

# Surface analysis of lanthanum strontium cobalt oxides under cathodic polarization at high temperature through *operando* total-reflection X-ray absorption and X-ray fluorescence spectroscopy

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## ABSTRACT

The surface chemistry of perovskite-type lanthanum transition metal oxides is an important parameter in high-temperature electrochemical devices such as solid oxide fuel cells (SOFCs). Strontium segregation on the surface of lanthanum strontium cobalt oxides (LSC) may lead to poor oxygen reduction kinetics. However, previous studies have not sufficiently addressed the oxide surface chemistry, including the oxidation change of LSC under cathodic polarization. In this study, *operando* total-reflection X-ray fluorescence absorption and X-ray fluorescence spectroscopy were performed to investigate cathodic polarization at high temperature (773 K) and at 1 atm pressure. Dense thin-film electrodes of  $\text{La}_{0.6}\text{Sr}_{0.4}\text{CoO}_3$  (LSC<sub>113</sub>) on  $\text{Zr}_{0.92}\text{Y}_{0.08}\text{O}_{1.96}$  electrolytes with a gadolinium-doped ceria interlayer were examined as model SOFC cathodes. This may be the first time that predominant surface reduction of cobalt is directly observed by using LSC<sub>113</sub> under cathodic polarization. This phenomenon is related to strontium enrichment on the LSC surface.  $(\text{La}_{0.5}\text{Sr}_{0.5})_2\text{CoO}_4$  decoration on LSC<sub>113</sub> enhanced the oxygen reduction kinetics compared with that of bare LSC<sub>113</sub>. This is due to the suppression of cobalt surface reduction and Sr enrichment in the LSC<sub>113</sub> phase.

## 1. Introduction

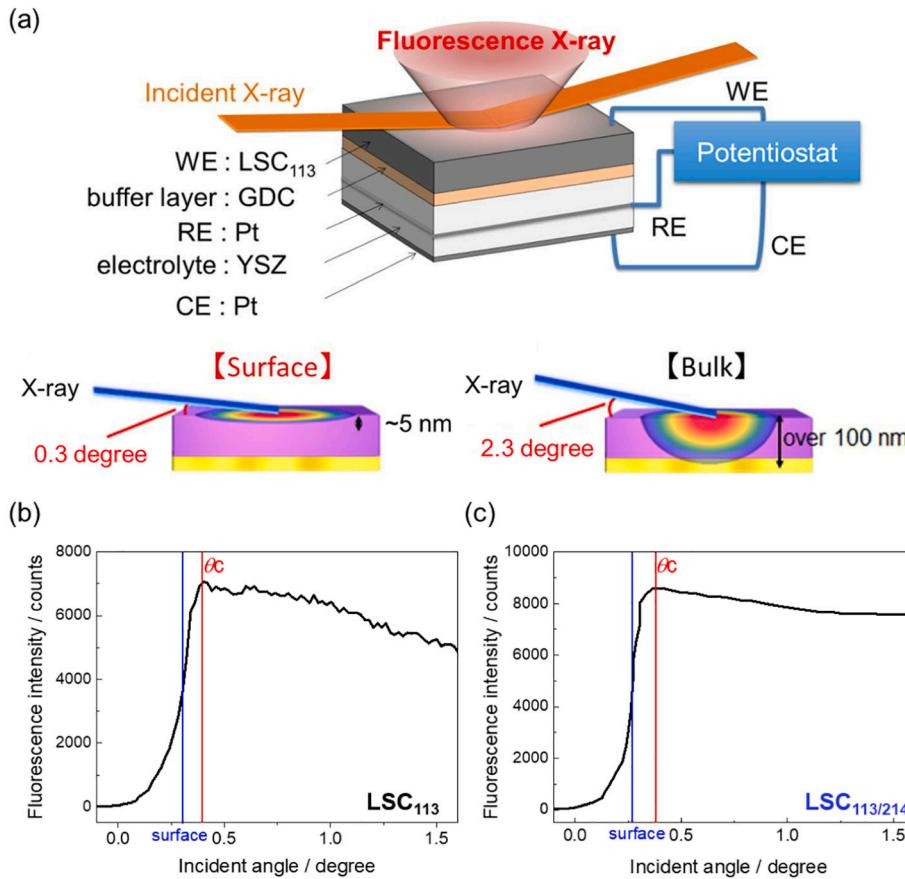
Reducing the operating temperature of solid oxide fuel cells (SOFCs) is advantageous because it reduces the cell operating cost and improves durability, which can broaden the practical applications of SOFCs [1]. However, kinetics of the cathode reaction considering the conversion of the oxygen gas to the oxide ion, and the subsequent ion transport into the electrolyte, is diminished at low temperature [2]. In order to enhance cathode kinetics, the cathode reaction mechanism at high temperature must be elucidated. Oxygen reduction on the electrode surface is one of the primary reactions that occur in the SOFC cathodes [3]. For cathode reactions using typical perovskite or Ruddelesden–Popper oxide cathodes, except for lanthanum manganese oxides, oxygen reduction including adsorption and ionization is the rate-limiting step [4–8]. The importance of cathode surface reactions in

SOFCs motivated us to investigate the surface chemistry of cathode materials. Because such surface phenomena are considered to be specific under high temperature conditions, *in situ* or *operando* observation of the nanosurface can lead to a novel approach for understanding cathode reactions [9].

