Measurement of the neutrino magnetic moment at the NOvA experiment

Technical note

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Abstract

This is the abstract

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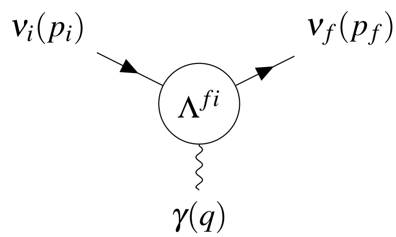

Figure 1: Effective coupling of neutrinos with one photon electromagnetic field.

1 Introduction

(TO DO: *Describe the main motivations for the analysis. Briefly mention that there was a previous study by Biao, what were the results there and what limitations (or maybe talk about this in the Experimental overview?)*)

2 Theoretical overview

(TO DO: *Go over the theoretical overview and make it better*)

2.1 Electromagnetic properties of the neutrino

In the standard model, neutrino can have electromagnetic interaction only at a higher order of the perturbative expansion of the interaction - from loop diagrams. (TO DO: *To be more specific, neutrinos can only have charge radius in standard model. Any other electromagnetic interactions must come from beyond standard model theories*)

In the one photon approximation, electromagnetic interactions of a neutrino field $\nu_k(x), k \in \{1, \ldots, N\}$, for N neutrino mass states, can be described by an effective interaction Hamiltonian [1]

$$\mathscr{H}_{em}^{(\nu)}(x) = \sum_{k,j=1}^{N} \overline{\nu}_k(x) \Lambda_\mu^{kj} \nu_j(x) A^\mu(x). \tag{1}$$

(TO DO: *decribe here what is lambda and everything else in the equation. Also mention figure 1*)

The amplitude of neutrino-to-neutrino interaction for **Dirac** neutrinos is

$$\langle \nu_f(p_f) | j_\mu^{(\nu)}(x) | \nu_i(p_i) \rangle = e^{i(p_f - p_i)x} \overline{u}_f(p_f) \Lambda_\mu^{fi}(p_f, p_i) u_i(p_i), \tag{2}$$

where p_f and p_i are the final and initial four momentums respectively and $u/\overline{u}$ are the solutions to the Dirac equation for a free particle. We take into account possible transitions between different mass states ν_i and ν_f [1]. (TO DO: *also describe what is j*)

The vertex function $\Lambda_\mu^{fi}(p_f, p_i)$ is generally a matrix and in the most general case it can be written in terms of linearly independent products of Dirac matrices (γ) and four momentum of the photon $(q = p_f - p_i)$:

$$\begin{aligned}\Lambda_{\mu}^{fi}(p_f, p_i) =& \mathbb{F}_1^{fi}(q^2)\, q_\mu + \mathbb{F}_2^{fi}(q^2)\, q_\mu \gamma_5 + \mathbb{F}_3^{fi}(q^2)\, \gamma_\mu + \mathbb{F}_4^{fi}(q^2)\, \gamma_\mu \gamma_5 + \\ & \mathbb{F}_5^{fi}(q^2)\, \sigma_{\mu\nu} q^\nu + \mathbb{F}_6^{fi}(q^2)\, \varepsilon_{\mu\nu\rho\gamma} q^\nu \sigma^{\rho\gamma}, \end{aligned} \tag{3}$$

where $\mathbb{F}_i^{fi}(q^2)$ are six Lorentz invariant form factors. For $f = i$ they are called "diagonal" and for $f \neq i$ "transition form factors" [1].

Applying conditions of hermiticity $\left(\mathscr{H}_{em}^{(\nu)\dagger} = \mathscr{H}_{em}^{(\nu)}\right)$ and of conservation of the current (continuity equation: $\partial^\mu j_\mu^{(\nu)}(x) = 0$), we can rewrite the vertex function as (TO DO: *maybe describe why we would want to apply these conditions*)

$$\Lambda_{\mu}^{fi}(q) = \left(\gamma_\mu - q_\mu \not{q}/q^2\right)\left[\mathbb{F}_Q^{fi}(q^2) + \mathbb{F}_A^{fi}(q^2)\, q^2 \gamma_5\right] - i\sigma_{\mu\nu} q^\nu \left[\mathbb{F}_M^{fi}(q^2) + i\mathbb{F}_E^{fi}(q^2)\, \gamma_5\right], \tag{4}$$

where $\mathbb{F}_Q^{fi}, \mathbb{F}_M^{fi}, \mathbb{F}_E^{fi}$ and $\mathbb{F}_A^{fi}$ are hermitian matrices representing the real charge, dipole magnetic, dipole electric and anapole neutrino form factors. In coupling with a real photon $(q^2 = 0)$ these become the neutrino charge and magnetic, electric and anapole moments [1].

For antineutrinos the form factors are transformed as:

$$\overline{\mathbb{F}}_\Omega^{fi} = -\mathbb{F}_\Omega^{if} = -\left(\mathbb{F}_\Omega^{fi}\right)^\star \quad \Omega = Q, M, E, \tag{5}$$

$$\overline{\mathbb{F}}_A^{fi} = \mathbb{F}_A^{if} = \left(\mathbb{F}_A^{fi}\right)^\star. \tag{6}$$

(TO DO: *maybe describe what does this mean?*)

In case of **Majorana neutrinos**, the general expression for the vertex function in terms of charge, magnetic, electric and anapole form factors looks the same as for Dirac neutrinos. (TO DO: *so does that mean that the interaction amplitude can be written in the same way for both Dirac and Majorana neutrinos?*) However, since Majorana antineutrinos are the same particle as Majorana neutrinos, from eq.5,6 we can see that:

$$\mathbb{F}_\Omega^M = -\left(\mathbb{F}_\Omega^M\right)^T \quad \Omega = Q, M, E, \tag{7}$$

$$\mathbb{F}_A^M = \left(\mathbb{F}_A^M\right)^T. \tag{8}$$

Therefore the Majorana charge, magnetic and electric form factor matrices are antisymmetric and the anapole form factor matrix is symmetric. This means that Majorana neutrino doesn't have any diagonal charge and dipole magnetic and electric moments, but it can have transition charge and magnetic and electric moment [1]. (TO DO: *Explain why is this worth mentioning or remove it if it's not*)

2.1.1 Neutrino electric and magnetic dipole moments

The size and effect of the neutrino electromagnetic properties depends on the specific theory beyond standard model.

Evaluating the one loop diagrams in the minimal extension of the standard model (TO DO: *should I mention here what actually do I mean by the minimal extension of SM? probably this: "introduction of three right-handed neutrinos"*) with right handed (Dirac) neutrinos gives us the first approximation of the electric and magnetic moments $\left(q^2 = 0\right)$:

$$\left.\begin{matrix} \mu_{kj}^D \\ i\varepsilon_{kj}^D \end{matrix}\right\} \simeq \frac{3eG_F}{16\sqrt{2}\pi^2}\left(m_k \pm m_j\right)\left(\delta_{kj} - \frac{1}{2}\sum_{l=e,\mu,\tau} U_{lk}^{\star} U_{lj}\frac{m_l^2}{m_W^2}\right), \tag{9}$$

where m_k, m_j are the neutrino masses and m_l are the masses of charged leptons which appear in the loop diagrams. Higher order electromagnetic corrections were neglected, but those can also have a significant contribution [1]. (TO DO: *ok, so what does that mean? Does that mean that this equation is not actually correct or that it's the higher/lower limit? or something else?*)

(It can be seen that) There are no diagonal electric moments $\left(\varepsilon_{kk}^D = 0\right)$ and the diagonal magnetic moments are approximately

$$\mu_{kk}^D \simeq \frac{3eG_F m_k}{8\sqrt{2}\pi^2} \simeq 3.2 \times 10^{-19}\left(\frac{m_k}{\mathrm{eV}}\right)\mu_B, \tag{10}$$

where μ_B is the Bohr magneton [1].

