

7/31/18 Lecture Notes: Back to Green's Functions

Last time: Associate function $f(x)$ with $\phi \mapsto \int f(x)\phi(x)dx$, a continuous, linear rule.

- Any such rule is a distribution
- The derivative of a distribution always exists and is a distribution

E.g. $\delta(x) \leftrightarrow \phi \mapsto \phi(0) (= \int \delta(x)\phi(x)dx)$

Green's Functions (in 3D)

Claim: $\Delta v = \delta(x)$, where $v = \frac{1}{4\pi r}$

Pf Green's 2nd Identity:

$$\iiint_D (u \Delta v - v \Delta u) d\vec{x} = \iint_{\partial D} \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS$$

\downarrow \downarrow \downarrow In representation formula proof
 $u(0)$ 0 if Du rep formula

Even if $Du \neq 0$, $\iiint_D u \Delta v d\vec{x} = u(0)$, so $\Delta v = \delta(x)$

Claim: $\Delta G = \delta(\vec{x} - \vec{x}_0) = \delta_{\vec{x}_0}$ in D solution is Green's function
 $G = 0$ on ∂D

Check: ii) obvious $G = 0$ on ∂D

i) obvious $\Delta G = 0$ away from \vec{x}_0

ii) $\Delta(G - v) = \Delta G - \Delta v = 0$ (If v centered at \vec{x}_0)

Alternatively, plug u, G into Green's 2nd with $Du = 0$ in D
 $u = h$ on ∂D

$$\iiint_D u \Delta G - G \Delta u d\vec{x} = \iint_{\partial D} u \frac{\partial G}{\partial n} - G \frac{\partial u}{\partial n} dS$$

$$u(\vec{x}_0) = \iint_{\partial D} h \frac{\partial G}{\partial n} dS$$

Heat Equation Claim: If $S(\vec{x}, t) = \frac{1}{(4\pi kt)^{d/2}} e^{-|\vec{x}|^2/4kt}$, then

a) $S_t = k \Delta S$ in $-\infty < x_i, t < \infty, t > 0$

b) $S(\vec{x}, 0) = \delta(\vec{x})$ at $t = 0$

Pf a) Already done

b) We showed as $t \rightarrow 0$,

$$\iiint S(\vec{x} - \vec{x}_0, t) \phi(\vec{x}_0) d\vec{x}_0 \rightarrow \phi(\vec{x}) = \iiint \delta(\vec{x} - \vec{x}_0) \phi(\vec{x}_0) d\vec{x}_0$$

Wave Equation Let $S(\vec{x}, t)$ be solution to

$$S_{tt} = c^2 \Delta S$$

$$S(\vec{x}, 0) = 0$$

$$S_t(\vec{x}, 0) = \delta(\vec{x})$$

Why $\delta(\vec{x})$ not ψ not ϕ ?

look

Claim: Solution to $u_{tt} = c^2 \Delta u$ is $u(\vec{x}, t) = \int S(\vec{x} - \vec{y}, t) \psi(\vec{y}) d\vec{y}$ familiar?
 $u(\vec{x}, 0) = \psi(\vec{x})$
 $u_t(\vec{x}, 0) = \psi(\vec{x}) + \frac{\partial}{\partial t} \int S(\vec{x} - \vec{y}, t) \phi(\vec{y}) d\vec{y}$

Pf $u_{tt} = c^2 \Delta u$ since u made from solutions

$$u(\vec{x}, 0) = \int \underbrace{S(\vec{x} - \vec{y}, 0)}_{\delta(\vec{x} - \vec{y})} \psi(\vec{y}) d\vec{y} + \int \underbrace{S_t(\vec{x} - \vec{y}, 0)}_{\delta(\vec{x} - \vec{y})} \phi(\vec{y}) d\vec{y} = \psi(\vec{x})$$

$$u_t(\vec{x}, 0) = \int \underbrace{S_t(\vec{x} - \vec{y}, 0)}_{\delta(\vec{x} - \vec{y})} \psi(\vec{y}) d\vec{y} + \int \underbrace{S_{tt}(\vec{x} - \vec{y}, 0)}_{\Delta S(\vec{x} - \vec{y}, 0)} \phi(\vec{y}) d\vec{y} = \psi(\vec{x})$$

What is S ?

$$1D: u(x, t) = \frac{1}{2} \left[\underbrace{\phi(x+ct)}_{\circ} + \underbrace{\phi(x-ct)}_{\circ} \right] + \frac{1}{2c} \int_{x-ct}^{x+ct} \underbrace{\psi(s)}_{\int(s)} ds$$

$$S(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \delta(s) ds = \begin{cases} \frac{1}{2c} & \text{if } -ct < x < ct \\ 0 & \text{else} \end{cases}$$

sgn pronounced "sighum"

$$= \frac{1}{2c} H(c^2 t^2 - x^2) \operatorname{sgn}(t), \quad \operatorname{sgn}(t) = \begin{cases} 1 & t > 0 \\ -1 & t < 0 \end{cases}$$

3D: $u(\vec{x}, t) = \frac{1}{4\pi c^2 t_0} \iint_S \psi(\vec{x}) dS + \frac{\partial}{\partial t_0} \left[\frac{1}{4\pi c^2 t_0} \iint_S \phi(\vec{x}) dS \right]$ Sorry for multiple uses of S !

$$S(\vec{x}, t) = \frac{1}{4\pi c^2 t} \delta(ct - |\vec{x}|) \text{ sign}(t) \text{ (almost matches 1D result)}$$

Q: Is $\delta(ct - |\vec{x}|) = \delta(c^2 t^2 - |\vec{x}|^2)$? Both δ 's supported where $|\vec{x}| = ct$.

Example: $\delta(2x)$: $\int \delta(2x) \phi(x) dx = \int \delta(r) \phi\left(\frac{r}{2}\right) \frac{dr}{2} = \frac{1}{2} \phi(0)$
 $y = 2x$
 $dy = 2dx$

By change of coordinates,

$$S(\vec{x}, t) = \frac{1}{2\pi c} \delta(c^2 t^2 - |\vec{x}|^2) \text{ Notice - pattern now?}$$

$$\underline{2D}: S(\vec{x}, t) = \begin{cases} \frac{1}{2\pi c} (c^2 t^2 - |\vec{x}|^2)^{-1/2} & |\vec{x}| < ct \\ 0 & |\vec{x}| > ct \end{cases}$$

- Can read off Huygen's principle (or lack thereof) by support of distribution with $c^2 t^2 - |\vec{x}|^2$

- This distribution has degree $\frac{1-d}{2}$ in dimension d , explaining difference between even and odd dimensions.

Next time: No more rephrasing. Solve PDE faster.