Bubble Nucleation

Vacuum Decay in the Early Universe

Mario L. Gutierrez Abed

School of Mathematical Sciences Rochester Institute of Technology

December 1st, 2020

Overview

- 1 Introduction
- 2 The Model
- 3 Numerical Methods
- 4 Results

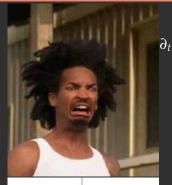


Introduction



MGA (SMS)

Equations... Equations anyone?





$$\partial_{t}\bar{A}_{ij} = \chi[\alpha(R_{ij} + KK_{ij} - 2K_{ik}K^{k}_{j})$$

$$-8\pi\alpha(S_{ij} - \frac{1}{2}\underbrace{\gamma_{ij}}(S - \rho)) - D_{i}D_{j}\alpha$$

$$+ \mathcal{L}_{\vec{\beta}}K_{ij}]^{TF} + \chi^{-1}\bar{A}_{ij}\left[\frac{2}{3}\chi(\alpha K - \partial_{i}\beta^{i})\right]$$

$$+ \beta^{i}\partial_{i}\chi\right]$$

$$= \left[\alpha\chi R_{ij} + \alpha\chi\chi^{-1}K\left(\bar{A}_{ij} + \frac{1}{3}\bar{\gamma}_{ij}K\right)\right]$$

$$-2\alpha\chi\chi^{-1}\left(\bar{A}_{ik} + \frac{1}{3}\bar{\gamma}_{ik}K\right)\left(\bar{A}^{k}_{j} + \delta^{k}_{j}K\right)...$$

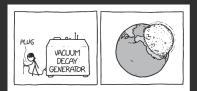
$$... - 8\pi\alpha\chi S_{ij} - \chi D_{i}D_{j}\alpha\right]^{TF} + ...$$

Dramatic much?

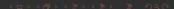




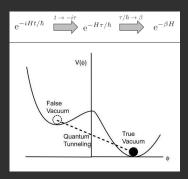




The Model

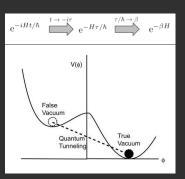


Semiclassical Vacuum Decay



$$\Gamma/V = Ae^{-B/\hbar}[1 + \mathcal{O}(\hbar)]$$

Semiclassical Vacuum Decay

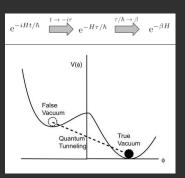


$$\Gamma/V = Ae^{-B/\hbar}[1 + \mathcal{O}(\hbar)]$$

$$B = S_E = \int \left[\frac{1}{2}\mathring{\nabla}^a \phi \mathring{\nabla}_a \phi + V(\phi)\right] d\tau d\vec{x}$$

$$V(\phi) = \frac{1}{2}\lambda^2 \sin^2 \phi - \cos \phi - 1$$

Semiclassical Vacuum Decay

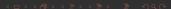


$$\Gamma/V = Ae^{-B/\hbar}[1 + \mathcal{O}(\hbar)]$$

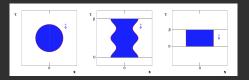
$$B = S_E = \int \left[\frac{1}{2}\mathring{\nabla}^a \phi \mathring{\nabla}_a \phi + V(\phi)\right] d\tau d\vec{x}$$

$$V(\phi) = \frac{1}{2}\lambda^2 \sin^2 \phi - \cos \phi - 1$$

Numerical Methods

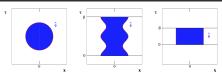


Good Ole shooting methods?



Shooting method very effective under O(D) assumption

Good Ole shooting methods?

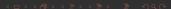




- Shooting method very effective under O(D) assumption
- ...but utterly useless otherwise ...

Introduce auxiliary scalar $\Phi(s,\tau,x)$ (so that $\lim_{s \to \infty} \Phi = \phi_b$) and set

$$\mathscr{F} \equiv \mathring{\nabla}^2 \Phi - \partial_{\Phi} V$$



MGA (SMS) Bubble Nucleation Dec

Introduce auxiliary scalar $\Phi(s,\tau,x)$ (so that $\lim_{s\to\infty}\Phi=\phi_b$) and set

$$\mathscr{F} \equiv \mathring{\nabla}^2 \Phi - \partial_{\Phi} V$$

First attempt:

$$\frac{\mathrm{d}\Phi}{\mathrm{d}s} = \mathcal{O}\mathscr{F}$$

Introduce auxiliary scalar $\Phi(s, au,x)$ (so that $\lim_{s o\infty}\Phi=\phi_b$) and set

$$\mathscr{F} \equiv \mathring{\nabla}^2 \Phi - \partial_{\Phi} V$$

First attempt:

$$\begin{split} \frac{\mathrm{d}\Phi}{\mathrm{d}s} &= \mathscr{F} \\ \Phi_{i,j}^{n+1} &= \varsigma \left[\Phi_{i+1,j}^n + \Phi_{i-1,j}^n + \Phi_{i,j+1}^n + \Phi_{i,j-1}^n \right] + \Phi_{i,j}^n \left[1 - 4\varsigma \right] \\ &- \Delta s \left[\frac{\lambda^2}{2} \sin \left(2\Phi_{i,j}^n \right) + \sin \Phi_{i,j}^n \right] \\ \varsigma &\equiv \frac{\Delta s}{\mathrm{h}^2} \end{split}$$

