#### GravLens

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History

Gravitational Lensing Theories

....

Applications

# Gravitational Lensing Theories, Questions and Applications

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#### Overview

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History

**Gravitational Lensing Theories** 

**Related Questions** 

### Predictions from Newtonian and GR

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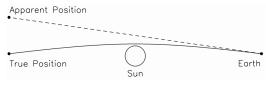
related Questi

- First proposed by Soldner (1801) using Newtonian theory, given a deflection angle (Appendix-1 (Page 66) for derivation).  $\alpha = \frac{2GM}{c^2r}, \ 0.85 \ \text{arcsec for the Sun}$
- ► Einstein (1911) derived the same result using Equivalence principle and Euclidean metric
- ► Einstein (1915) derived the new result using General Relativity.
  - $\alpha = \frac{4GM}{c^2r}$ , 1.7 arcsec for the Sun

## Eddington's observation of the Solar Eclipse

In 1919, Eddington measured a value close to GR's prediction using the data collected during an eclipse, stars with apparent position near the sun become visible.





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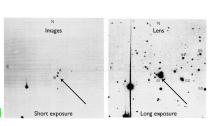
### The first example for GravLens: 0957+561

**Eddington** (1920): Multiple light paths o multi images

 $\blacktriangleright$  Walsh et al., (1979) quasar QSO 957+561 A,B found at  $z\sim1.4$ , two seen images separated by  $6^{''}$ 

Lens: evidence

- 1. Lensing galaxy at  $z\sim0.36$
- 2. Similar spectra
- 3. Ratio of optical and radio flux
- 4. VLBI imging: small scale features



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Walsh, Carswell & Weymann 1979, 0957+561 A, B: twin quasistellar objects or gravitational lens? https://ui.adsabs.harvard.edu/abs/1979Natur.279..381W/abstract

► Einstein Field Equations:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$$

Geodesic equations:

$$\frac{d^2x^{\beta}}{d\lambda^2} + \Gamma^{\beta}_{\mu\nu} \frac{dx^{\nu}}{d\lambda} \frac{dx^{\nu}}{d\lambda} = 0$$

We can calculate how gravity bends light by solving geodesic eqution.

▶ To compute the Christoffel symbols  $\Gamma^{\beta}_{\mu\nu}$ , requires solving for the metric tensor  $g_{\mu\nu}$ , which requires solving the curvature equations  $R_{\mu\nu}=0$ ,  $\leftarrow$  ten nonlinear partial differential equations.

### General Relaticity and light deflection

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Or, the velocity of the photon from the Schwarzshild metric,

 $ds^{2} = -(1+2\Phi)dt^{2} + (1-2\Phi)(dx^{2} + dy^{2} + dz^{2}),$ 

- ▶ and Poisson Equation,  $\nabla^2 \Phi = 4\pi G \rho$ ,
- ▶ and light interval:  $g_{\mu\nu} \frac{dx^{\mu}}{d\lambda} \frac{dx^{\nu}}{d\lambda} = 0$
- which gives,

$$v = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{dt} = \sqrt{\frac{1 + 2\Phi}{1 - 2\Phi}} \approx 1 + 2\Phi$$

The gravitational field decreases the speed of propagation

 $https://web.stanford.edu/\ oas/SI/SRGR/notes/SchwarzschildSolution.pdf,\ The\ Schwarzschild\ Solution$ 

# General Relaticity and light deflection

$$v=\frac{\sqrt{dx^2+dy^2+dz^2}}{dt}=\sqrt{\frac{1+2\Phi}{1-2\Phi}}\approx 1+2\Phi$$
 (natural units)  $\rightarrow v=c\left(1+\frac{2}{c^2}\Phi\right)$  (SI)

- $\blacktriangleright$  define refraction index:  $n=1-\frac{2}{c^2}\Phi=1+\frac{2}{c^2}|\Phi|\geq 1$
- ▶ deflection angle:  $\vec{\hat{\alpha}} = -\int \vec{\nabla}_{\perp} n dl = \frac{2}{c^2} \int \vec{\nabla}_{\perp} \Phi dl$
- which is twice of the Newtonian prediction,
- for point mass lens,  $\hat{\alpha} = \frac{4GM}{bc^2}$

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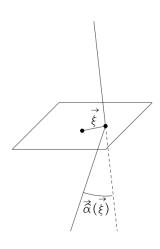
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# Thin screen approximation

- Most deflection occurs near to the lens  $(|z| \sim b)$   $\rightarrow$  treat all deflection as in the lens plane
- Projected surface density:  $\Sigma(\vec{\xi}) = \int \rho(\vec{\xi}, z) dz$
- ▶ Deflection angle:  $\hat{\alpha}(\vec{\xi}) = \frac{4G}{c^2} \int \frac{(\vec{\xi} \vec{\xi'})\Sigma(\vec{\xi'})}{|\vec{\xi} \vec{\xi'}|^2} d^2 \vec{\xi'}$
- In circular symmetry cases:  $\hat{\alpha}(\xi) = \frac{4GM(\xi)}{c^2 \xi}$   $M(\xi) = 2\pi \int_0^{\xi} \Sigma(\xi') \xi' d\xi'$



