# 1 The $\chi^2$ in the presence of correlations

Suppose to have an ensamble of $n$ measurements having the following structure:

$$m_i \pm \sigma_{i,\mathrm{stat}} \pm \sigma_{i,\mathrm{unc}} \pm \sigma^{(1)}_{i,\mathrm{corr}} \pm \cdots \pm \sigma^{(k)}_{i,\mathrm{corr}} \,, \tag{1}$$

where $m_i$, with $i = 1, \ldots, n$, is the central value of the $i$-th measurement, $\sigma_{i,\mathrm{stat}}$ its (uncorrelated) statistical uncertainty, $\sigma_{i,\mathrm{unc}}$ its uncorrelated systematic uncertainty[1], and $\sigma^{(l)}_{i,\mathrm{corr}}$, with $l = 1, \ldots, k$, its correlated systematic uncertainties. With this information at hand, one can construct the full covariance matrix $V_{ij}$ as follows (see for example Ref. [3]):

$$V_{ij} = \left(\sigma^2_{i,\mathrm{stat}} + \sigma^2_{i,\mathrm{unc}}\right) \delta_{ij} + \sum_{l=1}^{k} \sigma^{(l)}_{i,\mathrm{corr}} \sigma^{(l)}_{j,\mathrm{corr}} \,. \tag{2}$$

This is a clearly symmetric matrix. Given a set of predictions $t_i$ corresponding to the $n$ measurements of the ensamble, the $\chi^2$ takes the form:

$$\chi^2 = \sum_{i,j=1}^{n} (m_i - t_i)\, V^{-1}_{ij}\, (m_j - t_j) = \mathbf{y}^T \cdot \mathbf{V}^{-1} \cdot \mathbf{y} \,, \tag{3}$$

where in the second equality we have used the matricial notation and defined $y_i = m_i - t_i$. A convenient way to compute the $\chi^2$ relies on the Cholesky decomposition of the covariance matrix $\mathbf{V}$. In particular, it can be proven that any symmetric and positive definite matrix $\mathbf{V}$ can be decomposed as:

$$\mathbf{V} = \mathbf{L} \cdot \mathbf{L}^T \,, \tag{4}$$

where $\mathbf{L}$ is a lower triangular matrix whose entries are related recursively to those of $\mathbf{V}$ as follows:

$$\begin{aligned} L_{kk} &= \sqrt{V_{kk} - \sum_{j=1}^{k-1} L^2_{kj}} \,, \\ L_{ik} &= \frac{1}{L_{kk}} \left( V_{ik} - \sum_{j=1}^{k-1} L_{ij} L_{kj} \right) \,, \quad k < i \,, \\ L_{ik} &= 0 \,, \quad k > i \,. \end{aligned} \tag{5}$$

It is then easy to see that the $\chi^2$ can be written as:

$$\chi^2 = \left|\mathbf{L}^{-1} \cdot \mathbf{y}\right|^2 \,. \tag{6}$$

But the vector $\mathbf{x} \equiv \mathbf{L}^{-1} \cdot \mathbf{y}$ is the solution of the linear system:

$$\mathbf{L} \cdot \mathbf{x} = \mathbf{y} \,, \tag{7}$$

that can be efficiently solved by forward substitution, so that:

$$\chi^2 = |\mathbf{x}|^2 \,. \tag{8}$$

Following this procedure, one does not need to compute explicitly the inverse of the covariance matrix $\mathbf{V}$, simplifying significantly the computation of the $\chi^2$.

[1] There could be more than one uncorrelated systematic uncertainty. In this case, $\sigma_{i,\mathrm{unc}}$ is just the square root of the sum in quadrature of all the uncorrelated systematic uncertainties.

## 2 Additive and multiplicative uncertainties

The correlated systematic uncertainties $\sigma^{(l)}_{i,\text{corr}}$ may be either *additive* or *multiplicative*. The nature of the single uncertainties is typically provided by the experiments that release the measurements. A typical example of multiplicative uncertainty is the luminosity uncertainty but there can be others.