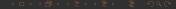
Introduce auxiliary scalar $\Phi(s, \tau, x)$ (so that $\lim_{s\to\infty} \Phi = \phi_b$) and set

$$\mathscr{F} \equiv \mathring{\nabla}^2 \Phi - \partial_{\Phi} V$$

First attempt:

$$\begin{split} \frac{\mathrm{d}\Phi}{\mathrm{d}s} &= \mathscr{F} \\ \Phi_{i,j}^{n+1} &= \varsigma \left[\Phi_{i+1,j}^n + \Phi_{i-1,j}^n + \Phi_{i,j+1}^n + \Phi_{i,j-1}^n \right] + \Phi_{i,j}^n \left[1 - 4\varsigma \right] \\ &- \Delta s \left[\frac{\lambda^2}{2} \sin \left(2\Phi_{i,j}^n \right) + \sin \Phi_{i,j}^n \right] \\ \varsigma &\equiv \frac{\Delta s}{\mathrm{h}^2} \end{split}$$

Unstable algorithm...try gain ...©



Second attempt: Examine the response of the field near the solution Φ_b under the effect of a slight perturbation $\delta\Phi$:

$$\frac{d\Phi}{ds} = \frac{d}{ds} \left(\Phi_b + \delta \Phi \right) = \underbrace{\frac{d\Phi_b}{ds}}_{\bullet \bullet} + \underbrace{\frac{d(\delta \Phi)}{ds}}_{\bullet \bullet} = \frac{d(\delta \Phi)}{ds}.$$

Hence the behavior of the system close to the bubble solution Φ_b is governed by a second-order operator.

Second attempt: Make \mathcal{O} a 2nd-order operator Δ^{\dagger} so that

$$\Delta^{\dagger}\mathcal{F} \equiv (\mathcal{F}')^{\dagger}\mathcal{F} = -\mathring{\nabla}^2\mathcal{F} + \partial_{\Phi}^2V\mathcal{F}.$$

Thus, we have

$$\frac{\mathrm{d}\Phi}{\mathrm{d}s} = \Delta^{\dagger} \mathscr{F}.$$

Second attempt: Make \mathcal{O} a 2nd-order operator Δ^{\dagger} so that

$$\Delta^{\dagger}\mathcal{F} \equiv (\mathcal{F}')^{\dagger}\mathcal{F} = -\mathring{\nabla}^2\mathcal{F} + \partial_{\Phi}^2V\mathcal{F}.$$

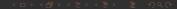
Thus, we have

$$\frac{\mathrm{d}\Phi}{\mathrm{d}s} = \Delta^{\dagger} \mathscr{F}.$$

Much better this time! ...But not good enough ... \odot (To have stability, vonNeumann shows $\Delta s \sim O(h^4)$)

Third time's the charm!: Put

$$\frac{\mathrm{d}^2\Phi}{\mathrm{d}s^2} + k\frac{\mathrm{d}\Phi}{\mathrm{d}s} = \Delta^\dagger \mathcal{F}$$



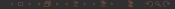
Third time's the charm!: Put

$$\frac{\mathrm{d}^2\Phi}{\mathrm{d}s^2} + k\frac{\mathrm{d}\Phi}{\mathrm{d}s} = \Delta^\dagger \mathcal{F}$$

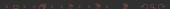
Then

$$\Phi_{i,j}^{n+1} = \frac{1}{1 + \frac{k}{2}\Delta s} \left\{ 2\Phi_{i,j}^{n} - \Phi_{i,j}^{n-1} \left[1 - \frac{k}{2}\Delta s \right] + \Delta^{\dagger} \mathscr{F}(\Phi_{i,j}^{n}) \Delta s^{2} \right\}$$

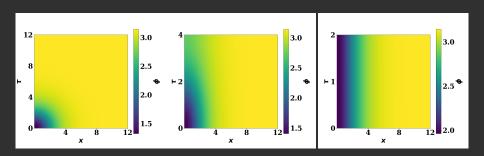
This is it! igotimes (To have stability, vonNeumann shows $\Delta s \sim O({
m h}^2)$)



Results

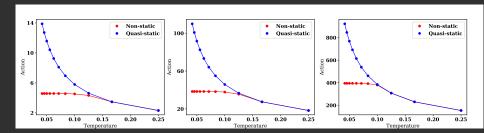


Transition from vacuum to thermal state

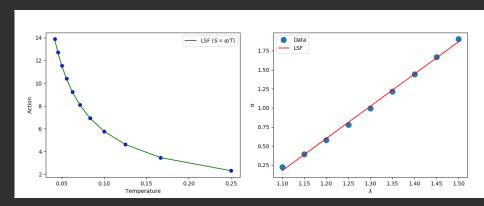


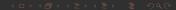


Action v Temperature (all three D cases)

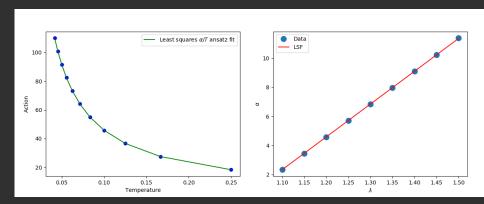


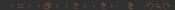
The most intriguing result!



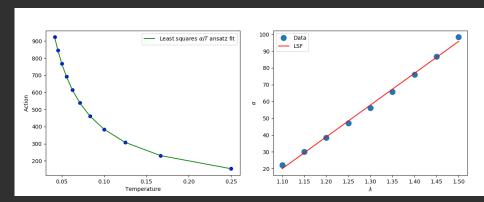


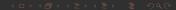
The most intriguing result!





The most intriguing result!





THANK YOU!