

Properties of the FT

L09

Goal: To survey some of the properties of the Fourier transform that are useful in image processing (beyond the convolution theorem).

But first, a quick look at what a 2D Fourier coefficient means.

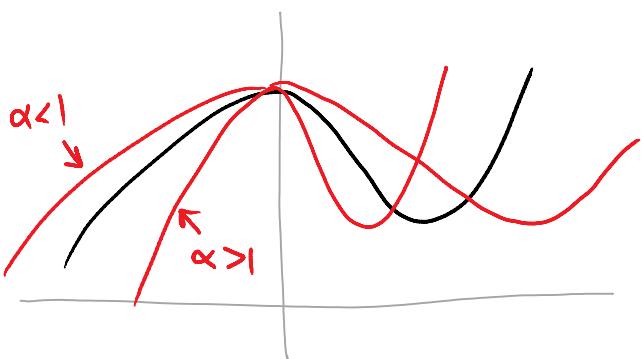
Demo: L09_wavefront.m

Scale Property

What happens to $F(\omega)$ if we scale f ?

Let $\bar{f}(x) = f(\alpha x)$

$$\begin{aligned}\bar{F}(\omega) &= \int_{-\infty}^{\infty} \bar{f}(x) e^{-2\pi i \omega x} dx \\ &= \int_{-\infty}^{\infty} f(\alpha x) e^{-2\pi i \omega x} dx\end{aligned}$$



Change of variables: $y = \alpha x \Rightarrow dy = \alpha dx$

Note: Lots of Fourier proofs involve
a change of variables.

$$\frac{dy}{\alpha} = dx$$

$$\bar{F}(\omega) = \frac{1}{\alpha} \int_{-\infty}^{\infty} f(u) e^{-2\pi i \frac{\omega u}{\alpha}} du = \frac{1}{\alpha} F\left(\frac{\omega}{\alpha}\right)$$

$$\widehat{F}(\omega) = \frac{1}{\alpha} \int_{-\infty}^{\infty} f(y) e^{-2\pi i \frac{\omega y}{\alpha}} dy = \frac{1}{\alpha} F\left(\frac{\omega}{\alpha}\right)$$

Thus, $\mathcal{F}\{f(\alpha x)\}(\omega) = \frac{1}{\alpha} \mathcal{F}\{f(x)\}\left(\frac{\omega}{\alpha}\right)$

So, as $f(x)$ stretches, $F(\omega)$ contracts & shrinks (and vice versa).

Example: Consider the Shah function with spacing Δx , $s(x)$. Recall the spacing of $S(\omega) = \mathcal{F}\{s(x)\}(\omega)$.

Exercise:

What is the spacing of $s(\alpha x)$, $\alpha \neq 0$?

What is the spacing of $S(\alpha x)$?

There is also a scale property for the DFT, though the interpretation is different. Suppose f_n and F_k are each an array of N numbers. What do those numbers mean? What points in the spatial/frequency domain do those values correspond to?

$$F_k = \sum_{n=0}^{N-1} f_n e^{-2\pi i \frac{nk}{N}}$$

One period... $0 \leq \frac{nk}{N} < N$ (*)

Notice we can introduce an arbitrary spatial scale, $0 \leq \frac{n}{N} < 1 \Rightarrow 0 \leq \frac{nL}{N} < L$
 $\Rightarrow 0 \leq x < L$

We compensate by scaling the frequency variable,

(*) $\Rightarrow 0 \leq \frac{nL}{N} \frac{k}{L} < N$

i.e. $x = \frac{nL}{N}$ for $n = 0, \dots, N-1$

$w = \frac{k}{L}$ for $k = 0, \dots, N-1$

Thus, if $x \in \left\{ \frac{0L}{N}, \frac{L}{N}, \frac{2L}{N}, \dots, \frac{(N-1)L}{N} \right\}$ ↑ "Field of view" (Fov) or period

then $w \in \left\{ \frac{0}{L}, \frac{1}{L}, \frac{2}{L}, \dots, \frac{N-1}{L} \right\}$ $\frac{N}{L}$

In general, $(FOV_x)(\Delta w) = (FOV_w)(\Delta x) = 1$

$$(FOV_x)(\Delta w) \rightarrow 1 \quad (FOV_w)(\Delta x) \rightarrow 1$$

$$\left(\frac{L}{N} \right) \left(\frac{1}{L} \right) = 1 \quad \left(\frac{N}{L} \right) \left(\frac{L}{N} \right) \rightarrow 1$$