Now let us express all the correlated systematic uncertainties $\sigma^{(l)}_{i,\text{corr}}$ as relative to the associate central value $m_i$, so that we define[2]:

$$\sigma^{(l)}_{i,\text{corr}} \equiv \delta^{(l)}_{i,\text{corr}} m_i \tag{9}$$

and let us also define $s_i^2 \equiv \sigma^2_{i,\text{stat}} + \sigma^2_{i,\text{unc}}$ so that Eq. (2) can be rewritten as:

$$V_{ij} = s_i^2 \delta_{ij} + \left( \sum_{l=1}^{k} \delta^{(l)}_{i,\text{corr}} \delta^{(l)}_{j,\text{corr}} \right) m_i m_j \,. \tag{10}$$

Now we split the correlated systematic uncertainties into $k_a$ additive uncertainties and $k_m$ multiplicative uncertainties, such that $k_a + k_m = k$. This way Eq. (10) takes the form:

$$V_{ij} = s_i^2 \delta_{ij} + \left( \sum_{l=1}^{k_a} \delta^{(l)}_{i,\text{add}} \delta^{(l)}_{j,\text{add}} + \sum_{l=1}^{k_m} \delta^{(l)}_{i,\text{mult}} \delta^{(l)}_{j,\text{mult}} \right) m_i m_j \,. \tag{11}$$

It is well known that this definition of the covariance matrix is problematic in that it results in the so-called D'Agostini bias of the multiplicative uncertainties [2]. A possible solution to this problem is the so-called $t_0$-prescription [1], where the experimental central value $m_i$ in the multiplicative term is replaced by a fixed theoretical predictions $t_i^{(0)}$, typically computed in a previous fit in which the "standard" definition of the covariance matrix in Eq. (2) (often referred to as *experimental* definition) is used. Applying the $t_0$ prescription, the covariance matrix takes the form:

$$V_{ij} = s_i^2 \delta_{ij} + \sum_{l=1}^{k_a} \delta^{(l)}_{i,\text{add}} \delta^{(l)}_{j,\text{add}} m_i m_j + \sum_{l=1}^{k_m} \delta^{(l)}_{i,\text{mult}} \delta^{(l)}_{j,\text{mult}} t_i^{(0)} t_j^{(0)} \,. \tag{12}$$

## 3 Artificial generation of correlated systematics

In order to implement the definition of the $\chi^2$ discussed above, it is necessary to have the experimental information in terms of the correlated systematic uncertainties $\sigma^{(l)}_{i,\text{corr}}$. This is what the experimental collaborations usually release. However, in some cases this information is given in terms of a covariance matrix. Therefore, one needs to find a workaround to generate correlated systematic uncertainties out of a covariance matrix.

Given a $n \times n$ symmetric matrix $\mathbf{C}$, it will have $n$ orthonormal eigenvectors $\mathbf{x}^{(i)}$, such that $\mathbf{x}^{(i)} \cdot \mathbf{x}^{(j)} = \delta_{ij}$, each of which will have a non-negative eigenvalue $\lambda_i$ associated:

$$\mathbf{C} \cdot \mathbf{x}^{(i)} = \lambda_i \mathbf{x}^{(i)} \,, \quad i = 1, \ldots, n \,. \tag{13}$$

[2]Note that this redefinition does not change the nature of the uncertainties, additive uncertainties remain additive as well as multiplicative uncertainties remain multiplicative.

If we define:

$$\sigma_{i,\mathrm{corr}}^{(l)} = \sqrt{\lambda_l} x_i^{(l)} \,, \quad i, l = 1, \ldots, n \,, \tag{14}$$

one can show that:

$$\sum_{l=1}^{n} \sigma_{i,\mathrm{corr}}^{(l)} \sigma_{j,\mathrm{corr}}^{(l)} = C_{ij} \,. \tag{15}$$

To prove this equality we start from the following matricial relation:

$$\mathbf{C} = \mathbf{Q} \cdot \mathbf{\Lambda} \cdot \mathbf{Q}^{-1} \,, \tag{16}$$

where $\mathbf{\Lambda}$ is a diagonal matrix with the eigenvalues $\lambda_i$ on the diagonal ($\Lambda_{ij} = \lambda_i \delta_{ij}$), while $\mathbf{Q}$ is a matrix whose columns are the eigenvectors $\mathbf{x}^{(i)}$ ($Q_{ij} = x_i^{(j)}$). In addition, since in this particular case $\mathbf{x}^{(i)} \cdot \mathbf{x}^{(j)} = \delta_{ij}$, this implies that:

$$\mathbf{Q}^T \cdot \mathbf{Q} = \mathbf{I} \quad \Rightarrow \quad \mathbf{Q}^{-1} = \mathbf{Q}^T \,, \tag{17}$$

so that:

$$\mathbf{C} = \mathbf{Q} \cdot \mathbf{\Lambda} \cdot \mathbf{Q}^T \,. \tag{18}$$

