

BME2 – Biomedical Ultrasonics

Lecture 6: An Introduction to Non-Linear Effects

Dr. Constantin-C. Coussios



Office: 43 Banbury Road

Tel: 01865-(2)74750

Email: constantin.coussios@eng.ox.ac.uk

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Department of Engineering Science

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6.1 Non-linear sound propagation

- Recall that for *linear* sound propagation, we must have

$$p(\mathbf{r}, t) = p_o + p'(\mathbf{r}, t) \quad p' \ll p_o$$

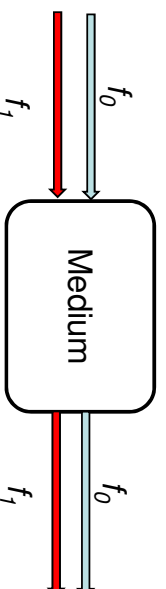
$$\rho(\mathbf{r}, t) = \rho_o + \rho'(\mathbf{r}, t) \quad \rho' \ll \rho_o$$

$$u(\mathbf{r}, t) = u'(\mathbf{r}, t) \quad |u'| \text{ small}$$

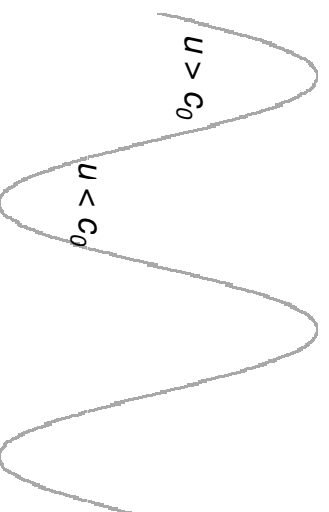
- In practice, this means that if we feed a wave at a single frequency f_o through a medium, we get



- Similarly, if the input consists of two or more signals of different frequencies:

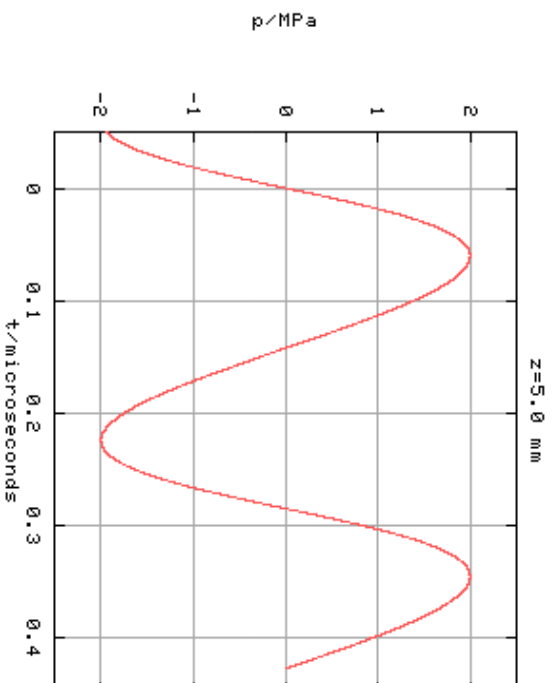


- What happens if the excitation amplitude is not small?
- The local fluid velocity is no longer small *compared to the speed of wave propagation*