The transition magnetic moments are suppressed with respect to the largest of the diagonal magnetic moments by at least a factor of 10^{-4} due to the m_W^2 in denominator and the transition electric moments are even smaller than that due to the mass difference [1]. Therefore an experimental observation of a magnetic moment larger than in eq.10 would indicate physics beyond the minimally extended standard model [2].

Majorana neutrinos can be obtained by either adding a $\mathrm{SU}(2)_L$ Higgs triplet, or right handed neutrinos together with a $\mathrm{SU}(2)_L$ Higgs singlet (TO DO: *is this also considered a minimally extended SM? or is this the only way to get Majorana neutrinos? Most likely the former - from nuElmagInt2015.pdf: the absence of Higgs triplets, without which it is not possible to have Majorana mass terms.*). If we neglect the Feynman diagrams which depend on the model of the scalar sector (TO DO: *what does that mean for the result?*), the magnetic and electric dipole moments are

$$\mu_{kj}^M \simeq -\frac{3ieG_F}{16\sqrt{2}\pi^2}\left(m_k + m_j\right)\sum_{l=e,\mu,\tau} \mathrm{Im}\left[U_{lk}^{\star} U_{lj}\right]\frac{m_l^2}{m_W^2}, \tag{11}$$

$$\varepsilon_{kj}^M \simeq \frac{3ieG_F}{16\sqrt{2}\pi^2}\left(m_k - m_j\right)\sum_{l=e,\mu,\tau} \mathrm{Re}\left[U_{lk}^{\star} U_{lj}\right]\frac{m_l^2}{m_W^2}. \tag{12}$$

These are difficult to compare to the Dirac case, due to possible presence of Majorana phases in the PMNS matrices, but it is clear that they have the same order of magnitude as Dirac transition dipole moments. However, the neglected model dependent contributions can enhance the transition dipole moments [1]. (TO DO: *but how much?*)

It is possible [2] to obtain "natural" upper limits on the size of neutrino magnetic moment by calculating its contribution to the neutrino mass by standard model radiative corrections. For Dirac neutrinos the radiative correction induced by neutrino magnetic moment, generated at an energy scale Λ, to the neutrino mass is generically

$$m_\nu^D \sim \frac{\mu_\nu^D}{3\times 10^{-15}\mu_B}\,[\Lambda\,(\mathsf{TeV})]^2\,\mathsf{eV}. \tag{13}$$

So for $\Lambda \simeq 1\mathsf{TeV}$ and $m_\nu \lesssim 0.3\mathsf{eV}$ the limit becomes $\mu_\nu^D \lesssim 10^{-15}\mu_B$. This applies only if the new physics is well above the electroweak scale ($\Lambda_{EW} \sim 100\mathsf{GeV}$). It is possible to get Dirac neutrino magnetic moment higher than this limit, for example in frameworks of minimal super-symmetric standard model, by adding more Higgs doublets, or by considering large extra dimensions [1].

The limit for Majorana neutrino magnetic moment is less stringent, due to the antisymmetry condition from eq.7 and considering $m_\nu \lesssim 0.3\mathsf{eV}$ can be expressed as

$$\mu_{\tau\mu}, \mu_{\tau e} \lesssim 10^{-9}\,[\Lambda\,(\mathsf{TeV})]^{-2} \tag{14}$$

$$\mu_{\mu e} \lesssim 3\times 10^{-7}\,[\Lambda\,(\mathsf{TeV})]^{-2} \tag{15}$$

which is shown in the flavour basis , which relates to the framework used previously as

$$\mu_{ij} = \sum_{\alpha\beta} \mu_{\alpha\beta} U^\star_{\alpha i} U_{\beta j}, \quad \alpha,\beta \in \{e,\mu,\tau\}\,. \tag{16}$$

This limits imply, that if a magnetic moment $\mu \gtrsim 10^{-15}\mu_B$ would be measured, it is plausible neutrinos are Majorana fermions and the scale of lepton violation would be well below the conventional see-saw scale [2].

2.1.2 Other neutrino electromagnetic properties

Neutrino electric charge is heavily constraint by the measurements on the neutrality of matter (since generally neutrinos having an electric charge would also mean that neutrons have charge which would affect all heavier nuclei). It is also constrained by the SN1987A, since neutrino having an effective charge would lengthen its path through the extragalactic magnetic fields and would arrive on earth later. It can also be obtained from nu-on-e scatter from the relationship between neutrino millicharge and magnetic moment. [nuElmagInt2015.pdf - sec. VIIA]

The neutrino charge radius is determined by the second term in the expansion of the neutrino charge form factor and can be interpreted using the Fourier transform of a spherically symmetric charge distribution. It can also be negative since the charge density is not a positively defined quantity. In the SM the charge radius has the form of (possible other definitions exist)

$$\langle r^2_{\nu_l}\rangle_{\mathsf{SM}} = \frac{G_\mathsf{F}}{4\sqrt{2}\pi^2}\left[3 - 2\log\left(\frac{m_l^2}{m_W^2}\right)\right]. \tag{17}$$

This corresponds to $\langle r^2_{\nu_\mu}\langle_{\mathsf{SM}} = 2.4\times 10^{-33}$ cm^2 and similar scale for other neutrino flavours. [nuElmagInt2015.pdf - sec. VIIB] [nuElmagInt2015.pdf - sec. VIIB] The effect of the neutrino charge radius on the neutrino-on-electron scattering cross section is through the following shift of the vector coupling constant (Grau and Grifols, 1986; Degrassi, Sirlin, and VMarciano, 1989; Vogel and Engel, 1989; Hagiwara et al., 1994):

$$g_V^{\nu_l} \to g_V^{\nu_l} + \frac{2}{3} m_W^2 \langle r_{\nu_l}^2 \rangle \sin^2 \theta_W \tag{18}$$

[nuElmagInt2015.pdf - sec. VIIB] The current experimental limits for muon neutrinos are from (TO DO: *check the current exp. limits*) Hirsch, Nardi, and Restrepo (2003) who obtained the following 90% C.L. bounds on $\langle r_{\nu_\mu}^2 \rangle$ from a reanalysis of CHARM-II (Vilain et al., 1995) and CCFR (McFarland et al.,1998) data:

$$-0.52 \times 10^{-32} < \langle r_{\nu_\mu}^2 \rangle < 0.68 \times 10^{-32}\ \mathrm{cm}^2 \tag{19}$$

In the Standard Model, the neutrino anapole moment is somehow coupled with the neutrino charge radii and is functionally identical. the phenomenology of neutrino anapole moments is similar to that of neutrino charge radii. Hence, the limits on the neutrino charge radii discussed in Sec. VII.B also apply to the neutrino anapole moments multiplied by 6. in the standard model the neutrino charge radius and the anapole moment are not defined separately and one can interpret arbitrarily the charge form factor as a charge radius or as an anapole moment. Therefore, the standard model values for the neutrino charge radii in Eqs. (7.35)–(7.38) can be interpreted also as values of the corresponding neutrino anapole moments. [nuElmagInt2015.pdf - sec. VIIC]