$\text{La}_{1-x}\text{Sr}_x\text{CoO}_{3-\delta}$  and their related perovskite oxides are representative cathode materials and have been considered as potential cathode materials for SOFCs because of their high ionic and electric conductivity properties [10,11]. Previous studies related to the surface analysis of perovskite oxides have identified the formation of a Sr-enriched phase (*i.e.*, Sr segregation) at high temperature [12–23], which induces compositional changes in cathode materials and impedes the kinetics of cathode reactions. During SOFC operation, the cathodic polarization of the electrodes also influences the surface chemistry. Sr segregation reduction is an efficient path to enhance the surface cathode reactions.

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**Fig. 1.** (a) Schematic of the setup for data collection in the *operando* total-reflection XAS cell. Fluorescence counts as a function of incident angle of (b) LSC<sub>113</sub> and (c) LSC<sub>113/214</sub>. The penetration depth was denoted as “surface”(blue font). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

Since the enhanced oxygen exchange reaction at the heterostructured oxide interface between La<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3-δ</sub> and La<sub>2-x</sub>Sr<sub>x</sub>CoO<sub>4±δ</sub> was found [24], the improved kinetics of the cathode reaction using La<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3-δ</sub>-La<sub>2-x</sub>Sr<sub>x</sub>CoO<sub>4±δ</sub> mixed composite perovskite [25,26], La<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3-δ</sub> thin films decorated with strontium oxide, and La<sub>2-x</sub>Sr<sub>x</sub>CoO<sub>4±δ</sub> have been reported [12–15,27]. Studies based on the thin-film interface model between La<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3-δ</sub> and La<sub>2-x</sub>Sr<sub>x</sub>CoO<sub>4±δ</sub> have reported the presence of a stable phase of La<sub>2-x</sub>Sr<sub>x</sub>CoO<sub>4±δ</sub> on the surface at high temperature and the observation of a large number of active sites for oxygen dissociation-incorporation under cathodic polarization [19]. These surface mechanisms are commonly explained by X-ray photoemission spectroscopy, because it is a surface-sensitive analytical technique. Considering the oxygen potential profile under cathodic polarization at high temperature [4], the oxygen non-stoichiometry and oxidation state of cations in the electrode bulk change significantly [28], which is considered to influence the electrode surface. Moreover, the chemical analysis of cations combined with surface analysis can elucidate the mechanism of the oxygen reduction reaction (ORR) at high temperature.

To determine the oxidation state of transition metals, X-ray absorption spectroscopy (XAS) is an effective technique that provides information on the electronic and local structures of absorbance atoms [9]. It is possible to analyze the chemical state of target atoms, because XAS data are related to specific atoms. In addition, spectra in the hard X-ray region can be obtained even in the presence of oxygen gas. Therefore, XAS has been used for *in situ* analysis of SOFC cathodes under high temperature and pressure conditions [6,7,29–35]. Although standard XAS reflects bulk information, the use of geometry setups including grazing incident X-ray diffraction [36] or low-angle X-ray fluorescence

[37,38] can provide surface-sensitive spectra. We have previously detected oxidation changes related to the surface segregation phenomena of La<sub>0.8</sub>Sr<sub>0.2</sub>CoO<sub>3</sub> thin films at high temperature by depth-resolved XAS [39].

In this study, the surface oxidation state of transition metals in lanthanum strontium cobalt oxides (LSC) was determined during cathodic polarization and compared with that of the bulk state. La<sub>0.6</sub>Sr<sub>0.4</sub>CoO<sub>3-δ</sub> thin-film electrodes on yttrium-stabilized zirconium (YSZ)-poly crystalline substrates with a gadolinium-doped ceria interlayer were prepared by pulse laser deposition (PLD), followed by decoration with (La<sub>0.5</sub>Sr<sub>0.5</sub>)<sub>2</sub>CoO<sub>4±δ</sub> of various thicknesses. Oxygen reduction kinetics for the electrodes were examined by electrochemical polarization and impedance measurements. The surface oxidation state was analyzed using *operando* total-reflection fluorescence X-ray absorption spectroscopy (TRF-XAS), which integrates the fluorescence yield obtained under total reflection and allows the investigation of surface reactions at the nanoscale under cathodic polarization. Furthermore, by conducting *operando* (TRF-XAS)-based XRF analysis, elemental analysis of LSC was performed under the operating conditions. Through the combination of these techniques, the electronic structure and elemental content change that occurred on the electrode surface of the LSC thin-film electrode were determined.