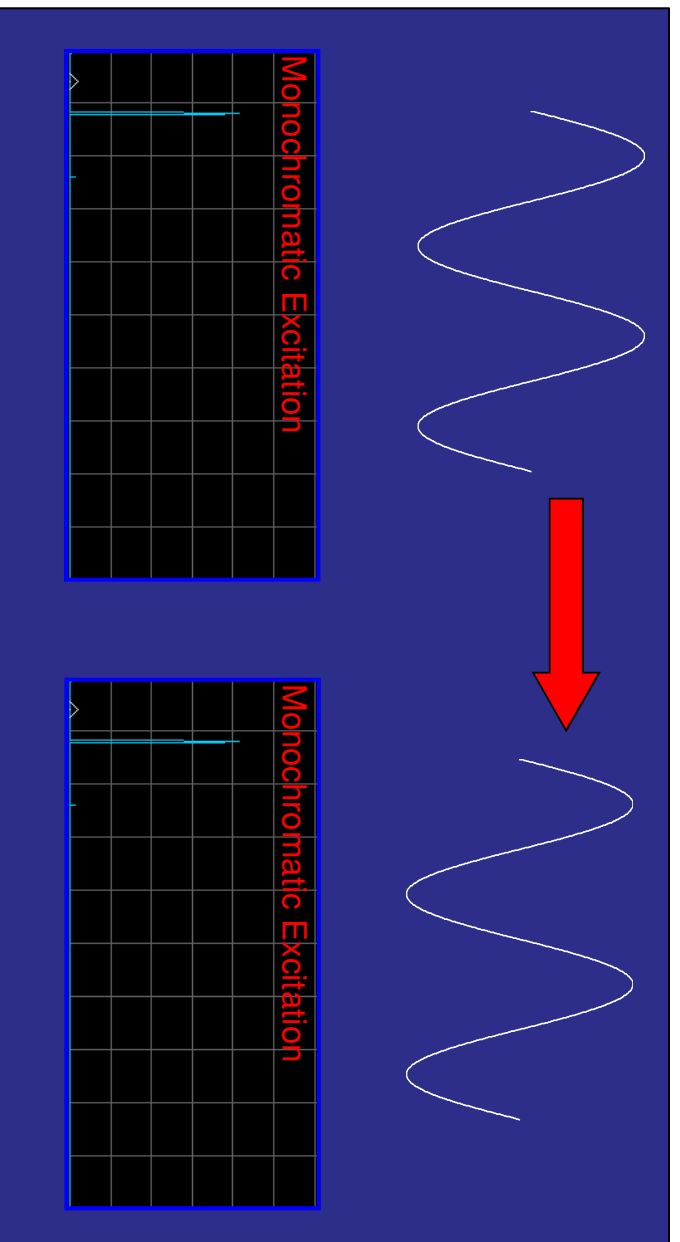
- Hence, sound in regions of positive u travels faster, whilst sound in regions of negative u travels slower.
- However, sound in regions of zero u still travels at speed c_o , meaning that the period is unchanged!

- This results in steepening of the wave and gradual shock formation:

