} + k_T N_{ext}),
                                                                                 G_{18} = \gamma_I
H_{11} = 0,
                                                                                  H_{12} = 0,
                                                                                  H_{14} = H^+(1),
H_{13} = 0,
                                                                                  H_{16} = H^{-}(1),
H_{15} = 0,
H_{17} = H^{+}(0),
                                                                                  H_{18} = -(H^{+}(1) + H^{-}(1) + \gamma + \gamma + \gamma_{I} + k_{T}D_{ext} + k_{T}N_{ext}).
```

Each row and column corresponds to a state in the order: (0,0,0), (0,0,1), (0,1,0), (0,1,1), (1,0,0), (1,0,1), (1,1,0), (1,1,1)

The matrix element $A_{j,i}$ represents the transition rate from state i to state j. The diagonal elements $A_{i,i}$ contain the negative sum of all outgoing rates from state i.

C Theoretical Justification

In this section, we need to show that in expectation, the model above converges to the following system of ODEs (adapted from Boareto et al., 2015):

$$\begin{split} \frac{dN}{dt} &= \frac{n_m I^2}{n_0^2 + I^2} - k_T N D_{ext} - \gamma N \\ \frac{dD}{dt} &= \frac{d_m d_0^2}{d_0^2 + I^2} - k_T D N_{ext} - \gamma D \\ \frac{dI}{dt} &= k_T N D_{ext} - \gamma_I I \end{split}$$

Setup and Notation

Let us define:

- $\langle N \rangle = \sum_{n.d.i} n \cdot P(n,d,i,t) =$ expected number of Notch receptors
- $\langle D \rangle = \sum_{n.d.i} d \cdot P(n,d,i,t) =$ expected number of Delta ligands
- $\langle I \rangle = \sum_{n,d,i} i \cdot P(n,d,i,t) =$ expected number of NICD molecules

Deriving ODEs for Expected Values

Now we derive the ODE for the expected value of Notch:

$$\frac{d\langle N \rangle}{dt} = \frac{d}{dt} \sum_{n,d,i} n \cdot P(n,d,i,t)$$
$$= \sum_{n,d,i} n \cdot \frac{\partial P(n,d,i,t)}{\partial t}$$

Substituting the Kolmogorov forward equation:

$$\frac{d\langle N \rangle}{dt} = \sum_{n,d,i} n \cdot \left[P(n-1,d,i,t) \cdot H^{+}(i) + P(n,d-1,i,t) \cdot H^{-}(i) + P(n,d-1,i,t) \cdot \gamma(n+1) + P(n,d+1,i,t) \cdot \gamma(d+1) + P(n,d,i+1,t) \cdot \gamma(d+1) + P(n,d,i+1,t) \cdot \gamma_{I}(i+1) + P(n+1,d,i-1,t) \cdot k_{T}(n+1) D_{ext} - P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}n D_{ext}] \right]$$

Now we use index shifting to simplify.

