

DARK MATTER THEORY II

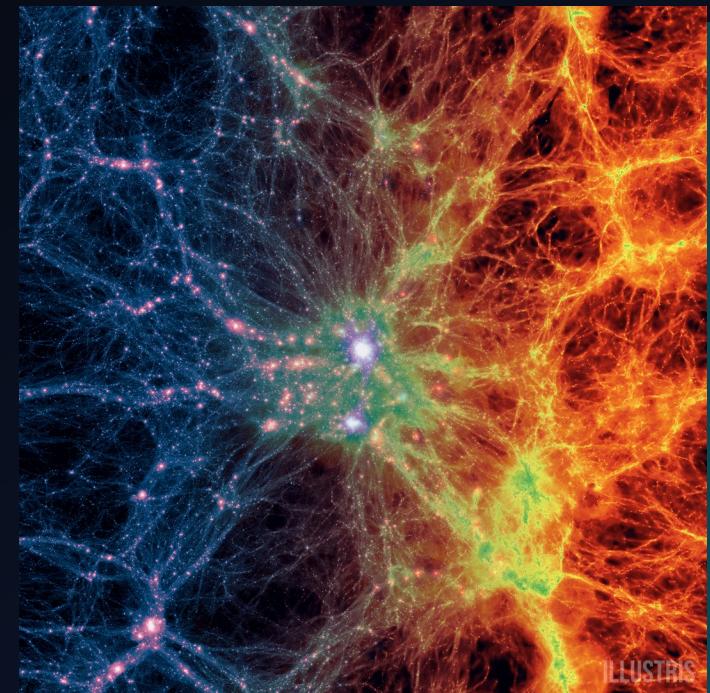
REBECCA LEANE
SLAC NATIONAL ACCELERATOR LABORATORY

SLAC SSI LECTURES
AUG 15-16th 2022

SLAC

Today's Outline

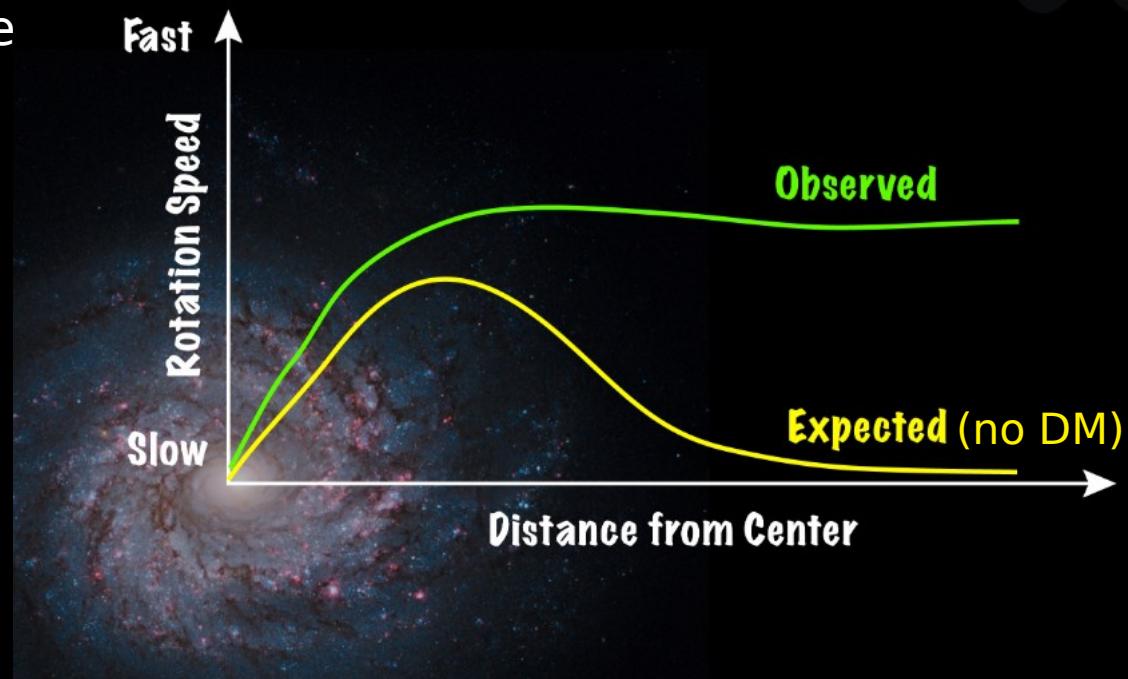
- WIMP recap and status
 - Searches and Constraints
- Newer Production Mechanisms
 - SIMPs, co-SIMPs
 - Freeze-in
- Axions
- MACHOs
 - Primordial Black Holes
- Outlook



Illustris Simulation

Recap: what we know about Dark Matter

- Makes up about 85% of matter in the Universe (~5x more mass than SM)
- Doesn't interact too much with light (or itself)
- Provides large scale structure of the Universe
- Structure tells us non-relativistic
- Forms halos around Galaxies
- Stable on cosmological timescales

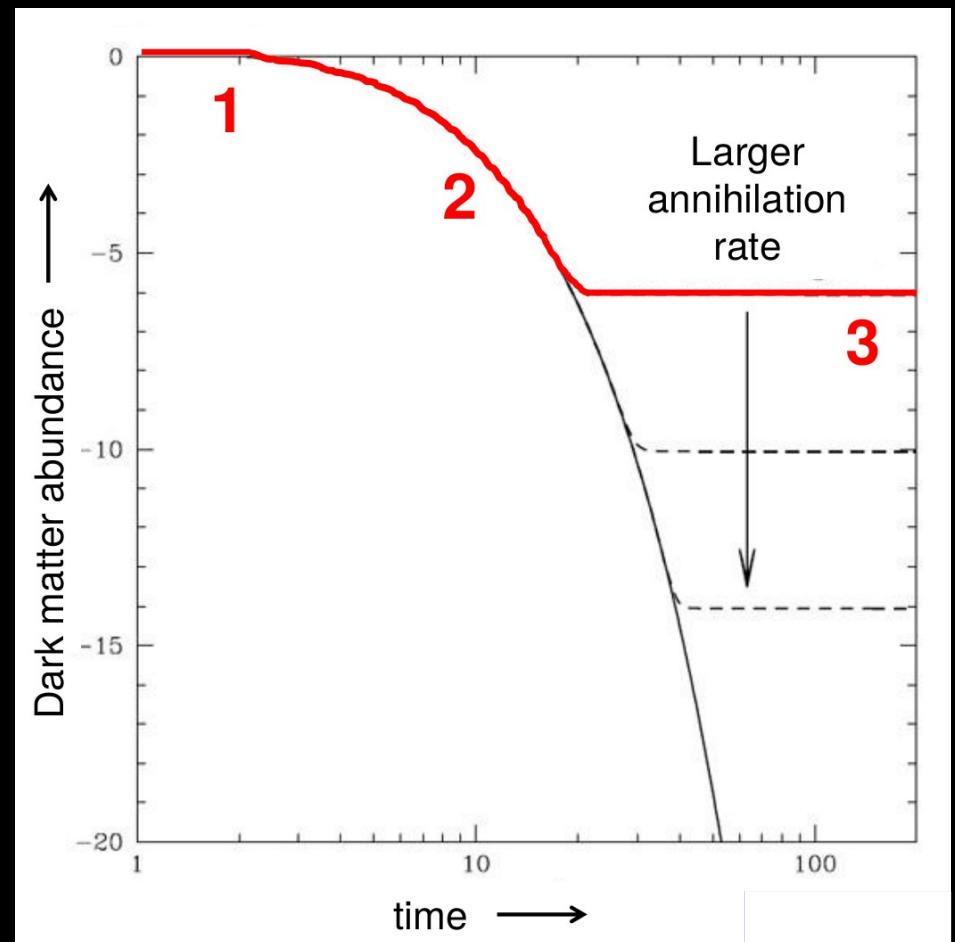


Cannot be explained by any known particles!

Recap: Dark Matter Abundance: WIMPs

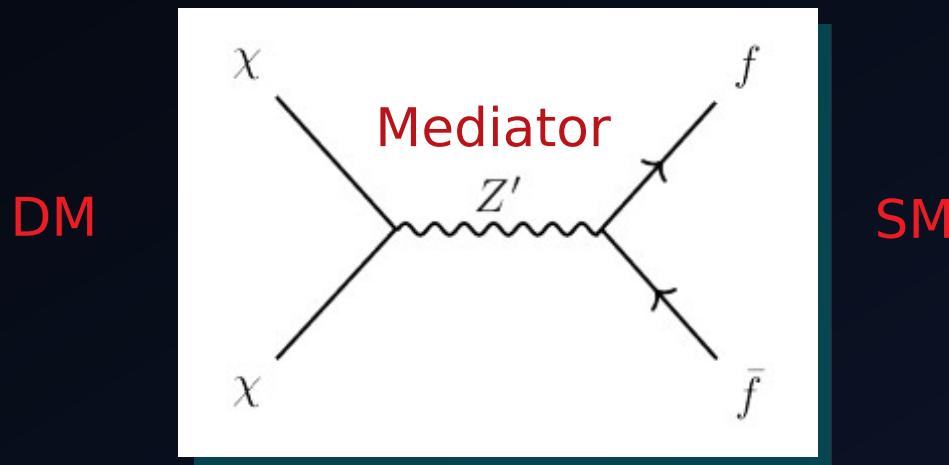
- 1)** Thermal equilibrium:
 $\text{DM} + \text{DM} \Rightarrow \text{visible particles}$
 $\text{Visible particles} \Rightarrow \text{DM} + \text{DM}$
- 2)** Universe cools, only
 $\text{DM} + \text{DM} \Rightarrow \text{visible particles}$
- 3)** Universe expands too fast.
No more annihilations.
DM abundance is set.

Predicts a particular annihilation rate
for dark matter.



Recap: Dark Sectors

- Standard Model (SM) mediated WIMP processes, e.g. with Z bosons, extremely constrained
- Need new fields beyond the SM!
- These can mediate interactions between the dark sector, and the SM



Dark Sectors

- Need new fields beyond the SM, consider extensions:
 - New scalar, pseudoscalar, fermion, or vector
- Symmetries of SM restrict how these fields interact with the SM

$$-\lambda_{HS}(S^\dagger S)(H^\dagger H)$$

Higgs Portal

$$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Axion Portal

$$-\frac{\epsilon}{2} B_{\mu\nu} Z'^{\mu\nu}$$

Hypercharge Portal

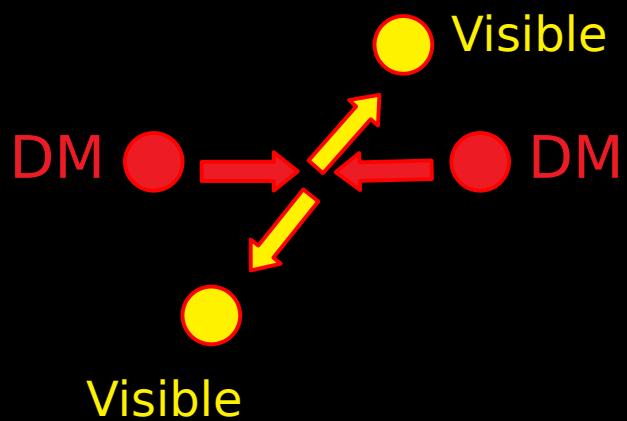
$$y_n LHN$$

Fermion Portal

Generally easiest to parameterize searches by generic masses and couplings

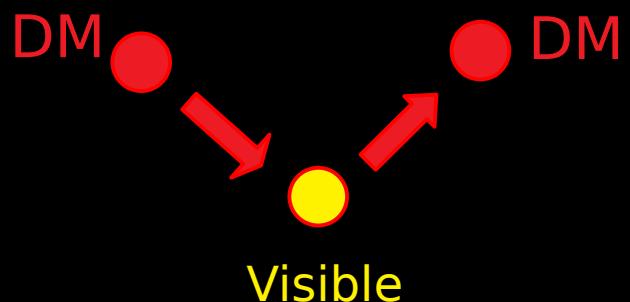
WIMP Dark Matter Search Program

Annihilation



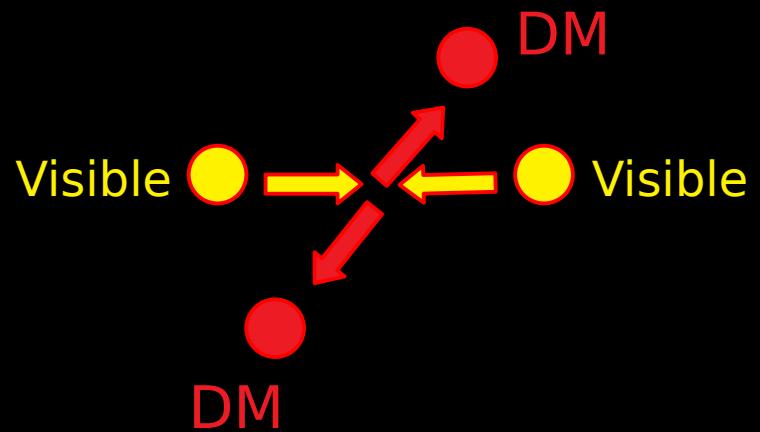
Astrophysics

Scattering



Direct detection
+ Astrophysics

Production

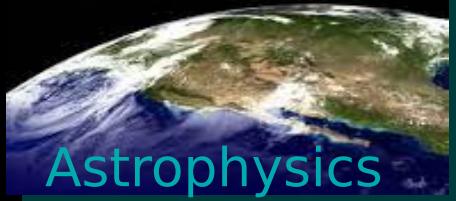
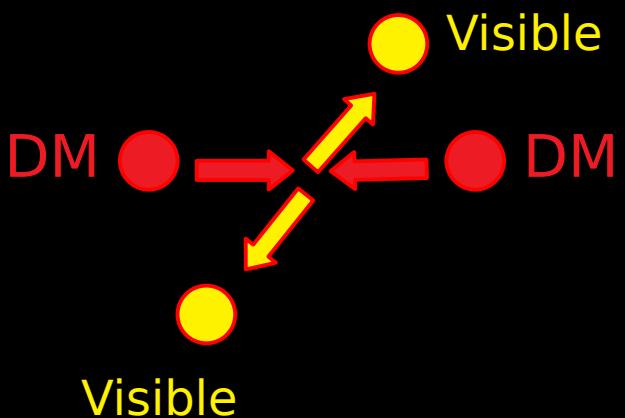


Colliders
+ Astrophysics

WIMP Dark Matter Search Program

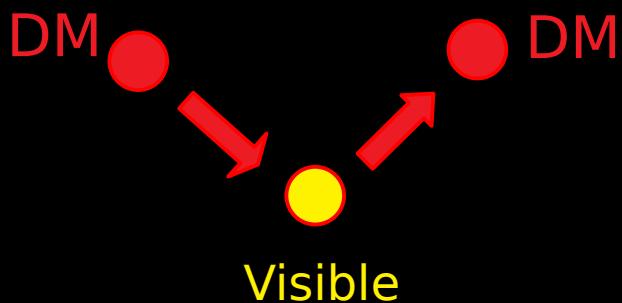
M.E. Mozani's lecture yesterday

Annihilation



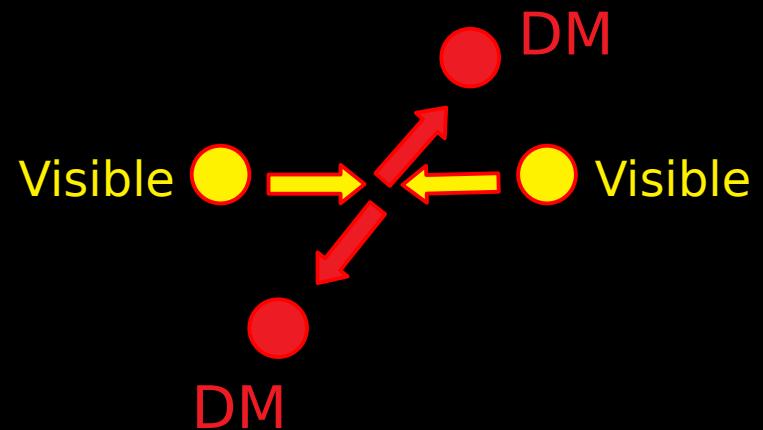
Astrophysics

Scattering



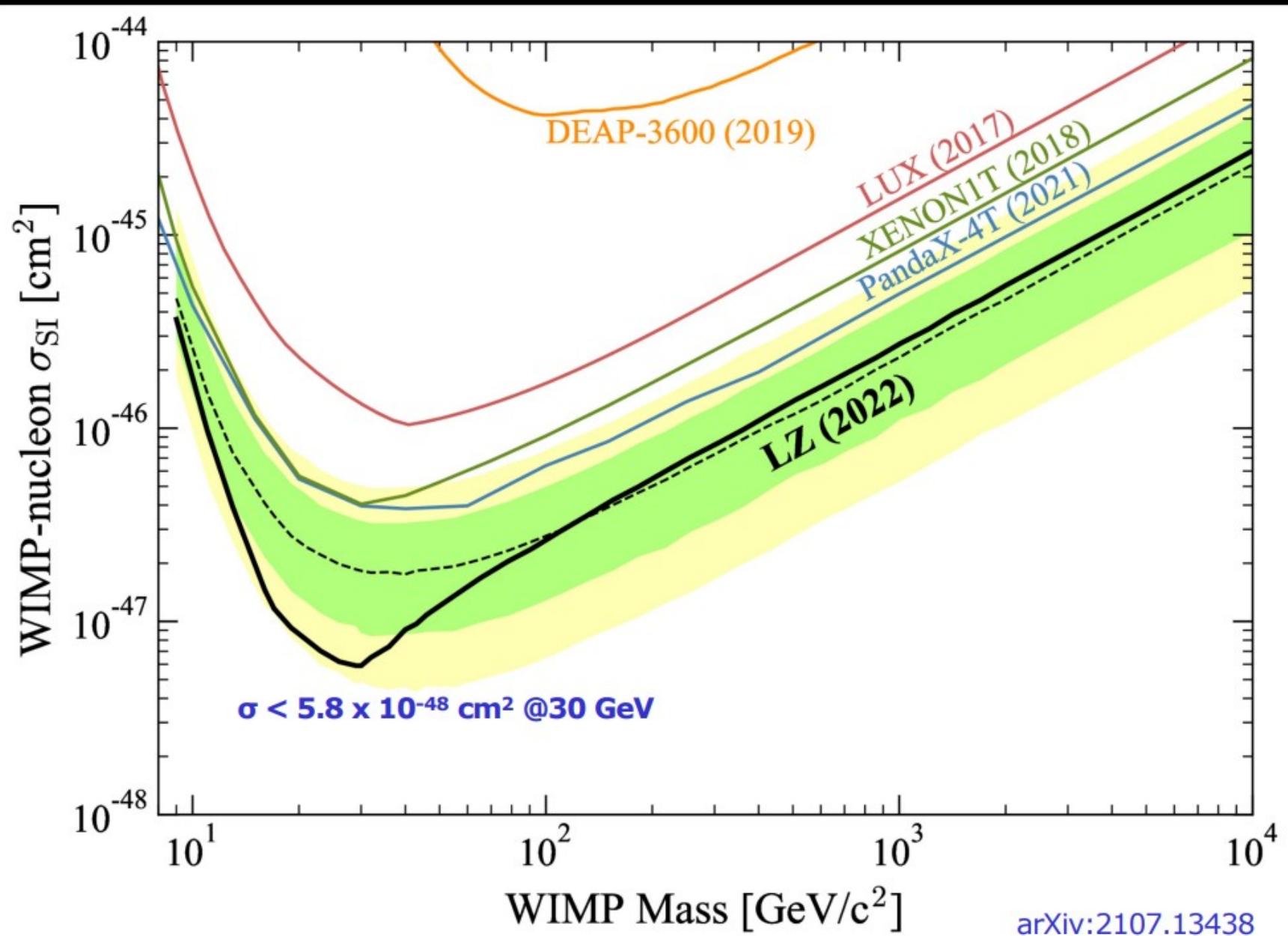
Direct detection
+ Astrophysics

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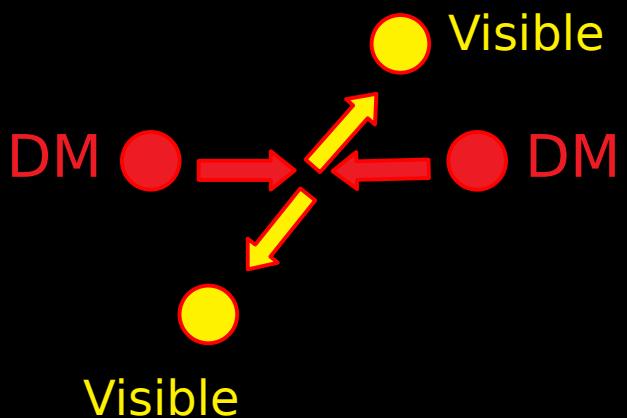
Colliders
+ Astrophysics