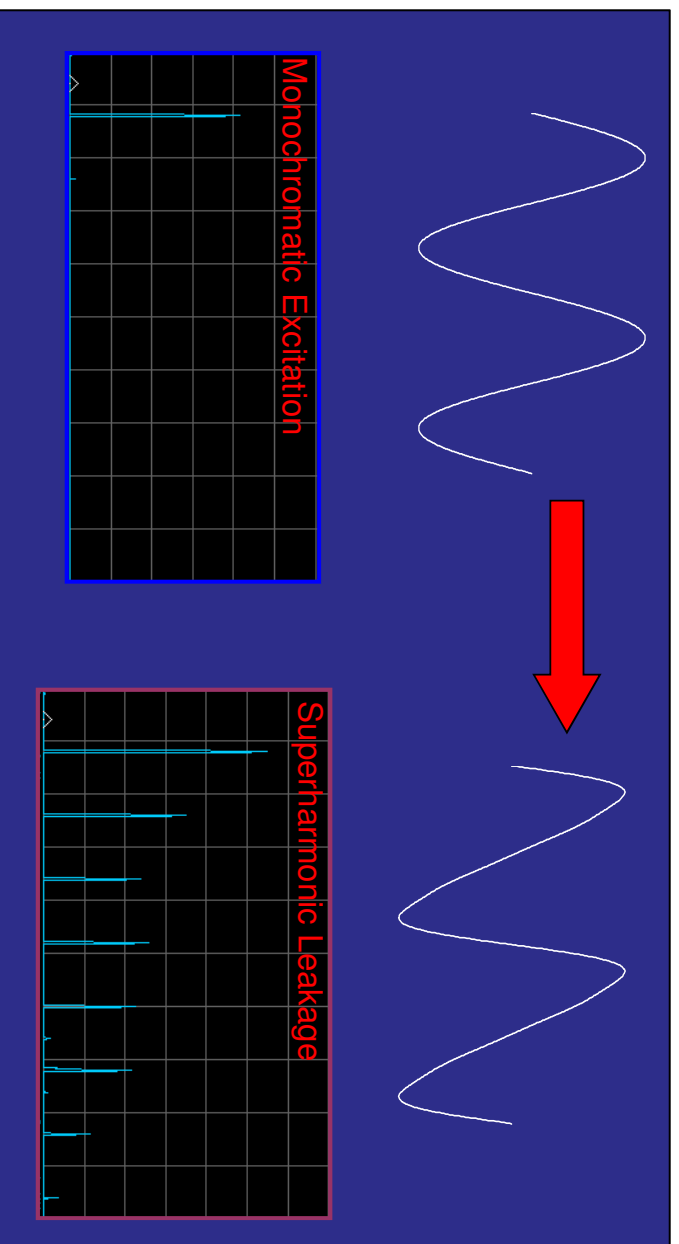


(recall that in most media attenuation is a power law of frequency).

Frequency response of a linear medium



Frequency response of a non-linear medium



- This now means that
-
- The degree of wave steepening / harmonic generation is primarily a property of the medium, described by the parameter B/A or the coefficient of non-linearity $\beta = 1 + B/2A$
- For comparison, note that
 - Liver: $B/A = 7.6$
 - Water: $B/A = 5.2$

- The propagation distance required for an initially-sinusoidal plane wave to form a shock wave is termed the shock formation distance

$$\bar{x} = \frac{\lambda}{2\pi\beta M}$$

This means that shocks form sooner at higher frequencies, in more non-linear media and at higher amplitudes M .

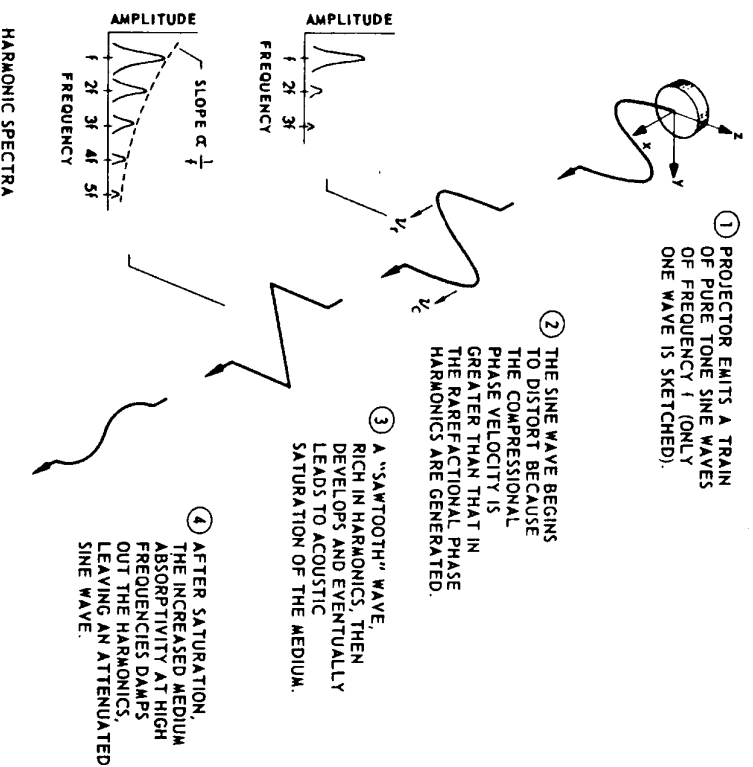
- The Westervelt equation* includes diffraction, absorption, nonlinearity and inhomogeneity

$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) p - \frac{1}{\rho} \nabla p \cdot \nabla \rho + \frac{\delta}{c^4} \frac{\partial^3 p}{\partial t^3} + \frac{\beta}{\rho c^4} \frac{\partial^2 p^2}{\partial t^2} = 0$$

where p is the propagating acoustic pressure perturbation, c is the local sound speed, ρ is the local density for the medium, δ the local acoustic [4.2] diffusivity, and β is a local coefficient of nonlinearity.