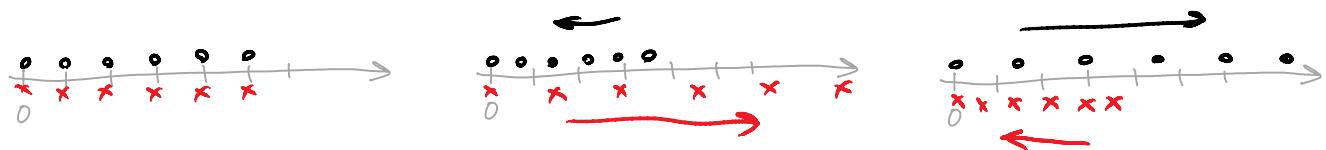
Equivalently, we could use

$$x \in \left\{ \frac{0}{L}, \frac{1}{L}, \frac{2}{L}, \dots, \frac{N-1}{L} \right\}$$

$$w \in \left\{ \frac{0L}{N}, \frac{L}{N}, \frac{2L}{N}, \dots, \frac{(N-1)L}{N} \right\}$$

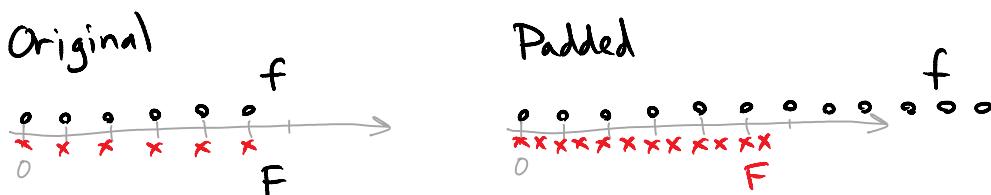
Exercise: Verify that this sampling obeys the discrete scaling rule.

Visual Examples: For 6-vectors ($N=6$) ...



Side-Effect: Scaling using padding & cropping, & the DFT
Consider this situation ...

- Pad f (add samples), then apply DFT.
Where will the samples be in the freq. domain?



In other words, padding in one domain can be thought of as **subsampling** in the other.

Demo: Image scaling using the DFT.

Shift Property

Consider f_n and its DFT F_k , $n, k = 0, \dots, N-1$.

Let $g_n = f_{n-d}$ (a shifted version of f_n).

$$G_k = \sum_{n=0}^{N-1} g_n e^{-2\pi i \frac{nk}{N}} \quad k = 0, \dots, N-1$$

$$= \sum_{n=0}^{N-1} f_{n-d} e^{-2\pi i \frac{nk}{N}}$$

Change of vars: let $m = n-d$

$$= \sum_{n=0}^{N-1} f_{n-d} e^{\frac{-2\pi i nk}{N}}$$

$$= \sum_{m=-d}^{N-1-d} f_m e^{\frac{-2\pi i (m+d)k}{N}}$$

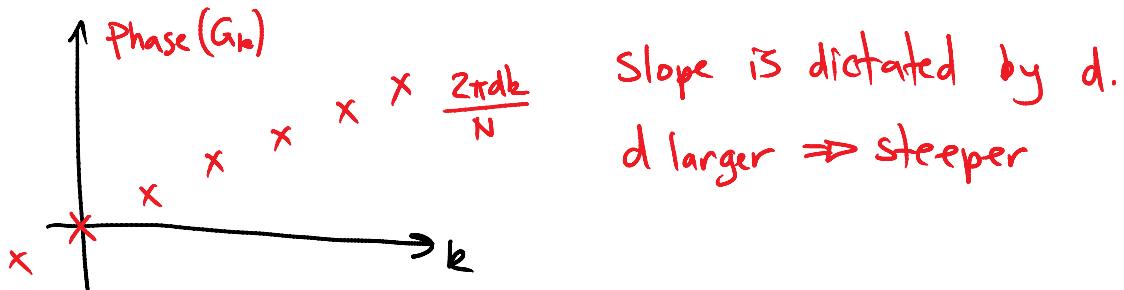
Change of vars: Let $m = n-d$
 $n = m+d$

Since f & e are N -periodic

$$= \sum_{m=0}^{N-1} f_m e^{\frac{-2\pi i mk}{N}} e^{\frac{-2\pi idk}{N}}$$

$$= e^{\frac{-2\pi idk}{N}} F_k$$

That is, shifting f by d samples is equivalent to multiplying F by a "phase ramp".



"Phase" because it only influences the phase of the Fourier coeffs.

This property is also used quite often in MRI.

Demo: Fourier Shift

Rotational Invariance

If R is a rotation matrix, then $\bar{f}(\vec{x}) \equiv f(R\vec{x})$ is a rotated version of f . What does $\mathcal{F}\{\bar{f}(\vec{x})\}(\omega)$ look like?

$\mathcal{F}\{f(\vec{x})\}(\omega)$ look like?