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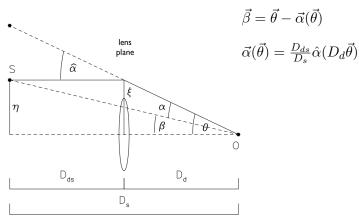
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reduced Question

### The Lens Equation

Connecting the position of images in the Lens plane and corresponding sources the Source plane.



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### Case 1: Point mass Lens

$$\begin{cases} \hat{\alpha}(\xi) = \frac{4GM(\xi)}{c^2 \xi} \\ \vec{\alpha}(\vec{\theta}) = \frac{D_{ds}}{D_s} \hat{\alpha}(D_d \vec{\theta}) \end{cases} \rightarrow \vec{\alpha}(\vec{\theta}) = \frac{4GM}{c^2} \frac{D_{ds}}{D_d D_s} \frac{\vec{\theta}}{|\vec{\theta}|^2}$$

Source

Lens

Observer

and  $\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$ , if  $\beta = 0$ , gives the Einstein Radius,

$$\theta_E = \sqrt{\frac{4GM}{c^2} \frac{D_{ds}}{D_d D_s}}$$

The Lens equation:  $\vec{eta} = \vec{ heta} - heta_E^2 \frac{\vec{ heta}}{|\vec{ heta}|^2}$ 

- if  $\beta > \theta_E \rightarrow$ , weakly lensed and weakly distorted image,
- if  $\beta < \theta_E \rightarrow$ , stronly lensed and multi images:

$$\theta_{\pm} = \frac{1}{2} \left( \beta \pm \sqrt{\beta^2 + 4\theta_E^2} \right)$$



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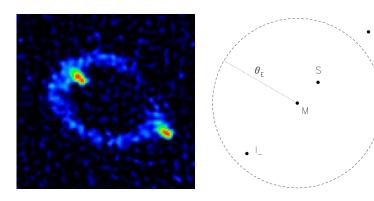
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#### Case 1: Point mass Lens

Hewitt+ 1987, First Einstein Ring discovered in Radio Band & VLA



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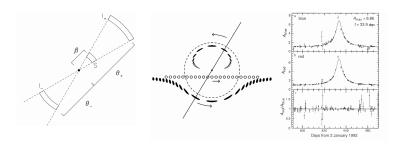
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Hewitt, J., Turner, E., Schneider, D. et al. Unusual radio source MG1131+0456: a possible Einstein ring. Nature 333, 537–540 (1988) doi:10.1038/333537a0

### Magnification

Gravitational lensing preserves surface brightness (Liouville's Theorem), but changes the apparent solid angle of the source  $\rightarrow magnification = \frac{Area_{image}}{Area_{source}}$ 



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Possible Gravitational Microlensing of a Star in the Large Magellanic Cloud https://arxiv.org/pdf/astro-ph/9309052v1.pdf

### Magnification

Local properties of the lens mapping, described by its Jacobian matrix A

$$A \equiv \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} - \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j}\right) = \left(\delta_{ij} - \frac{\partial^2 \Psi(\vec{\theta})}{\partial \theta_i \partial \theta_j}\right)$$
$$\mu = \left|\det\left(\frac{\partial \vec{\beta}}{\partial \vec{\theta}}\right)\right|^{-1} \equiv \left|\det\left(\frac{\partial \beta_i}{\partial \theta_j}\right)\right|^{-1}$$

If circularly symmetric,

$$\mu = \frac{\theta}{\beta} \frac{d\theta}{d\beta}$$

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# Case 1: Point mass Lens - Magnification

► Images:

$$\theta_{\pm} = \frac{1}{2} \left( \beta \pm \sqrt{\beta^2 + 4\theta_E^2} \right)$$

▶ Define  $u = \beta \theta_E^{-1}$ , Magnification:

$$\mu_{\pm} = \left[1 - \left(\frac{\theta_E}{\theta_{\pm}}\right)^4\right]^{-1}$$
$$= \frac{u^2 + 2}{2u\sqrt{u^2 + 4}} \pm \frac{1}{2}$$

► Total magnification:

$$\mu = |\mu_+| + |\mu_-| = \frac{u^2 + 2}{u\sqrt{u^2 + 4}}$$

e.g. for source on the Einstein Ring:

$$\beta=\theta_E, u=1 
ightarrow \mu=1.34 
ightarrow$$
 magnitude increase 0.32.