World-leading constraint on WIMPs



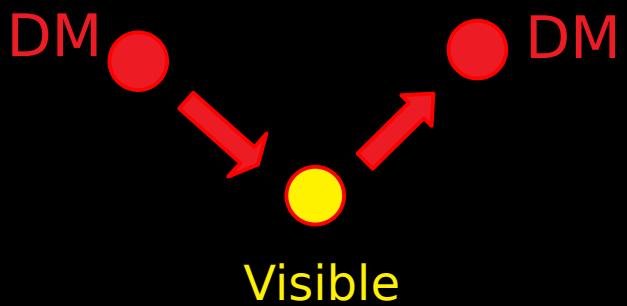
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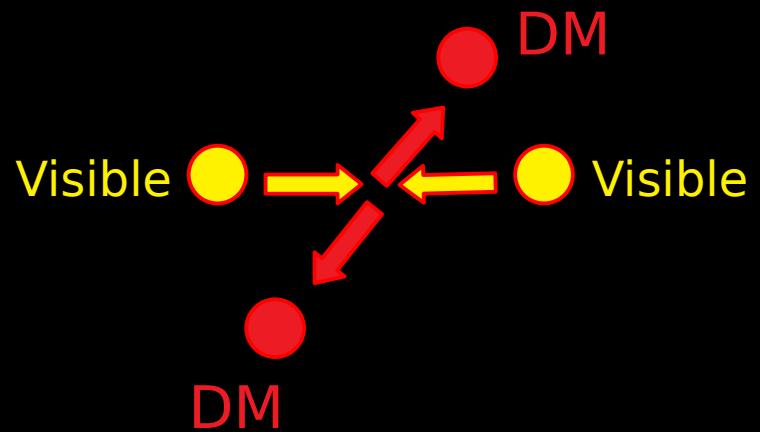
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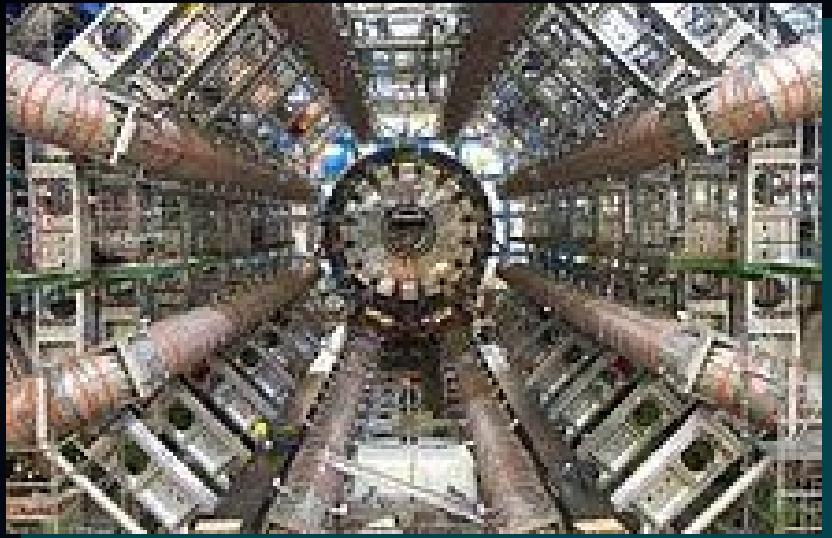
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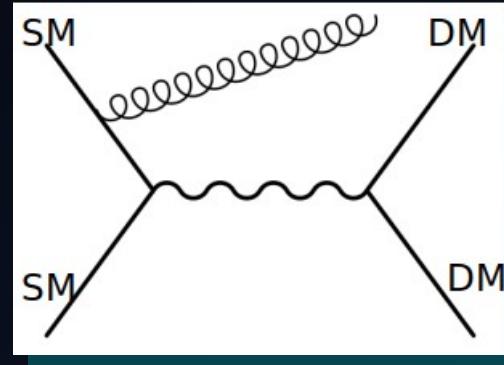
Colliders
+ Astrophysics

Dark matter at colliders



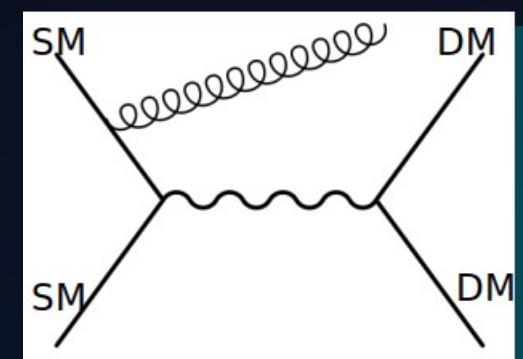
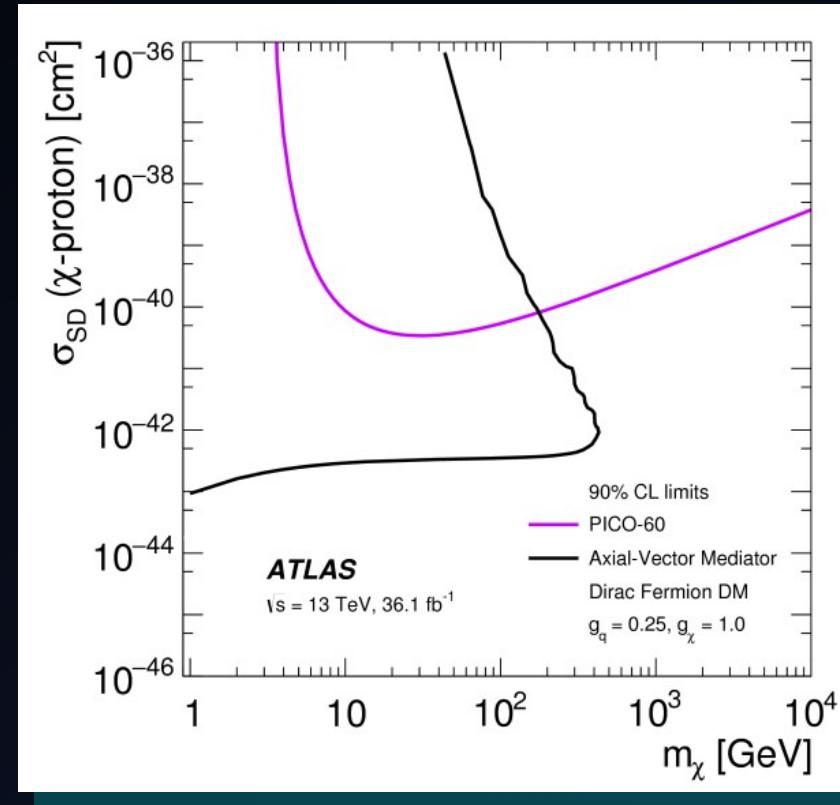
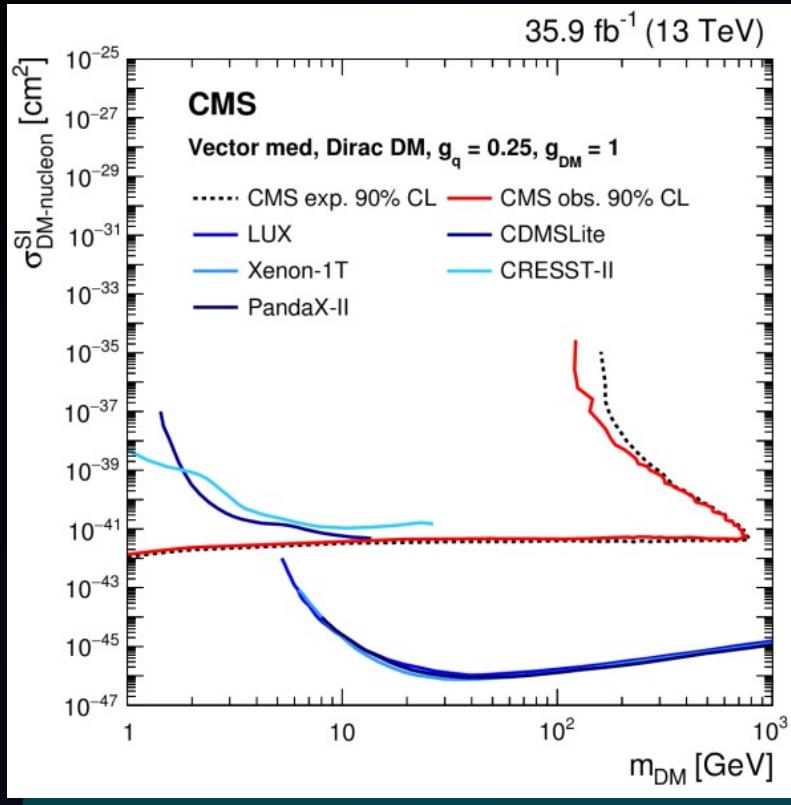
ATLAS

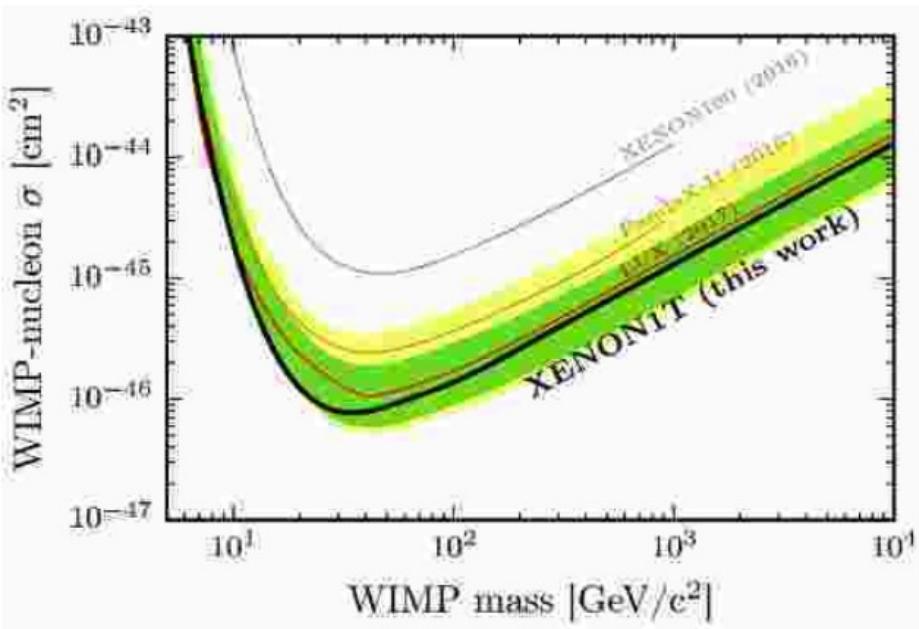
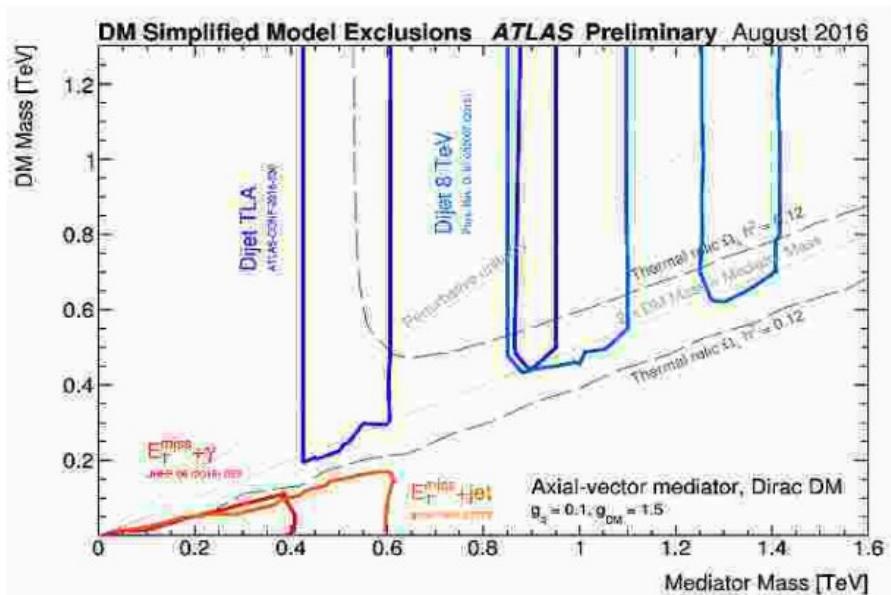
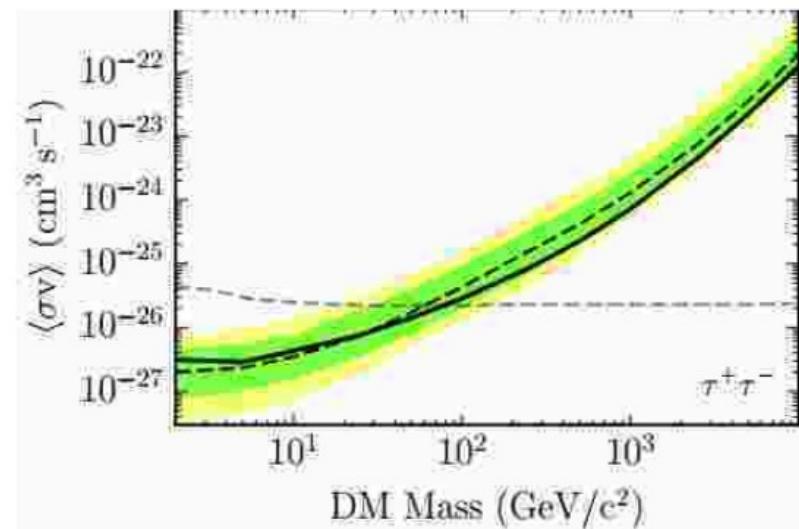
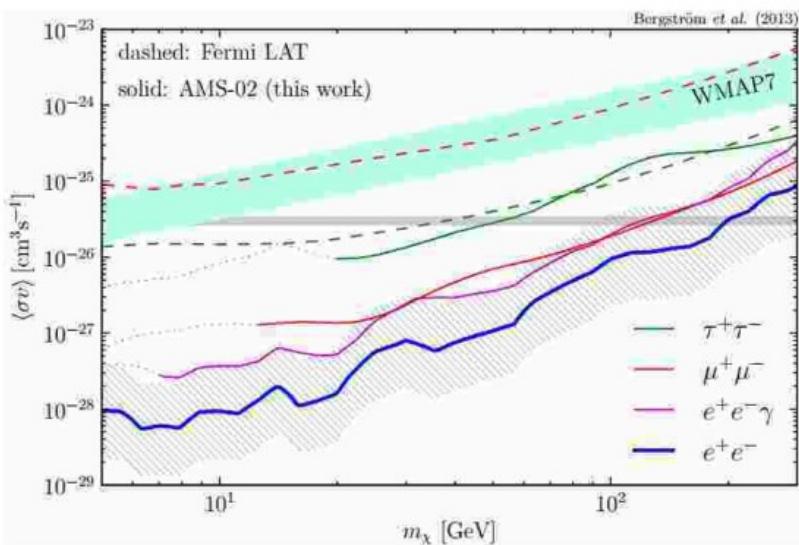
- Missing momentum searches test weak-scale dark matter

