

# Appendix A

## Derivations

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### A.1 DERIVATION OF EULER'S FORMULA

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Assume a sinusoidal function:

$$x = \cos \theta + j \sin \theta \quad (\text{A.1})$$

(where  $j = \sqrt{-1}$  as usual)

Differentiating with respect to  $\theta$  produces:

$$\frac{dx}{d\theta} = j(\cos \theta + j \sin \theta) = jx \quad (\text{A.2})$$

Separating the variables gives:

$$\frac{dx}{x} = j d\theta \quad (\text{A.3})$$

and integrating both sides gives:

$$\ln x = j\theta + K$$

where  $K$  is the constant of integration. To solve for this constant, note that in [Equation \(A.1\)](#):  $x = 1$  when  $\theta = 0$ . Applying this condition to [Equation \(A.3\)](#):

$$\ln 1 = 0 = 0 + K; \quad K = 0;$$

so [Equation \(A.3\)](#) becomes:  $\ln x = j\theta$ .

or

$$x = e^{j\theta} \quad (\text{A.4})$$

but since  $x$  is defined in [Equation \(A.1\)](#) as:  $\cos \theta + j \sin \theta$

$$e^{j\theta} = \cos \theta + j \sin \theta \quad (\text{A.5})$$

Alternatively,

$$e^{-j\theta} = \cos \theta - j \sin \theta \quad (\text{A.6})$$

### A.2 CONFIRMATION OF THE FOURIER SERIES

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Fourier showed that a periodic function of period  $T$  can be represented by a series, possibly infinite, of sinusoids, or sine and cosines:

$$x(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos n\omega_0 t + b_n \sin n\omega_0 t) \quad (\text{A.7})$$

where  $\omega_0 = 2\pi/T$  and  $a_n$  and  $b_n$  are the Fourier coefficients.

To derive the Fourier coefficients, let us begin with the  $a_0$  or DC term. Integrating both sides of Equation (A.7) over a full period:

$$\int_0^T x(t)dt = \int_0^T \frac{a_0}{2}dt + \sum_{n=1}^{\infty} \int_0^T (a_n \cos n\omega_0 t + b_n \sin n\omega_0 t)dt \quad (\text{A.8})$$

For  $n = 0$ , the second term on the right-hand side is zero since the sum begins at  $n = 1$ , and the equation becomes

$$\begin{aligned} \int_0^T x(t)dt &= \int_0^T \frac{a_0}{2}dt; & \int_0^T x(t)dt &= \frac{a_0 T}{2}; \\ a_0 &= \frac{2}{T} \int_0^T x(t)dt \end{aligned} \quad (\text{A.9})$$

To find the other coefficients, multiply both sides of Equation (A.7) by  $\cos(m\omega_0 t)$ , where  $m$  is an integer, and again integrate both sides.

$$\begin{aligned} \int_0^T x(t)\cos(m\omega_0 t)dt &= \int_0^T \frac{a_0}{2}\cos(m\omega_0 t)dt + \sum_{n=1}^{\infty} \int_0^T (a_n \cos(m\omega_0 t)\cos n\omega_0 t \\ &\quad + b_n \cos(m\omega_0 t)\sin n\omega_0 t)dt \end{aligned} \quad (\text{A.10})$$

Rearranging:

$$\begin{aligned} \int_0^T x(t)\cos(m\omega_0 t)dt &= \int_0^T \frac{a_0}{2}\cos(m\omega_0 t)dt + \sum_{n=1}^{\infty} a_n \int_0^T \cos(m\omega_0 t)\cos n\omega_0 t dt \\ &\quad + \sum_{n=1}^{\infty} b_n \int_0^T \cos(m\omega_0 t)\sin n\omega_0 t dt \end{aligned} \quad (\text{A.11})$$

Since  $m$  is an integer, the first and third terms on the right-hand side integrate to zero for all  $m$ . The second term also integrates to zero for all  $m$  except  $m = n$ . At  $m = n$ , the second term becomes:

$$an \int_0^T \cos^2(n\omega_0 t)dt = \frac{\pi}{\omega_0} a_n = \frac{T}{2} a_n; \quad (\text{A.12})$$

so that:

$$\begin{aligned} \frac{T}{2}a_n &= \int_0^T x(t)dt \\ a_n &= \int_0^T \cos(n\omega_0 t)dt; \quad m = 1, 2, 3, \dots \end{aligned} \quad (\text{A.13})$$

The  $b_n$  coefficients are found in a similar fashion except Equation (A.7) is multiplied by  $\sin(m\omega_0 t)$ , then integrated. In this case, all but the third term integrate to zero and the third term is nonzero only for  $m = n$ .

$$bn \int_0^T \sin^2(n\omega_0 t)dt = \frac{\pi}{\omega_0}b_n = \frac{T}{2}b_n;$$

so that:

$$\begin{aligned} \frac{T}{2}b_n &= \int_0^T x(t)dt \\ b_n &= \int_0^T \sin(n\omega_0 t)dt; \quad m = 1, 2, 3, \dots \end{aligned}$$

### A.3 DERIVATION OF THE TRANSFER FUNCTION OF A SECOND-ORDER OP AMP FILTER

The op amp circuit for a second-order low-pass filter is shown in Figure A.1.

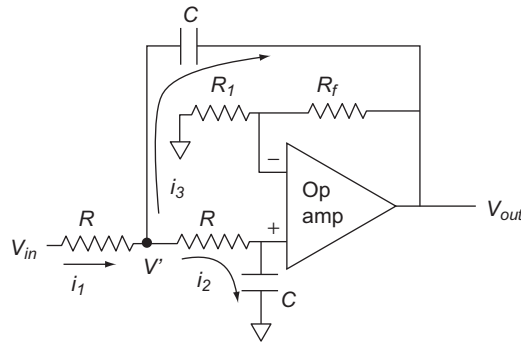


FIGURE A.1 A second-order op amp filter circuit.