- Note that a major limitation of this equation is that it does not accurately account for the frequency dependence of attenuation.

Imaging Ultrasound is not Necessarily Linear



Imaging Ultrasound is not Necessarily Linear

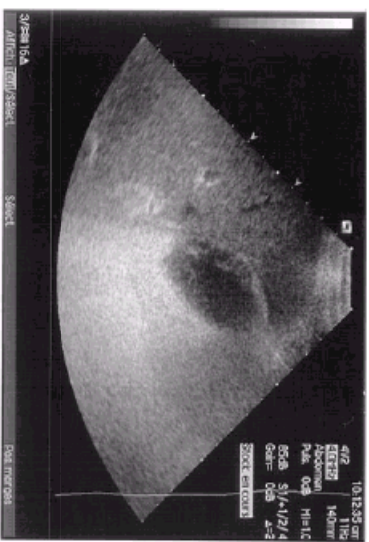
Harmonic Imaging

- Finite amplitude sound waves propagate nonlinearly through real media
- This results in the transfer of energy from the fundamental frequency to higher harmonics
- This conversion occurs further away from the source, this near field aberrating layers do not impact the formation of these harmonics.
 - **Nearfield aberration not as important at the 2nd harmonic**
- Since the conversion to harmonics is second order in the pressure amplitude, this higher frequency energy is generated along a narrower beam
 - **Better lateral resolution at the 2nd harmonic**