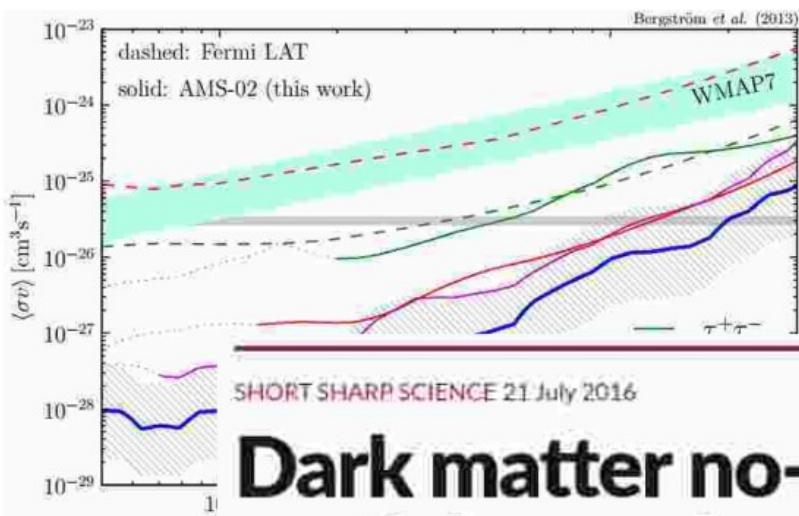


- Strong limits for mediator searches directly as well

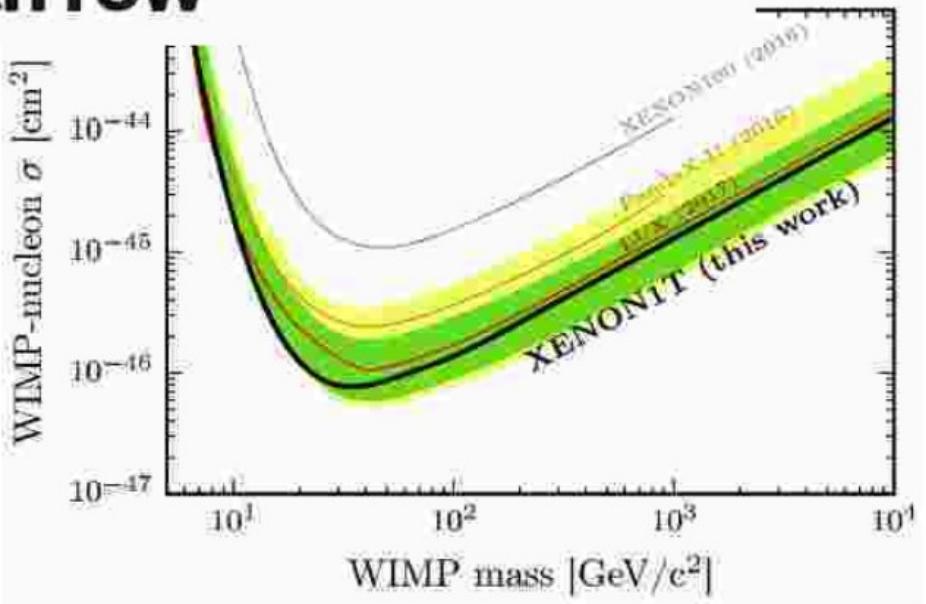
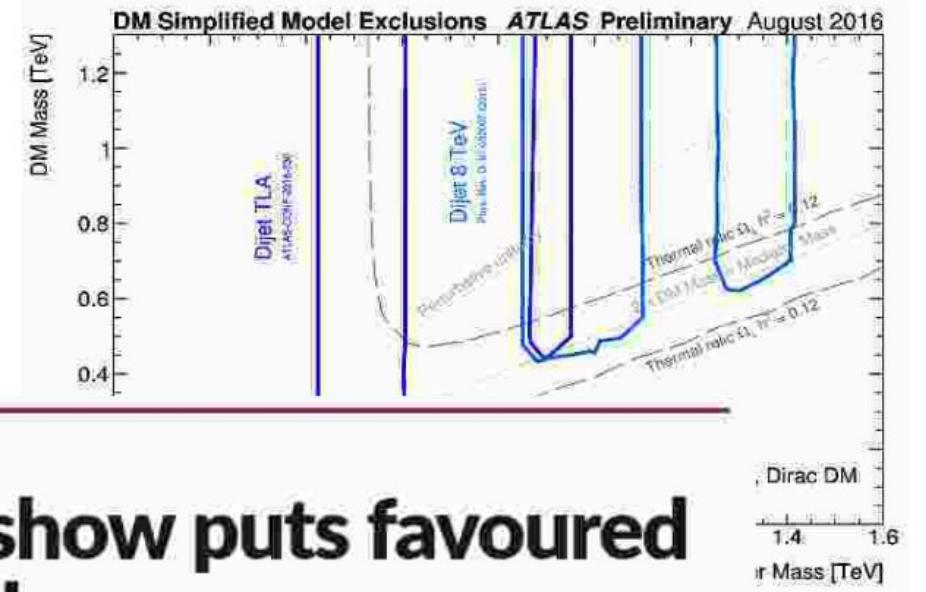
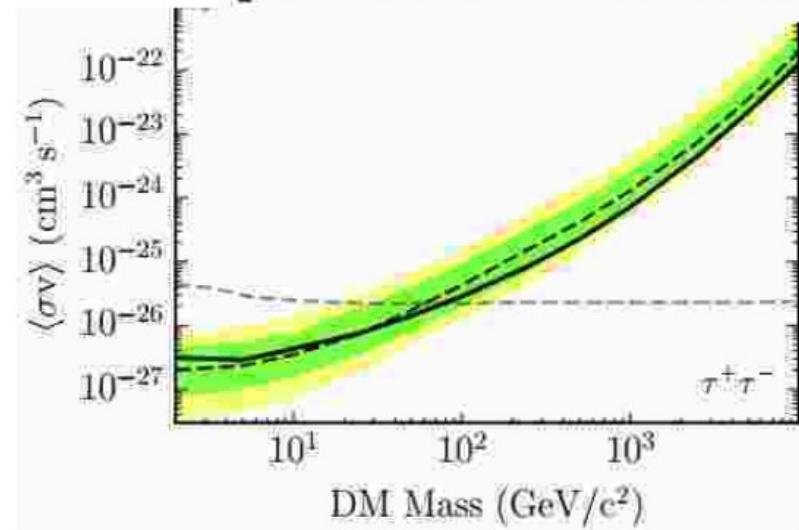
Mono-jet searches

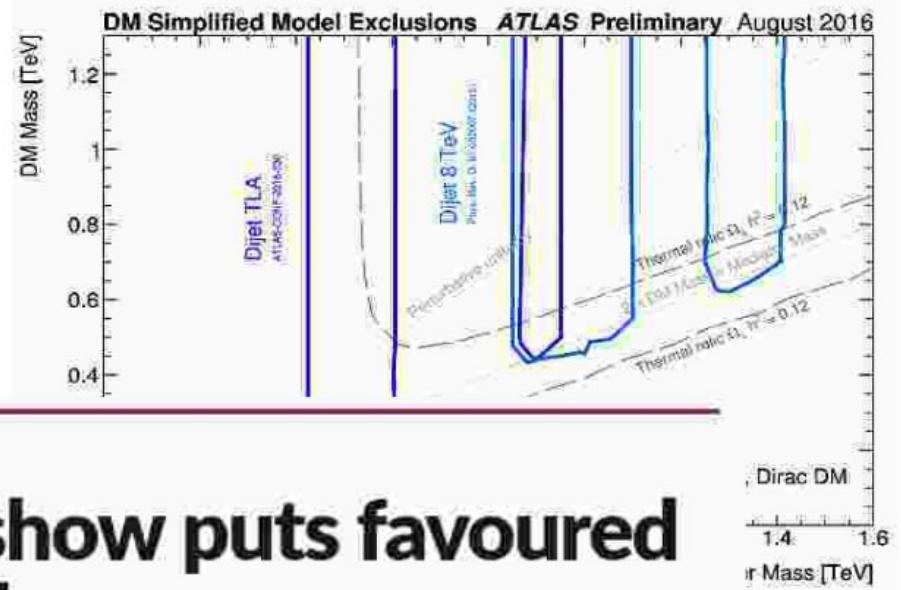
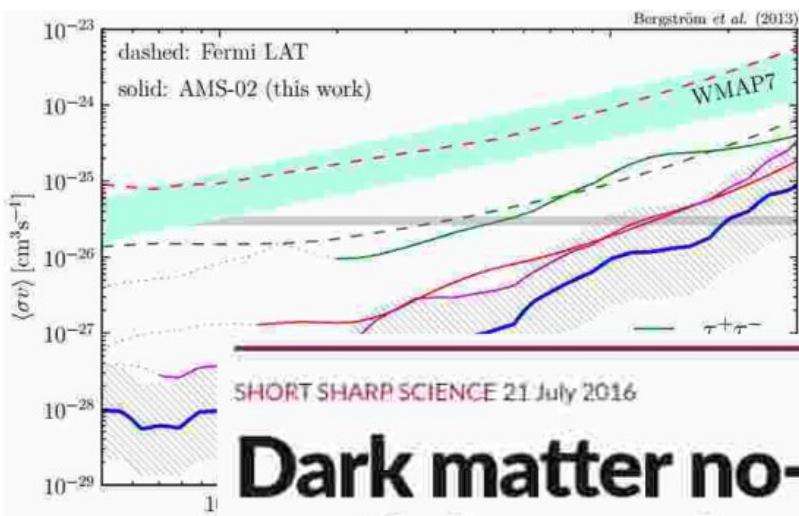




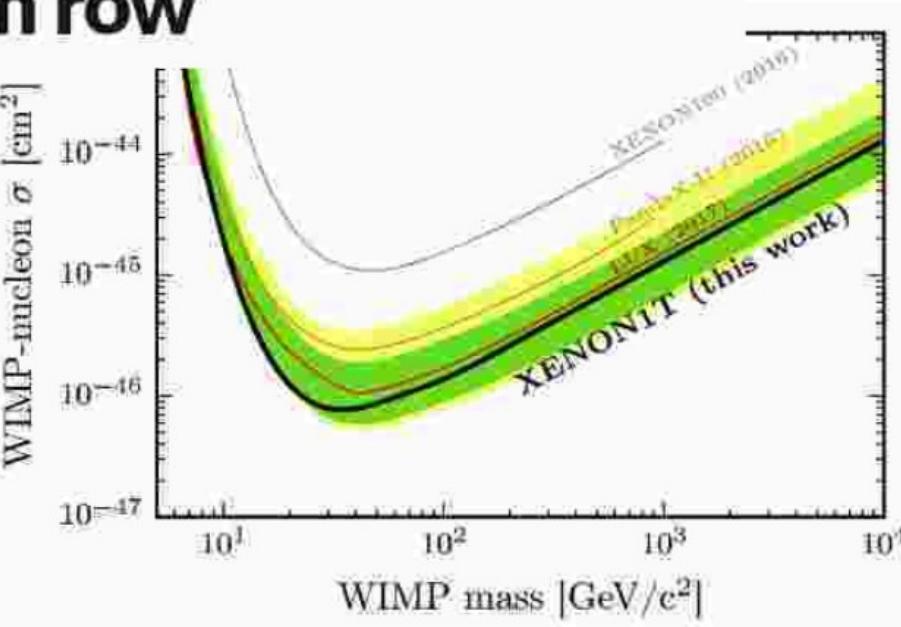
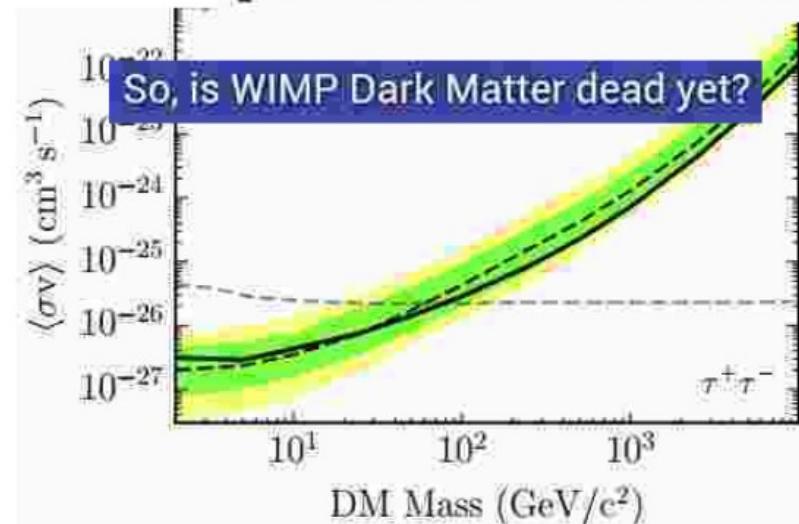


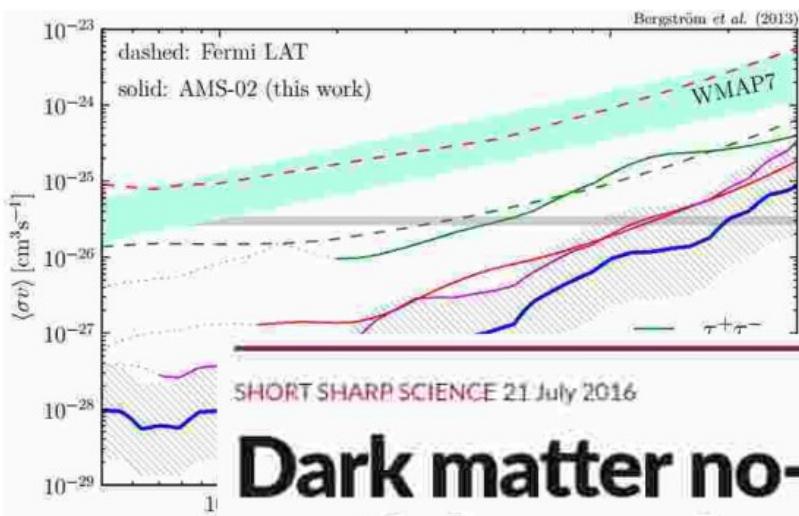
Dark matter no-show puts favoured particles on death row



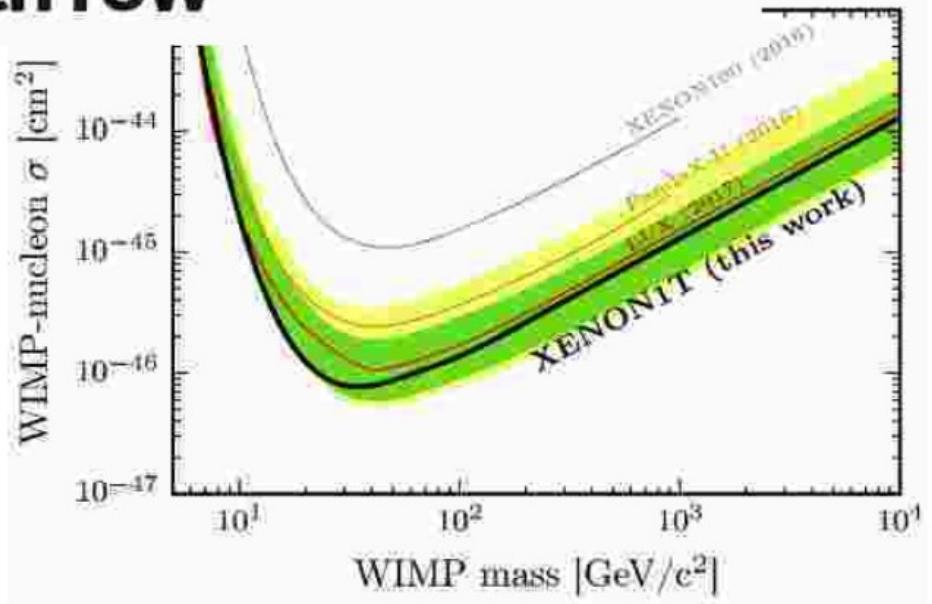
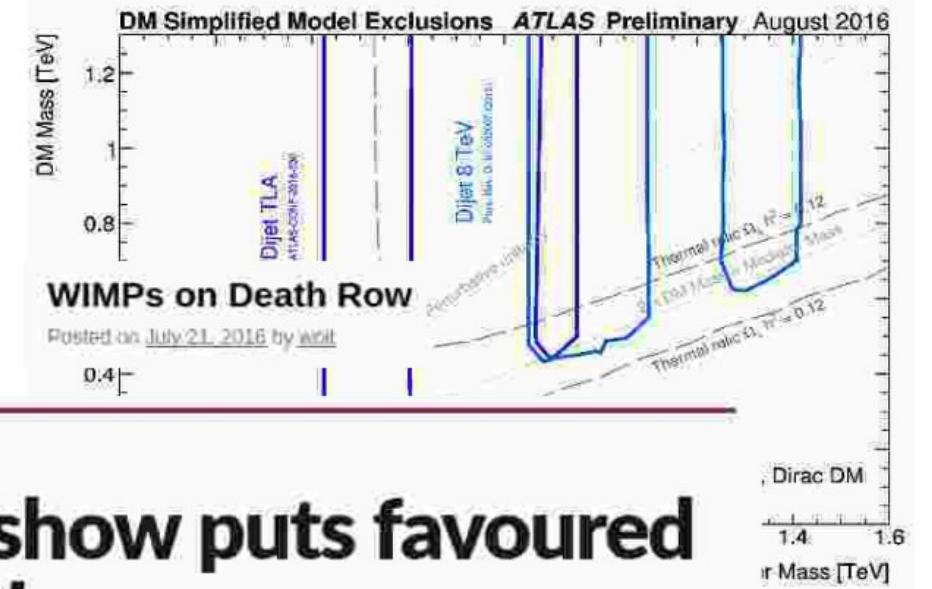
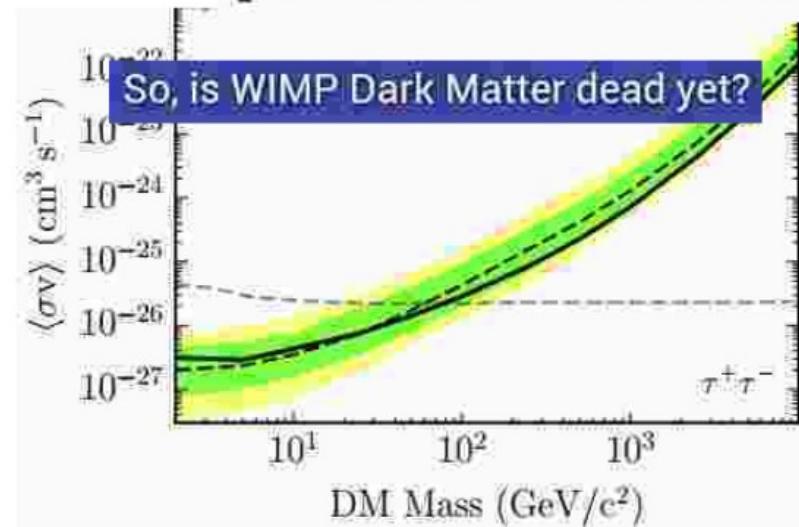