As it turns out

$$\mathcal{F}\{f(R\vec{x})\}(\omega) = F(R\vec{\omega})$$

We won't prove this here, but it's quite simple. Thus, rotating f is equivalent to **rotating F by the same amount**.

This works for the FT, and also for the DFT with a few minor caveats.

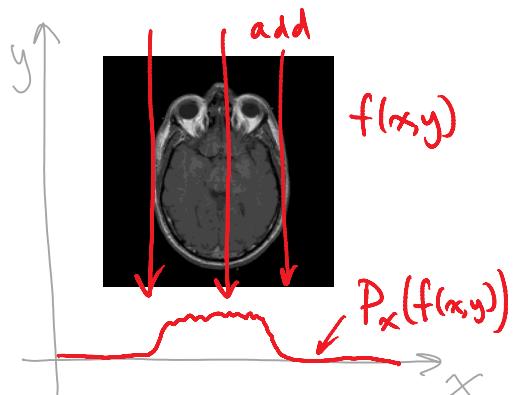
Demo: Fourier rotation

Fourier Projection Theorem

Suppose you have $f: \mathbb{C}^2 \rightarrow \mathbb{C}$, and $F(\omega, \lambda) = \mathcal{F}\{f(x, y)\}(\omega, \lambda)$.

Consider the Radon projection of f onto the x -axis,

$$P_x\{f(x, y)\}(x) = \int_{-\infty}^{\infty} f(x, y) dy$$



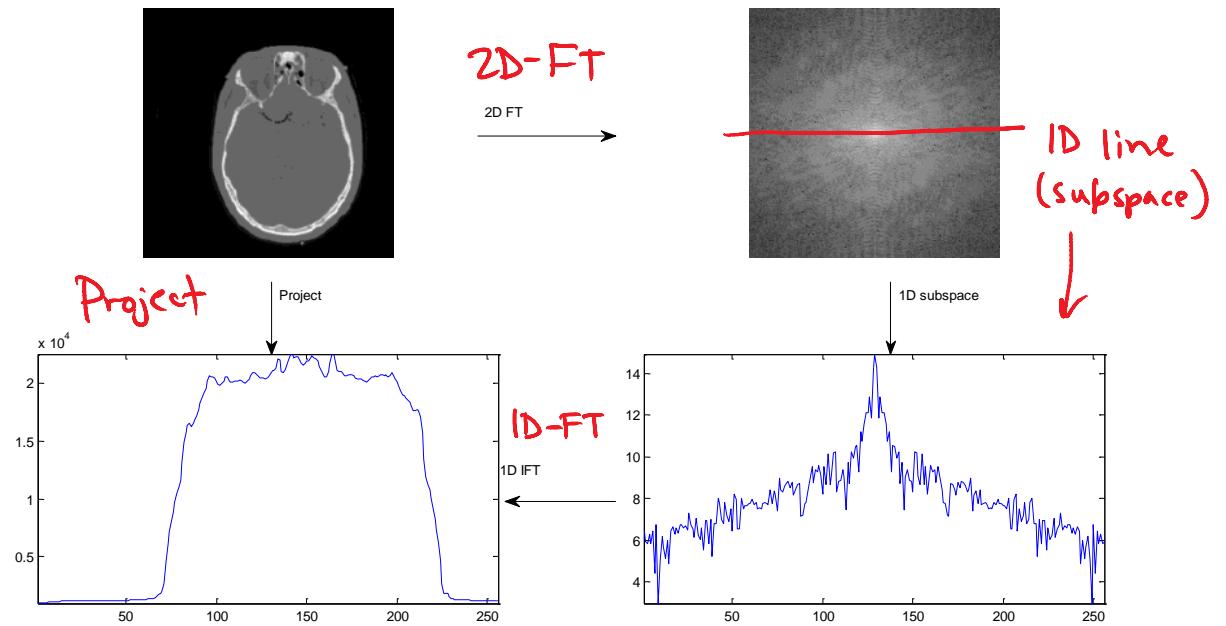
For no apparent reason, let's take the 1D FT of P_x .

$$\int_{-\infty}^{\infty} P_x\{f(x, y)\}(x) e^{-2\pi i \omega x} dx = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) e^{-2\pi i \omega x} dy dx$$

$$\begin{aligned}
 \int_{-\infty}^{\infty} \mathcal{P}_x \{f(x,y)\}(x) e^{-j\omega x} dx &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) e^{-j\omega x} dy dx \\
 &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) e^{-2\pi i(\omega x + \lambda y)} dy dx \quad \text{but } \lambda = 0. \\
 &= F(\omega, 0)
 \end{aligned}$$

Thus, the 1D-FT of the projection is the same as the 1D line of Fourier coeffs taken from the 2D-FT.

In a picture...



This works for a projection along any direction, as long as the subspace in the freq. domain has the same orientation.