## 2. Experimental

### 2.1. Sample preparation

LSC [La<sub>0.6</sub>Sr<sub>0.4</sub>CoO<sub>3-δ</sub> and (La<sub>0.5</sub>Sr<sub>0.5</sub>)<sub>2</sub>CoO<sub>4±δ</sub>] thin films (the oxides are defined as LSC<sub>113</sub> and LSC<sub>214</sub>, respectively) were prepared by PLD

with a YAG laser ( $\lambda = 256$  nm). The target samples were synthesized through a conventional solid-state reaction and a complex polymerization process. To prepare LSC<sub>113</sub>, stoichiometric amounts of La<sub>2</sub>O<sub>3</sub>, CoO, and Sr(NO<sub>3</sub>)<sub>2</sub> (Wako Pure Chemical) were mixed in a mortar. Subsequently, calcination was performed at 1423 K for 5 h. The resulting LSC<sub>113</sub> powder was pressed into pellets and sintered at 1523 K for 10 h under air. To prepare LSC<sub>214</sub>, La<sub>2</sub>O<sub>3</sub> (Kojundo Chemical) was pre-dried at 1273 K under air and dissolved in nitric acid. Then, stoichiometric amounts of Co(NO<sub>3</sub>)<sub>2</sub>·6H<sub>2</sub>O and Sr(NO<sub>3</sub>)<sub>2</sub> (Wako Pure Chemical) were added until dissolved. Citric acid and ethylene glycol were added in excess to the solution. The solution was gradually fired at 423 K to obtain a gel. Organic residuals were removed by heating at 673 K, and this precursor was calcined at 1273 K for 10 h. Gadolinium-doped ceria (GDC, Rhodia, Japan) was used as the buffer layer. GDC powder was pressed into pellets and sintered at 1473 K for 10 h under air. Separately, YSZ powder was pressed into pellets and sintered at 1773 K for 10 h under air. The sintered YSZ pellets were mirror-polished and used as the substrate; the base pressure in the chamber was  $<10^{-5}$  Pa. Deposition under 10 Hz repetition was performed for 5 min using a GDC target and for 15 min for the LSC<sub>113</sub> thin film at a substrate temperature of 1173 K. Decorated deposition of LSC<sub>214</sub> was performed on LSC<sub>113</sub> through numerous pulses (25–2700 pulses) at a substrate temperature of 1173 K. All film samples were characterized by X-ray diffraction (XRD, Rigaku RINT2200) with Cu-K $\alpha$  (40 kV, 40 mA). To perform *operando* TRF-XAS and XRF, an (110)-oriented YSZ single crystal was used as a substrate for PLD because a thin film sample with a flat and smooth surface was required. The PLD preparation conditions were similar for each thin film when using a polycrystalline YSZ substrate, and the LSC<sub>214</sub> thin film obtained with 900 pulses was selected because it showed the highest ORR activity. The thin films prepared on (110)-oriented YSZ single-crystal substrates were characterized by XRD, atomic force microscopy (AFM), and transmission electron microscopy (TEM).

## 2.2. Electrochemical testing

A three-electrode cell was constructed to determine the electrochemical performance (Fig. S1) at 1073 K and under the open circuit condition (OCV) at oxygen partial pressure [ $p(O_2)$ ] of  $10^5$  Pa. The working electrode (WE) consisted of LSC<sub>113</sub> decorated with LSC<sub>214</sub>. The counter electrode (CE) and reference electrode were porous platinum. Electrochemical impedance spectroscopy (EIS) data and direct current (DC) polarization of the film electrodes were respectively conducted using an impedance analyzer and a potentiostat (Solartron 1255 and 1287, respectively). Impedance measurements were conducted over the frequency range from 1 MHz to 0.1 Hz with an amplitude alternating current voltage of 10 mV.