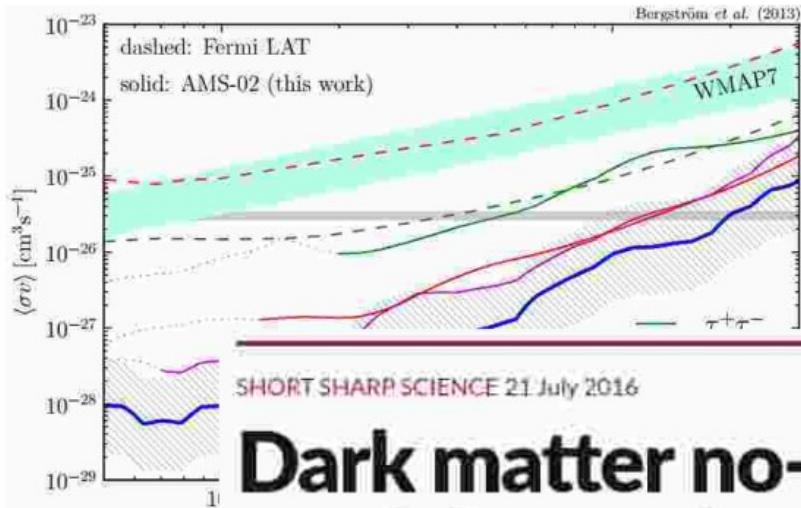
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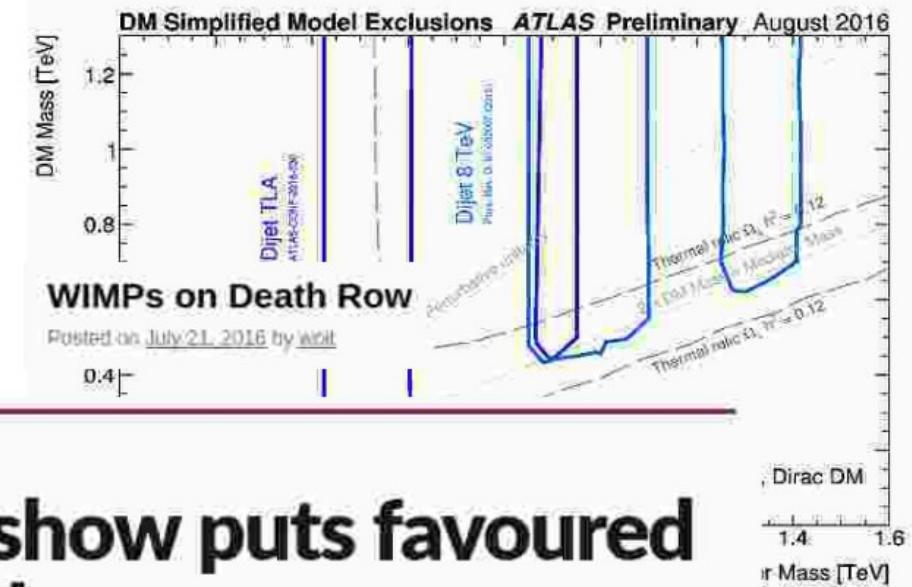
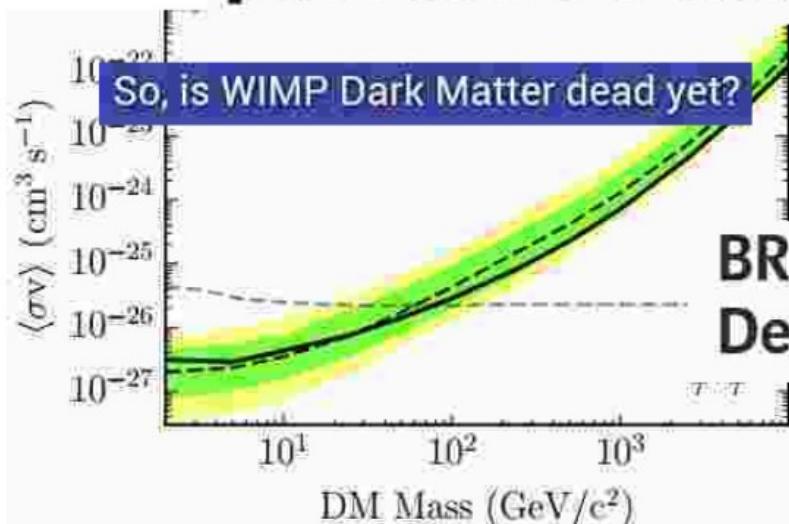


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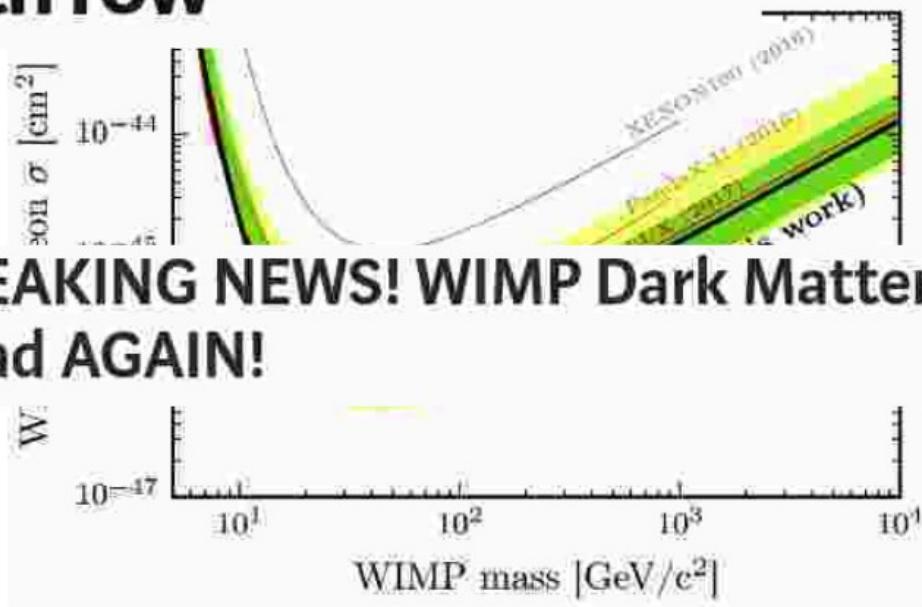




Dark matter no-show puts favoured particles on death row



BREAKING NEWS! WIMP Dark Matter Dead AGAIN!

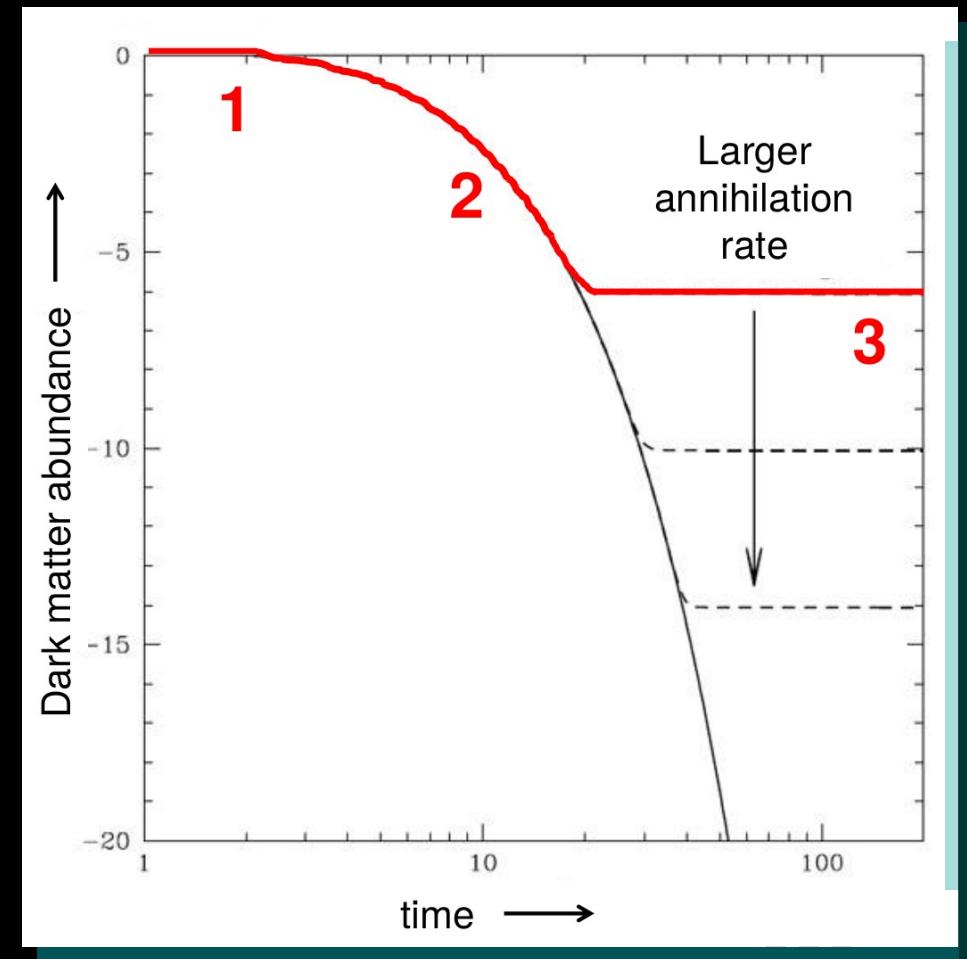


Only annihilation provides the target rate

Scattering (direct detection) and production (colliders):

No well-defined scale because only some aspects of the interaction are being considered.

Those constraints tell you which WIMP models are excluded, but nothing about the remaining possibilities.



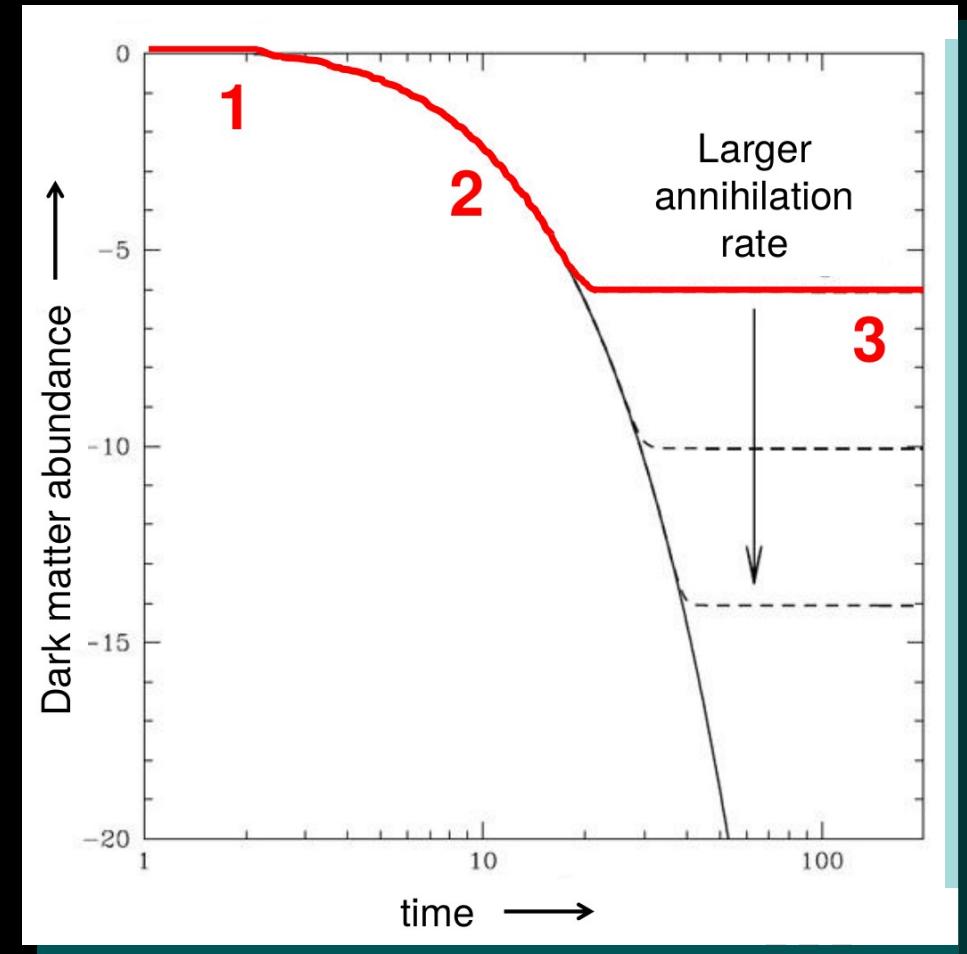
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Scattering (direct detection) and production (colliders):

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The most decisive way to test thermal WIMPs is through their annihilation products, as this exactly goes to their most fundamental feature: being annihilation relics, which sets a well-defined scale for the total cross section.



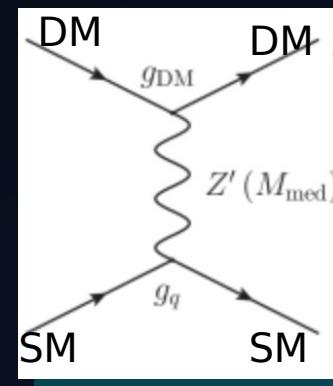
Hidden dark sectors:

- New fields that have **small couplings** to the SM
- Naturally occurs in many model scenarios; symmetries, or Higgs may have mass coupling dependence, etc
- If zero coupling to the SM, can't detect the DM!
- Need to **interact at least very weakly** with the SM to probe at particle experiments

Searching for hidden dark sectors

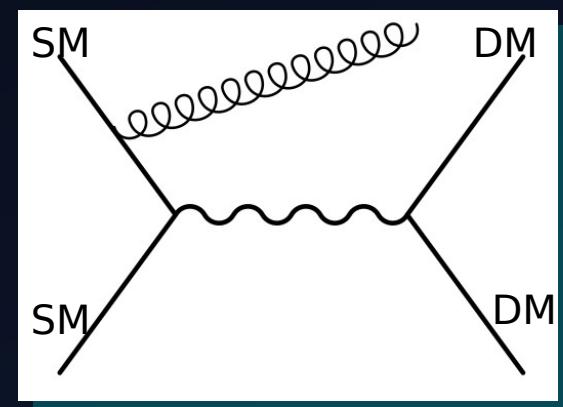
- Small coupling to SM minimizes constraints from scattering or production
 - Difficult but not impossible to probe!

Scattering



Suppressed :(
Rebecca Leane

Production

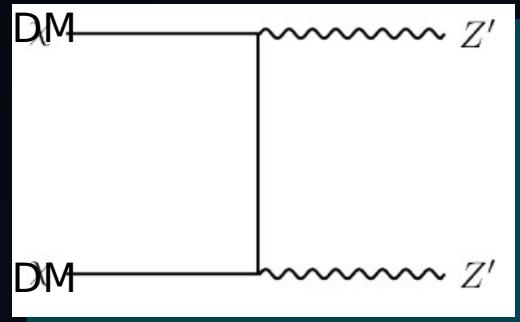


Suppressed :(

Searching for hidden dark sectors

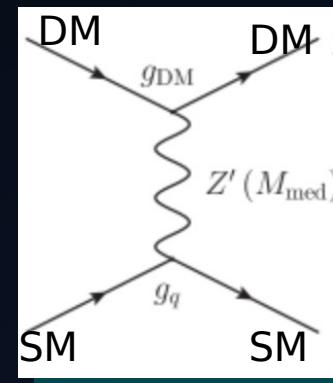
- Small coupling to SM minimizes constraints from scattering or production
 - Difficult but not impossible to probe!
- Large signals still possible for annihilation!
 - Cross section independent of SM coupling if particles produced on shell
 - Only then multiplied by branching fraction

Annihilation



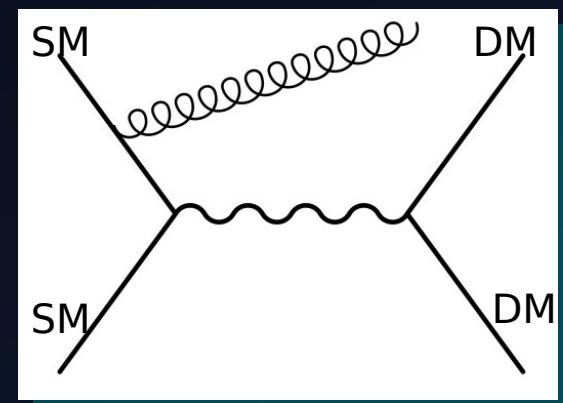
Not suppressed :)