It is possible to consider the toroidal dipole moment as a characteristic of the neutrino which is more convenient and transparent than the anapole moment for the description of T-invariant interactions with nonconservation of the P and C symmetries. the toroidal and anapole moments coincide in the static limit when the masses of the initial and final neutrino states are equal to each other. The toroidal (anapole) interactions of a Majorana as well as a Dirac neutrino are expected to contribute to the total cross section of neutrino elastic scattering off electrons, quarks, and nuclei. Because of the fact that the toroidal (anapole) interactions contribute to the helicity preserving part of the scattering of neutrinos on electrons, quarks, and nuclei, its contributions to cross sections are similar to those of the neutrino charge radius. In principle, these contributions can be probed and information about toroidal moments can be extracted in low-energy scattering experiments in the future. Different effects of the neutrino toroidal moment are discussed by Ginzburg and Tsytovich (1985), Bukina, Dubovik, and Kuznetsov (1998a, 1998b), and Dubovik and Kuznetsov (1998). In particular, it has been shown that the neutrino toroidal electromagnetic interactions can produce Cherenkov radiation of neutrinos propagating in a medium. [nuElmagInt2015.pdf - sec. VIIC]

2.2 Measuring neutrino magnetic moment

(TO DO: *Need to add some general description of the measurements*)

2.2.1 Effective neutrino magnetic moment

(TO DO: *Describe why experiments measure an effective magnetic moment*) What we measure in experiments is an effective "flavour" magnetic moment, which is influenced by mixing of "mass" magnetic moments (and electric moments) and oscillations. In the ultrarelativistic limit this is

$$\mu^2_{\nu_l}(L,E_\nu) = \sum_j \left| \sum_k U^\star_{lk} e^{-i\Delta m^2_{kj} L/2E_\nu} \left(\mu_{jk} - i\varepsilon_{jk}\right) \right|^2 . \tag{20}$$

What is called the effective magnetic moment (often just magnetic moment) therefore contains contributions from both the neutrino magnetic and electric moment [1].

For antineutrinos, the effective magnetic moment is (TO DO: *maybe I should just combine these two equations to avoid overcrowding*)

$$\mu^2_{\bar{\nu}_l}(L,E_\nu) = \sum_j \left| \sum_k U^\star_{lk} e^{+i\Delta m^2_{kj} L/2E_\nu} \left(\mu_{jk} - i\varepsilon_{jk}\right) \right|^2 . \tag{21}$$

So the only difference is in the phase induced by neutrino oscillations.

For experiments with baselines short enough for neutrino oscillations to not develop ($\frac{\Delta m^2 L}{2E_\nu} \ll 1$), such as the NOvA ND, the effective magnetic moment can be expressed as

$$\mu^2_{\nu_l} \simeq \mu^2_{\bar{\nu}_l} \simeq \sum_j \left| \sum_k U^\star_{lk} \left(\mu_{jk} - i\varepsilon_{jk}\right) \right|^2 = \left[U\left(\mu^2 + \varepsilon^2\right) U^\dagger + 2\,\mathrm{Im}\left(U\mu\varepsilon U^\dagger\right) \right]_{ll'}, \tag{22}$$

which is independent of the neutrino energy and of the source to detector distance.

It is important to mention, that since the effective magnetic moment depends on the flavour of the studied neutrino, it is different for different types of neutrino experiment. Also the solar neutrino experiments need to include the effect of the solar matter on the neutrino oscillations. Therefore the reports on the value (or upper limit) of the effective neutrino magnetic moment are not directly comparable between different types of neutrino experiments.

2.2.2 Neutrino-on-electron elastic scattering

The most sensitive method to measure neutrino magnetic moment is the low energy elastic scattering of (anti)neutrinos on electrons [1]. This interaction has two observables, the recoil electron's kinetic energy (T_e) and the recoil angle with respect to the incoming neutrino beam (θ). From simple $2 \to 2$ kinematics we can get

$$\left(P_\nu - P_{e'}\right)^2 = \left(P_{\nu'} - P_e\right)^2, \tag{23}$$

$$m^2_\nu + m^2_e - 2E_\nu E_{e'} + 2E_\nu p_{e'} \cos\theta = m^2_\nu + m^2_e - 2E_{\nu'} m_e . \tag{24}$$

Using the energy conservation

$$E_\nu + m_e = E_{\nu'} + E_{e'} = E_{\nu'} + T_e + m_e \Rightarrow E_{\nu'} = E_\nu - T_e \tag{25}$$

we get

$$E_\nu p_{e'} \cos\theta = E_\nu E_{e'} - E_{\nu'} m_e = E_\nu\left(T_e + m_e\right) - \left(E_\nu - T_e\right) m_e = T_e\left(E_\nu + m_e\right), \tag{26}$$

$$\cos\theta = \frac{E_\nu + m_e}{E_\nu}\sqrt{\frac{T_e^2}{E_{e'}^2 - m_e^2}} = \frac{E_\nu + m_e}{E_\nu}\sqrt{\frac{T_e^2}{T_e^2 + 2T_e m_e}}. \tag{27}$$

And finally we get

$$\cos\theta = \frac{E_\nu + m_e}{E_\nu}\sqrt{\frac{T_e}{T_e + 2m_e}}. \tag{28}$$

Electron's kinetic energy is kinematically constrained by

$$T_e \leq \frac{2E_\nu^2}{2E_\nu + m_e}. \tag{29}$$

Considering $E_\nu \sim \mathsf{GeV}$, we can approximate $\frac{m_e^2}{E_\nu^2} \to 0$ and in the small angle approximation we get from eq.28

$$T\theta^2 \cong 2m_e\left(1 - \frac{T_e}{E_\nu}\right) < 2m_e. \tag{30}$$

In the ultrarelativistic limit, the neutrino magnetic moment changes the neutrino helicity, turning active neutrinos into sterile (TO DO: *this is a very strong statement and it probably need a bit more backing up*). Since the SM weak interaction conserves helicity we can add the two contribution to the neutrino on electron cross section incoherently [1]:

$$\frac{d\sigma_{\nu_l e^-}}{dT_e} = \left(\frac{d\sigma_{\nu_l e^-}}{dT_e}\right)_{\mathsf{SM}} + \left(\frac{d\sigma_{\nu_l e^-}}{dT_e}\right)_{\mathsf{MAG}}. \tag{31}$$

The standard model contribution can be expressed as [1]:

$$\left(\frac{d\sigma_{\nu_l e^-}}{dT_e}\right)_{\mathsf{SM}} = \frac{G_F^2 m_e}{2\pi}\left\{\left(g_V^{\nu_l} + g_A^{\nu_l}\right)^2 + \left(g_V^{\nu_l} - g_A^{\nu_l}\right)^2\left(1 - \frac{T_e}{E_\nu}\right)^2 + \left(\left(g_A^{\nu_l}\right)^2 - \left(g_V^{\nu_l}\right)^2\right)\frac{m_e T_e}{E_\nu^2}\right\}, \tag{32}$$

where the coupling constants g_V and g_A are different for different neutrino flavours and for antineutrinos. Their values are:

$$g_V^{\nu_e} = 2\sin^2\theta_W + 1/2, \qquad g_A^{\nu_e} = 1/2, \tag{33}$$

$$g_V^{\nu_{\mu,\tau}} = 2\sin^2\theta_W - 1/2, \qquad g_A^{\nu_{\mu,\tau}} = -1/2. \tag{34}$$

For antineutrinos $g_A \to -g_A$.