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Trajectories

Time (tm)

- ▶ Passage through potential also leads to time delay
- lacktriangle without potential:  $t_0 = \int rac{dl}{c}$
- with potential:

$$t_{1} = \int_{src}^{obs} \frac{dl}{v} = \int_{src}^{obs} \frac{dl}{c - \frac{2}{c}|\Phi|}$$
$$= \int_{src}^{obs} \frac{dl/c}{1 - \frac{2}{c^{2}}|\Phi|} = \int_{src}^{obs} \frac{dl}{c} [1 + \frac{2}{c^{2}}|\Phi|]$$

- $lackbox{ so, } \Delta t = \int_{src}^{obs} rac{2}{c^3} |\Phi| dl 
  ightarrow {
  m Shapiro \ delay} \ {
  m (1964)}$
- Total time delay is the sum of the extra path length from the deflection and the gravitational time delay

$$t(\vec{\theta}) = \frac{1 + z_d}{c} \frac{D_d D_s}{D_{ds}} \left[ \frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \psi \vec{\theta} \right] = t_{geom} + t_{grav}$$

- Galaxy lenses, the distributed nature of mass
- lackbox Simple model assumes that mass ightarrow particles of ideal gas
- ► ideal gas:
  - equation of state:  $p = \frac{\rho kT}{m}$
  - In thermal equilibrium, T is related to the 1-d velocity dispersion:  $m\sigma_v^2=kT$
  - In hydrostatic equilibrium,

$$\frac{p'}{\rho} = -\frac{GM(r)}{r^2}, M'(r) = 4\pi r^2 \rho$$

- > solve the EOS  $\rightarrow$  density profile:  $\rho(r) = \frac{\sigma_v^2}{2\pi G} \frac{1}{r^2}$
- ightharpoonup so mass profile:  $M(r)=rac{2\sigma_v^2}{G}$

# Case 2: Singular Isothemal Sphere (SIS)

For ideal gas, rotational velocity in circular orbit:

$$\frac{p'}{\rho} = -\frac{GM(r)}{r^2}, M'(r) = 4\pi r^2 \rho$$

Surface mass density:

$$\begin{split} v_{rot}^2 &= GM/r \rightarrow \\ dM &= \frac{v_{rot}^2}{G} dr = 4\pi r^2 \rho(r) \rightarrow \\ \rho(r) &= \frac{v_{rot}^2}{4\pi G} \frac{dr}{r^2} \rightarrow \\ \Sigma(\xi) &= \frac{v_{rot}^2}{4\pi G} \int_{-\infty}^{\infty} \frac{dz}{(z+\xi)^2} = \frac{v_{rot}^2}{4G\xi} = \frac{\sigma_v^2}{2G\xi} \end{split}$$

where, 
$$M(\xi)=2\pi\int_0^\xi \Sigma(\xi')\xi'd\xi'$$
 , and  $\hat{\alpha}(\xi)=\frac{4GM(\xi)}{c^2\xi}$  which gives:  $\hat{\alpha}=4\pi\frac{\sigma_v^2}{c^2}$ 

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# Case 2: Singular Isothemal Sphere (SIS)

using 
$$\begin{cases} \hat{\alpha} = 4\pi\sigma_v^2/c^2 \\ \vec{\alpha}(\vec{\theta}) = \hat{\alpha}(D_d\vec{\theta})D_{ds}/D_s \end{cases}$$
 and lens equation  $\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$ 

we can get:

$$\alpha = \hat{\alpha} \frac{D_{ds}}{D_s} = 4\pi \frac{\sigma_v^2}{c^2} \frac{D_{ds}}{D_s} = \theta_E$$

for strong lensing, get two images as for point mass:

$$\vec{\beta} = \vec{\theta} - \theta_E \frac{\vec{\theta}}{|\vec{\theta}|}, \frac{\vec{\theta}}{|\vec{\theta}|} = \pm 1 \rightarrow \theta_{\pm} = \beta \pm \theta_E$$

Magnification can be very large for source aligned with lens, from  $\mu=\frac{\theta}{\beta}\frac{d\theta}{d\beta}\to$ :

$$\mu_{\pm} = \frac{\theta_{\pm}}{\beta} = 1 \pm \frac{\theta_E}{\beta} = \left(1 \mp \frac{\theta_E}{\theta_{\pm}}\right)^{-1}$$

Separation of the two images is typically  $\sim$  arcsec for galaxy lenses:

$$\theta_E = 1''.6 \left(\frac{\sigma_v}{200kms^{-1}}\right)^2 \left(\frac{D_{ds}}{D_s}\right)$$

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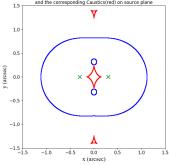
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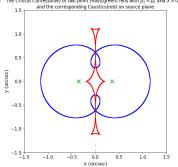
- ▶ Grav Lens changes the observed brightness of the source, determined by the Jacobian matrix from the lens equation:  $A \equiv \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j}\right) = \left(\delta_{ij} \frac{\partial^2 \Psi(\vec{\theta})}{\partial \theta_i \partial \theta_j}\right)$   $\mu(\theta) = \frac{1}{\det A(\theta)}$
- $\blacktriangleright$  Image at  $\vec{\theta}$  is magnified by a factor of  $|\mu(\vec{\theta})|$
- Notice that  $|\mu(\vec{\theta})|$  diverge at  $det A(\theta) = 0 \to$  these points in the image plane form closed curves, which is so called <u>critical lines</u>.
- Corresponding curves in the source plane obtained via the lens equation are called caustics.

