Tartícula en una caja

$$H = \frac{\hat{p}^2}{2m} + V(x)$$

$$\hat{H}|\phi_n\rangle = E_n|\phi_n\rangle$$

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$$\hat{H} | \phi_n \rangle = E_n | \phi_n \rangle$$

$$E_n = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 n^2, \quad n = 2, ..., \infty$$

$$\phi_n(x) = \langle x | \phi_n \rangle = \begin{cases} \frac{2}{L} \sin\left(\frac{n\pi x}{L}\right) & 0 \le x \le L \\ 0 & x \le 0, x \ge L \end{cases}$$

$$|\Psi\rangle$$
 arbitraria $\Rightarrow \Psi(x) = \langle x | \Psi \rangle$, $\langle \Psi | \Psi \rangle = 1$

$$\psi(x) = \sum_{n=1}^{\infty} \langle x | \phi_n \rangle \langle \phi_n | \psi \rangle = \sum_{n=1}^{\infty} \phi_n(x) a_n = \sum_{n=1}^{\infty} a_n \phi_n(x)$$

$$a_n = \langle \phi_n | \psi \rangle^{\ell}$$

$$a_n = \int_{-\infty}^{\infty} \langle \phi_n | \mathbf{x} \rangle \langle \mathbf{x} | \psi \rangle d\mathbf{x}$$

$$\int_{-\infty}^{\infty} \phi_n^*(x) \psi(x) dx = \int_{0}^{L} \phi_n^*(x) \psi(x) dx$$

| > |an|2: prob. de hallar el eigenvalor del sistema en 14) en de la enegía 4(x)

$$\langle \psi | \psi \rangle = \int_{-\infty}^{\infty} \langle \psi | x \rangle \langle x | \psi \rangle dx = \int_{-\infty}^{\infty} | \psi (x) |^2 dx = 1$$

$$= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} \alpha_m^* \phi_m^*(x) a_n \phi_n(x) dx$$

$$= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} a_m^* a_n \int_{-\infty}^{\infty} \oint_{m}^{*} (x) \oint_{n} (x) dx$$

$$= \sum_{n=1}^{\infty} a_n^* a_n = 1$$

$$|a_n|^2 \rightarrow \text{probabilidad}$$

$$t=0 | \psi \rangle \rightarrow | \psi(t) \rangle = \dot{c}?$$

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \hat{H} |\Psi(t)\rangle$$

$$|\Psi\rangle = \sum_{n=1}^{\infty} a_n |\phi_n\rangle, \quad a_n = \langle \phi_n | \Psi \rangle$$

$$|\psi(t)\rangle = \sum_{n=1}^{\infty} a_n e^{-iE_nt/\hbar} |\phi_n\rangle$$

$$\langle \Psi(t) | \hat{A}(\hat{x}, \hat{p}) | \Psi(t) \rangle = \sqrt{$$

 $\langle \psi(t) | \hat{\rho} | \psi(t) \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) \frac{\hbar}{i} \frac{\partial}{\partial x} \psi(x,t) dx - \frac{\lambda}{i}$

$$\psi(x,t) = \sum_{n=1}^{\infty} a_n e^{-i E_n t/\hbar} \phi_n(x)$$

$$\Rightarrow \frac{a_{\ell} = 1}{a_{n} = 0} \quad \forall n \neq \ell \quad \Rightarrow | \psi(t) \rangle = e^{-i E_{\ell} t / \hbar} | \phi_{\ell} \rangle$$

 $\langle \Psi(t) | \hat{A} | \Psi(t) | \Psi(t) \rangle = e^{i E_{\ell} t/\hbar} \langle \phi_{\ell} | \hat{A} | \phi_{\ell} \rangle e^{-i E_{\ell} t/\hbar} = \langle \phi_{\ell} | \hat{A} | \phi_{\ell} \rangle$

donde 1917 son los ESTADOS ESTACIONARIOS.

$$|a_n e^{-iE_n t/\hbar}|^2 = |a_n|^2$$
 Dado un edo. arbitrario $|\Psi(t)\rangle = \sum_{n=1}^{\infty} (a_n e^{-iE_n t/\hbar}) |\phi_n\rangle$

$$\hat{A}|\chi_n\rangle = \alpha_n|\chi_n\rangle$$
, entonces

$$|\langle \chi_{K} | \psi(t) \rangle|^{2} = \left| \sum_{n=1}^{\infty} a_{m} e^{-iE_{n}t/\hbar} \langle \chi_{K} | \phi_{n} \rangle \right|^{2}$$

$$= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} a_{m}^{*} e^{iE_{m}t/\hbar} a_{n} e^{-iE_{n}t/\hbar}$$

(\$\phi_n | x_K)(x_K | \$\phi_n\$) cambia on of tiempo

$$f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right), \quad a_n = \int_0^L \phi_n^*(x) f(x) dx$$

... i Métalo a Mothematica!

$$|\phi_n\rangle = \int_{-\infty}^{\infty} |\rho\rangle\langle\rho|\phi_n\rangle d\rho$$

$$\int_{0}^{L} \frac{1}{\sqrt{2\pi\hbar}} e^{-ipx/\hbar} \phi_{n}(x) dx$$

$$=\frac{1}{2}\sqrt{\frac{L}{\pi\hbar}}\left[(-1)^{n}e^{-i\rho L/\hbar}-1\right]\left[\frac{1}{\rho\frac{L}{\hbar}-n\hbar}-\frac{1}{\rho\frac{L}{\hbar}+n\hbar}\right]$$

$$\phi_n(x) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{ipx/\hbar} \widetilde{\phi}_n(p) dp$$