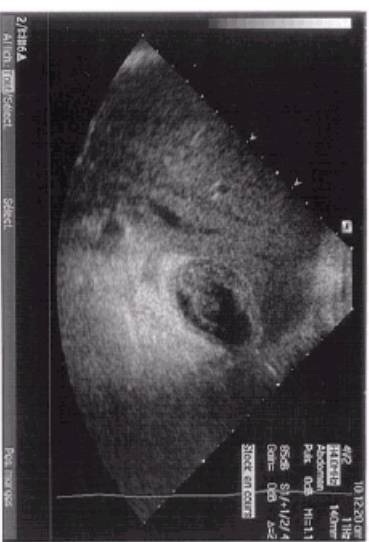


2nd Harmonic Imaging Mode

B-Mode Image of
Chronic Cholecystitis



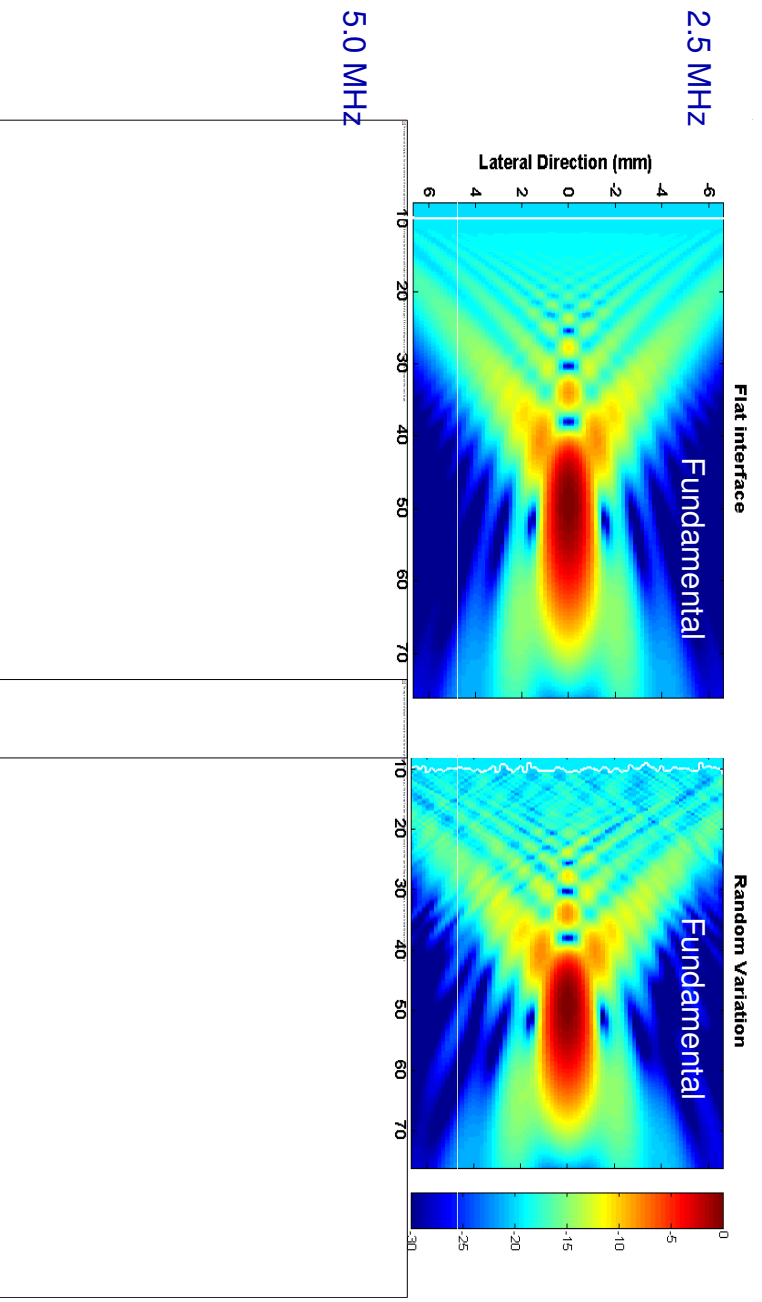
2nd Harmonic Image
Improved Detail and
Contrast



Can Resolve Contents and
Wall of the Gall Bladder