#### Caustics and Critical lines - Point sources

### Separation 0.6'', 0.7'':

The Critical Curves(blue) of two point mass(green) lens with  $\mu_1 = \mu_2$  and X = 0.30 The Critical Curves(blue) of two point mass(green) lens with  $\mu_1 = \mu_2$  and X = 0.35 and the corresponding Caustics(red) on source plane





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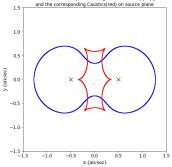
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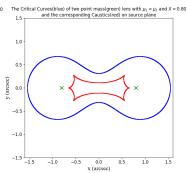
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### Caustics and Critical lines - Point sources

# Separation 1.0'', 1.6'':







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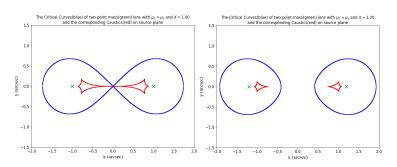
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### Caustics and Critical lines - Point sources

### Separation 2.0'', 2.4'':



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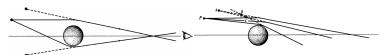
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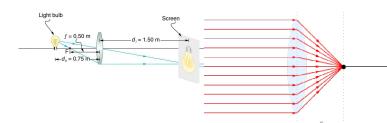
#### Gravitational Lens vs Genuine Lens

• Grav lens: the observer sees the source at two distinct locations.  $\alpha \propto b^{-1}$ 



Grav lens has no well-defined focal length and cannot produce genuine images, the "images" are corresponds merely to a direction of incidence of light on the observer, not a genuine image in the sky

▶ Genuine lens,  $\alpha \propto b$ 



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A galaxy at redshift 0.5 can be modelled as a singular isothermal sphere; its dispersion is 200km/s. A background source at redshift 2 is lensed by the foreground galaxy into two images with a brightness ratio of 3:1, what are the angular separation and time delay between the two images? You can assume the usual cosmology.

First assume  $H_0=69.6kms^{-1}, \Omega_m=0.286, \Omega_{vac}=0.714$  then the distance of z = 0.5 and z = 2 can be computed, using theories of SIS model and the time delay function we can estimate the angular separation  $\sim 1.2$  arcsec and time delay  $\sim 45$  days

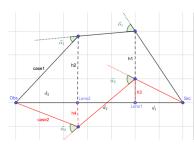
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Two (N=2) galaxies are aligned perfectly with the Earth and a distant quasar. Each galaxy can be modelled as a singular isothermal sphere. How many Einstein rings are formed as a result? How will your results generalize when you have N>2 galaxies?



For N galaxies, the quasar will render  $2^{N-1}$  Einstein rings, and if we consider galaxy lensed by galaxy, will generate (at most)  $\sum_{i=1}^{N-1} 2^{i-1}$  more, so the total number of Einstein rings rendered by N galaxies and a quasar will be:  $\sum_{i=1}^{N} 2^{i-1} = 2^{N} - 1$ , at most.

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Related Questions

A background source is multiply-imaged, can the brightest image arrive last?

Using the relation between effective lensing potential and magnification  $\&\ \text{time}\ \text{delay}$ 

Effective lensing potential:

$$\Psi(\vec{\theta}) = \frac{D_{ds}}{D_d D_s} \frac{2}{c^2} \int \Phi(D_d \vec{\theta}, z) dz, \ \vec{\nabla}_{\theta} \Psi = \vec{\alpha}$$

Jacobian matrix:
$$A = \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \delta_{ij} - \frac{\partial^2 \Psi(\vec{\theta})}{\partial \vec{\theta_i} \partial \vec{\theta_j}} \equiv \delta_{ij} - \Psi_{ij}$$

Magnification: 
$$\mu = |det A|^{-1}$$

time delay: 
$$t(\vec{\theta}) = \frac{1+z_d}{c} \frac{D_d D_s}{D_{ds}} \left[ \frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \Psi(\vec{\theta}) \right]$$

$$\mu \leftarrow \Psi \rightarrow t$$

Simulation, under a potential, given a  $\vec{\beta}$ , solve the lens equation for  $\vec{\theta_i}$  see if  $\vec{\theta_k}$  which has maximum magnification also have largest time delay

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Study the caustics and critical curves of two singular isothermal spheres lensing a background quasar. Study two cases 1) Both galaxies are at the same redshift and, 2) these two galaxies are at different redshifts.