(A) Notch Production

Let's substitute n' = n - 1 (so n = n' + 1):

$$\sum_{n,d,i} n \cdot P(n-1,d,i,t) \cdot H^{+}(i) = \sum_{n',d,i} (n'+1) \cdot P(n',d,i,t) \cdot H^{+}(i)$$

$$= \sum_{n',d,i} n' \cdot P(n',d,i,t) \cdot H^{+}(i) + \sum_{n',d,i} P(n',d,i,t) \cdot H^{+}(i)$$

$$= \sum_{n,d,i} n \cdot P(n,d,i,t) \cdot H^{+}(i) + \sum_{n,d,i} P(n,d,i,t) \cdot H^{+}(i)$$

The first term is $\langle N \cdot H^+(I) \rangle$. If we assume that $H^+(I)$ is approximately $H^+(\langle I \rangle)$, then this term becomes $\langle N \rangle \cdot H^+(\langle I \rangle)$. The second term is $\langle H^+(I) \rangle \approx H^+(\langle I \rangle)$. So the contribution of this term is:

$$\sum_{n,d,i} n \cdot P(n-1,d,i,t) \cdot H^{+}(i) \approx \langle N \rangle \cdot H^{+}(\langle I \rangle) + H^{+}(\langle I \rangle)$$

Since we are indexing from 0, and n-1 could be negative, the sum should start from n=1:

$$\sum_{n\geq 1,d,i} n \cdot P(n-1,d,i,t) \cdot H^{+}(i) = \sum_{n'\geq 0,d,i} (n'+1) \cdot P(n',d,i,t) \cdot H^{+}(i)$$

$$= \sum_{n'\geq 0,d,i} n' \cdot P(n',d,i,t) \cdot H^{+}(i) + \sum_{n'\geq 0,d,i} P(n',d,i,t) \cdot H^{+}(i)$$

$$= \langle N \rangle \cdot H^{+}(\langle I \rangle) + H^{+}(\langle I \rangle)$$

$$= H^{+}(\langle I \rangle) \cdot (1 + \langle N \rangle)$$

(B) Delta Production

This term does not affect N, so the prefactor n remains with the original probability:

$$\begin{split} \sum_{n,d,i} n \cdot P(n,d-1,i,t) \cdot H^-(i) &= \sum_{n,d',i} n \cdot P(n,d',i,t) \cdot H^-(i) \\ &= \langle N \cdot H^-(I) \rangle \\ &\approx \langle N \rangle \cdot H^-(\langle I \rangle) \end{split}$$

(C) Notch Decay

$$\begin{split} \sum_{n,d,i} n \cdot P(n+1,d,i,t) \cdot \gamma(n+1) &= \sum_{n',d,i} (n'-1) \cdot P(n',d,i,t) \cdot \gamma n' \\ &= \gamma \sum_{n',d,i} (n'-1)n' \cdot P(n',d,i,t) \\ &= \gamma \sum_{n',d,i} n'^2 \cdot P(n',d,i,t) - \gamma \sum_{n',d,i} n' \cdot P(n',d,i,t) \\ &= \gamma (\langle N^2 \rangle - \langle N \rangle) \end{split}$$

If we assume that the variance of N is small compared to its mean (which is reasonable for high molecule counts), then $\langle N^2 \rangle \approx \langle N \rangle^2$. So:

$$\gamma(\langle N^2 \rangle - \langle N \rangle) \approx \gamma(\langle N \rangle^2 - \langle N \rangle)$$
$$= \gamma\langle N \rangle (\langle N \rangle - 1)$$

(D) Delta Decay

$$\begin{split} \sum_{n,d,i} n \cdot P(n,d+1,i,t) \cdot \gamma(d+1) &= \gamma \sum_{n,d',i} n \cdot P(n,d',i,t) \cdot d' \\ &= \gamma \langle N \cdot D \rangle \\ &\approx \gamma \langle N \rangle \langle D \rangle \end{split}$$

(E) NICD Decay

$$\sum_{n,d,i} n \cdot P(n,d,i+1,t) \cdot \gamma_I(i+1) = \gamma_I \sum_{n,d,i'} n \cdot P(n,d,i',t) \cdot i'$$
$$= \gamma_I \langle N \cdot I \rangle$$
$$\approx \gamma_I \langle N \rangle \langle I \rangle$$

(F) Trans-activation

$$\sum_{n,d,i} n \cdot P(n+1,d,i-1,t) \cdot k_T(n+1) D_{ext} = k_T D_{ext} \sum_{n',d,i'} (n'-1) \cdot P(n',d,i',t) \cdot n'$$

$$= k_T D_{ext} \sum_{n',d,i'} (n'^2 - n') \cdot P(n',d,i',t)$$

$$= k_T D_{ext} (\langle N^2 \rangle - \langle N \rangle)$$