Scattering



Suppressed :(
Rebecca Leane

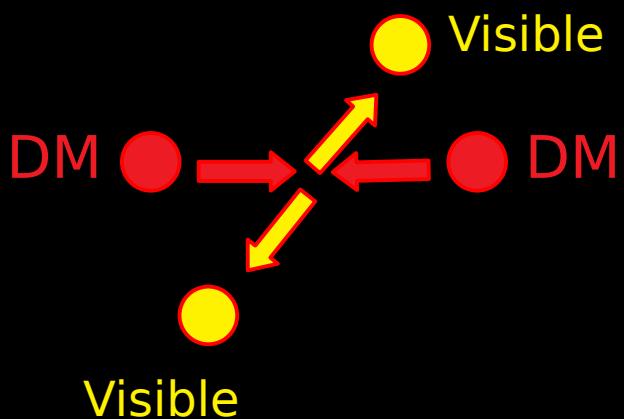
Production



Suppressed :(

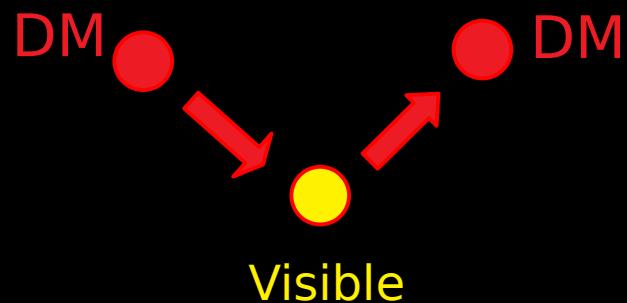
WIMP Dark Matter Search Program

Annihilation



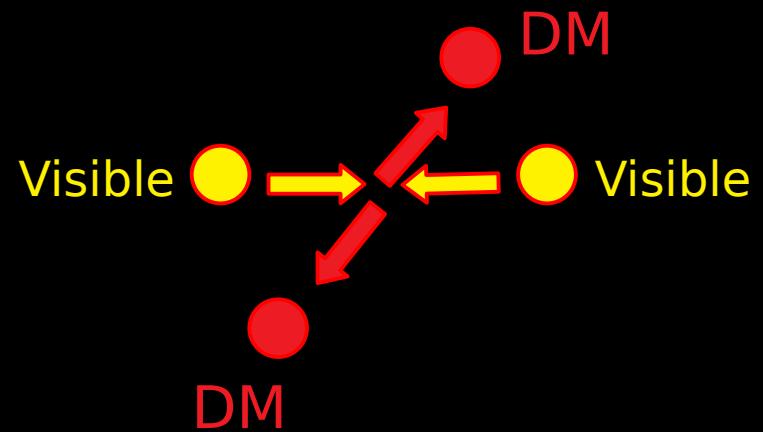
Astrophysics

Scattering



Direct detection
+ Astrophysics

Production



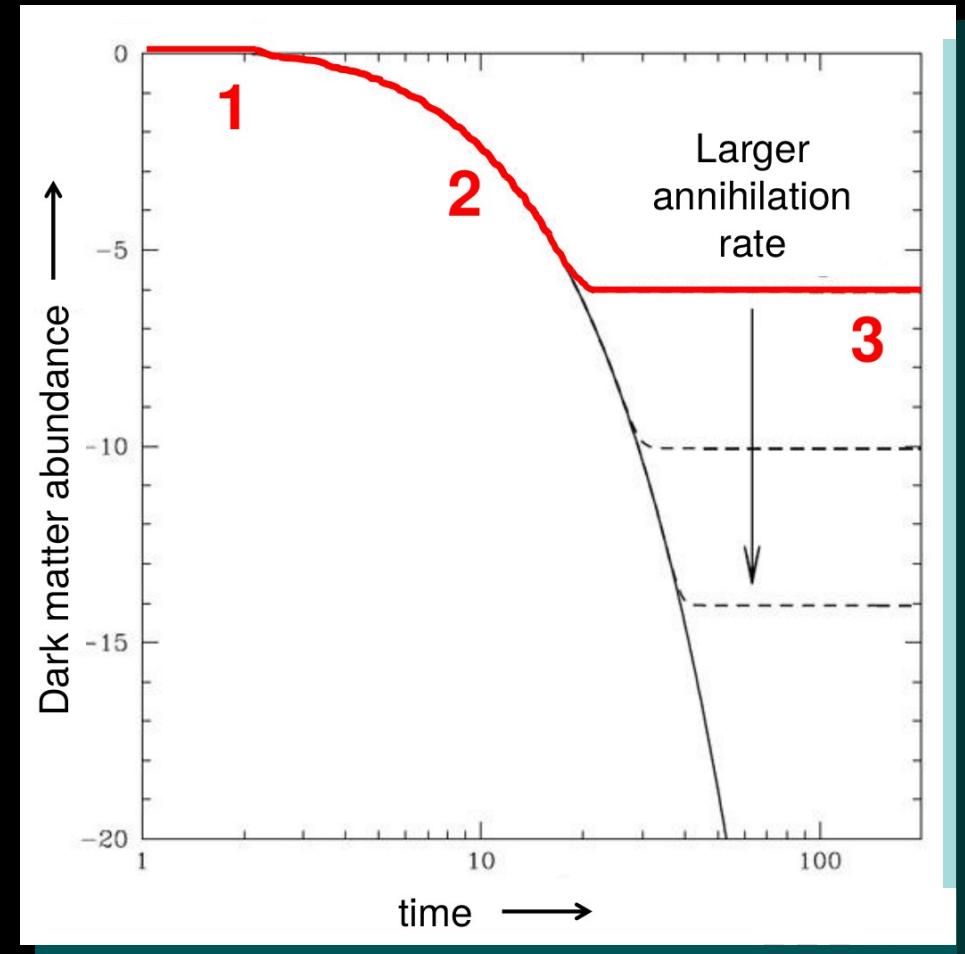
Colliders
+ Astrophysics

INDIRECT DETECTION

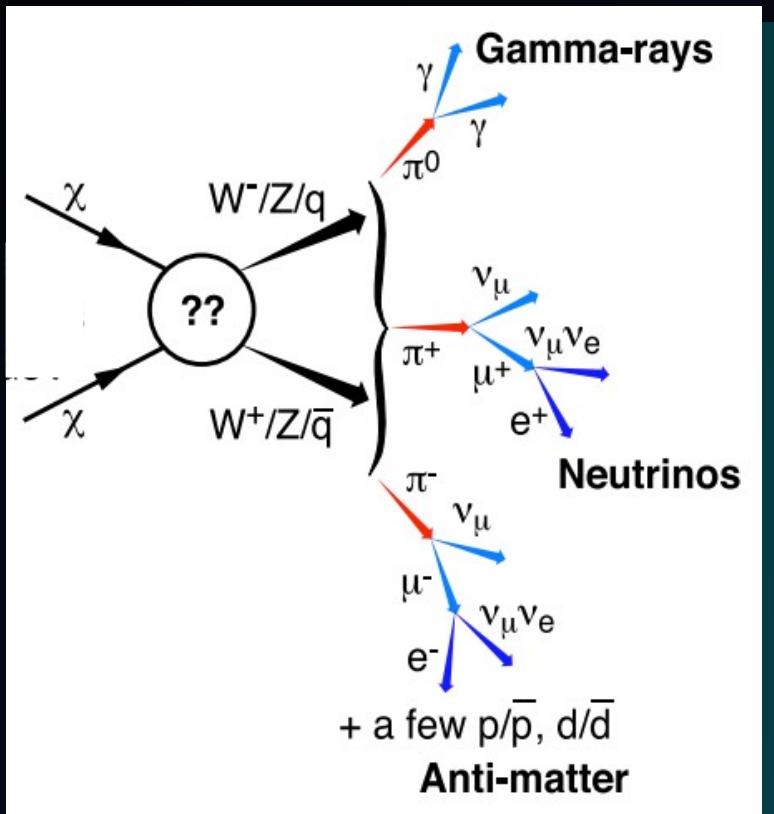
Ingredient #1: DM Interaction Rate

- 1) Thermal equilibrium:
 $\text{DM} + \text{DM} \Rightarrow \text{visible particles}$
 $\text{Visible particles} \Rightarrow \text{DM} + \text{DM}$
- 2) Universe cools, only
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- 3) Universe expands too fast.
No more annihilations.
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Predicts a particular annihilation rate for dark matter.

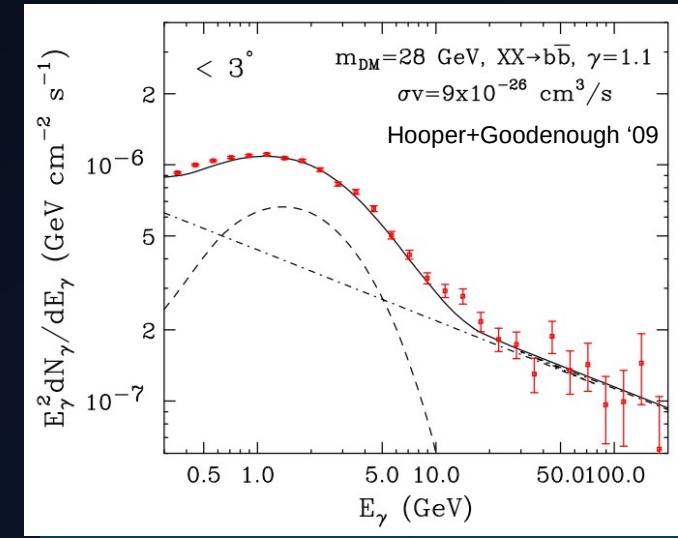


Ingredient #2: Energy Spectrum



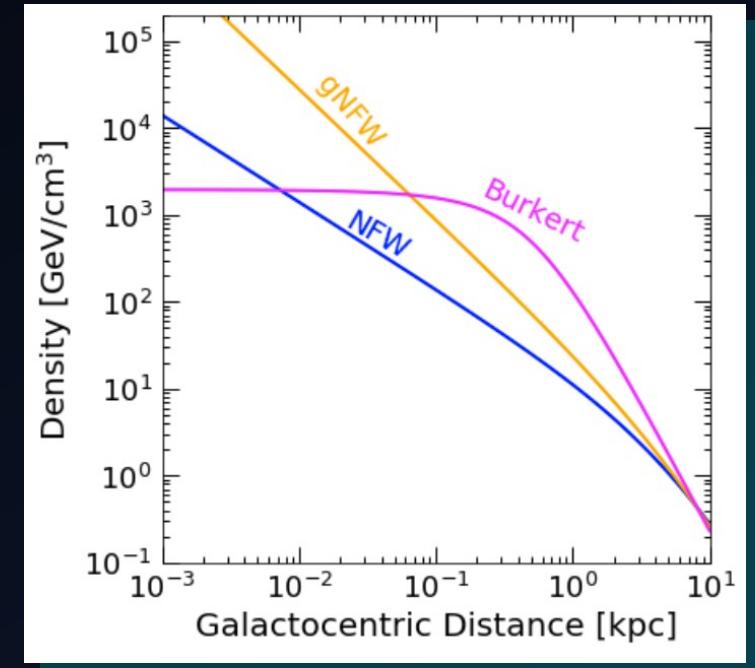
Baltz et al 0806.2911

- Also driven by particle physics model
- Shape depends on:
 - branching ratios to final SM states
 - boosts of particles



Ingredient #3: DM Density+Distribution

- Line of sight integral over DM density
 - J-factor (annihilation)
 - D-factor (decay)
- DM density profiles not well-known
 - large uncertainties



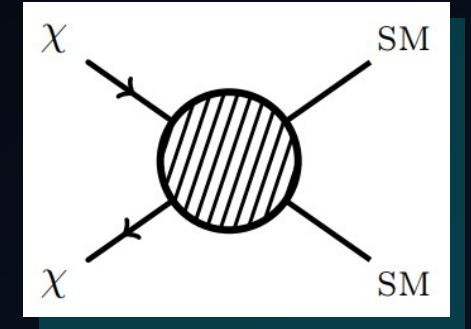
Indirect Detection Ingredients

(Neutral
particles)

Particle Physics

Astrophysics

$$\Phi(E, \phi) = \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell.$$



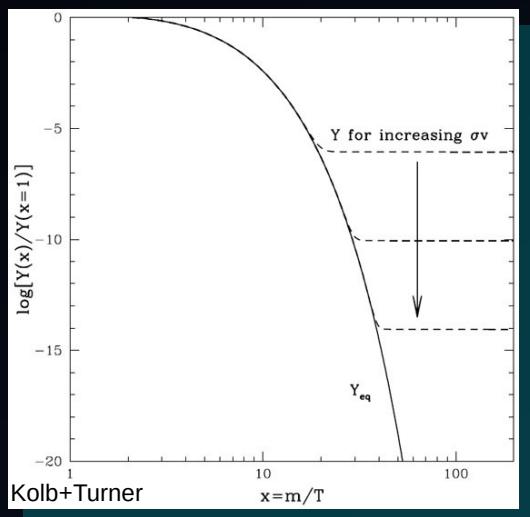
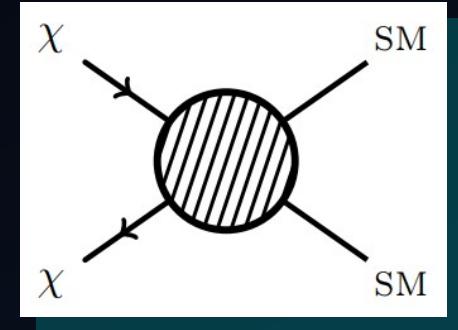
Indirect Detection Ingredients

Particle Physics Astrophysics

$$\Phi(E, \phi) \xrightarrow{\Gamma} \frac{1}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell.$$

(Neutral particles)

Annihilation cross section



Indirect Detection Ingredients

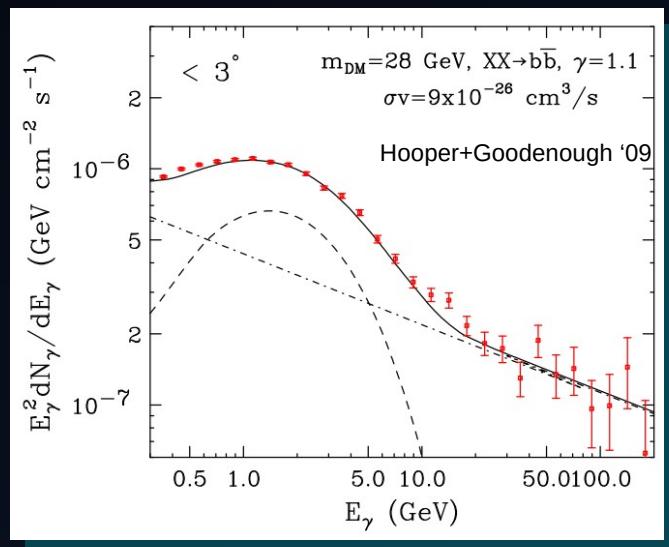
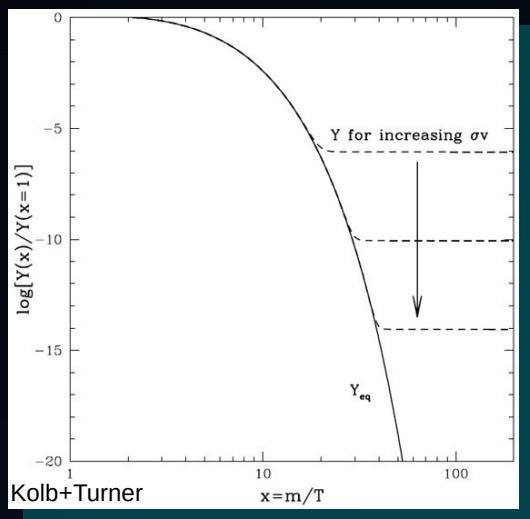
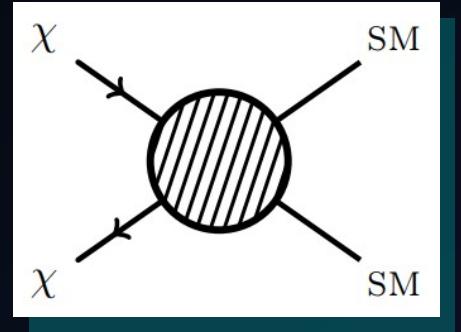
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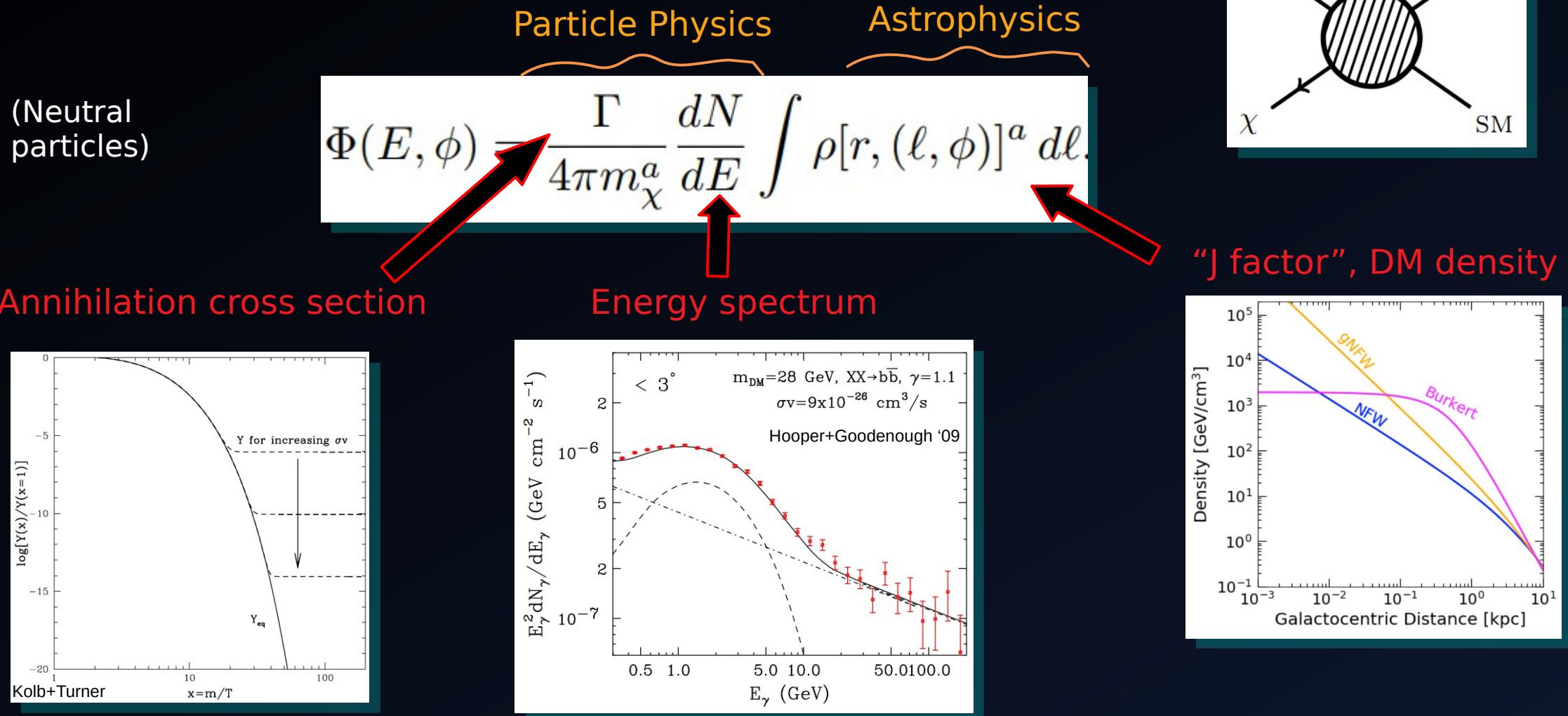
(Neutral particles)

Annihilation cross section

Energy spectrum



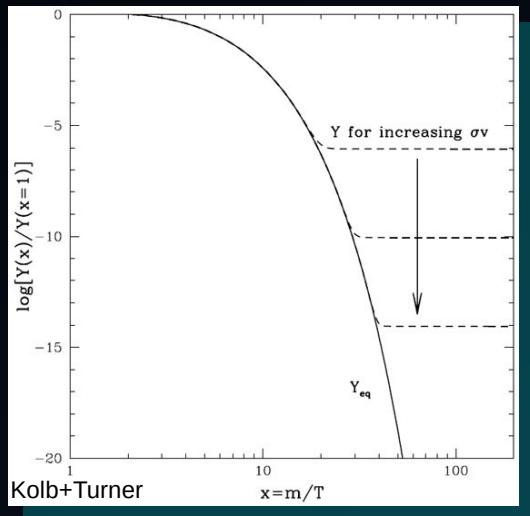
Indirect Detection Ingredients



Indirect Detection Ingredients

(Neutral particles)

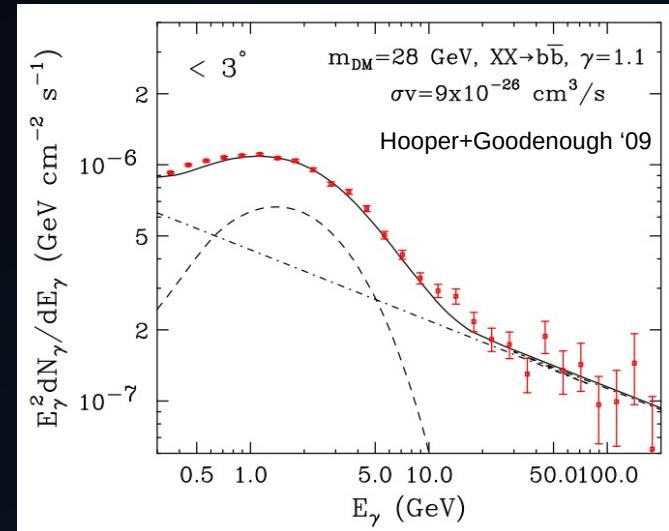
Annihilation cross section



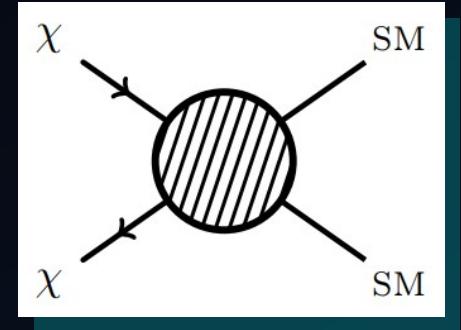
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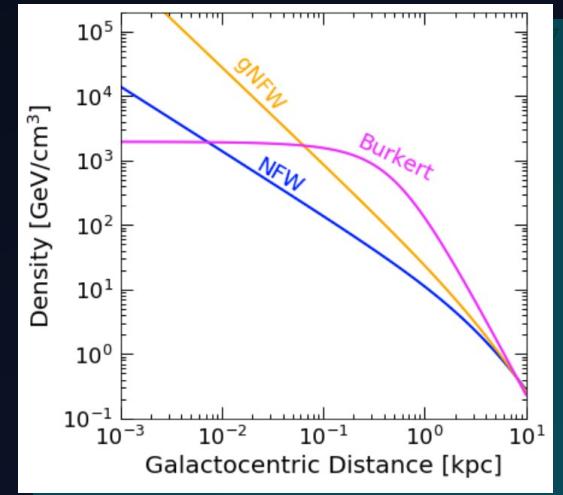
Energy spectrum



Astrophysics



"J factor", DM density

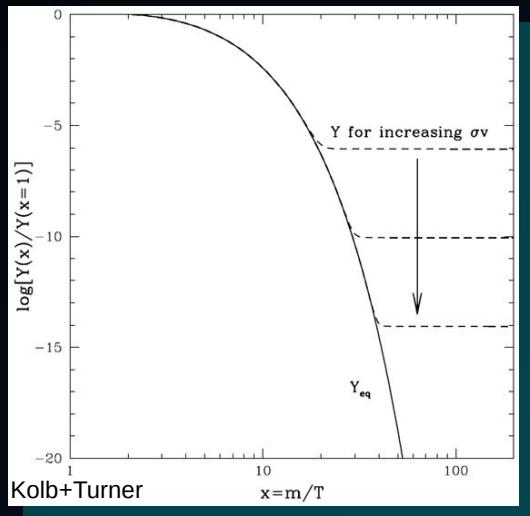


Look where this is large!