Using expressions 28 and 30 we can also derive [3] cross sections with respect to $\cos\theta$, θ^2 and $T\theta^2$: (TO DO: *I don't think these equations are actually valid. We've been using some approximations and therefore I don't think these equations work for any theta dependence*)

$$\left(\frac{d\sigma_{\nu_l e^-}}{d\cos\theta}\right)_{\mathsf{SM}} = \frac{2G_F^2 E_\nu^2 m_e^2 \cos\theta\,(E_\nu+m_e)^2}{\pi\left((E_\nu+m_e)^2 - E_\nu^2\cos^2\theta\right)^2}$$
$$\left\{(g_V^{\nu_l}+g_A^{\nu_l})^2 + (g_V^{\nu_l}-g_A^{\nu_l})^2\left(1-\frac{2m_eE_\nu\cos^2\theta}{(E_\nu+m_e)^2-E_\nu^2\cos^2\theta}\right)^2 + \right.$$
$$\left.\left((g_A^{\nu_l})^2-(g_V^{\nu_l})^2\right)\frac{2m_e^2\cos^2\theta}{\left((E_\nu+m_e)^2-E_\nu^2\cos^2\theta\right)}\right\}, \tag{35}$$

$$\left(\frac{d\sigma_{\nu_l e^-}}{d\theta^2}\right)_{\mathsf{SM}} = \frac{G_F^2 m_e^2}{\pi\left(\theta^2+\frac{2m_e}{E_\nu}\right)^2}\left\{(g_V^{\nu_l}+g_A^{\nu_l})^2 + (g_V^{\nu_l}-g_A^{\nu_l})^2\left(\frac{\theta^2}{\theta^2+\frac{2m_e}{e_\nu}}\right)^2 + \right.$$
$$\left.\left((g_A^{\nu_l})^2-(g_V^{\nu_l})^2\right)\frac{2m_e^2}{E_\nu^2\left(\theta^2+\frac{2m_e}{E_\nu}\right)}\right\}, \tag{36}$$

$$\left(\frac{d\sigma_{\nu_l e^-}}{dT\theta^2}\right)_{\mathsf{SM}} = \frac{G_F^2 E_\nu}{4\pi}\left\{(g_V^{\nu_l}+g_A^{\nu_l})^2 + (g_V^{\nu_l}-g_A^{\nu_l})^2\left(\frac{T\theta^2}{2m_e}\right)^2 + \right.$$
$$\left.\left((g_A^{\nu_l})^2-(g_V^{\nu_l})^2\right)\frac{m_e}{E_\nu}\left(1-\frac{T\theta^2}{2m_e}\right)\right\}. \tag{37}$$

The neutrino magnetic moment contribution is (TO DO: *include derivation from [4]*) [1]:

$$\left(\frac{d\sigma_{\nu_l e^-}}{dT_e}\right)_{\mathsf{MAG}} = \frac{\pi\alpha^2}{m_e^2}\left(\frac{1}{T_e}-\frac{1}{E_\nu}\right)\left(\frac{\mu_{\nu_l}}{\mu_B}\right)^2, \tag{38}$$

147 where α is the fine structure constant.

Analogically to previous, we can also express this cross section in $\cos\theta$, θ^2 and $T\theta^2$:

$$\left(\frac{d\sigma_{\nu_l e^-}}{d\cos\theta}\right)_{\mathsf{MAG}} = \frac{2\pi\alpha^2(E_\nu+m_e)^2}{m_e^2\cos\theta}\frac{(E_\nu+m_e)^2-E_\nu^2\cos^2\theta-2m_eE_\nu\cos^2\theta}{\left((E_\nu+m_e)^2-E_\nu^2\cos^2\theta\right)^2}\left(\frac{\mu_{\nu_l}}{\mu_B}\right)^2, \tag{39}$$

$$\left(\frac{d\sigma_{\nu_l e^-}}{d\theta^2}\right)_{\mathsf{MAG}} = \frac{\pi\alpha^2}{m_e^2}\frac{\theta^2}{\left(\theta^2+\frac{2m_e}{E_\nu}\right)}\left(\frac{\mu_{\nu_l}}{\mu_B}\right)^2, \tag{40}$$

$$\left(\frac{d\sigma_{\nu_l e^-}}{dT\theta^2}\right)_{\mathsf{MAG}} = \frac{\pi\alpha^2}{4m_e^4}\frac{T\theta^2}{\left(1-\frac{T\theta^2}{2m_e}\right)}\left(\frac{\mu_{\nu_l}}{\mu_B}\right)^2. \tag{41}$$

The magnetic moment contribution exceeds the standard model contribution for low enough T_e [1]:

$$T_e \lesssim \frac{\pi^2 \alpha^2}{G_F^2 m_e^3} \left(\frac{\mu_\nu}{\mu_B}\right)^2 \simeq 2.9 \times 10^{19} \left(\frac{\mu_\nu}{\mu_B}\right)^2 [\mathsf{MeV}]\,, \tag{42}$$

which does not depend on the neutrino energy and makes neutrino experiment sensitive to lower energetic neutrinos more sensitive to the neutrino magnetic moment. (TO DO: *this is probably not about sensitivity to lower energetic neutrinos but to lower energetic electron, isn't it?*)

2.2.3 Neutrino on nucleus scattering

2.2.4 Cosmological effects

[NuMMBasicsAndAstro_2022.pdf] One of the most important astrophysical consequences of neutrino non-zero effective magnetic moments is the neutrino helicity change $\nu_l \to \nu_R$ with the appearance of nearly sterile right-handed neutrinos ν_R. In general, this phenomena can proceed in three different mechanisms:

1. the helicity change in the neutrino magnetic moment scattering on electrons (or protons and neutrons),
2. the neutrino spin and spin-flavour precession in an external magnetic field, and
3. the neutrino spin and spin-flavour precession in the transversally moving matter currents or in the transversally polarized matter at rest

For completeness note that the important astrophysical consequence of nonzero neutrino millicharges is the neutrino deviation from the rectilinear trajectory. (TO DO: *find out if this is the same thing that IceCube describes in their paper*)

3 Experimental overview

(TO DO: *Create a story for the experimental overview. Point out what is the hole in the current knowledge that NOvA can fill up*)

3.1 Direct muon (anti)neutrino magnetic moment measurements

3.1.1 NOvA (Biao's thesis)

- ν_μ only
- Only comparing total event counts - 25 events observed and 23.78 expected
- Put an upper limit (90% C.L.) of $\mu_{\nu_\mu} < 1.58 \times 10^{-9} \mu_B$ with 10.9% systematic uncertainty on the standard model background

- Used 3.62×10^{20} POT of data (6.74×10^{23} POT for MC) with $T\theta^2 < 0.003\mathsf{GeV} \times \mathsf{Rad}^2$, $0.3 < T < 0.9\mathsf{GeV}$

3.1.2 MiniBooNE

- ν_μ only
- Observed excess of events (seems a bit too high)

3.1.3 E734 at the Alternating Gradient Synchrotron (AGS) of the Brookhaven National Laboratory

- **Both ν_μ and $\overline{\nu}_\mu$**
- $\mu_{\nu_\mu} < 8.5 \times 10^{-10} \mu_B$

3.1.4 LSND

3.2 Direct electron (anti)neutrino magnetic moment measurements

3.3 Solar neutrino magnetic moment measurements

3.3.1 XENONnT

First results published in arXiv:2207.11330[5] on 22 July 2022.