It follows that:

$$C_{ij} = \sum_{k,l=1}^{n} Q_{ik} \Lambda_{kl} Q_{jl} = \sum_{k,l=1}^{n} x_i^{(k)} \lambda_k \delta_{kl} x_j^{(l)} = \sum_{l=1}^{n} \lambda_l x_i^{(l)} x_j^{(l)} = \sum_{l=1}^{n} \sigma_{i,\mathrm{corr}}^{(l)} \sigma_{j,\mathrm{corr}}^{(l)} \,, \tag{19}$$

as required.

The matrix $\mathbf{C}$ can be regarded as the correlated contribution to the full covariance matrix $\mathbf{V}$. In particular, considering Eqs. (2) and (10), one can write:

$$\mathbf{V} = \mathbf{U} + \mathbf{C} \,, \tag{20}$$

where $\mathbf{U}$ is a diagonal matrix of uncorrelated uncertainties:

$$U_{ij} = s_i^2 \delta_{ij} \,. \tag{21}$$

This defines the matrix $\mathbf{C}$ as:

$$\mathbf{C} = \mathbf{V} - \mathbf{U} \,, \tag{22}$$

such that, given a $n \times n$ covariance matrix $\mathbf{V}$ along with its uncorrelated contribution $\mathbf{U}$, one can generate a set of $n$ *artificial* correlated systematics according to Eq. (14), where $\mathbf{C}$ is given in Eq. (22), for each of the $n$ measurements. This allows us to implement Eq. (12) for the construction of the covariance matrix.

# 4 Determining the systematic shifts

In order to visualise the effect of systematic uncertainties, it is instructive to compute the *systematic shift* generated by the systematic uncertainties. To do so, we need to write the $\chi^2$ in terms of the so-called "nuisance parameters" $\lambda_\alpha$. One can show that the definition of the $\chi^2$ in Eq. (3) is equivalent to [3]:

$$\chi^2 = \sum_{i=1}^{n} \frac{1}{s_i^2} \left( m_i - t_i - \sum_{\alpha=1}^{k} \lambda_\alpha \sigma_{i,\mathrm{corr}}^{(\alpha)} \right)^2 + \sum_{\alpha=1}^{k} \lambda_\alpha^2 \,. \tag{23}$$

The optimal value of the nuisance parameters can be computed by minimising the $\chi^2$ with respect to them, that is imposing:

$$\frac{\partial \chi^2}{\partial \lambda_\beta} = 0\,. \tag{24}$$

This yields the system:

$$\sum_{\beta=1}^{k} A_{\alpha\beta}\lambda_\beta = \rho_\alpha\,, \tag{25}$$

with:

$$A_{\alpha\beta} = \delta_{\alpha\beta} + \sum_{i=1}^{n} \frac{\sigma_{i,\mathrm{corr}}^{(\alpha)}\sigma_{i,\mathrm{corr}}^{(\beta)}}{s_i^2} \quad \text{and} \quad \rho_\alpha = \sum_{i=1}^{n} \frac{m_i - t_i}{s_i^2}\sigma_{i,\mathrm{corr}}^{(\alpha)}\,, \tag{26}$$

that determines the values of $\lambda_\beta$. The quantity:

$$d_i = \sum_{\alpha=1}^{k} \lambda_\alpha \sigma_{i,\mathrm{corr}}^{(\alpha)} \tag{27}$$

in Eq. (23) can be interpreted as a shift caused by the correlated systematic uncertainties. Defining the shifted predictions as:

$$\bar{t}_i = t_i + d_i\,, \tag{28}$$

the $\chi^2$ reads:

$$\chi^2 = \sum_{i=1}^{n} \left(\frac{m_i - \bar{t}_i}{s_i}\right)^2 + \sum_{\alpha=1}^{k} \lambda_\alpha^2\,. \tag{29}$$

Therefore, up to a penalty term given by the sum of the square of the nuisance parameters, the $\chi^2$ takes the form of the uncorrelated definition. In order to achieve a visual assessment of the agreement between data and theory, it appears natural to compare the central experimental values $m_i$ to the shifted theoretical predictions $\bar{t}_i$ in units of the uncorrelated uncertainty $s_i$.

# References


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