## 2.3. Operando XAS and XRF analysis

**A three-electrode cell using thin films was prepared to perform *operando* TRF-XAS and XRF at the BL01B1 beamline (SPring-8, JASRI, Japan).** Co K-edge X-ray absorption near-edge spectroscopy (XANES) was conducted under the fluorescence mode using a 19-element Ge solid-state detector (SSD). A schematic of the TRF-XAS setup of the cell is shown in Fig. 1a. XAS and XRF measurements were performed at 773 K and  $p(O_2)$  of  $10^5$  Pa under OCV. In addition, the measurements under constant polarizations of  $-100$  and  $-900$  mV (WE–CE) were conducted. Figs. 1b and c show the fluorescence detector counts as a function of the incident angle for LSC<sub>113</sub> and LSC<sub>113</sub> with 3 nm of LSC<sub>214</sub> decoration (denoted as LSC<sub>113/214</sub>). The calculated critical angle of total reflection ( $\theta_c$ ) for LSC<sub>113</sub> was  $0.375^\circ$ . The highest count was observed at  $0.40^\circ$  and  $0.38^\circ$  for LSC<sub>113</sub> and LSC<sub>113/214</sub>, respectively, which was consistent with the calculated angle. Because the incident angle for the surface region of the sample was required to be smaller than  $\theta_c$ , we set half of the SSD counts of the highest counts as the surface measurement angle. When the incident angle was set at  $0.30^\circ$ , the penetration depth was estimated to

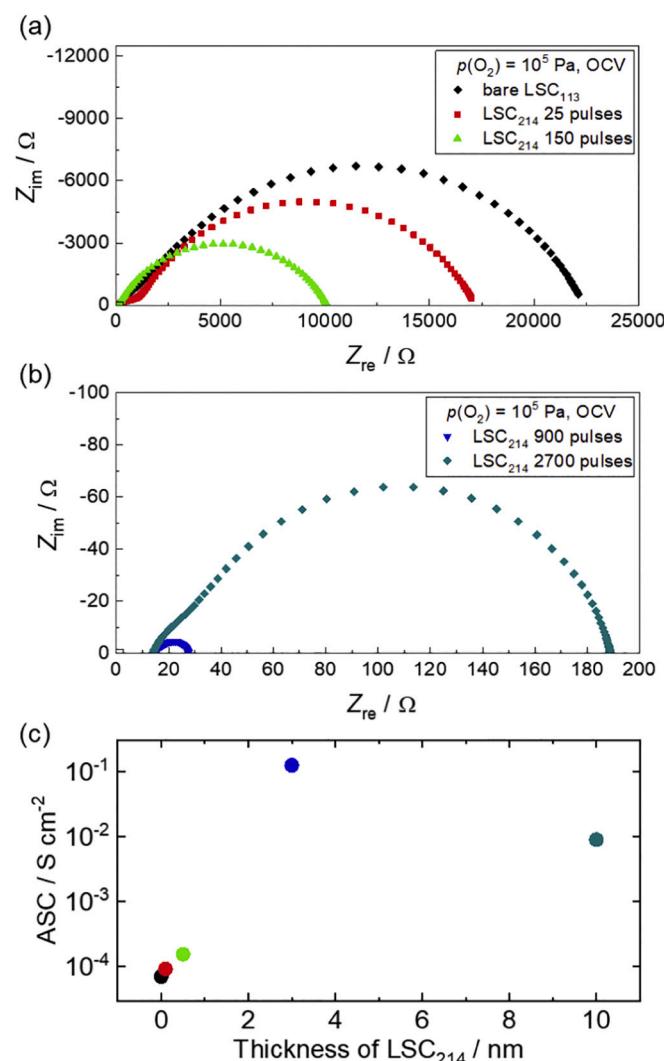
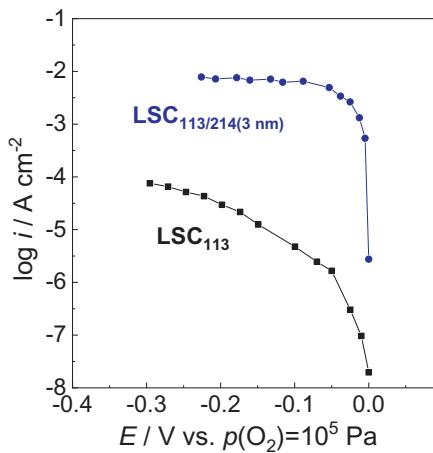


Fig. 2. Nyquist plot from EIS at 1073 K and  $p(O_2)$  of  $10^5$  Pa for (a) bare LSC<sub>113</sub> and LSC<sub>214</sub>-decorated LSC<sub>113</sub> (25 and 150 pulses); (b) LSC<sub>214</sub>-decorated LSC<sub>113</sub> (900 and 2700 pulses). (c) Area specific interfacial conductivity as a function of LSC<sub>214</sub> layer thickness.

be  $\sim 5$  nm (denoted as “surface” in Figs. 1b and c). The bulk information was obtained by setting the incident angle at  $2.3^\circ$ ; under that condition, the penetration depth was estimated to be  $>100$  nm. Considering that the thickness of LSC<sub>113</sub> was  $\sim 60$  nm, information regarding this film could be obtained. XRF spectra were also collected at an incident energy of 16.5 eV.