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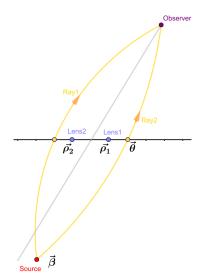
Applications

Mass density  $\rightarrow$  Surfacce mass density  $\rightarrow$  Effective lensing potential  $\rightarrow$  Jacobian matrix  $\rightarrow$  Magnification draw those points which make magnification largest  $\rightarrow$  critical lens  $\rightarrow$  caustics

### case 1) Both galaxies are at the same redshift

$$\vec{\alpha}(\vec{\theta}) = \theta_{E1} \frac{\vec{\theta} - \vec{\rho_1}}{|\vec{\theta} - \vec{\rho_1}|} + \theta_{E2} \frac{\vec{\theta} - \vec{\rho_2}}{|\vec{\theta} - \vec{\rho_2}|}$$
 where  $\vec{\rho_1}, \vec{\rho_2}$  are positions of Lens1, Lens2 on the lens plane.  $\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$ 

$$A = \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} - \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j}\right)$$
$$\mu(\vec{\theta}) = \frac{1}{\det A(\vec{\theta})}$$



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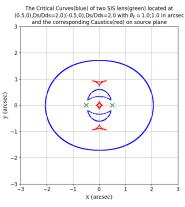
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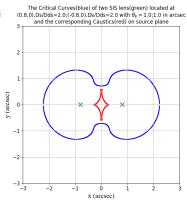
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### case 1) Both galaxies are at the same redshift

# Two SIS lenses ( $\theta_E=1.0^{''}$ ), separation $1.0^{''}, 1.6^{''}$





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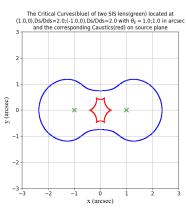
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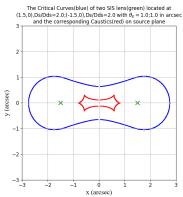
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### case 1) Both galaxies are at the same redshift

# Two SIS lenses ( $\theta_E=1.0^{''}$ ), separation $2.0^{''},3.0^{''}$





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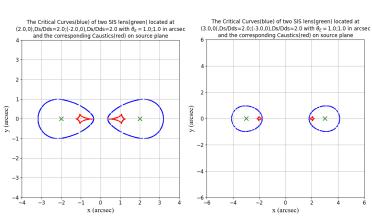
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### case 1) Both galaxies are at the same redshift

# Two SIS lenses ( $\theta_E=1.0^{''}$ ), separation $4.0^{''}, 6.0^{''}$



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### case 2) these two galaxies are at different redshifts

Lens equations of two lens planes system:

$$\vec{\eta} = \frac{D_s}{D_1} \vec{\xi}^{(1)} - D_{1s} \vec{\alpha}^{(1)} \left( \vec{\xi}^{(1)} \right) - D_{2s} \vec{\alpha}^{(2)} \left( \vec{\xi}^{(2)} \right)$$

$$\vec{\xi}^{(2)} = \frac{D_2}{D_1} \vec{\xi}^{(1)} - D_{12} \vec{\alpha}^{(1)} \left( \vec{\xi}^{(1)} \right)$$

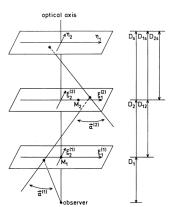
For two point mass lenses:

$$\vec{\alpha}^{(1)}\left(\vec{\xi}^{(1)}\right) = \frac{4GM_1}{c^2} \frac{\vec{\xi}^{(1)}}{|\vec{\xi}^{(1)}|^2},$$

$$\vec{\alpha}^{(2)}\left(\vec{\xi}^{(2)}\right) = \frac{4GM_2}{c^2} \frac{\vec{\xi}^{(2)} - \vec{\xi}_{M_2}^{(2)}}{|\vec{\xi}^{(2)} - \vec{\xi}_{M_2}^{(2)}|^2}$$

For two SIS lenses: 
$$\vec{\alpha}^{(1)} \left( \vec{\xi}^{(1)} \right) = \frac{4\pi\sigma_{v1}^2}{c^2} \frac{\vec{\xi}^{(1)}}{|\vec{\xi}^{(1)}|},$$

$$\vec{\alpha}^{(2)}\left(\vec{\xi}^{(2)}\right) = \frac{4\pi\sigma_{v2}^2}{c^2} \frac{\vec{\xi}^{(2)} - \vec{\xi}_{M_2}^{(2)}}{|\vec{\xi}^{(2)} - \vec{\xi}_{M_2}^{(2)}|}$$