- 5.9 tonne dual-phase liquid xenon TPC dark matter detector
- Region Of Interest is (1,140) keV
- Very low background (5 times lower than XENON1T)
- Tritium excluded as the potential background (also in XENON1T)
- No excess found - XENON1T excess excluded with 4σ
- The 90% C.L. upper limit on solar neutrinos with an "enhanced" magnetic moment is $\mu_{\nu_{sol}} < 6.3 \times 10^{-12} \mu_B$, the strongest non-astronomical limit so far (see fig.2)

Amir Khan used[6] XENONnT's results and derived limits on electromagnetic properties for the three SM neutrino flavours (see fig.3). For ν_μ they

3.3.2 XENON1T

3.3.3 BOREXINO

Should be $\mu_{\nu_e} < 2.8 \times 10^{-11} \mu_B$ [BorexinoLimit2017.pdf]

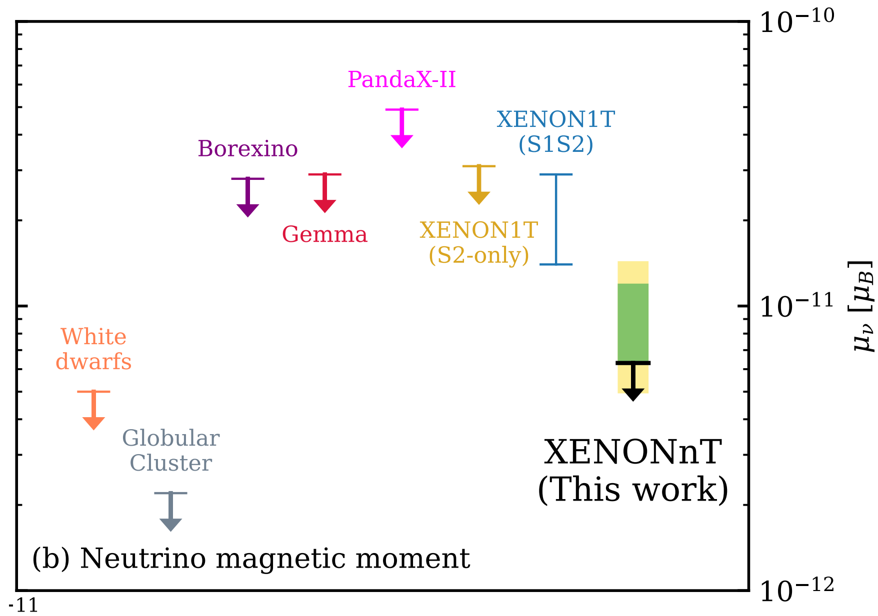

Figure 2: 90% C.L. upper limit on solar neutrinos with an enhanced magnetic moment.

3.3.4 GEMMA

Should be $\mu_{\nu\ Eff} < 2.9 \times 10^{-11} \mu_B$. [GemmaLimits2013.pdf]

3.4 Other

3.4.1 LHC Forward Physics Facilities

Preliminary sensitivity studies for future experiments (namely for FLArE and FASERν2)

- LHC's Forward Physics Facilities study high energy (TeV) neutrinos of all flavours from the ATLAS interaction point.
- Large opportunity to study tau neutrinos in more detail

3.5 Astrophysics

[NuMMBasicsAndAstro_2022.pdf] Neutrino electromagnetic processes that could be studied/observed in astrophysics

- Neutrino radiative decay
  - Decay of heavier neutrino flavour into a lighter neutrino and a photon
  - "The neutrino radiative decay has been constrained from the absence of decay photons in studies of the solar, supernova and reactor (anti)neutrino fluxes, as well as of the spectral distortions of the cosmic microwave background radiation."

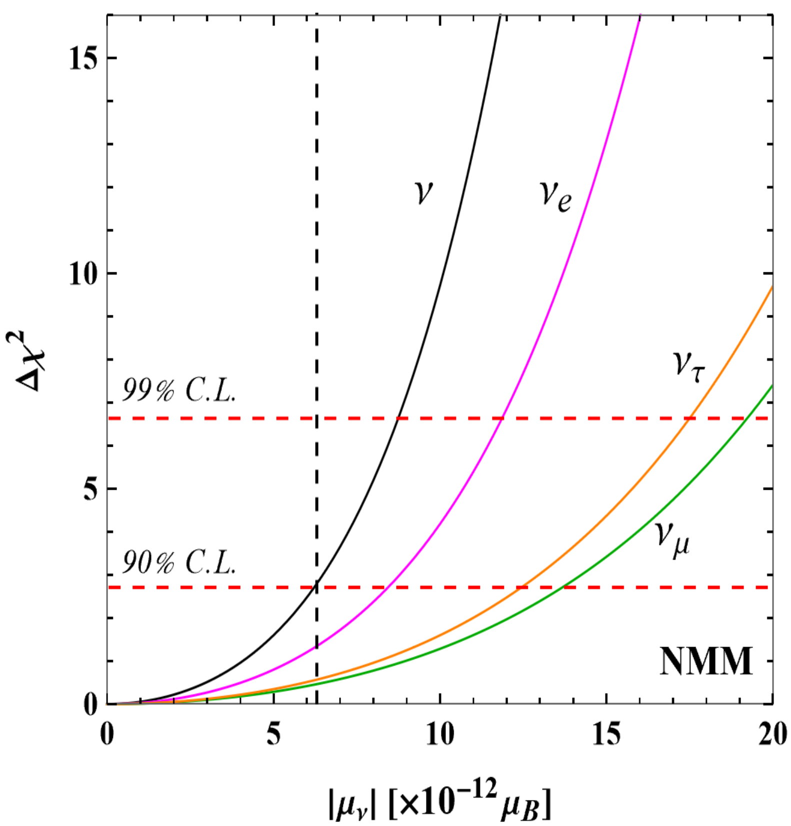

Figure 3: One-dimensional $\Delta\chi^2$ distribution with 90% and 99% C.L. boundaries of neutrino magnetic moments. The distribution in black corresponds to the effective flavor independent magnetic moment

    - Less stringent than the plasmon decay into a nu-antinu pairs

- Plasmon decay to neutrino-antineutrino pair
    - "For constraining neutrino electromagnetic properties, and obtaining upper bounds on neutrino magnetic moments in particular, the most interesting process is the plasmon decay into a neutrino-antineutrino pair [11]"
    - Plasmon decay frees the energy from the stars plasma in form of neutrinos that escape and therefore speeds up the star cooling
    - "observed properties of globular cluster stars provides new upper bounds on the effective neutrino magnetic moment $\mu_{ef} \leq (1.2-2.6)\times 10^{-12}\mu_B$ that is valid for both cases of Dirac and Majorana neutrinos."

- Transition of neutrino helicities $\nu_L \to \nu_R$ from active to sterile neutrinos
    - Supernovas would cool much faster - not observed for 1987A by Kamioka II and IMB, constraining Dirac neutrino mag. moment

4 Analysis overview

(TO DO: *Describe the motivations for this analysis*) What are we trying to achieve? Are we aiming for purity of the final sample or efficiency of the selection?

We are trying to select nu-on-e events with low electron recoil energies.

(TO DO: *Briefly introduce what are we going to do with these events after the selection (fitting)*) Are we just going to compare the event counts of signal and background (and possibly correct the background based on some other "sideband" selection?), or are we doing a fit to some spectra - either electron energy, angle or ETh2.

(TO DO: *Describe what I'm talking about in this section (datasets, weights, selection, resolution, fitting framework).*)

(TO DO: *Describe already here that we're dividing the signal/background into four due to ...*) Here on forward I'm going to describe the differences between these (definitions, weights, signal def, systematics. What is the same: event selection and binning. They're joint together in the fitting framework, where the ν_eCC MEC and the other backgrounds are simply summed together and scaled together. The ν-on-e background (also called the irreducible background by the LDM analysis) is treated/scaled separately.

Should I describe the NOvA Near Detector here? Specifically its capabilities for detecting electrons?

4.1 Datasets and Event Reconstruction details

For this analysis we are using the near detector samples with a standard Production 5.1 reconstruction. To tackle low number of ν-on-e and ν_eCC MEC events (especially after full selection) in the nominal simulation sample, and to increase speed and lower the computation costs of each study, we are using the following samples for the signal and background components, for the nominal prediction as well as systematically shifted.