Indirect Detection Ingredients

(Neutral particles)

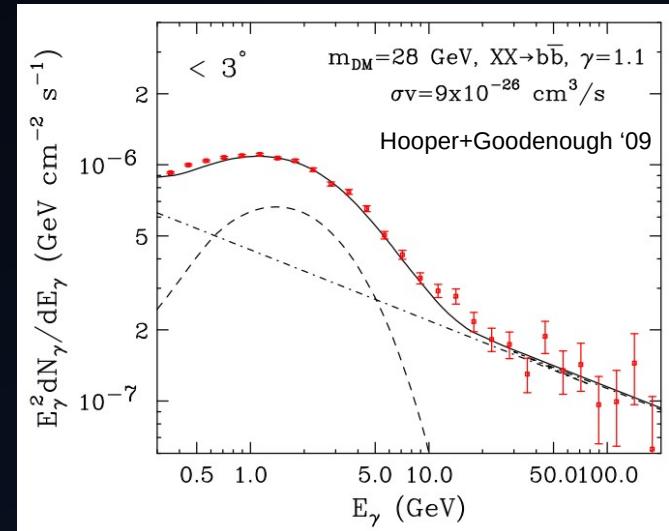
Annihilation cross section



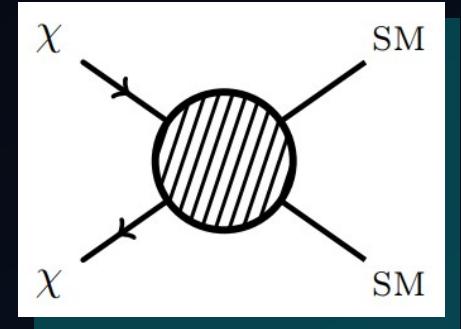
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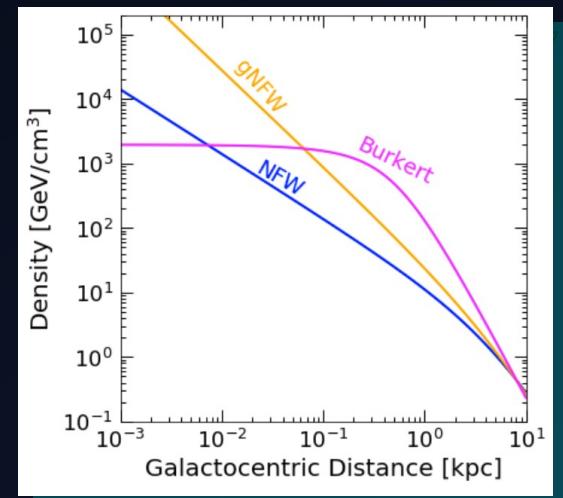
Energy spectrum



Astrophysics



"J factor", DM density

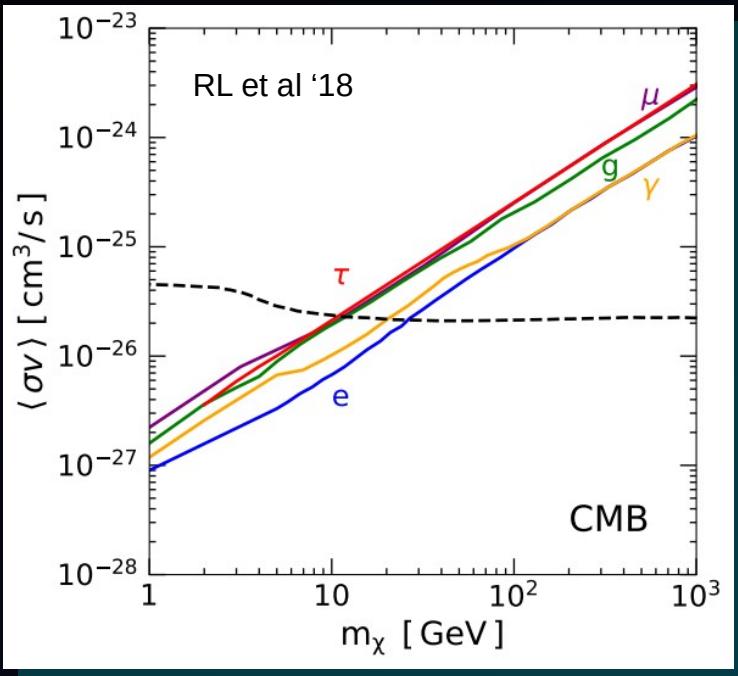


Look where this is large!

...or places with low background!

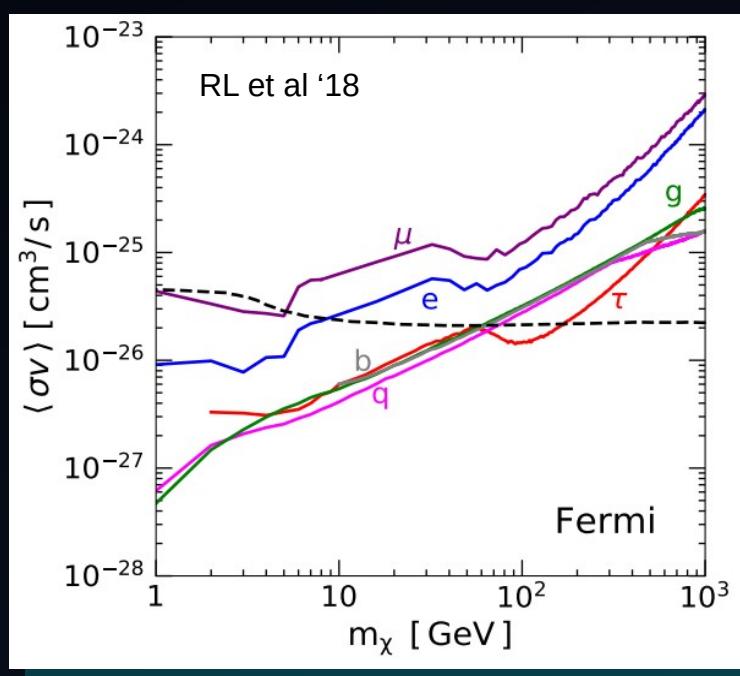
Complementarity: cornering WIMPs

Strongest low mass



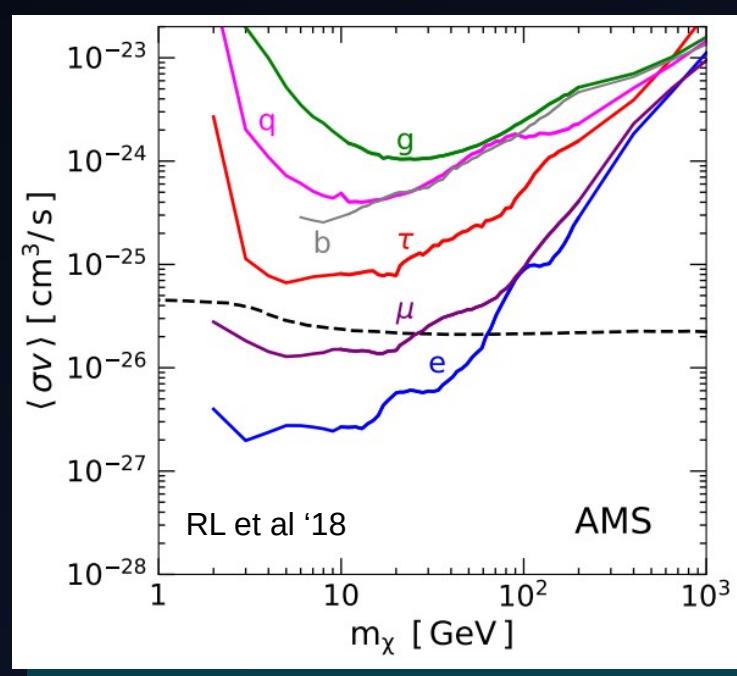
Also see Slatyer '15

Strongest for hadrons



Also see Fermi Collab '16

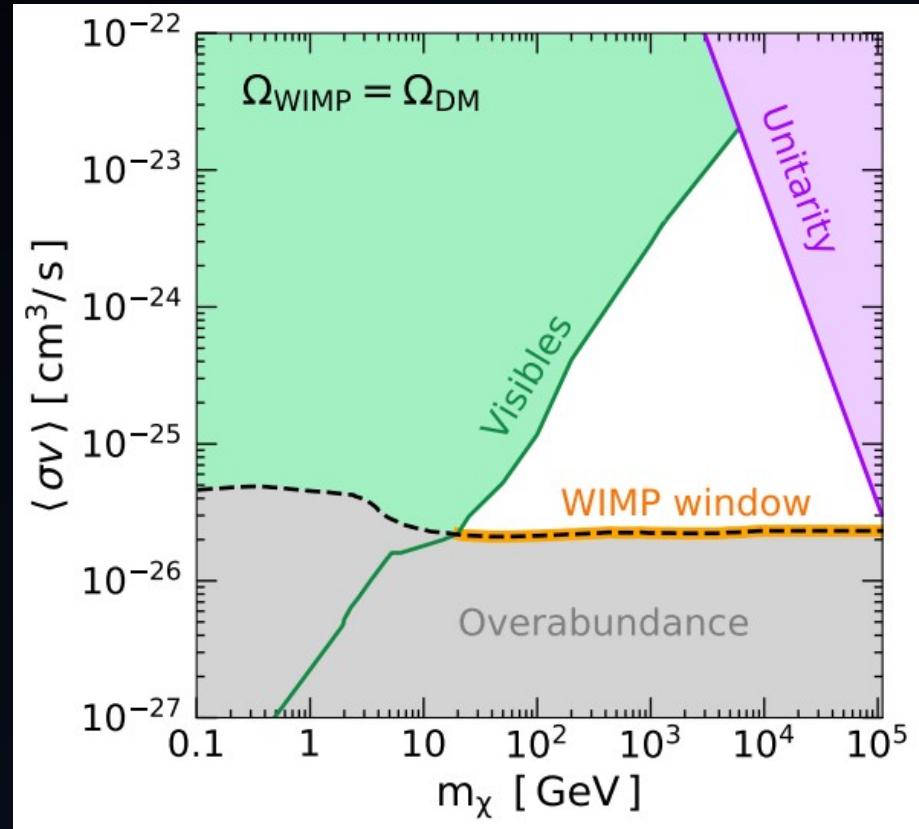
Strongest for leptons



Also see AMS collab '14

(strongest *and most robust* bounds)

Complementarity: cornering WIMPs



WIMP is not dead!

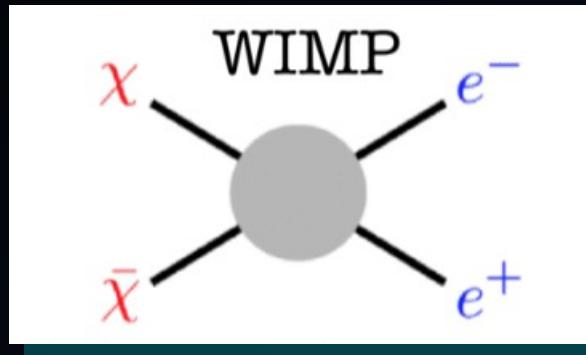
Use all possible final states, combine strongest limits
S-wave $2 \rightarrow 2$ thermal DM to visible states: mass greater than ~ 20 GeV

Still need to push through this window

OTHER TYPES OF FREEZEOUT

Range of freezeout options!

Zeldovich, Lee,
Weinberg, Steigman,
Turner,...+

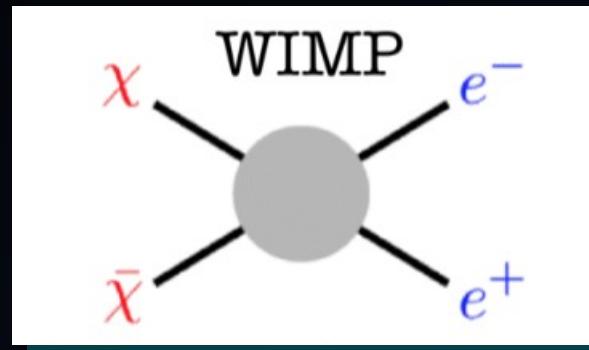


$$\Gamma_{\text{DM}} = \langle \sigma v_{\text{rel.}} \rangle n_{\text{DM}} > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \frac{0.12}{\langle \sigma v_{\text{rel.}} \rangle [25 \text{TeV}]^2}$$

Range of freezeout options!

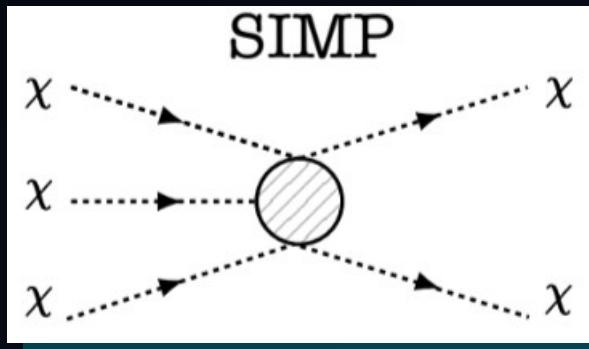
Zeldovich, Lee,
Weinberg, Steigman,
Turner,...+



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Hochberg, Kuflik,
Volansky, Wacker, 2014

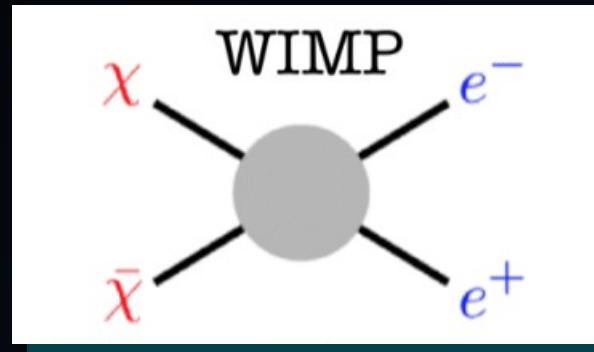


$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}}^2 > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \left(\frac{\text{MeV}}{m_{\text{DM}}} \right) \frac{0.12}{\sqrt{\langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle [3 \text{MeV}]^5}}$$

Range of freezeout options!

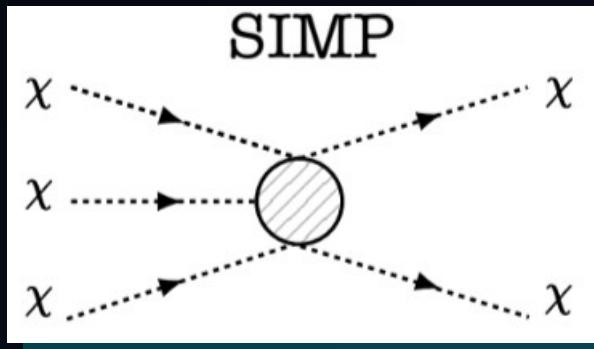
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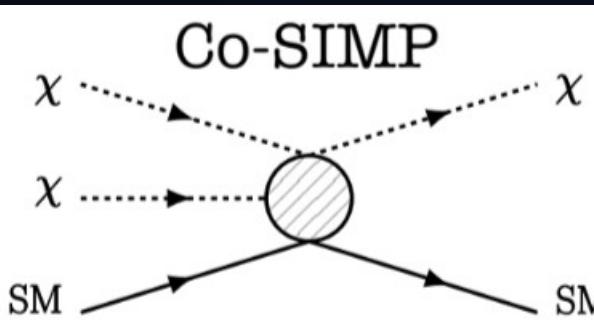
Hochberg, Kuflik,
Volansky, Wacker, 2014



$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}}^2 > H(T)$$

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Beacom, Smirnov, 2020



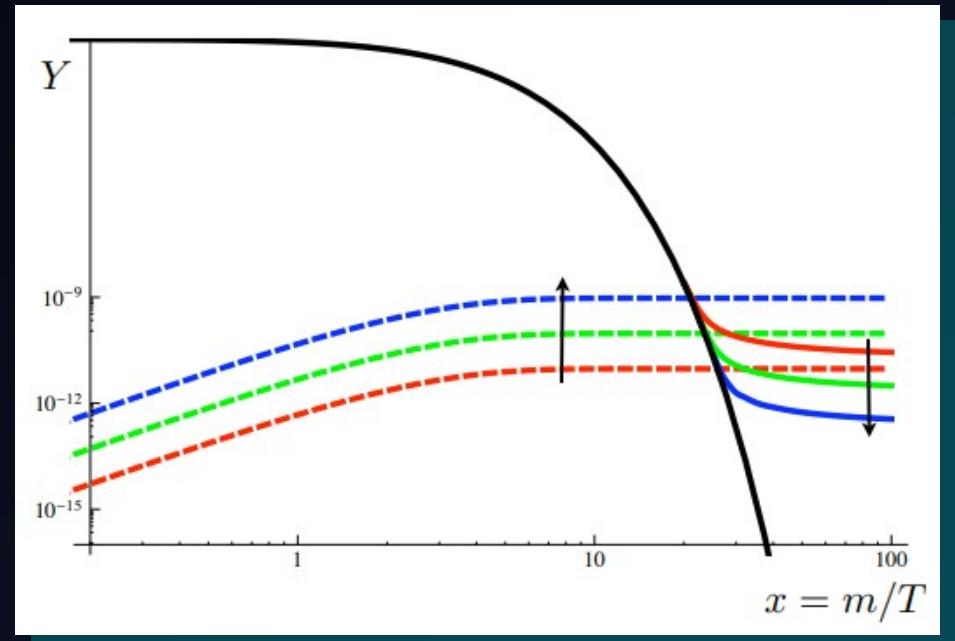
$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}} n_{\text{SM}} > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \left(\frac{\text{MeV}}{m_{\text{DM}}} \right)^3 \frac{0.12}{\langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle [100 \text{ MeV}]^5}$$

FREEZEOUT? FREEZE IN!

