

## Noninteracting, Zero temperature

$\hat{h}_j$  operator:  $\hat{h}_j = -J(\hat{c}_j^\dagger \hat{c}_{j+1} + \hat{c}_{j+1}^\dagger \hat{c}_j)$

Initial Hamiltonian:  $\hat{P} = \sum_{j=0}^{L-1} j \hat{h}_j = -J \sum_{j=0}^{L-1} j(\hat{c}_j^\dagger \hat{c}_{j+1} + \hat{c}_{j+1}^\dagger \hat{c}_j)$

Post-quench Hamiltonian:  $\hat{H} = \sum_{j=0}^{L-1} \hat{h}_j = \sum_{j=0}^{L-1} -J(\hat{c}_j^\dagger \hat{c}_{j+1} + \hat{c}_{j+1}^\dagger \hat{c}_j)$

Emergent Hamiltonian:  $\mathcal{H}(t) = \hat{P} + t\hat{Q} = \sum_{j=0}^{L-1} \left[ j \hat{c}_j^\dagger \hat{c}_j - t(i\hat{c}_{j+1}^\dagger \hat{c}_j - i\hat{c}_j^\dagger \hat{c}_{j+1}) \right]$

Analytical single particle correlation solution ( $m \neq n$ ):  $C_{mn} = i^{n-m} \frac{t(J_m(2t)J_{n+1}(2t) - J_{m+1}(2t)J_n(2t))}{n-m}$

Analytical single particle correlation solution ( $m = n$ ):

Effective single particle correlation solution:  $C_{mn} = \sum_{j=1}^N \phi_{q_j}^*(m) \phi_{q_j}(n)$  ( $\{\phi_q\}$  is the set of  $N = 1$  energy eigenstates)

## Noninteracting, Infinite Temperature

Initial Hamiltonian, current operator, emergent Hamiltonian are the same as in the  $T = 0$  case.

Initial (mixed) state:

The full set of computational basis states contains each arrangement of  $N$  particles along length  $L = 4N$  lattice. Extract basis states in which all  $N$  particles are confined to the middle half of the lattice. Initialize initial state  $|\psi_0\rangle$  as an empty list. Convert these binary strings into normalized lists and append each to  $|\psi_0\rangle$ .

Effective (mixed) state:

Diagonalize emergent Hamiltonian. For each normalized computational basis state in initial state  $|\psi_0\rangle$ , for each site entry in this basis state, if the entry is not zero, append the eigenstate (list) from the emergent Hamiltonian corresponding to that entry to a list. Append that 2d list to the mixed state list.

## Interacting ( $t$ - $V$ )

$\hat{h}_j$  operator:  $\hat{h}_j = -J(c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + V(c_j^\dagger c_j - 1/2)(c_{j+1}^\dagger c_{j+1} - 1/2)$

Initial Hamiltonian:

$$\begin{aligned}\hat{P} &= \sum_{j=0}^{L-1} j \hat{h}_j \\ &= \sum_{j=0}^{L-1} j \left( J(c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + V(c_j^\dagger c_j - 1/2)(c_{j+1}^\dagger c_{j+1} - 1/2) \right)\end{aligned}$$

Current operator:

$$\hat{Q} = \sum_j \left[ iJ^2(c_{j+2}^\dagger c_j - c_j^\dagger c_{j+2}) - iV \left( (c_{j+1}^\dagger c_j - c_j^\dagger c_{j+1})(n_{j+2} - 1/2) + (c_{j+2}^\dagger c_{j+1} - c_{j+1}^\dagger c_{j+2})(n_j - 1/2) \right) \right]$$

Post-quench Hamiltonian:

$$\begin{aligned}\hat{H} &= \sum_{j=0}^{L-1} \hat{h}_j \\ &= \sum_{j=0}^{L-1} \left[ -J(c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + V(c_j^\dagger c_j - 1/2)(c_{j+1}^\dagger c_{j+1} - 1/2) \right]\end{aligned}$$

Emergent Hamiltonian:  $\mathcal{H}(t) = \hat{P} + t\hat{Q}$

Exact single particle correlation solution:  $C_{ij} = \sum_{\alpha, \beta} \sum_{m, n}^{\binom{L}{N}} e^{it(E_\alpha - E_\beta)} \langle \alpha | \psi_0 \rangle \langle \psi_0 | \beta \rangle \alpha^*(m) \beta(n) \langle m | \hat{c}_i^\dagger \hat{c}_j | n \rangle$

$|\alpha\rangle, |\beta\rangle$  denote energy eigenstates obtained via exact diagonalization of the post-quench Hamiltonian

Effective single particle correlation solution:  $C_{ij} = \sum_{m, n}^{\binom{L}{N}} \phi(m) \phi^*(n) \langle m | \hat{c}_i^\dagger \hat{c}_j | n \rangle$