(TO DO: *Find out what data sample we're using and write out the data POT*) For data we are only planning to unveil them after fully approved by the collaboration and we will be using the following data sample...

Yiwen has already looked at data for the following samples and the results are here...

Should I mention the POT counting here or somewhere else? - I think I should mention with each separate sample if the POT has to be rescaled or not.

The use of the samples can be briefly summarised as follows:

Signal	Enhanced ν-on-e sample
ν-on-e background	Enhanced ν-on-e sample
ν_eCC MEC background	Enhanced ν_eCC MEC sample
Other background	Nominal ND CAF sample

Table 1: Overview of simulation samples used.

Enhanced ν-on-e sample

(TO DO: *Describe the nuone sample*) Created by Wenjie Wu (was it just him or also Yiwen?) to do ... and fully described in the technote [7]. Using the overlayed and filematched samples for consistency.

(TO DO: *Find a reference and reasoning for why Wenjie hasn't created the other systematics samples*) We only have the selected few systematics definitions because ...

(TO DO: *Describe the differences*)

- Missing cross section parameters - unable to use cross section weights or so
- Special mode for nu-on-e elastic scattering 10005

The list of the nu-on-e sample definitions is in table 2.

Enhance ν_eCC MEC sample

Created by Yiwen Xiao [7] to tackle the low statistics of the ν_eCC MEC background events and subsequently large and unphysical cross section weights.

(TO DO: *List all the nueCC MEC sample definitions used. Do this after creating the filematched definitions maybe?*)

Near Detector filematched CAF sample

(TO DO: *describe and list all the ND nominal CAF samples*)

Nominal:
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_v1_nuone_overlay`
Systematically shifted samples:
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_calibup_v1_nuone_overlay`
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_calibdown_v1_nuone_overlay`
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_ckvup_v1_nuone_overlay`
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_ckvdown_v1_nuone_overlay`
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_lightlevelup_v1_nuone_overlay`
`prod_caf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_lightleveldown_v1_nuone_overlay`

Table 2: SAMWEB definitions for the ν-on-e samples.

Nominal:
`prod_flatsumdecaf_R20-11-25-prod5.1reco.g_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_v1_g4rwgt_respin_batch2_filematchedSystematics`
Systematically shifted samples:
`prod_flatsumdecaf_R20-11-25-prod5.1reco.e_nd_genie_N1810j0211a_nonswap_fhc_nova_v08_full_calibdown_v1_batch2_filematchedSystematics_calibdown_v1`

Table 3: SAMWEB definitions of the other background samples. THIS IS JUST A PLACE-HOLDER!

4.2 Analysis weights

(TO DO: *Describe why do we use weights*) What are the weights we are using and why?

To correct for known deficiencies in simulation of neutrino flux or cross sections we apply weights calculated for each event.

Table 4 shows what CAFAna weights are used to simulate what signal/background sample.

PPFX weight

`ana::kPPFXFluxCVWgt` [8] (TO DO: *What does this do (one sentence ish).*)

Signal	Flux and neutrino magnetic moment weights
ν-on-e background	Flux and radiative correction weights
ν_eCC MEC background	Flux and cross section weights
Other background	Flux and cross section weights

Table 4: Overview of CAFAna weights applied to each analysis sample.

Prod5.1 GSF XSec weight

`ana::kXSecCVWgt2020GSFProd51` (TO DO: *Find the reference: possibly Maria's docdb:53336 together with the official 2020 XSec tuning technote docdb:43962.*)

(TO DO: *Briefly describe what does this do. Also mention Yiwen's talk/technote about the large XSec weights that made her create an enhanced nueCC MEC sample.*)

We are only using the for the background since we assume that the cross section for the signal is perfect. Also there are not weights for this kind of interaction.

Radiative correction weight

(TO DO: *Why are we doing this? (reference Yiwen's talk/technote).*)

Mention here where did I get the original GENIE cross section from (reference Yiwen's talk or technote, plus the original paper that was used).

(TO DO: *Write out the actual version of the weight. Including the original and the corrected XSec constants*)

Say that we are not using the third part of the correction because it is tiny and it makes no difference. (tried and tested)

(TO DO: *correct the equation*) Calculated as

$$weight_{\text{Radiative Corr.}} = \left.\frac{d\sigma_{\nu-on-e}}{dy}\right|_{\text{Radiative Corr.}} \Big/ \left.\frac{\mathrm{d}\sigma_{\nu-on-e}}{\mathrm{d}y}\right|_{\text{GENIE 3}} ; y = \frac{E_e - m_e}{E_\nu} \tag{43}$$

4.2.1 Neutrino magnetic moment signal as a weight

(TO DO: *What does this do and why does it work? Reference the theory part as to why is the magnetic moment signal simply a rescaling of the GENIE cross section.*)

Using the same tree-level cross section from GENIE as in the rad. corr. weight.

(TO DO: *Write the name of the weight in CAFAna/nuone namespace and where it is located*)

(TO DO: *correct the equation*) Calculated as

$$weight_{\nu\text{ Mag. Moment}} = \left.\frac{d\sigma_{\nu-on-e}}{dy}\right|_{\nu\text{ Mag. Moment}} \Big/ \left.\frac{\mathrm{d}\sigma_{\nu-on-e}}{\mathrm{d}y}\right|_{\text{GENIE 3}} ; y = \frac{E_e - m_e}{E_\nu} \tag{44}$$

4.3 Event selection

Should this be a separate section or is it all right to keep it here? It will have a lot of plots...

(TO DO: *Add the link to the LDM group's technote and say what's different (or maybe do this after we discuss the cuts?*) Currently we are using the exact same selection as is used by the ND group [7] and very similar to the Light Dark Matter analysis (cite their technote).

Signal definition

(TO DO: *Define the signal of the NuMM. Reference the NuMM weight description above*) The signal of the neutrino magnetic moment analysis is just a reweighted signal of the ν-on-e analysis from the near detector group. We are using the same event selection as the near detector group.

(TO DO: *Decide and explain what signal definitions we're using (kIsVtxContained VS Fiducial volume*) What is the signal and all the background samples definition? Difference between using kIsVtxCont and the fiducial volume. Is there a fundamental difference or preference? Or does it just depend on me? The results/counts are quite different...

Signal	`kMode==` 10005 && `NDNuoneFiducial`
ν-on-e background	`kMode==` 10005 && `NDNuoneFiducial`
ν_eCC MEC background	!(`kMode== 5 && kElInFinState && NDNuoneFiducial`) && (`kIsCC` && `kIsNue` && `kMode == 10`)
Other background	!(`kMode== 5 && kElInFinState && NDNuoneFiducial`) || !(`kIsCC` && `kIsNue` && `kMode` == 10)

Table 5: Overview of signal and background definitions. Mode 10005 denotes ν-on-e events, while mode 5 denotes all electron scattering events, including inverse muon decay interactions. That is why we had to add a requirement of an electron in the final state. Mode 10 denotes all MEC events.(TO DO: *Check that the definitions are correct from the code.*)

Pre-selection

The pre-selection cuts have been kept from the ν_eCC analysis with loosened cut values (TO DO: *find a reference for this analysis*). Pre-selection cuts include basic quality cuts (TO DO: *describe the basic quality cuts that are implied from the preselection cuts*). They also remove the obvious νCC interactions by requiring that the length of the longest prong is < 800 cm, number of planes crossed by the longest prong is < 120, and the summed number of cells for all prongs in the slice is < 600. In pre-selection we also include a cut on the time difference between the mean times of the "current" slice and of the slice closest in time, which should be > 25 ns. This ensures that ... (TO DO: *describe why do we need the closest slice cut with reference to Yiwen's talk and technote*).