$$\approx k_T D_{ext} \langle N \rangle (\langle N \rangle - 1)$$

(G) Negative Terms

$$-\sum_{n,d,i} n \cdot P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext}]$$

$$= -\langle N \cdot H^{+}(I) \rangle - \langle N \cdot H^{-}(I) \rangle - \gamma \langle N^{2} \rangle - \gamma \langle N \cdot D \rangle - \gamma_{I} \langle N \cdot I \rangle - k_{T}D_{ext} \langle N^{2} \rangle$$

$$\approx -\langle N \rangle \cdot H^{+}(\langle I \rangle) - \langle N \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle N \rangle^{2} - \gamma \langle N \rangle \langle D \rangle - \gamma_{I} \langle N \rangle \langle I \rangle - k_{T}D_{ext} \langle N \rangle^{2}$$

(H) Final Synthesis of $\langle N \rangle$

Combining all these terms:

$$\frac{d\langle N \rangle}{dt} = H^{+}(\langle I \rangle) \cdot (1 + \langle N \rangle) + \langle N \rangle \cdot H^{-}(\langle I \rangle) + \gamma \langle N \rangle (\langle N \rangle - 1) + \gamma \langle N \rangle \langle D \rangle
+ \gamma_{I} \langle N \rangle \langle I \rangle + k_{T} D_{ext} \langle N \rangle (\langle N \rangle - 1)
- \langle N \rangle \cdot H^{+}(\langle I \rangle) - \langle N \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle N \rangle^{2} - \gamma \langle N \rangle \langle D \rangle - \gamma_{I} \langle N \rangle \langle I \rangle - k_{T} D_{ext} \langle N \rangle^{2}$$

After cancellations:

$$\frac{d\langle N \rangle}{dt} = H^{+}(\langle I \rangle) + \langle N \rangle \cdot H^{+}(\langle I \rangle) + \gamma \langle N \rangle (\langle N \rangle - 1) + k_{T} D_{ext} \langle N \rangle (\langle N \rangle - 1)
- \langle N \rangle \cdot H^{+}(\langle I \rangle) - \gamma \langle N \rangle^{2} - k_{T} D_{ext} \langle N \rangle^{2}
= H^{+}(\langle I \rangle) + \gamma \langle N \rangle (\langle N \rangle - 1) - \gamma \langle N \rangle^{2} + k_{T} D_{ext} \langle N \rangle (\langle N \rangle - 1) - k_{T} D_{ext} \langle N \rangle^{2}
= H^{+}(\langle I \rangle) + \gamma \langle N \rangle (\langle N \rangle - 1 - \langle N \rangle) + k_{T} D_{ext} \langle N \rangle (\langle N \rangle - 1 - \langle N \rangle)
= H^{+}(\langle I \rangle) - \gamma \langle N \rangle - k_{T} D_{ext} \langle N \rangle$$

Therefore, the ODE for $\langle N \rangle$ is:

$$\frac{d\langle N\rangle}{dt} = H^{+}(\langle I\rangle) - \gamma\langle N\rangle - k_T D_{ext}\langle N\rangle$$

Substituting the definition of $H^+(\langle I \rangle)$:

$$\frac{d\langle N\rangle}{dt} = \frac{n_m \langle I\rangle^2}{n_0^2 + \langle I\rangle^2} - \gamma \langle N\rangle - k_T D_{ext} \langle N\rangle$$

Which can be written as:

$$\frac{dN}{dt} = \frac{n_m I^2}{n_0^2 + I^2} - \gamma N - k_T D_{ext} N$$

where we use N, D, and I as shorthand for the expected values $\langle N \rangle$, $\langle D \rangle$, and $\langle I \rangle$.

Similarly, we can derive the ODE for the expected value of Delta:

$$\frac{d\langle D\rangle}{dt} = \frac{d}{dt} \sum_{n,d,i} d \cdot P(n,d,i,t)$$
$$= \sum_{n,d,i} d \cdot \frac{\partial P(n,d,i,t)}{\partial t}$$

We will analyse each term of the Kolmogorov equation, multiplied by d, as follows.