This derivation applies not only to the following low-pass version, but also to the high-pass version where the positions of  $R$  and  $C$  are reversed.

Note that at node  $V'$  by KCL:  $i_1 - i_2 - i_3 = 0$ .

This allows us to write a nodal equation around that node:

$$\frac{V_{in} - V'}{R} - \frac{V' - V_{out}}{\frac{1}{CS}} - i_2 = 0$$

where:

$$i_2 = \frac{V^+}{\frac{1}{CS}} = V^+ CS$$

Since the two terminals of an op amp must be at the same voltage, the voltage  $V^+$  must be equal to  $V^-$ . Applying the voltage divider equation to the two feedback resistors,  $V^-$ , and hence  $V^+$ , can be found in terms of  $V_{out}$ :

$$V^+ = V^- = V_{out} G$$

where

$$G = \frac{R_1}{R_f + R_1}$$

So  $i_2$  becomes:  $V_{out}(Cs)/G$ . Substituting  $i_2$  into the nodal equation at  $V'$ :

$$\frac{V_{in} - V'}{R} - (V_{out} - V')Cs - \frac{V_{out}(Cs)}{G} = 0$$

$$\frac{V_{in}}{R} - \frac{V'}{R} - V'Cs + V_{out}Cs - \frac{V_{out}Cs}{G} = 0$$

$$\frac{V_{in}}{R} - V' \left( \frac{1}{R} - V'Cs \right) + V_{out}Cs \left( 1 - \frac{1}{G} \right) = 0$$

Note that  $V'$  can also be written in terms of just  $i_2$ :

$$V' = i_3 \left( R + \frac{1}{CS} \right) = \frac{V_{out}}{G} \left( R + \frac{1}{CS} \right)$$

Substituting this for  $V'$  in the nodal equation:

$$\frac{V_{in}}{R} - \frac{V_{out}Cs}{G} \left( \frac{1}{R} + Cs \right) \left( R + \frac{1}{Cs} \right) + V_{out}Cs \left( 1 - \frac{1}{G} \right) = 0$$

$$\frac{V_{in}}{R} = \frac{V_{out}Cs}{G} \left( 3 + \frac{1}{RCs} + RCs - G \right)$$

Solving for  $V_{out}/V_{in}$ :

$$\frac{V_{out}}{V_{in}} = \frac{G}{RCs \left( 1 + \frac{1}{RC} + RCs - g \right)} = \frac{G}{(RCs)^2 + (3 - G)RCs + 1}$$

$$\frac{V_{out}}{V_{in}} = \frac{\frac{G}{(RC)^2}}{s^2 + \frac{3 - G}{RC}s + \frac{1}{(RC)^2}}$$

#### A.4 DERIVATION OF THE TRANSFER FUNCTION OF AN INSTRUMENTATION AMPLIFIER

The classic circuit for a three-op amp instrumentation amplifier is shown in [Figure A.2](#).

To determine the transfer function, note that the voltage  $V_{in1}$  appears on both terminals of op amp 1, whereas  $V_{in2}$  appears on both terminals of op amp 2. The voltage out of op amp 1 will be equal to  $V_{in2}$  plus the voltage drop across the two resistors,  $R_2$  and  $R_1$ :

$$V_{out1} = i_{12}(R_1 + R_2) + V_{in2};$$

but

$$i_{12} = \frac{V_{in1} - V_{in2}}{R_1}$$

$$V_{out1} = \frac{V_{in1} - V_{in2}}{R_1}(R_1 + R_2) + V_{in2}$$

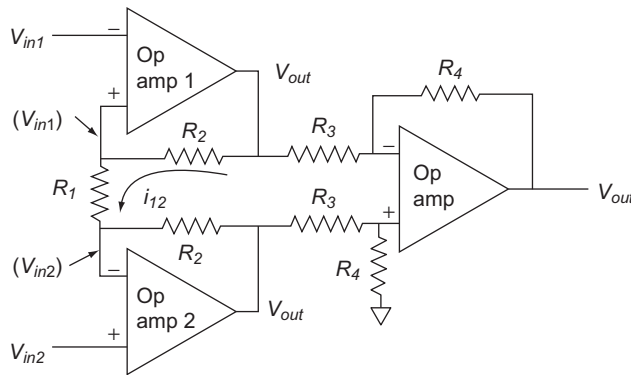


FIGURE A.2 Instrumentation amplifier circuit.

Applying the same logic to op amp 2, its output,  $V_{out2}$ , can be written as:

$$\begin{aligned}
 V_{out2} - V_{out1} &= \frac{V_{in1} - V_{in2}}{R_1} (R_1 + R_2) 2 - (V_{in1} - V_{in2}) \\
 &= \left[ \frac{2(R_1 + R_2)}{R_1} - 1 \right] (V_{in1} - V_{in2}) = \left[ \frac{2R_1 + 2R_2 - R_1}{R_1} \right] (V_{in1} - V_{in2}) \\
 V_{out2} - V_{out1} &= \left( \frac{R_1 + 2R_2}{R_1} \right) (V_{in1} - V_{in2}) \\
 V_{out} &= \left( \frac{R_4}{R_3} \right) (V_{out2} - V_{out1}) = \left( \frac{R_4}{R_3} \right) \left( \frac{R_1 + 2R_2}{R_1} \right) (V_{in1} - V_{in2})
 \end{aligned}$$

The overall output,  $V_{out}$ , is equal to the difference of  $V_{out2} - V_{out1}$  times the gain of the differential amplifier:  $R_4/R_3$ .