Fiducial and containment cuts

(TO DO: *Describe what does the fiducial cut do*) We require that the reconstructed vertex is contained within the following volume: $-185 < \mathsf{Vtx}_X < 175, -175 < \mathsf{Vtx}_Y < 175, 95 < \mathsf{Vtx}_Z < 1095$ cm.

To ensure all the energy is contained within the detector and to remove events originating outside of the detector (rock muons), we require that the extreme positions of hits for all prongs in the slice are within the following volume: $-190 < \mathsf{min}_X, \mathsf{max}_X < 180, -180 < \mathsf{min}_Y, \mathsf{max}_Y < 190, 105 < \mathsf{min}_Z, \mathsf{max}_Z < 1275$ cm

Single particle requirement

To selection events with a single particle we require that the fraction of energy contained in the most energetic shower is > 0.8, that the summed energy of all cells (above threshold and within ± 8 planes from the vertex) outside of the most energetic shower is < 0.02 GeV, and that the distance between the vertex and the start of the primary shower is < 20 cm.

Shower energy cut

(TO DO: *discuss the energy cut, should this be removed? What is the effect on the event count? Why was this included in the first place (the identifiers are not as strong for lowere energies - is this true though? - also there are further unexplored backgrounds that would need to be further studied and explore. Maybe depends on where would we move the cut...)*) The calorimetric energy of the primary shower is required to be within $0.5 < E_{cal} < 5$ GeV.

Event classifiers

We are using two event classifiers based on convolution neural network that were developed specifically to identify ν-on-e interactions. The first one (`NuoneID`) is trained to select ν-on-e events and the second one (`Epi0ID`) is trained on the events passing the `NuoneID` to reject the π^0 background. Our selection requires that `NuoneID`> 0.73 and that `Epi0ID`> 0.92.

(TO DO: *reference theory for the kinematics of nuone scattering*) We require that the product of reconstructed energy of the primary shower and the square of its angle from the Z axis is $E_{cal}\theta^2 < 0.005$ GeV $\times$ rad^2.

(TO DO: *Add plots of distributions of the event selection variables with two columns. LHS shows no cuts applied and RHS shows all previous cuts applied*)

Using the many plots below that show the effect of each of the cuts on the signal and all background events. (For signal we are showing NuMM=...)

(TO DO: *Describe the cutflow tables below*) The final event count and efficiency of each of the cuts is shown on the table 6. Table 4 shows the dissemination of background into the individual components.

(TO DO: *Add a discussion of possible improvements on the event selection on its limitations - mostly for the analysis review committee*) From here we can see that ... Maybe what can be improved is... This can likely be improved upon by specifically selection low energy events and removing the cut on the reconstructed shower energy.

4.4 Resolution and binning

(TO DO: *Add the energy resolution and binning plots*) The electron energy and angle distributions and resolutions. Are we going to fit in E, Th, or ETh2? Is there something else?

Show plots of Reco V True for both energy and angle. (Should I show it with or without the energy cut?). Also show the resolution plots.

4.5 Systematic uncertainties

(TO DO: *Describe the main systematic uncertainties. Add plots showing their effect on the NuMM events. Possibly with different event selection variables as X axes. Also show the final table with the percentages summed*) Plots showing combined uncertainties for signal and backgrounds. Maybe also some interpolations. Table of systematic uncertainties on the event count.

Normalization systematics

(TO DO: *Describe the normalization systematics (or just remove this if not using them in the end*) Should we include normalization systematics? Would that make any difference? There's a POT

Selection	ν Mag. Moment signal			ν-on-e background			Other background		
	N_{sig}	ε^{N-1}	ε (%)	N_{IBkg}	ε^{N-1}	ε (%)	N_{Bkg}	ε^{N-1}	ε (%)
No Cut	269.77	100	100	3.43×10^3	100	100	2.96×10^8	100	100
Closest Slc	262.20	97.19	97.19	3.35×10^3	97.50	97.50	2.74×10^8	92.66	92.66
Png Length	169.82	64.77	62.95	3.14×10^3	93.72	91.38	7.19×10^7	26.24	24.31
N Planes	169.82	100.00	62.95	3.14×10^3	99.98	91.37	7.19×10^7	99.98	24.31
N Cells	169.82	100.00	62.95	3.14×10^3	99.98	91.35	6.95×10^7	96.66	23.50
Fiducial	167.72	98.76	62.17	3.09×10^3	98.41	89.89	3.59×10^7	51.71	12.15
Cont.	159.37	95.02	59.08	2.48×10^3	80.43	72.30	1.38×10^7	38.35	4.66
ShwE Frac.	150.37	94.35	55.74	2.42×10^3	97.59	70.56	8.82×10^6	63.97	2.98
Vtx E	142.29	94.63	52.74	2.18×10^3	90.16	63.62	4.15×10^6	47.07	1.40
Shw Gap	137.96	96.96	51.14	2.09×10^3	95.58	60.80	3.25×10^6	78.34	1.10
Shw E	37.13	26.92	13.76	1.36×10^3	65.10	39.58	6.25×10^5	19.21	0.21
Nuoneid	29.48	79.39	10.93	940.21	69.18	27.38	2.42×10^4	3.88	8.19×10^{-3}
Epi0id	22.51	76.35	8.34	749.93	79.76	21.84	1.47×10^4	60.75	4.97×10^{-3}
$E\theta^2$	19.74	87.73	7.32	675.02	90.01	19.66	84.15	0.57	2.84×10^{-5}
$E\theta^2$ **(sb)**	2.74	-	1.01	74.30	-	2.16	1.01×10^3	-	3.43×10^{-4}
No ShwE	37.62	-	13.94	782.67	-	22.79	238.79	-	8.07E-05

Table 6: Event selection cutflow table

scaling uncertainty which is very small (find out exactly how small).

In the fitting experiment normalization uncertainties would probably not make any difference whatsoever, but in the counting experiment they might be important?

Neutrino flux systematics

(TO DO: *Describe the flux uncertainties. Describe the PCA. Describe the difference between what the ND group is doing and what we're doing*) Using the PCA vs using the PPFX universes+beam transport separately. Plots of energy showing shifts for signal and backgrounds separately

(TO DO: *understand differences with ND and 3F methods*)

This is mainly a normalization. Discuss how to use the fact that ν-on-e events can be used (and are used) to constraint the beam uncertainty. Would the counting experiment still be valid then? Maybe if we made another sideband sample...

Detector systematics

(TO DO: *Make plots of energy showing shifts for signal and backgrounds separately*)

Cross section systematics

(TO DO: *Describe the XSec systs*) Only for the non nu-on-e background. Assuming the nu-on-e events (including the signal events) are precisely known.

Plots of energy showing shifts for signal and backgrounds separately

4.6 Fitting framework

(TO DO: *Describe the fitting framework, ideally with an example plot, maybe with some arrows*) How does the fitting framework work? It's based on the framework developed by Mu Wei for the Light Dark Matter analysis (ref.) which was developed together (is this fair?). Basic description of the framework.

Also this framework is used for both LDM and NuMM together. It is trivial to simply switch between including the NuMM or LDM in it. This was done to save space in creating predictions since our backgrounds are exactly the same (or at least they should be...). Theoretically this could be separated into two difference frameworks.

- <NDPredictionSingleElectron> Prediction class which holds the LDM as a special 2-D spectrum (not used for NuMM), and NuMM, ν-on-e background, ν_eCC MEC background and other background as simple 1-D spectra. Also scaling each spectra by...

- <NDPredictionSystSingleElectron> class derived from PredictionInterp that takes in the NDPredictionSingleElectron and applies systematic shifts to it. Includes the interpolation/extrapolation between the systematic shifts.