Computed Nonlinear Pressure Fields FDTD Simulations

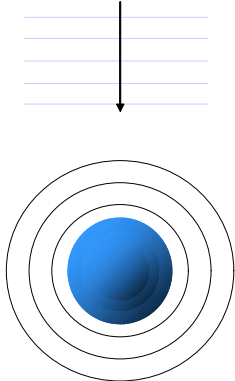


6.3 Acoustic cavitation

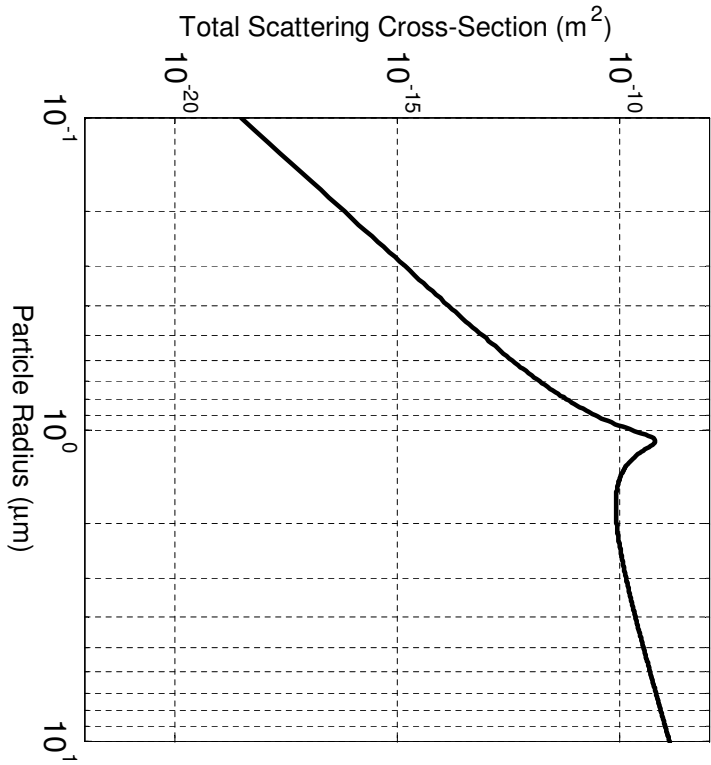
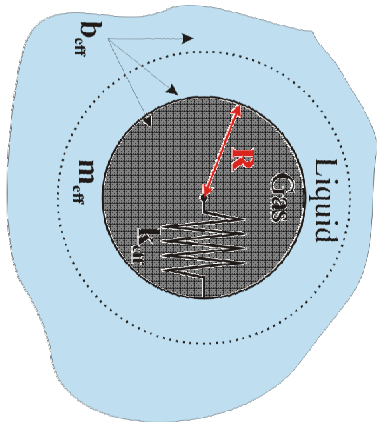
➤ Linear Bubble Behaviour