Classification of the multiple deflection two point-mass gravitational lens models and application of catastrophe theory in lensing. Erdl, Helmut and Schneider, Peter https://ui.adsabs.harvard.edu/abs/1993A%26A...268..453E/abstract

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### case 2) these two galaxies are at different redshifts

For two SIS lenses: 
$$\vec{\alpha}^{(1)} \left( \vec{\xi}^{(1)} \right) = \frac{4\pi\sigma_{v1}^2}{c^2} \frac{\vec{\xi}^{(1)}}{|\vec{\xi}^{(1)}|},$$

$$\vec{\alpha}^{(1)}\left(\vec{\xi}^{(1)}\right) = \frac{4\pi\sigma_{v1}^2}{c^2} \frac{\vec{\xi}^{(1)}}{|\vec{\xi}^{(1)}|},$$

$$\vec{\alpha}^{(2)}\left(\vec{\xi}^{(2)}\right) = \frac{4\pi\sigma_{v2}^2}{c^2} \frac{\vec{\xi}^{(2)} - \vec{\xi}_{M_2}^{(2)}}{|\vec{\xi}^{(2)} - \vec{\xi}^{(2)}|}$$

Rearrange these equations:

$$\vec{y} = \vec{x} - m_1 \frac{\vec{x}}{|\vec{x}|} - m_2 \frac{\vec{t} - \vec{\rho}}{|\vec{t} - \vec{\rho}|}$$
 
$$\vec{t} = \vec{x} - m_1 \beta \frac{\vec{x}}{|\vec{x}|}$$
 where  $m_1 = D_1$   $\theta_1$ ,  $m_2 = 0$ 

where 
$$m_1 = D_{1s}\theta_1$$
,  $m_2 = D_{2s}\theta_2$ ,  $\beta = D_{12}/(D_2D_{1s})$ ,  $\vec{\xi}^{(2)} = D_2\vec{t}$ ,  $\vec{\xi}^{(2)}_{M_2} = D_2\vec{\rho}$ ,  $\vec{\xi}^{(1)} = D_1\vec{x}$ ,  $\vec{\eta} = D_s\vec{y}$ ,

$$\theta_i = \frac{4\pi\sigma_{vi}^2}{c^2} = 1^{"}.6\left(\frac{\sigma_{vi}}{200km/s}\right)^2$$

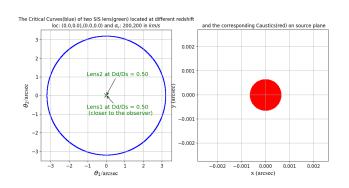
Thus the Jacobian matrix of this two lenses system:

$$A = \frac{\partial \vec{y}}{\partial \vec{x}}$$

$$\mu(\vec{x}) = \frac{1}{\det A(\vec{x})}$$

### case 2) these two galaxies are at different redshifts

▶ Obesever, Lens1, Lens2 are on a line, e.g.  $\vec{\rho}=(0,0)$ , and  $D_1=1, D_2=1, D_s=2$ 



#### GravLens

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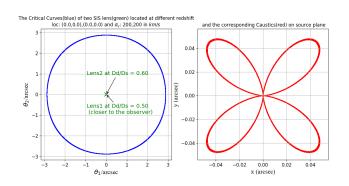
History

Gravitational

#### Related Questions

### case 2) these two galaxies are at different redshifts

▶ Obesever, Lens1, Lens2 are on a line, e.g.  $\vec{\rho}=(0,0)$ , and  $D_1=1, D_2=1.2, D_s=2$ 



#### GravLens

#### Renkun Kuang

History

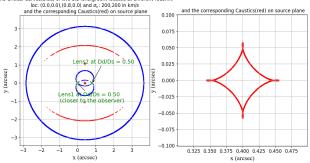
Gravitational

#### Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (0.8, 0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



GravLens

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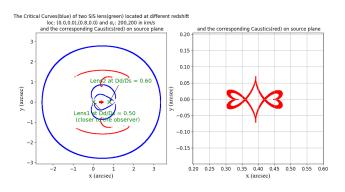
Histor

Gravitational Lensing Theories

Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (0.8,0)$$
, and  $D_1 = 1, D_2 = 1.2, D_s = 2$ 



#### GravLens

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Histor

Gravitational

### Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (0.8,0)$$
, and  $D_1 = 1, D_2 = 1.6, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (0.8.0.0) and a.: 200.200 in km/s and the corresponding Caustics(red) on source plane and the corresponding Caustics(red) on source plane 0.6 0.4 ens2 at Dd/Ds = 0.80 0.2 y (arcsec) y (arcsec) 0.0 -0.2Lens1 at Dd/Ds (closer to the observer -1-0.4-0.6 -0.6 -0.4 -0.2 0.0 0.2 0.4 0.6 0.8 x (arcsec) x (arcsec)