- FitVariables and what do they do

- Fitter which does exactly what... What are the parameters of the fit? What are the results/outputs?

5 Results / Fake data studies

[to do]: talk about the shape-only versus normalisation-only strategies for the differential versus single-bin analyses.

I think it might be a good idea to talk about this already in the analysis overview section.

I need to also discuss fake data studies. Should I talk about this now? Probably

5.1 Counting experiment

Single bin analysis - total numbers of signal and backgrounds predicted. Possible background corrections and its effect on the background uncertainties.

5.2 Binned experiment

Plots showing delta chi2 for stats and systs separately. Plot showing the chi2 with limits. Maybe talk about using different binning and variables and the effect.

Also need to somehow quantify the various fitter settings, like different seeding values, different range of profiling. Should I talk about Feldman-Cousins?

5.2.1 Sensitivities and limits

6 Conclusion

(TO DO: *Report the limit with its uncertainty*)

(TO DO: *Very briefly discuss differences with current world limit and how the techniques differ*)

(TO DO: *Very briefly summarise expectations for future measurements.*)

References

[1] Carlo Giunti and Alexander Studenikin. Neutrino electromagnetic interactions: A window to new physics. *Rev. Mod. Phys.*, 87:531–591, Jun 2015. URL: `https://link.aps.org/doi/10.1103/RevModPhys.87.531`, `doi:10.1103/RevModPhys.87.531`.

[2] Nicole F. Bell, Mikhail Gorchtein, Michael J. Ramsey-Musolf, Petr Vogel, and Peng Wang. Model independent bounds on magnetic moments of Majorana neutrinos. *Phys. Lett. B*, 642:377–383, 2006. `arXiv:hep-ph/0606248`, `doi:10.1016/j.physletb.2006.09.055`.

[3] A. O. Barut and Z. Z. Aydin. Angular distribution in electron-neutrino scattering and the anomalous magnetic moment of the neutrino. *Nuovo Cim. A*, 101:677–682, 1989. doi:10.1007/BF02848090.

[4] P. Vogel and J. Engel. Neutrino electromagnetic form factors. *Phys. Rev. D*, 39:3378–3383, Jun 1989. URL: https://link.aps.org/doi/10.1103/PhysRevD.39.3378, doi:10.1103/PhysRevD.39.3378.

[5] E. Aprile et al. Search for New Physics in Electronic Recoil Data from XENONnT. *Phys. Rev. Lett.*, 129:161805, Oct 2022. URL: https://link.aps.org/doi/10.1103/PhysRevLett.129.161805, arXiv:2207.11330, doi:10.1103/PhysRevLett.129.161805.

[6] Amir N. Khan. Light new physics and neutrino electromagnetic interactions in XENONnT. *Phys. Lett. B*, 837:137650, 2023. arXiv:2208.02144, doi:10.1016/j.physletb.2022.137650.

[7] Wenjie Wu and Yiwen Xiao. Neutrino-Electron Elastic Scattering in the NOvA Near Detector - Technote. NOVA Document 56383, October 2023. NOvA technical note. URL: https://nova-docdb.fnal.gov/cgi-bin/sso/ShowDocument?docid=56383.

[8] Leonidas Aliaga Soplin. PPFX tech-note for the 2017 analysis. NOVA Document 23441, November 2017. NOvA technical note. URL: https://nova-docdb.fnal.gov/cgi-bin/sso/ShowDocument?docid=23441.

Selection	ν_eCC MEC			ν_eCC Other			ν_μCC			NC			Other		
	N	ε^{N-1}	$\varepsilon\,(\%)$	N	ε^{N-1}	$\varepsilon\,(\%)$	N	ε^{N-1}	$\varepsilon\,(\%)$	N	ε^{N-1}	$\varepsilon\,(\%)$	N	ε^{N-1}	$\varepsilon\,(\%)$
No Cut	3.50×10^{4}	100	100	3.23×10^{6}	100	100	2.24×10^{8}	100	100	3.40×10^{7}	100.	100	3.49×10^{7}	100	100
Closest Slc	3.49×10^{4}	99.78	99.78	2.95×10^{6}	91.07	91.07	2.14×10^{8}	95.45	95.45	3.16×10^{7}	92.81	92.81	2.61×10^{7}	74.77	74.77
Png Length	2.73×10^{4}	78.23	78.06	1.30×10^{6}	43.99	40.06	5.51×10^{7}	25.81	24.64	1.41×10^{7}	44.53	41.33	1.43×10^{6}	5.49	4.10
N Planes	2.73×10^{4}	99.99	78.05	1.30×10^{6}	99.99	40.06	5.51×10^{7}	99.98	24.63	1.41×10^{7}	100.00	41.33	1.43×10^{6}	100.00	4.10
N Cells	2.73×10^{4}	99.99	78.05	1.21×10^{6}	93.29	37.37	5.33×10^{7}	96.76	23.84	1.35×10^{7}	96.24	39.77	1.43×10^{6}	100.00	4.10
Fiducial	1.39×10^{4}	51.12	39.90	6.30×10^{5}	52.10	19.47	2.60×10^{7}	48.77	11.62	8.25×10^{6}	60.99	24.26	1.05×10^{6}	73.53	3.02
Cont.	9.32×10^{3}	66.82	26.66	2.63×10^{5}	41.72	8.12	7.64×10^{6}	29.38	3.42	4.96×10^{6}	60.15	14.59	9.12×10^{5}	86.62	2.61
ShwE Frac.	9.20×10^{3}	98.70	26.32	1.95×10^{5}	74.39	6.04	4.82×10^{6}	63.10	2.15	2.97×10^{6}	59.78	8.72	8.28×10^{5}	90.81	2.37
Vtx E	5.92×10^{3}	64.33	16.93	6.05×10^{4}	30.96	1.87	1.97×10^{6}	40.79	0.88	1.36×10^{6}	45.75	3.99	7.62×10^{5}	92.03	2.18
Shw Gap	5.50×10^{3}	92.91	15.73	4.62×10^{4}	76.39	1.43	1.58×10^{6}	80.18	0.70	1.06×10^{6}	77.78	3.10	5.69×10^{5}	74.61	1.63
Shw E	3.62×10^{3}	65.81	10.35	1.12×10^{4}	24.15	0.35	4.38×10^{5}	27.80	0.20	1.71×10^{5}	16.15	0.50	1.28×10^{3}	0.23	3.68×10^{-3}
Nuoneid	1.40×10^{3}	38.63	4.00	2.11×10^{3}	18.89	0.065	1.17×10^{4}	2.66	5.21×10^{-3}	8.99×10^{3}	5.27	0.026	66.43	5.17	1.90×10^{-4}
Epi0id	1.14×10^{3}	81.78	3.27	1.61×10^{3}	76.40	0.050	7.17×10^{3}	61.52	3.20×10^{-3}	4.76×10^{3}	52.94	0.014	29.47	44.36	8.45×10^{-5}
$E\theta^2$	15.13	1.32	0.043	39.00	2.42	1.21×10^{-3}	8.62	0.12	3.85×10^{-6}	20.91	0.44	6.15×10^{-5}	0.50	1.69	1.43×10^{-6}
$E\theta^2$ **(sb)**	386.16	-	1.10	306.55	-	9.48×10^{-3}	165.59	-	7.40×10^{-5}	149.93	-	4.41×10^{-4}	6.24	-	1.79×10^{-5}
No ShwE	15.54	-	0.044	69.61	-	2.15×10^{-3}	68.48	-	3.06×10^{-5}	75.67	-	2.22×10^{-4}	9.49	-	2.72E-05

Figure 4: Event selection cutflow table for background components