Incident

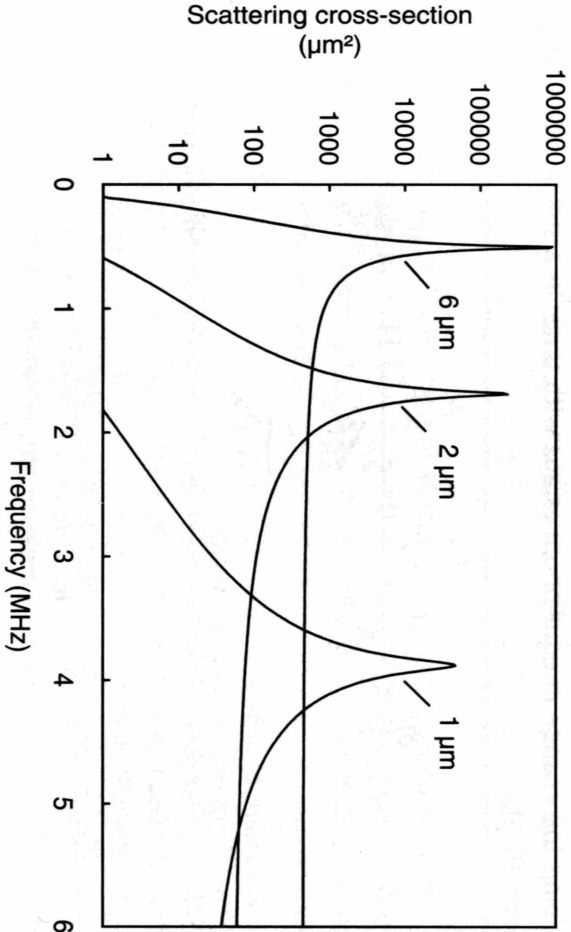
ultrasound



$$m_{eff} \ddot{r} + b_{eff} \dot{r} + k_{eff} r = F(t)$$

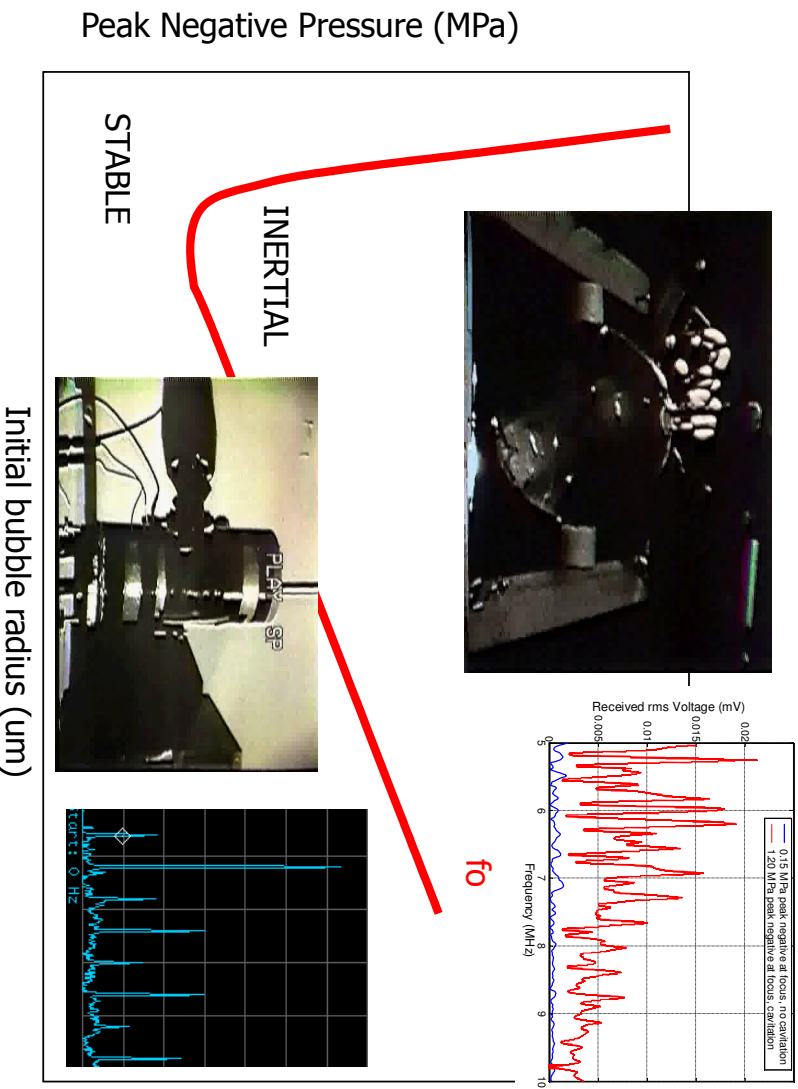


Linear bubble resonance



$$f_o R_o = \frac{1}{2\pi} \sqrt{\frac{3\rho_o}{\rho}} \approx 3.25 \text{ Hz} - m$$
$$\left\{ \begin{array}{l} 1 \text{ mm} \Rightarrow 3.25 \text{ kHz} \\ 1 \mu m \Rightarrow 3.25 \text{ MHz} \end{array} \right.$$

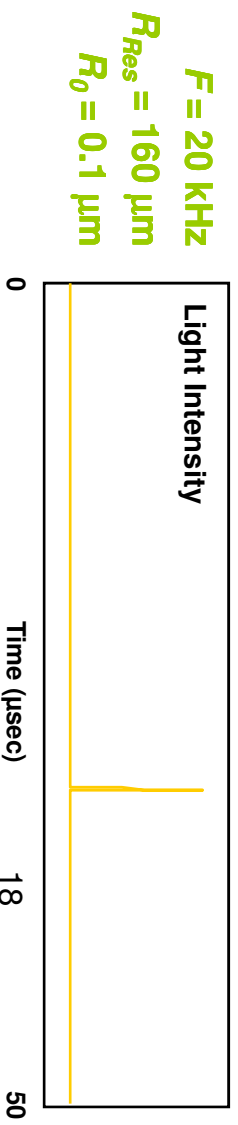
➤ Non-linear bubble behaviour: stable and inertial cavitation



Physical Effects: *Inertial Cavitation*

Inertial collapse of an acoustically forced bubble in a liquid

- Isothermal expansion
- Rapid adiabatic collapse
- Chemical dissociation
- Light production



$F = 20 \text{ KHz}$
 $R_{Res} = 160 \text{ } \mu\text{m}$
 $R_0 = 0.1 \text{ } \mu\text{m}$