## 3. Results and discussion

The crystal structures of each prepared thin film on a SiO<sub>2</sub> glass substrate were characterized. XRD patterns of GDC (the buffer layer) and LSC<sub>113</sub> and LSC<sub>214</sub> thin films on SiO<sub>2</sub> are shown in Fig. S2. All peaks were identified on the basis of the reported crystal structures for GDC, LSC<sub>113</sub>, and LSC<sub>214</sub>. The lattice constants of the thin films were consistent with the database (Table S1). We confirmed the formation of an LSC<sub>113</sub> phase for the LSC<sub>113</sub> thin film at the (110) YSZ plane, as shown in Fig. S3. AFM analysis demonstrated that LSC<sub>113</sub> and LSC<sub>113/214</sub> films at the (110) YSZ plane had a small surface roughness with root-mean-square values of  $\sim 1.32$  and  $\sim 3.01$  nm, respectively (Fig. S4). The cross-sectional TEM images of the thin films are shown in Fig. S5; dense film layers with thicknesses of 60, 3, and 25 nm were obtained for LSC<sub>113</sub>, LSC<sub>214</sub>, and GDC, respectively. These results confirmed that the



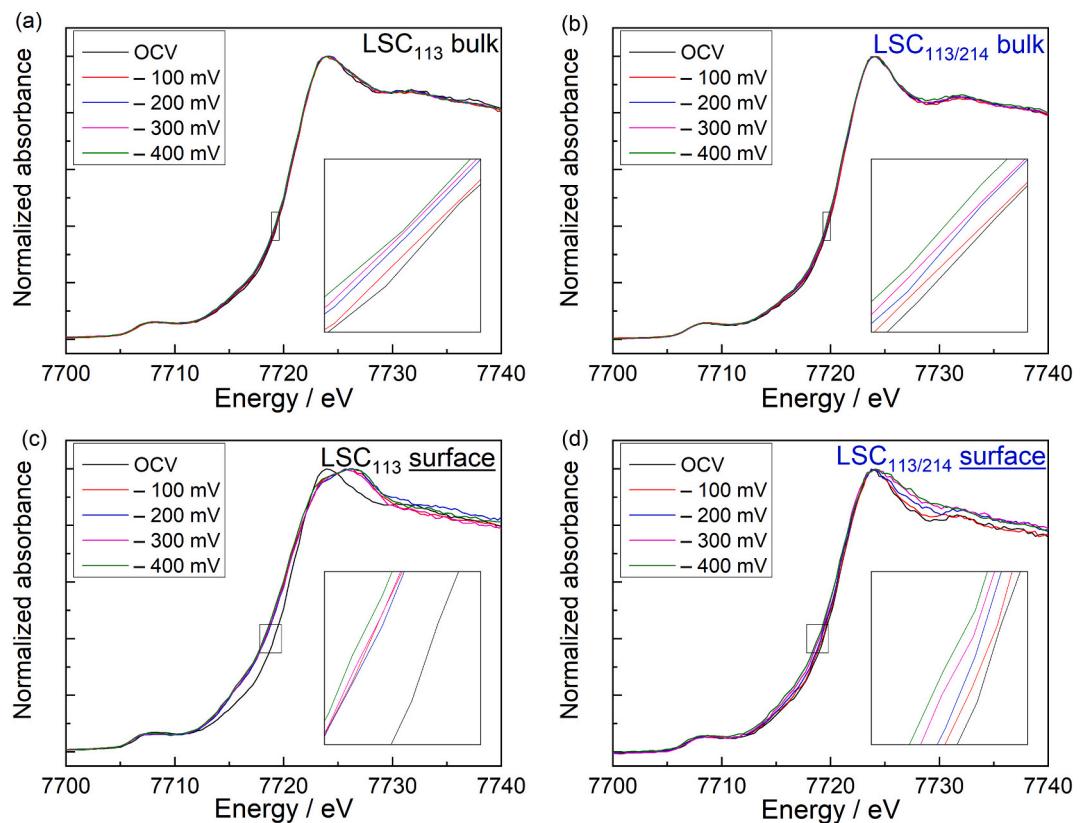
**Fig. 3.** Current–potential curve for  $\text{LSC}_{113}$  and  $\text{LSC}_{113/214}$  on YSZ at 1073 K and  $p(\text{O}_2)$  of  $10^5$  Pa.

prepared thin films could be used as model electrodes to investigate their surface phenomena *via operando* TRF-XAS.