$$\sum_{n,d,i} d \cdot P(n-1,d,i,t) \cdot H^{+}(i) = \langle D \cdot H^{+}(I) \rangle \approx \langle D \rangle \cdot H^{+}(\langle I \rangle)$$

$$\sum_{n,d,i} d \cdot P(n,d-1,i,t) \cdot H^{-}(i) = \langle D \cdot H^{-}(I) \rangle + \langle H^{-}(I) \rangle \approx \langle D \rangle \cdot H^{-}(\langle I \rangle) + H^{-}(\langle I \rangle)$$

$$\sum_{n,d,i} d \cdot P(n+1,d,i,t) \cdot \gamma(n+1) \approx \gamma \langle D \rangle \langle N \rangle$$

$$\sum_{n,d,i} d \cdot P(n,d+1,i,t) \cdot \gamma(d+1) \approx \gamma \langle D \rangle (\langle D \rangle - 1)$$

$$\sum_{n,d,i} d \cdot P(n,d,i+1,t) \cdot \gamma_{I}(i+1) \approx \gamma_{I} \langle D \rangle \langle I \rangle$$

$$\sum_{n,d,i} d \cdot P(n+1,d,i-1,t) \cdot k_{T}(n+1) D_{ext} \approx k_{T} D_{ext} \langle D \rangle \langle N \rangle$$

Now we consider this term which appears in the ODEs but was not explicitly stated in our Kolmogorov equation setup. This is a special case where we need to account for Delta from our cell binding to Notch in neighbouring cells:

$$-\sum_{n,d,i} d \cdot P(n,d,i,t) \cdot k_T dN_{ext} = -k_T N_{ext} \sum_{n,d,i} d^2 \cdot P(n,d,i,t)$$
$$= -k_T N_{ext} \langle D^2 \rangle$$
$$\approx -k_T N_{ext} \langle D \rangle^2$$

Now we consider negative terms from original equation.

$$-\sum_{n,d,i} d \cdot P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext}]$$

$$\approx -\langle D \rangle \cdot H^{+}(\langle I \rangle) - \langle D \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle D \rangle \langle N \rangle - \gamma \langle D \rangle^{2} - \gamma_{I}\langle D \rangle \langle I \rangle - k_{T}D_{ext}\langle D \rangle \langle N \rangle$$

Let's synthesise all terms of $\langle D \rangle$

$$\frac{d\langle D \rangle}{dt} = \langle D \rangle \cdot H^{+}(\langle I \rangle) + \langle D \rangle \cdot H^{-}(\langle I \rangle) + H^{-}(\langle I \rangle) + \gamma \langle D \rangle \langle N \rangle
+ \gamma \langle D \rangle (\langle D \rangle - 1) + \gamma_{I} \langle D \rangle \langle I \rangle + k_{T} D_{ext} \langle D \rangle \langle N \rangle - k_{T} N_{ext} \langle D \rangle^{2}
- \langle D \rangle \cdot H^{+}(\langle I \rangle) - \langle D \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle D \rangle \langle N \rangle - \gamma \langle D \rangle^{2} - \gamma_{I} \langle D \rangle \langle I \rangle - k_{T} D_{ext} \langle D \rangle \langle N \rangle
= H^{-}(\langle I \rangle) - \gamma \langle D \rangle - k_{T} N_{ext} \langle D \rangle^{2}$$

Making a final approximation that $\langle D \rangle^2 \approx \langle D \rangle$, we get:

$$\frac{d\langle D\rangle}{dt} = H^{-}(\langle I\rangle) - \gamma\langle D\rangle - k_T N_{ext}\langle D\rangle$$

Substituting the definition of $H^-(\langle I \rangle)$:

$$\frac{d\langle D\rangle}{dt} = \frac{d_m}{d_0^2 + \langle I\rangle^2} - \gamma \langle D\rangle - k_T N_{ext} \langle D\rangle$$

Which can be written as:

$$\frac{dD}{dt} = \frac{d_m}{d_0^2 + I^2} - \gamma D - k_T N_{ext} D$$

Finally, we derive the ODE for NICD:

$$\frac{d\langle I \rangle}{dt} = \frac{d}{dt} \sum_{n,d,i} i \cdot P(n,d,i,t)$$
$$= \sum_{n,d,i} i \cdot \frac{\partial P(n,d,i,t)}{\partial t}$$

Let's analyse each term:

$$\sum_{n,d,i} i \cdot P(n-1,d,i,t) \cdot H^{+}(i) = \langle I \cdot H^{+}(I) \rangle \approx \langle I \rangle \cdot H^{+}(\langle I \rangle)$$