Bioeffects caused by cavitation

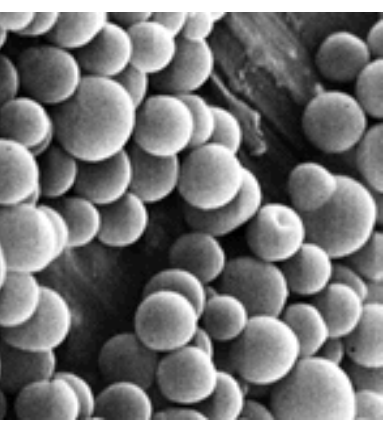


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Bubbles are excellent scatterers of sound

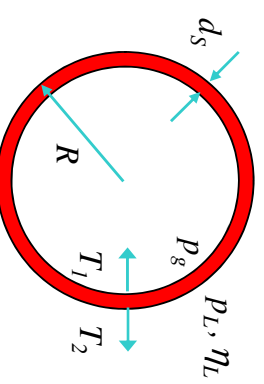
Ultrasound Contrast Agents (UCAs)

- Micron-sized bubbles injected
- Sizes range from 0.2 - 3 μm
- Must be stabilized against dissolution
 - Insoluble gases
 - Surface skins



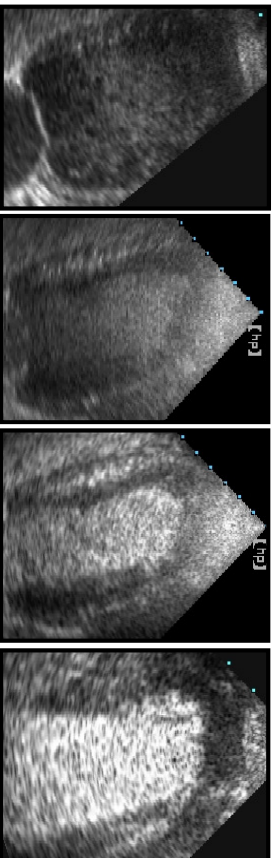
UCA's enhance backscatter from blood

- Image remote regions
- Enhance Doppler flow imaging
- Highlight vascularized regions (tumors!)
- Measure perfusion into tissues



The Role of Bubbles in Ultrasound Imaging

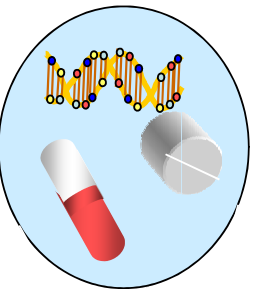
- **Bubble Contrast Imaging**
 - Adding bubbles to blood enhances echogenicity because **everyone knows** that bubbles are the most amazing scatterers of sound!
 - There are many different manifestations of bubble contrast agents -- all offer a measure of diffusional stability.
 - The stabilization mechanisms inhibit mechanical response -- regardless, they are still highly effective scattering sites.
 - The inherently nonlinear response of cavitation bubbles make them a natural for enhance harmonic imaging.



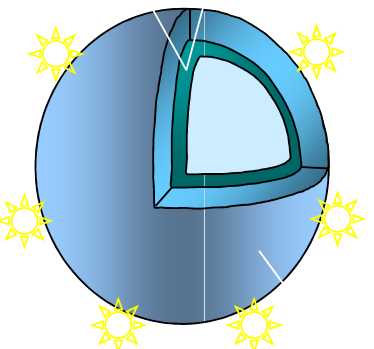
Applications

Targeted Drug Delivery

Gas core



Surfactant or polymer shell.



Targeting species may be attached to the shell surface.

- Drugs or even DNA can be incorporated into the shell and released by destroying the bubble with high intensity ultrasound at the required location.

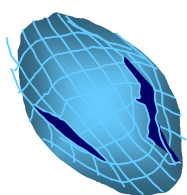
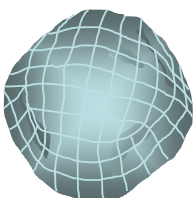


The Encapsulating Shell

Mechanisms of Destruction

Buckling

Rupture & fragmentation



Modification for Drug Delivery

- (i) Asymmetric shell
- (ii) Additional oil layer
- (iii) Stiffened phospholipid coating