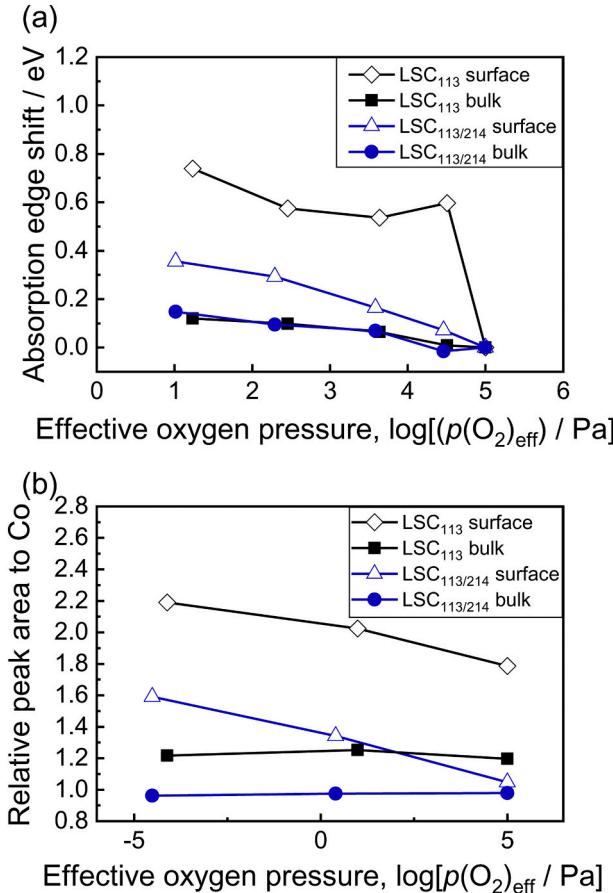
Previous studies have reported enhanced ORR kinetics through the decoration of  $\text{LSC}_{113}$  with  $\text{LSC}_{214}$  [13,16]. Similarly, in our thin film study, this kinetics enhancement using  $\text{LSC}_{214}$  decoration was also observed through EIS and DC polarization measurements. Fig. 2 shows Nyquist plots and the calculated area specific conductivity obtained from EIS measured at 1073 K and  $p(\text{O}_2)$  of  $10^5$  Pa. It can be observed that  $\text{LSC}_{214}$  decoration decreased the resistance related to the semicircle up to 3 nm (900 pulses of  $\text{LSC}_{214}$  deposition) and increased at 10 nm (2700 pulses). This represents an enhancement approximately three orders of magnitude higher for the decoration of 3 nm  $\text{LSC}_{214}$  compared to that for

bare  $\text{LSC}_{113}$ . However, electrode kinetics slightly decreased to 10 nm for  $\text{LSC}_{214}$ . The enhanced kinetics could also be confirmed through the DC polarization measurements for the bare  $\text{LSC}_{113}$  and the  $\text{LSC}_{113}$  with a 3 nm  $\text{LSC}_{214}$  deposition, at 1073 K and  $p(\text{O}_2)$  of  $10^5$  Pa (Fig. 3). As shown in Fig. 3, the current density for  $\text{LSC}_{113/214}$  significantly increased. These trends were consistent with those reported in previous studies [13,16]. We will discuss the surface chemistry from the results obtained by TRF-XAS and XRF using the bare  $\text{LSC}_{113}$  and the 3 nm  $\text{LSC}_{214}$  decoration on  $\text{LSC}_{113}$  sample ( $\text{LSC}_{113/214}$ ).

*Operando* Co K-edge XANES of LSC thin films enables us to determine the oxidation state obtained from electrochemical cathodic polarization. In addition, the TRF setup provides surface-sensitive information, which is compared with the bulk information. Fig. 4 shows the normalized Co K-edge XANES spectra for  $\text{LSC}_{113}$  and  $\text{LSC}_{113/214}$  under different applied potentials at 773 K and  $p(\text{O}_2)$  of  $10^5$  Pa. Bulk information for  $\text{LSC}_{113}$  and  $\text{LSC}_{113/214}$  from the standard fluorescence XAS spectra is depicted in Figs. 4a and b, respectively. When cathodic polarization increased, the absorption edge observed at 7720 eV shifted toward the lower energy side. The absorption energy of Co K-edge XANES in LSC corresponded to the oxidation state of cobalt in the electrodes, which was determined by oxygen nonstoichiometry [40]. The low energy shift observed in  $\text{LSC}_{113}$  and  $\text{LSC}_{113/214}$  indicated that cathodic polarization induced oxygen vacancy in LSC and decreased the oxygen chemical potential, which will be further discussed quantitatively in the following section. The surface-sensitive XANES spectra for  $\text{LSC}_{113}$  and  $\text{LSC}_{113/214}$  from the TRF measurements are illustrated in Figs. 4c and d, respectively. It should be noticed that even at low cathodic polarization, a drastic profile change was induced in the Co K-edge XANES spectra for  $\text{LSC}_{113}$ , while a continuous energy shift was observed in the bulk XANES data. Compared with the Co K-edge XANES spectra under OCV, the XANES spectra measured at 100 mV broadened and located at lower energy. This structural change is related to the Sr enrichment in the LSC phase



**Fig. 4.** Co K-edge XANES spectra from *operando* TRF-XAS for (a, c)  $\text{LSC}_{113}$  and (b, d)  $\text{LSC}_{113/214}$ , measured under cathodic polarization at 773 K and  $p(\text{O}_2)$  of  $10^5$  Pa. (a, b) Bulk information from the standard fluorescence XAS and (c, d) surface information under the TRF condition.