#### GravLens

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Histor

Gravitational Lensing Theories

### Related Questions

## case 2) these two galaxies are at different redshifts

 $ightharpoonup ec{
ho} = (0.8,0)$ , and  $D_1 = 1, D_2 = 1.8, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (0.8.0.0) and a.: 200.200 in km/s and the corresponding Caustics(red) on source plane and the corresponding Caustics(red) on source plane 2.0 1.5 1.0 0.1 Lens2 at Dd/Ds = 0.00 y (arcsec) (arcsec) 0.0 Lens1 at Dd/Ds = 0 -0.1-0.5(closer to the observer -0.2-1.0-1.5 -0.3-1.5 -1.0 -0.50.0 0.5 1.0 1.5 2.0 -0.10.0 0.1 0.3 0.4 0.5 x (arcsec) x (arcsec)

### GravLens

### Renkun Kuang

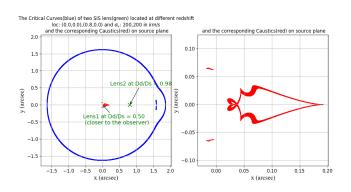
Histor

Gravitational
Lensing Theories

### Related Questions

## case 2) these two galaxies are at different redshifts

 $ightharpoonup ec{
ho} = (0.8, 0)$ , and  $D_1 = 1, D_2 = 1.96, D_s = 2$ 



GravLens

Renkun Kuang

Histor

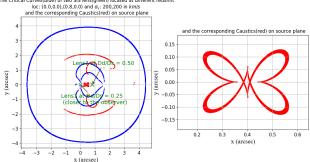
Gravitational Lensing Theories

Related Questions

## case 2) these two galaxies are at different redshifts

$$ightarrow$$
  $ec{
ho} = (0.8,0)$ , and  $D_1 = 0.5, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (0.8.0.0) and a.: 200.200 in km/s



#### Gravl ens

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## case 2) these two galaxies are at different redshifts

$$ightarrow$$
  $ec{
ho} = (0.8,0)$ , and  $D_1 = 0.2, D_2 = 1.6, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0,0.0),(0.8,0.0) and σ<sub>v</sub>: 200,200 in km/s and the corresponding Caustics(red) on source plane and the corresponding Caustics(red) on source plane 3 1.5 1.0 Lens2 at Dd/Ds = 0.800.5 y (arcsec) y (arcsec) × 0.0 -1-0.5-1.0-1.5 -3 -1.0-0.5 0.0 0.5 1.0 1.5 2.0 x (arcsec) x (arcsec)

#### GravLens

### Renkun Kuang

Histor

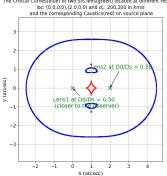
Gravitational

### Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(2.0.0.0) and \u03c3.: 200.200 in km/s



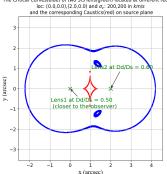
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### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1.2, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(2.0.0.0) and \u03c3.: 200.200 in km/s and the corresponding Caustics(red) on source plane



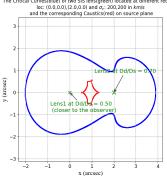
#### Gravl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1.4, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(2.0.0.0) and \u03c3.: 200.200 in km/s



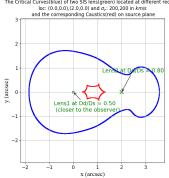
#### GravLens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1.6, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(2.0.0.0) and \u03c3.: 200.200 in km/s



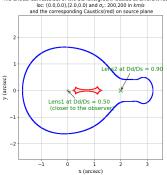
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1.8, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(2.0.0.0) and \u03c3.: 200.200 in km/s and the corresponding Caustics(red) on source plane



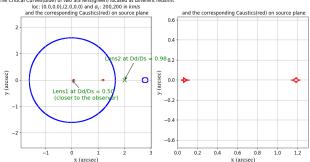
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 1, D_2 = 1.96, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



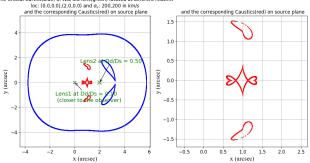
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 0.2, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



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History

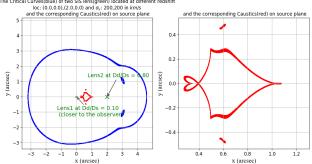
Gravitational

### Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 0.2, D_2 = 1.6, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



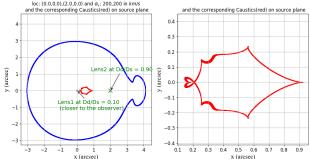
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (2,0)$$
, and  $D_1 = 0.2, D_2 = 1.8, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (2.0.0.0) and a.: 200.200 in km/s and the corresponding Caustics(red) on source plane