$$\sum_{n,d,i} i \cdot P(n,d-1,i,t) \cdot H^{-}(i) \approx \langle I \rangle \cdot H^{-}(\langle I \rangle)$$

$$\sum_{n,d,i} i \cdot P(n+1,d,i,t) \cdot \gamma(n+1) \approx \gamma \langle I \rangle \langle N \rangle$$

$$\sum_{n,d,i} i \cdot P(n,d+1,i,t) \cdot \gamma(d+1) \approx \gamma \langle I \rangle \langle D \rangle$$

$$\sum_{n,d,i} i \cdot P(n,d,i+1,t) \cdot \gamma_{I}(i+1) \approx \gamma_{I} \langle I \rangle \langle \langle I \rangle - 1)$$

$$\sum_{n,d,i} i \cdot P(n+1,d,i-1,t) \cdot k_{T}(n+1) D_{ext} \approx k_{T} D_{ext} \langle I \rangle \langle N \rangle + k_{T} D_{ext} \langle N \rangle$$

Negative terms can be rewritten as:

$$-\sum_{n,d,i} i \cdot P(n,d,i,t) \cdot [H^{+}(i) + H^{-}(i) + \gamma n + \gamma d + \gamma_{I}i + k_{T}nD_{ext}]$$

$$\approx -\langle I \rangle \cdot H^{+}(\langle I \rangle) - \langle I \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle I \rangle \langle N \rangle - \gamma \langle I \rangle \langle D \rangle - \gamma_{I} \langle I \rangle^{2} - k_{T}D_{ext} \langle I \rangle \langle N \rangle$$

Let's combine all terms for $\langle I \rangle$

$$\frac{d\langle I \rangle}{dt} = \langle I \rangle \cdot H^{+}(\langle I \rangle) + \langle I \rangle \cdot H^{-}(\langle I \rangle) + \gamma \langle I \rangle \langle N \rangle + \gamma \langle I \rangle \langle D \rangle
+ \gamma_{I} \langle I \rangle (\langle I \rangle - 1) + k_{T} D_{ext} \langle I \rangle \langle N \rangle + k_{T} D_{ext} \langle N \rangle
- \langle I \rangle \cdot H^{+}(\langle I \rangle) - \langle I \rangle \cdot H^{-}(\langle I \rangle) - \gamma \langle I \rangle \langle N \rangle - \gamma \langle I \rangle \langle D \rangle - \gamma_{I} \langle I \rangle^{2} - k_{T} D_{ext} \langle I \rangle \langle N \rangle$$

Therefore, the ODE for $\langle I \rangle$ is:

$$\frac{d\langle I\rangle}{dt} = k_T D_{ext} \langle N\rangle - \gamma_I \langle I\rangle$$

Which can be written as:

$$\frac{dI}{dt} = k_T D_{ext} N - \gamma_I I$$

Final System of ODEs

Combining our three derived equations, we get the following system of ODEs:

$$\begin{split} \frac{dN}{dt} &= \frac{n_m I^2}{n_0^2 + I^2} - \gamma N - k_T D_{ext} N \\ \frac{dD}{dt} &= \frac{d_m}{d_0^2 + I^2} - \gamma D - k_T N_{ext} D \\ \frac{dI}{dt} &= k_T D_{ext} N - \gamma_I I \end{split}$$

This matches the system presented in section 3.6 of the document:

$$\frac{dN}{dt} = \frac{n_m I^2}{n_0^2 + I^2} - k_T N D_{ext} - \gamma N$$

$$\frac{dD}{dt} = \frac{d_m}{d_0^2 + I^2} - k_T D N_{ext} - \gamma D$$

$$\frac{dI}{dt} = k_T N D_{ext} - \gamma_I I$$

Mapping Between Probabilities and Rates

To relate the probabilistic formulation to the deterministic ODE system, we need to understand how reaction probabilities map to reaction rates. In the stochastic formulation, the probability of a reaction occurring in a small time interval Δt is:

$$P(\text{reaction}) \approx \text{rate} \times \Delta t$$

For example:

$$P(\emptyset \to N_A) = H^+(I)\Delta t = \frac{n_m I^2}{n_0^2 + I^2} \Delta t$$

In the limit as $\Delta t \to 0$, these probabilities per unit time become rates, which appear directly in the deterministic ODEs. The ODE system essentially tracks the expected values of these stochastic processes under the assumption that fluctuations are small compared to the means.