**Fig. 5.** (a) Absorption edge shift at the Co K-edge as a function of effective oxygen partial pressure for the *operando* TRF-XAS. (b) Relative peak area changes of Sr-K normalized by the peak area of Co-K for the *operando* TRF-XRF. Open squares denote the shift measured under various  $p(O_2)$  under OCV. Cross symbols represent the shift measured when applying different DC biases at  $p(O_2)$  of  $10^5$  Pa.

because of the similarity with the corresponding XANES spectra for Sr<sub>6</sub>Co<sub>5</sub>O<sub>15</sub>, which is a Sr-rich and non-stoichiometric compound [39]. In addition, the XANES spectra for LSC<sub>113/214</sub> showed a gradual increase in edge inclination, implying that Sr-rich phase formation was suppressed at low cathodic polarization. The surface XANES spectra significantly changed upon potential application. As the bulk XANES spectra showed negligible differences between LSC<sub>113</sub> and LSC<sub>113/214</sub> upon potential application, the structural difference is the nanoscale ordered

phenomena at the surface, which cannot be explained by conventional techniques in order to analyze bulk materials. The estimated edge energy shift as a function of the applied voltage is provided in Fig. S6 in the supporting information.

The edge energy shift, which is an indicator of the oxidation state, was further analyzed as a function of the effective oxygen partial pressure,  $p(O_2)_{eff}$ . The electrode polarization potentials,  $\eta$ , were calculated using the following expression:

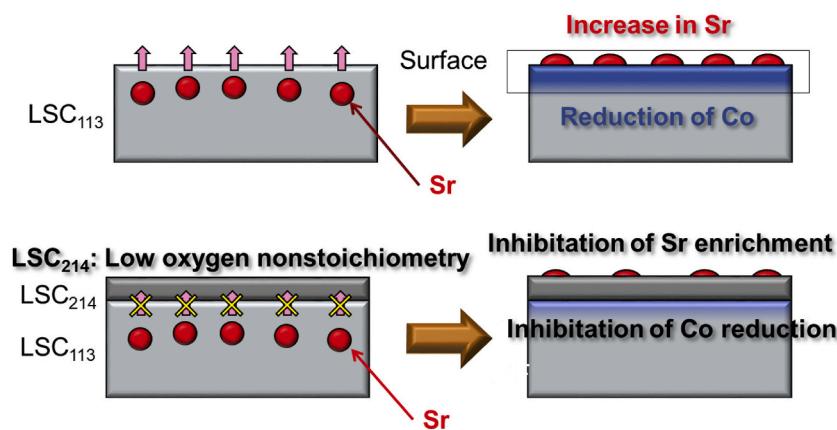
$$\eta = V_{apply} - R_{bulk}I \quad (1)$$

where  $V_{apply}$ ,  $R_{bulk}$ , and  $I$  are the voltage applied to the cell, electrolyte resistance, and current, respectively. Under polarization conditions,  $p(O_2)_{eff}$  is defined in terms of  $\eta$  as follows:

$$p(O_2)_{eff} = \exp\left\{\frac{2\mu_{O_{eff}}}{RT}\right\} = p(O_2)\exp\left\{\frac{4F\eta}{RT}\right\} \quad (2)$$

where  $R$  is the gas constant and  $F$  is the Faraday constant. Under OCV, ( $\eta = 0$ ),  $p(O_2)_{eff}$  is equivalent to  $p(O_2)$ . In this study, the application of  $-400$  mV between the two electrodes (WE–CE) corresponded to  $p(O_2)_{eff}$  that was approximately four orders of magnitude high. The energy shift calculated from Co K-edge XANES spectra as a function of  $p(O_2)_{eff}$  is shown in Fig. 5a. For bulk XANES, there was no difference between the spectral results for LSC<sub>113</sub> and LSC<sub>113/214</sub>. This demonstrated that the oxygen potential change induced by cathodic polarization is similar for LSC<sub>113</sub> and LSC<sub>113/214</sub>; thus, these LSC had the same rate-determining step in the reaction mechanism (probably surface reaction limitation as in previously reported studies [4]). The surface oxidation state exhibited an evident difference between LSC<sub>113</sub> and LSC<sub>113/214</sub>. While the energy shift on the LSC<sub>113</sub> surface was induced under small polarization, a gradual energy shift was observed for LSC<sub>113/214</sub>. The energy shift for LSC<sub>113/214</sub> even at  $p(O_2)_{eff}$  of  $10$  Pa was lower than that for LSC<sub>113</sub> at  $p(O_2)_{eff}$  of  $10^4$  Pa. As mentioned previously, the energy shift was related to Sr enrichment in the LSC phase. The difference in the edge energy shifts confirmed the suppression of that Sr enrichment by LSC<sub>214</sub> decoration. Previous reports using XPS have not clearly distinguished the surface secondary phase and LSC phases [14,16]. Because the Co K-edge TRF-XAS results directly corresponded to the oxidation state of cobalt oxides, the *operando* XAS results confirmed the Sr enrichment in LSC phases and the suppression effect of LSC<sub>214</sub> decoration.