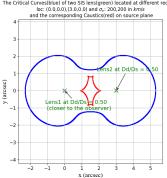
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(3.0.0.0) and \u03c3.: 200.200 in km/s



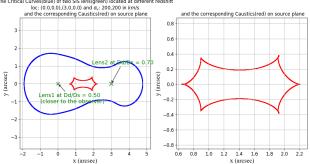
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 1, D_2 = 1.4, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (3.0.0.0) and a.: 200.200 in km/s

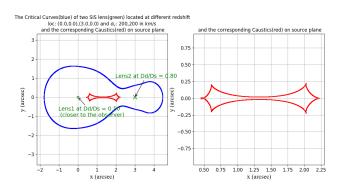


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## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 1, D_2 = 1.6, D_s = 2$ 



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History

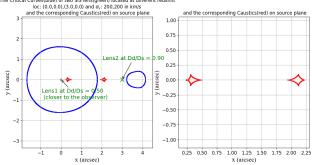
Gravitational

Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 1, D_2 = 1.8, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



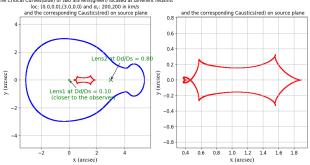
#### Grayl ens

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## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 0.2, D_2 = 1.6, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0) (3.0.0.0) and a.: 200.200 in km/s



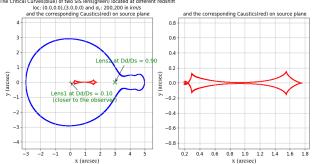
#### Grayl ens

### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 0.2, D_2 = 1.8, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



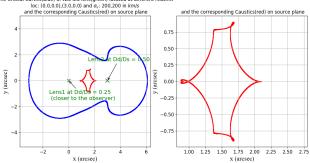
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## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (3,0)$$
, and  $D_1 = 0.5, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift



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Histor

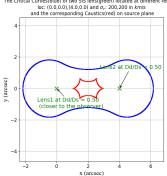
Gravitational

### Related Questions

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (4,0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(4.0.0.0) and \u03c3.: 200.200 in km/s



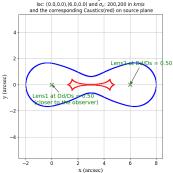
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### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (6,0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(6.0.0.0) and \u03c3.: 200.200 in km/s



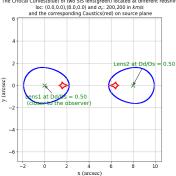
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### Renkun Kuang

## case 2) these two galaxies are at different redshifts

$$ightharpoonup ec{
ho} = (8,0)$$
, and  $D_1 = 1, D_2 = 1, D_s = 2$ 

The Critical Curves(blue) of two SIS lens(green) located at different redshift loc: (0.0.0.0),(8.0.0.0) and \u03c3.: 200.200 in km/s



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## **Gravitational Lensing Applications**

History

Gravitational Lensing Theories

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- ► Cosmic telescopes: distant, faint objects observation
- 2-d mass distribution of lenses, dark matter
- Hubble constant, cosmological constant, density parameter
- .....

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History

Gravitational Lensing Theories

Applications

# The End, Thanks!

- https://lacosmo.com/DeflectionOfLight/index.html
- ► The Mathematical Theory of Relativity, Arthur Stanley Eddington (P101)
- http://web.mit.edu/6.055/old/S2009/notes/bending-oflight.pdf
- https://www.mathpages.com/rr/s6-03/6-03.htm Gravitation and Spacetime, Hans C. Ohanian, Remo Ruffini. – 3rd ed
- https://en.wikipedia.org/wiki/Gauss%27s\_law\_for\_gravity
- Lectures On Gravitational Lensing, Narayan & Bartelmann
- https://www.cfa.harvard.edu/ dfabricant/huchra/ay202/lectures/lecture12.pdf
- https://web.stanford.edu/ oas/SI/SRGR/notes/ SchwarzschildSolution.pdf
- · .....

## Appendix-1, Newtonian prediction

$$\phi(b,z) = -\frac{GM}{(b^2 + z^2)^{1/2}}$$

$$\alpha = \frac{v_b}{c} = \frac{1}{c} \int \frac{d\Phi}{db} dt = \frac{1}{c^2} \int \frac{d\Phi}{db} dl$$

$$\approx \frac{1}{c^2} \int \frac{d\Phi}{db} dz$$

$$= \frac{GMb}{c^2} \int \frac{dz}{(b^2 + z^2)^{3/2}}$$

$$= \frac{GMb}{c^2} \left[ \frac{z}{b^2 \sqrt{b^2 + z^2}} \Big|_{-\infty}^{+\infty} \right]$$

$$= \frac{2GM}{c^2b}$$

b: impact factor

Gravitational potential simplify the calculation by integrating not along the deflected ray but along z axis

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History

Gravitational
Lensing Theories

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