

13.1 Scalar Scattering Theory

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August 2024

1 The Basic Integral Equation

From [1.2 EM Fields] :

$$\nabla^2 \vec{E} - \epsilon \mu \frac{\partial^2 \vec{E}}{\partial t^2} + (\nabla \ln \mu) \times (\nabla \times \vec{E}) + \nabla [\vec{E} \cdot (\nabla \ln \epsilon)] = 0$$

Specialize the above equation to a nonmagnetic, monochromatic case :
 $\nabla \ln \mu = 0$ and $\epsilon \mu = \frac{1}{v^2} = \frac{1}{c^2} \cdot \frac{c^2}{v^2} = \frac{n^2}{c^2}$:

$$\nabla^2 \vec{E}(\vec{r}, w) + \frac{n^2(\vec{r}, w)}{c^2} w^2 \vec{E}(\vec{r}, w) + \nabla [\vec{E}(\vec{r}, w) \cdot (\nabla \ln \epsilon(\vec{r}, w))] = 0$$

Denote $k \equiv \frac{w}{c}$:

$$\nabla^2 \vec{E}(\vec{r}, w) + k^2 n^2(\vec{r}, w) \vec{E}(\vec{r}, w) + \nabla [\vec{E}(\vec{r}, w) \cdot (\nabla \ln \epsilon(\vec{r}, w))] = 0$$

If the dielectric constant varies so slowly such that $\nabla \ln \epsilon(\vec{r}, w) = 0$ then

$$\nabla^2 \vec{E}(\vec{r}, w) + k^2 n^2(\vec{r}, w) \vec{E}(\vec{r}, w) = 0$$

A nonmagnetic, time-harmonic, monochromatic, $\epsilon(\vec{r}, w)$ slowly varying wave equation. Note that each Cartesian components of \vec{E} are decoupled now. So, just let $U(\vec{r}, w)$ one of the components :

$$\nabla^2 U(\vec{r}, w) + k^2 n^2(\vec{r}, w) U(\vec{r}, w) = 0$$

To solve the differential equation with homogeneous/inhomogeneous part, manipulate the equation :

$$\nabla^2 U(\vec{r}, w) + k^2 U(\vec{r}, w) = -k^2 [n^2(\vec{r}, w) - 1] U(\vec{r}, w)$$

Define the *scattering potential of the medium* $F(\vec{r}, w) = \frac{1}{4\pi} k^2 [n^2(\vec{r}, w) - 1]$:

$$\nabla^2 U(\vec{r}, w) + k^2 U(\vec{r}, w) = -4\pi F(\vec{r}, w) U(\vec{r}, w)$$

This is a differential form of the scalar scattering equation. We'll rewrite this into a integral form.

Express $U(\vec{r}, w) = U^{(i)}(\vec{r}, w) + U^{(s)}(\vec{r}, w)$ where $U^{(i)}$ is an incident field and $U^{(s)}$ is a scattered field.

The incident field is usually a plane wave so $(\nabla^2 + k^2)U^{(i)}(\vec{r}, w) = 0$:

$$(\nabla^2 + k^2)U^{(s)}(\vec{r}, w) = -4\pi F(\vec{r}, w)U(\vec{r}, w) \quad (1)$$

Denote $G(\vec{r} - \vec{r}')$: a Green's function of the Helmholtz operator $(\nabla^2 + k^2)$:

$$(\nabla^2 + k^2)G(\vec{r} - \vec{r}') = -4\pi\delta(\vec{r} - \vec{r}') \quad (2)$$

(1) $\times G(\vec{r} - \vec{r}', w) - (2) \times U^{(s)}(\vec{r}, w)$ then

$$G\nabla^2 U^{(s)} - U^{(s)}\nabla^2 G = -4\pi GFU + 4\pi U^{(s)}\delta(\vec{r} - \vec{r}')$$

Take volume integral over V_R (larger than the medium volume V and has a radius R) :

(1) LHS

With the Green's second identity,

$$\int_{V_R} d\tau' [G\nabla^2 U^{(s)} - U^{(s)}\nabla^2 G] = \int_{S_R} (G\frac{\partial U^{(s)}}{\partial \vec{n}} - U^{(s)}\frac{\partial G}{\partial \vec{n}}) \cdot d\vec{S}$$

(2) RHS

$$\begin{aligned} & \int_{V_R} d\tau' [-4\pi GFU + 4\pi U^{(s)}\delta(\vec{r} - \vec{r}')] \\ &= -4\pi \int_{V_R} d\tau' G(\vec{r} - \vec{r}') F(\vec{r}') U(\vec{r}') + 4\pi \int_{V_R} d\tau' U^{(s)}(\vec{r}') \delta(\vec{r} - \vec{r}') \\ &= 4\pi U^{(s)}(\vec{r}) - 4\pi \int_V d\tau' G(\vec{r} - \vec{r}') F(\vec{r}') U(\vec{r}') \end{aligned}$$

Then, (1)=(2)

$$U^{(s)}(\vec{r}) = \int_V d\tau' G(\vec{r} - \vec{r}') F(\vec{r}') U(\vec{r}') + \int_{S_R} (G\frac{\partial U^{(s)}}{\partial \vec{n}} - U^{(s)}\frac{\partial G}{\partial \vec{n}}) \cdot d\vec{S}$$

Now, we'll show the surface integral vanishes when $R \rightarrow \infty$. To do that, choose the Green function :

$$G(\vec{r} - \vec{r}', w) = \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|}$$

the *outgoing free-space Green's function*. Note that $G \rightarrow 0$ as $R \rightarrow \infty$ in this case. We can also assume when $R \rightarrow \infty$, $U^{(s)}$ behaves as an outgoing spherical wave. Then, $U^{(s)} \rightarrow 0$ as $R \rightarrow \infty$. With these choice and assumption, we get :

$$\boxed{U^{(s)}(\vec{r}, w) = \int_V d\tau' G(\vec{r} - \vec{r}', w) F(\vec{r}', w) U(\vec{r}', w)}$$

Or,

$$U^{(s)}(\vec{r}, w) = \int_V d\tau' F(\vec{r}', w) U(\vec{r}', w) \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|}$$

Using a plane incident wave $U^{(i)}(\vec{r}, w) = e^{iks_0 \cdot \vec{r}}$, we have

$$U(\vec{r}, w) = e^{iks_0 \cdot \vec{r}} + \int_V d\tau' F(\vec{r}', w) U(\vec{r}', w) \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|}$$

the *integral equation of potential scattering*.

References

- [1] Born, M., Wolf, E., Bhatia, A. B. (2019). Principles of Optics: 60th Anniversary Edition. United Kingdom: Cambridge University Press.