To investigate the phenomena related to Sr enrichment in LSC phases, the elemental analysis of for LSC<sub>113</sub> and LSC<sub>113/214</sub> was conducted through *operando* TRF-XRF. The XRF spectra are provided in Fig. S7 in the supporting information. The scattered X-ray was observed at 16.5 keV. The peaks observed at approximately 4.5, 6.9, and 14.1 keV can be assigned to La-L, Co-K, and Sr-K, respectively. In addition to these peaks from LSC phases, other fluorescence peaks were detected at 7.9 and 9.1 keV. These peaks can be assigned to Hf-L from the YSZ substrates



**Fig. 6.** Schematic of the surface mechanism of LSC<sub>113</sub> and LSC<sub>113/214</sub> cathodic polarization at high temperature.

because the YSZ single crystals contain a few weight percent of Hf impurity [41]. Even for the surface-sensitive TRF configuration, the peaks from Hf-L also appeared. This is because the scattered X-ray spreads in all directions, which induces the fluorescence from the substrates. To compare the changes of the relative intensity of Sr fluorescence for LSC<sub>113</sub> and LSC<sub>113/214</sub> on the surface and bulk as a function of  $p(O_2)_{\text{eff}}$ , the peak areas of Sr normalized by Co were plotted in Fig. 5b. We note that the higher Sr ratio of LSC<sub>113</sub> bulk than LSC<sub>113/214</sub> bulk is within the experimental error. However, compared to the experimental error, we observed a more than twofold change in the area ratio with the bias applied for the surface spectra. While the peak areas of Sr at the surface for LSC<sub>113</sub> and LSC<sub>113/214</sub> increased under cathodic polarization, the peaks of the bulk did not change under polarization. For LSC<sub>113</sub>, the amount of Sr on the surface remarkably increased, which demonstrated Sr enrichment in the LSC phase because XRF detects La and Co simultaneously. Although an increase in Sr fluorescence intensity at the surface was also observed for LSC<sub>113/214</sub>, the amount of enriched Sr significantly decreased, which is also in agreement with the energy shift of *operando* Co K-edge XANES spectra.

The experimental results provided insights into the surface mechanism under cathodic polarization at high temperature, as shown in Fig. 6. Although the *operando* XAS and XRF measurements were performed at 773 K due to the limitation of thermal insulation properties of the cell, the observed phenomena are expected to be also happen at higher temperature condition. The cathodic polarization of LSC<sub>113</sub> induced the predominant Sr enrichment on the surface, which decreased ORR activity due to deviation from the ideal electronic state, such as the O p-band center [42]. The driving force of Sr enrichment is the high temperature and/or electrochemically derived reducing conditions at the electrode surface. These essential factors cannot be avoided, and the subsequent Sr migration can be disturbed by surface modification of the secondary phase such LSC<sub>214</sub>. Compared to LSC<sub>113</sub>, the amount of oxygen nonstoichiometry for LSC<sub>214</sub> was considerably low [43,44], which hindered the atomic configuration change in the crystal. When the LSC<sub>214</sub> phase was decorated on the LSC<sub>113</sub> surface, this inertial property of LSC<sub>214</sub> prevented Sr from moving toward the surface of LSC<sub>113</sub>, which preserved the ideal band structure of LSC<sub>113</sub>.

#### 4. Conclusion

Dense LSC<sub>113</sub> and LSC<sub>113/214</sub> thin film electrodes on the YSZ substrate with a GDC interlayer were prepared. From the results, it can be concluded that LSC<sub>214</sub> decoration on LSC<sub>113</sub> enhanced ORR activity by more than two orders of magnitude at high temperature. *Operando* TRF-XAS and XRF under cathodic polarization were performed. The predominant surface reduction of Co ions with Sr enrichment under cathodic polarization was observed for LSC<sub>113</sub>. This was suppressed by LSC<sub>214</sub> decoration, which improved ORR activity of the as-prepared electrode.

#### Declaration of Competing Interest

The authors have no conflicts of interest directly relevant to the content of this article.

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#### Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.ssi.2020.115502>.

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