Dark matter freeze-in

- For freeze-out, the universe has a large population of DM that was in thermal equilibrium with the bath, that is depleted as the Universe cools
- For freeze-in, the universe starts with little DM and it never reaches thermal equilibrium with the bath
- At temperatures lower than the DM mass, the bath no longer has enough energy to produce it. Then, the DM is “frozen-in”

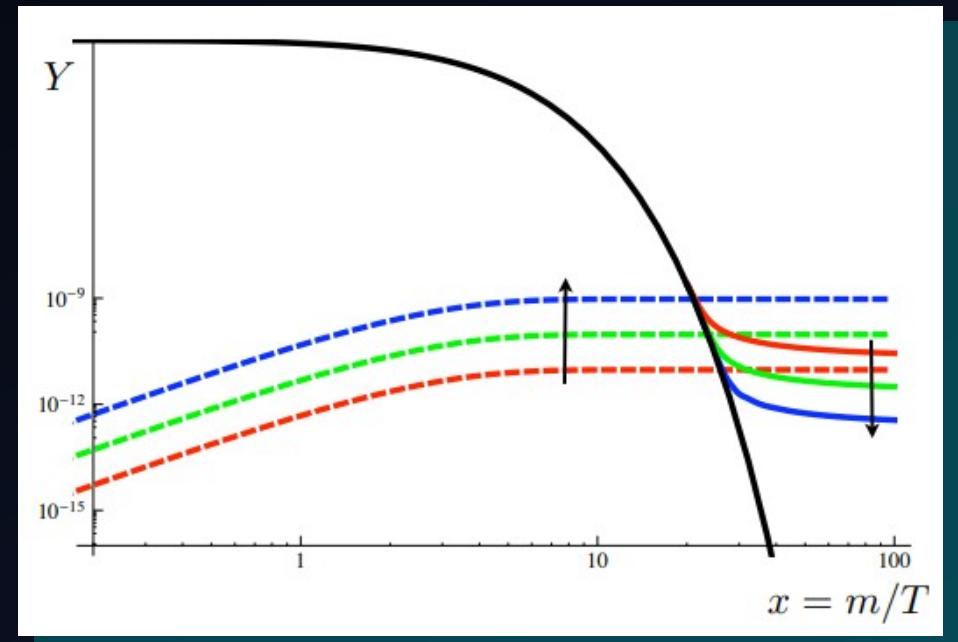


Arrows point in direction of increased DM-bath coupling

Hall, Jedamzik,
March-Russell, West
2009

Dark matter freeze-in

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- At temperatures lower than the DM mass, the bath no longer has enough energy to produce it. Then, the DM is “frozen-in”
- Still gets us right abundance, but applies to models with much smaller couplings -- “Feebly Interacting Massive Particles”, or FIMPs

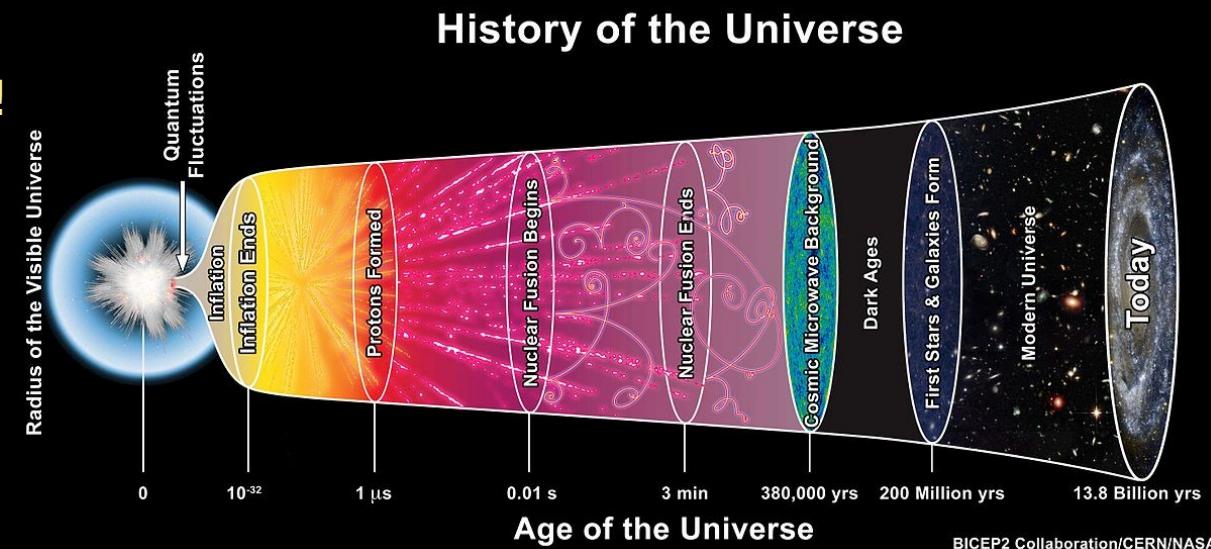


Arrows point in direction of increased DM-bath coupling

Hall, Jedamzik,
March-Russell, West
2009

Uncertainties for the early Universe?

- WIMP freezes out before BBN and CMB
 - A lot could happen between these times!
- Some particle could decay into DM
- Some particle could decay into visible matter, diluting the original amount of DM
- Maybe physics was just different to what we expect! Hard to know.
 - Finding DM may give us a new window further back into the Universe



AXIONS

(see also Jure Zupan's earlier lecture)

Strong CP problem

- If interactions are allowed by all symmetries, they generally are present
- One term is suspiciously silent

$$\mathcal{L}_\theta = \frac{\theta}{16\pi^2} G_{\mu\nu} \tilde{G}_{\mu\nu}$$

- Strange also because symmetry between matter and antimatter is broken in the weak interactions (CP violated), but not in strong interactions
- Term should induce neutron electric dipole moment, but experimentally it is measured to be at least **~10 ORDERS OF MAGNITUDE** below naive expectations ???

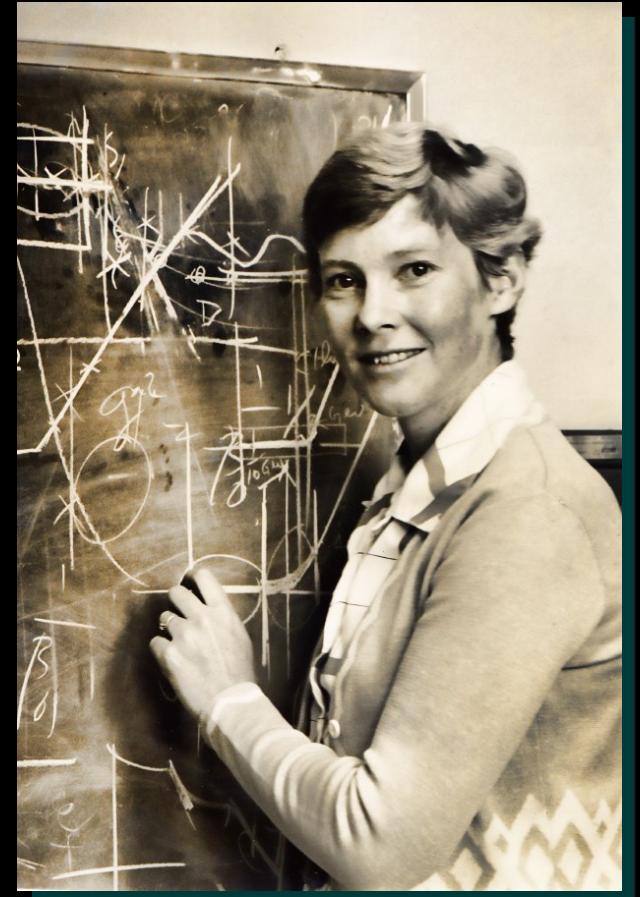
ENTER PECCEI + QUINN



Image: Quanta Magazine

Peccei-Quinn Symmetry

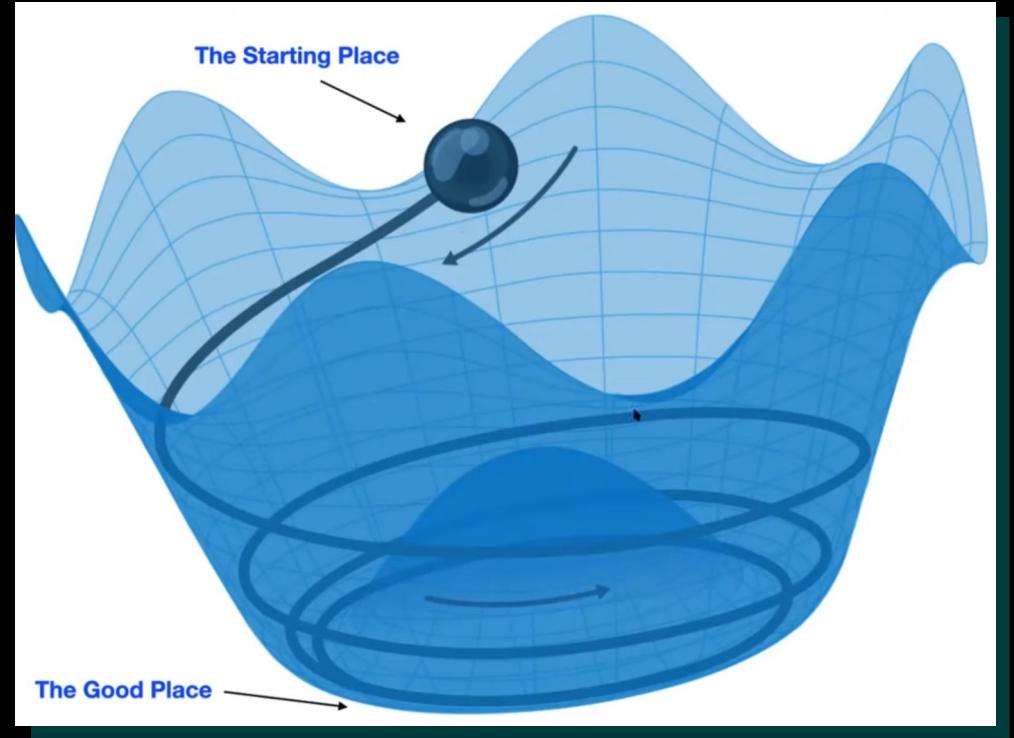
- If interactions are allowed by all symmetries, they generally are present
- So: Invoke new symmetry, Peccei-Quinn (PQ) symmetry, which is a minimal extension of the SM, which is spontaneously broken



Helen Quinn in her SLAC office, ~1977
Source: Helen Quinn

Dynamical field

- Promote the unwanted coupling (theta) to a dynamical field that relaxes to zero: not fixed for all time, but instead changes over time
- Value of the field is the co-efficient of the "unwanted" interaction
- Energetically preferred coupling is small/zero (so over time of the Universe it evolves to zero)
- Quanta of this field is the QCD axion



Peccei, Quinn 1977

Wilczek 1978, Weinberg 1978

Vafa, Witten 1984

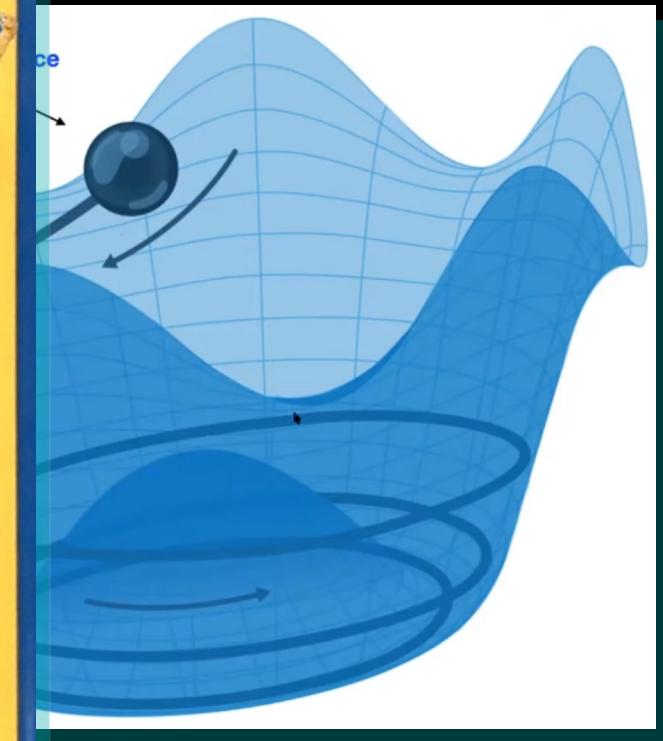
$$m_a = 5.7 \mu\text{eV} \left(\frac{10^{12}\text{GeV}}{f_a} \right)$$

Rebecca Leane

Wilczek

Dynamic

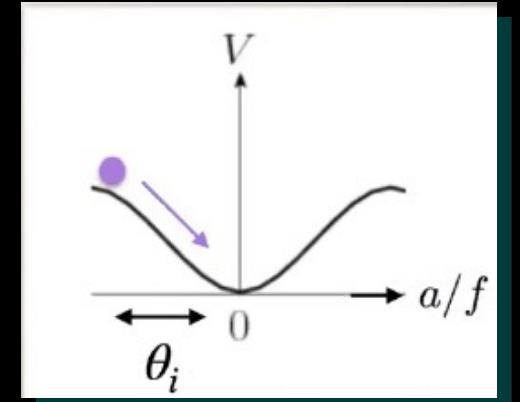
- Promote the unwanted interaction between the dynamical field that is present for all time, but interacts with matter.
- Value of the field is proportional to the "unwanted" interaction.
- Energetically preferred over static field (so over time of the universe)
- Quanta of this field



Wilczek

Can be dark matter

- Production occurs in the early universe, misalignment mechanism
 - Relic abundance ok for small masses, $\sim 10^{-8} - 10^{-3}$ eV
- Generally want mass less than about 20 eV for stability
- Can act as cold dark matter!



Zupan's slide



Increasingly popular, and can be DM!

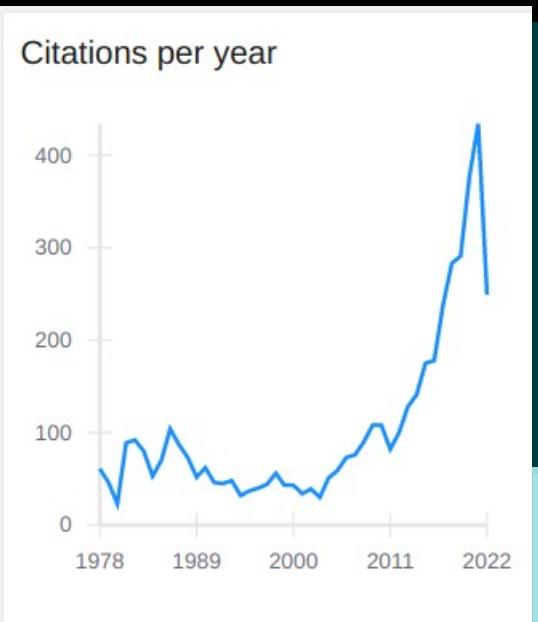
A New Light Boson?