D Bifurcation Analysis

One-Cell System

First, we consider the one-cell system under the influence of a second, static cell. In this model, the Notch and Delta concentrations in the second cell $(N_{ext}$ and $D_{ext})$ are treated as parameters. Equations 3, 4, and 5 define this system.

$$\frac{dN}{dt} = \frac{n_m I^2}{n_0^2 + I^2} - (k_T D_{ext} + \gamma)N \tag{3}$$

$$\frac{dD}{dt} = \frac{d_m d_0^2}{d_0^2 + I^2} - (k_T N_{ext} + \gamma)D \tag{4}$$

$$\frac{dI}{dt} = k_T N D_{ext} - \gamma_I I \tag{5}$$

Observe that dN/dt and dI/dt do not depend on D. Therefore, we can analyze the N-I phase plane to learn more about the system. First, we find the nullclines of Equations 3 and 5 as functions of I. The results of this calculation are shown in Equations 6 and 7.

$$f_N(I) = \frac{1}{\gamma + k_T D_{ext}} \left(\frac{n_m I^2}{n_0^2 + I^2} \right)$$
 (6)

$$f_I(I) = \frac{\gamma_I I}{k_T D_{ext}} \tag{7}$$

Plotting the nullclines in Equations 6 and 7 on the N-I plane for small and large values of D_{ext} (assuming the other parameters are suitably chosen) reveals a fold bifurcation. This is shown in Figure D.

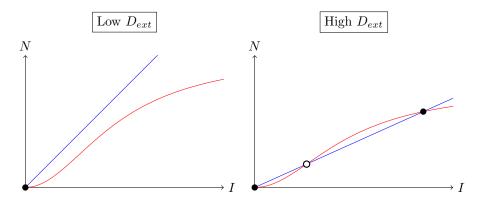


Figure 5: Approximate N-I phase planes for the one-cell ODE model under the influence of low and high Delta from a static neighbour. The N-nullcline is shown in red and the I-nullcline is shown in blue. Below a critical value of D_{ext} , the system undergoes a fold bifurcation which results in the annihilation of the high Notch steady state. The high Notch steady state is also absent for qualitatively high values of γ_I , low values of k_T , and low values of n_m .

Two-Cell System

Now, we will consider the coupled two-cell system in which both cells respond to the Delta-Notch signalling of the other. Let N_1, D_1, I_1 and N_2, D_2, I_2 denote the Notch, Delta, and NICD concentrations in the first and second cells respectively. Our previous analysis reveals that the $N_1 = I_1 = 0$ steady state exists for all values of D_2 . The steady state concentrations $\{N_1^*, D_1^*, I_1^*, N_2^*, D_2^*, I_2^*\}$ for this case are given in Equation 8.

$$N_{1}^{*} = 0 N_{2}^{*} = \frac{1}{\gamma + k_{T}D_{1}^{*}} \left(\frac{n_{m}(I_{2}^{*})^{2}}{n_{0}^{2} + (I_{2}^{*})^{2}} \right)$$

$$D_{1}^{*} = \frac{d_{m}}{\gamma + k_{T}N_{2}^{*}} D_{2}^{*} = \frac{1}{\gamma} \left(\frac{d_{m}d_{0}^{2}}{d_{0}^{2} + (I_{2}^{*})^{2}} \right)$$

$$I_{1}^{*} = 0 I_{2}^{*} = \frac{k_{T}N_{2}^{*}D_{1}^{*}}{\gamma_{I}}$$

$$(8)$$

There is a second steady state occurs when $N_1^* = N_2^*$, $D_1^* = D_2^*$, and $I_1^* = I_2^*$. While it is difficult to analytically determine the steady state concentrations, we can verify that this steady state exists and is unstable using computer simulations.