Steven Weinberg (Harvard U.)
Dec, 1977

12 pages
Published in: *Phys.Rev.Lett.* 40 (1978) 223-226
DOI: [10.1103/PhysRevLett.40.223](https://doi.org/10.1103/PhysRevLett.40.223)
Report number: HUTP-77/A074
View in: OSTI Information Bridge Server, ADS Abstract Service, KEK scanned document

[pdf](#) [cite](#) [claim](#)

4,670 citations

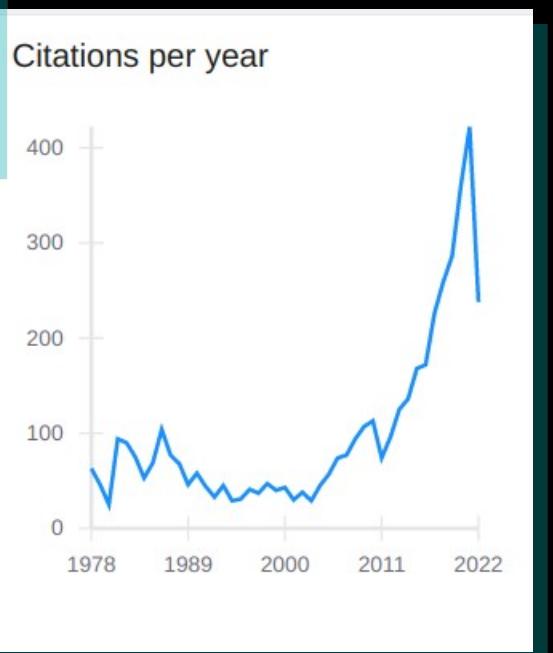


Frank Wilczek (Columbia U. and Princeton, Inst. Advanced Study)
Nov, 1977

12 pages
Published in: *Phys.Rev.Lett.* 40 (1978) 279-282
DOI: [10.1103/PhysRevLett.40.279](https://doi.org/10.1103/PhysRevLett.40.279)
Report number: Print-77-0939 (COLUMBIA)
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[pdf](#) [cite](#) [claim](#)

4,484 citations



Testable interactions

- Can couple to gluons, EM, leptons, quarks

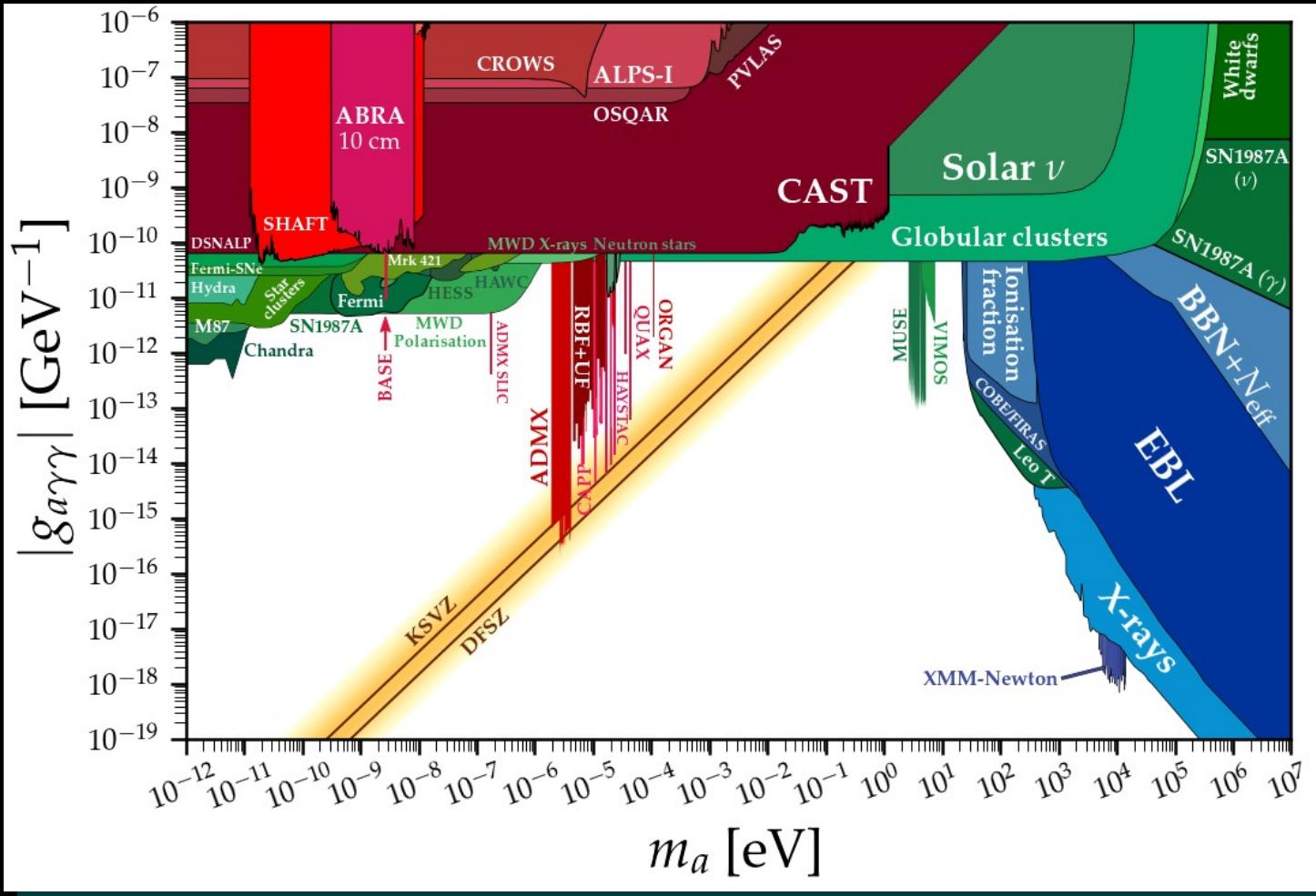
$$\mathcal{L}_{\text{int}} \sim -\frac{a}{F} (c_G \alpha_s G_{\mu\nu} \tilde{G}^{\mu\nu} + c_\gamma \alpha F_{\mu\nu} \tilde{F}^{\mu\nu} + d_q \sum_q m_q \bar{q} \gamma_5 q + d_l \sum_l m_l \bar{l} \gamma_5 l + \dots)$$

- Modifies Maxwell's equations:

$$\begin{aligned}\nabla \cdot E &= -\kappa \nabla a \cdot B \\ \nabla \times E &= -\frac{\partial B}{\partial t} \\ \nabla \cdot B &= 0 \\ \nabla \times B &= \frac{\partial E}{\partial t} + \kappa (\dot{a} B + \nabla a \times E)\end{aligned}$$

Some interactions
can also apply to
axion-like particles
(ALPs)

Axion limits

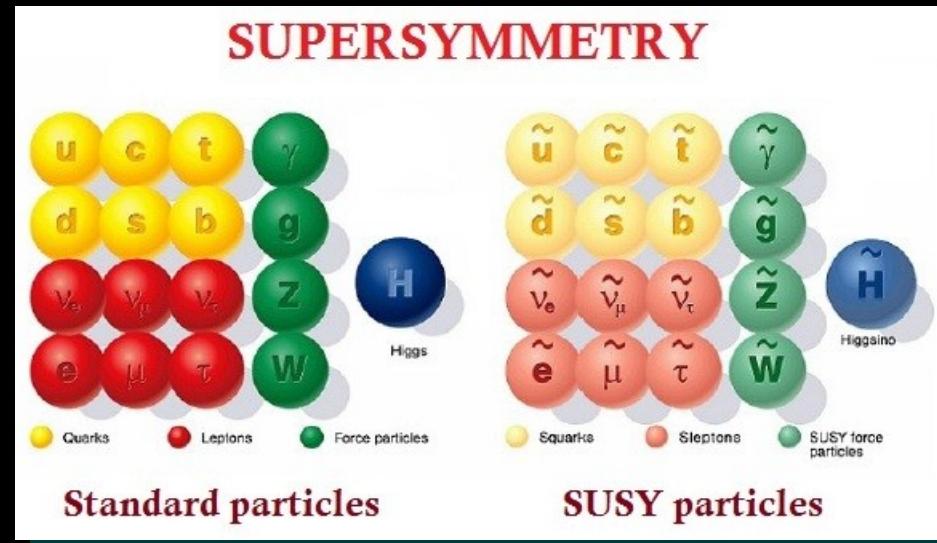


<https://cajohare.github.io/AxionLimits/>

COMPLETE DM MODELS

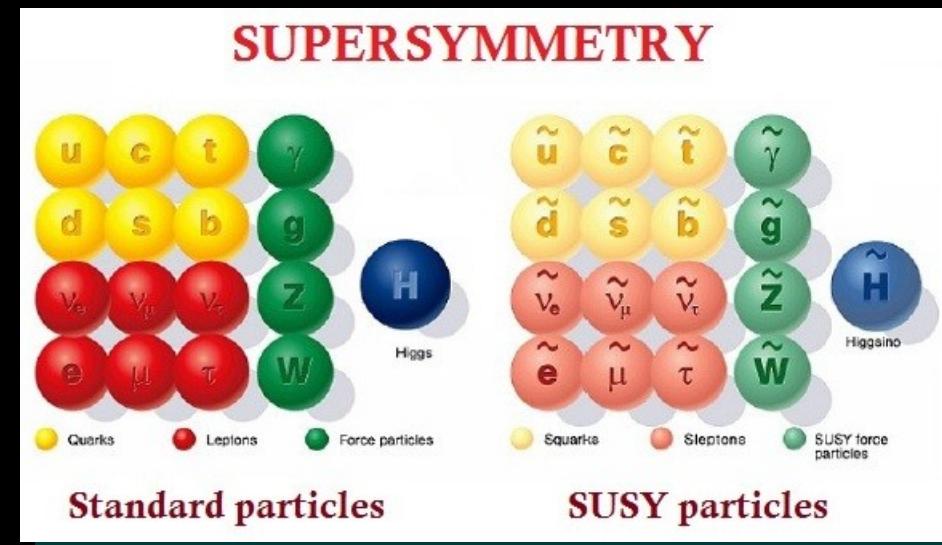
Complete DM models

- Axions:
naturally appear in string theory, GUTs
- Supersymmetry:
Prevailing new physics model for a long time,
naturally gave us:
 - Dark matter candidate (avoid proton decay
enforcing R-parity, so superpartners only
couple in pairs to the SM \rightarrow LSP)
 - Hierarchy problem solution, unification...



Complete DM models

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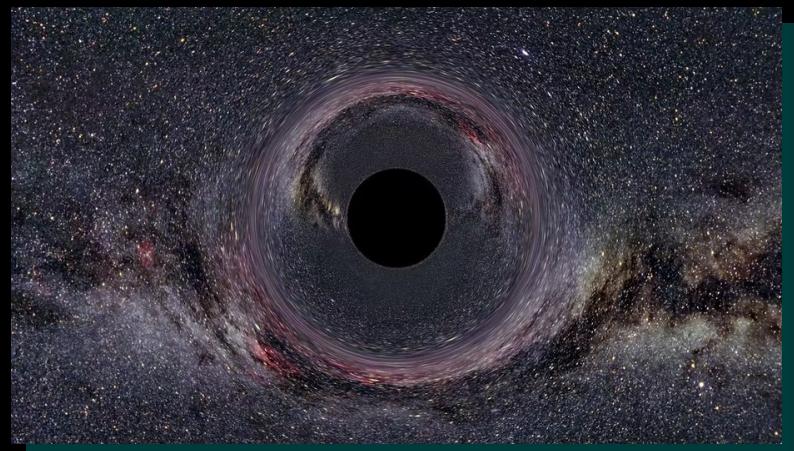
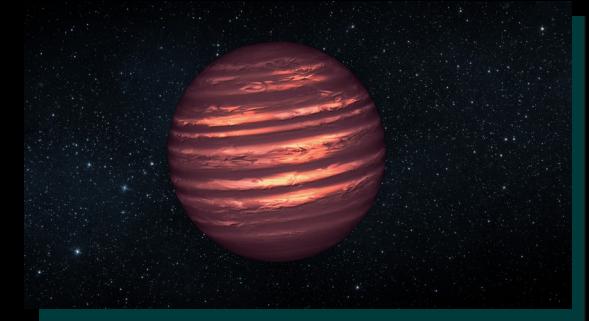


Complete models can guide us where to look for other new physics
(see unification talks later this week)

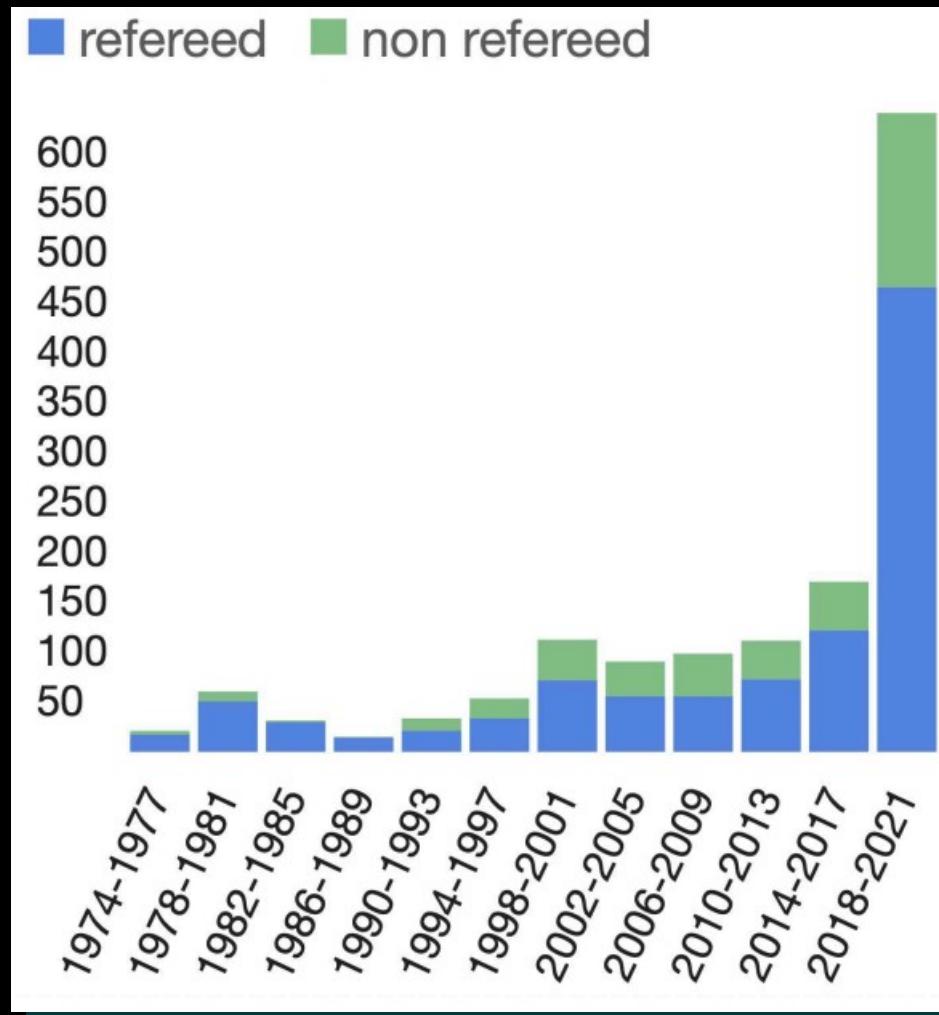
MACHOS: COMPOSITE DARK MATTER

MACHOs

- Compact objects like black holes, neutron stars, brown dwarfs...
- Originally, was just thought could be baryonic objects
 - Hard to reconcile with observations:
 - CMB / BBN
 - Microlensing surveys
- Increasingly popular candidate: primordial black holes



Primordial Black Hole Popularity



Carr + Kuhnel, 2022

Rebecca Leane

Primordial Black Holes

- First considered in 1966 by Zel'dovich and Novikov, in more detail by Hawking and Carr in the early 1970s
- What if DM is just black holes that formed before the epoch of BBN, AND with masses below the sensitivity range of microlensing surveys?
- Viable option, though we need (1) some production mechanism, and (2) some explanation of the current constraints from again CMB and lensing, among other things



Primordial Black Holes: Production

- Cosmological density time t after the Big Bang is

$$\rho \sim 1/(G t^2)$$

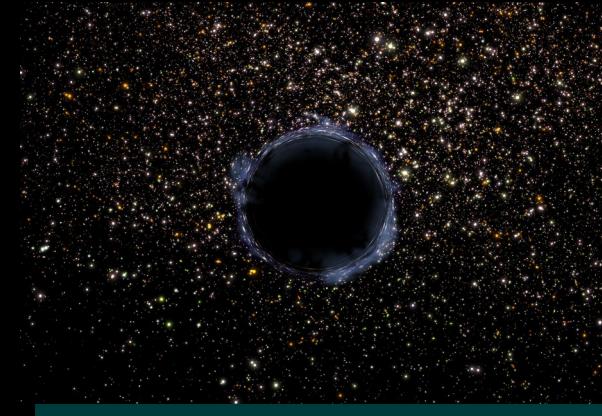
- Density needed for volume with mass M to fall within Schwarzschild radius is

$$\rho \sim c^6/(G^3 M^2)$$

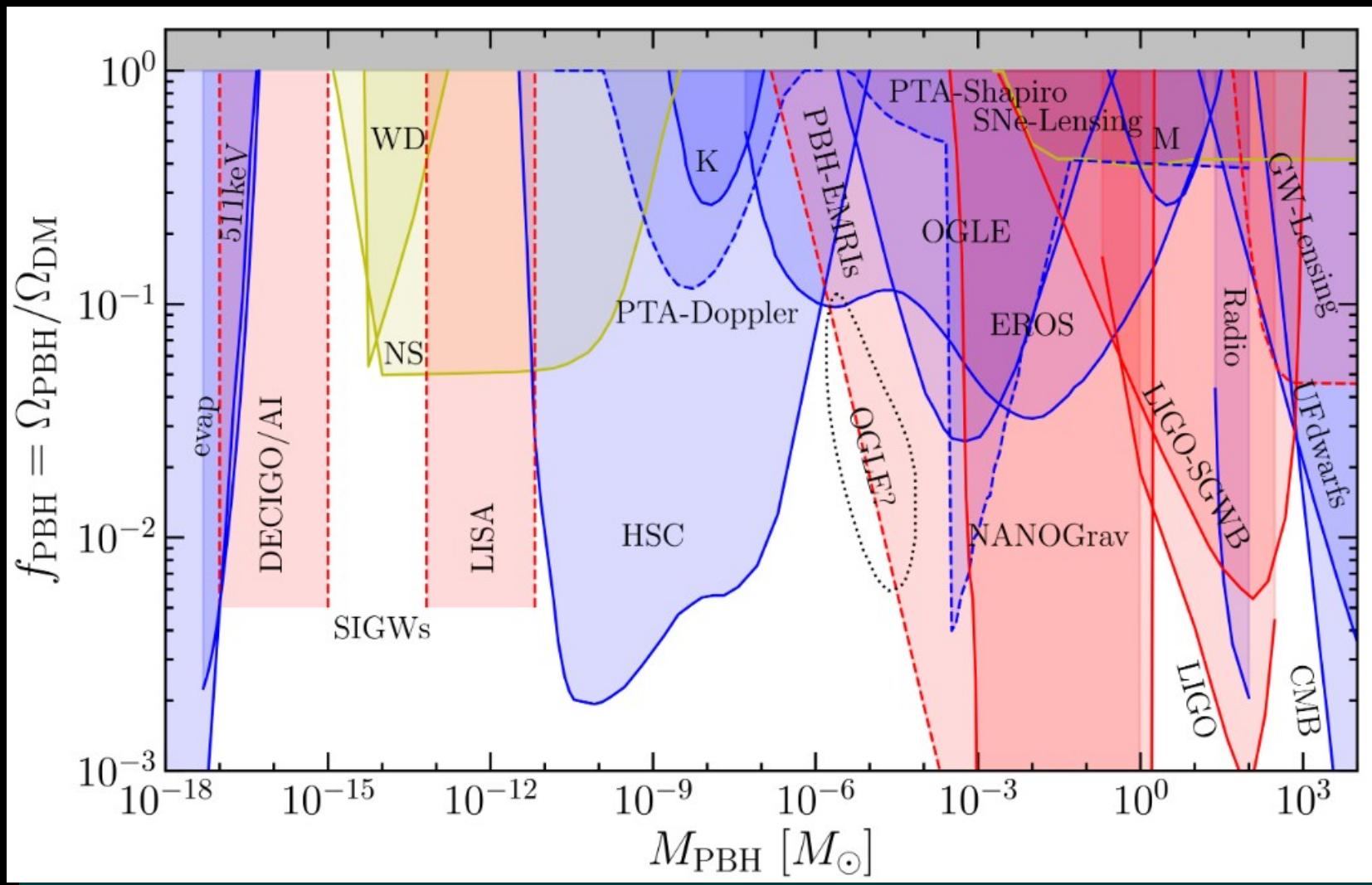
- Tells us that around the cosmological horizon, mass of PBHs would be

$$M \sim \frac{c^3 t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

- Therefore, depending on formation time, could span huge mass range!



Primordial Black Holes: Constraints



Outlook for dark matter

- Dark matter exists! Diverse range of evidence across many length scales
 - Finding its nature is a key goal of our community
- Hard to know the right dark matter theory direction.
No clear excesses to explain, no clear “correct” model.
- Experiment and theory have historically informed each other, and will continue to do so. Extensive new technological advances and new instruments on the horizon.
- Dark matter could be around the corner!

