

Non-Hermitian Photonic Topological Insulators

17th International Congress on Artificial Materials for Novel Wave Phenomena



Prof. Tatiana Rappoport

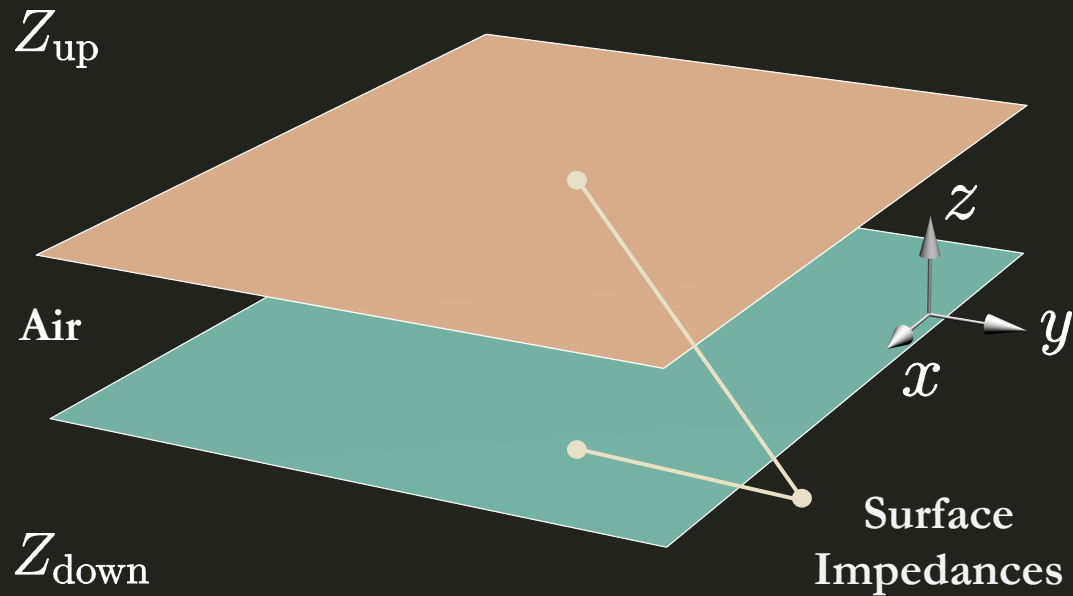


Prof. Mário Silveirinha



Reciprocal and \mathcal{PD} -Symmetric Waveguides

1



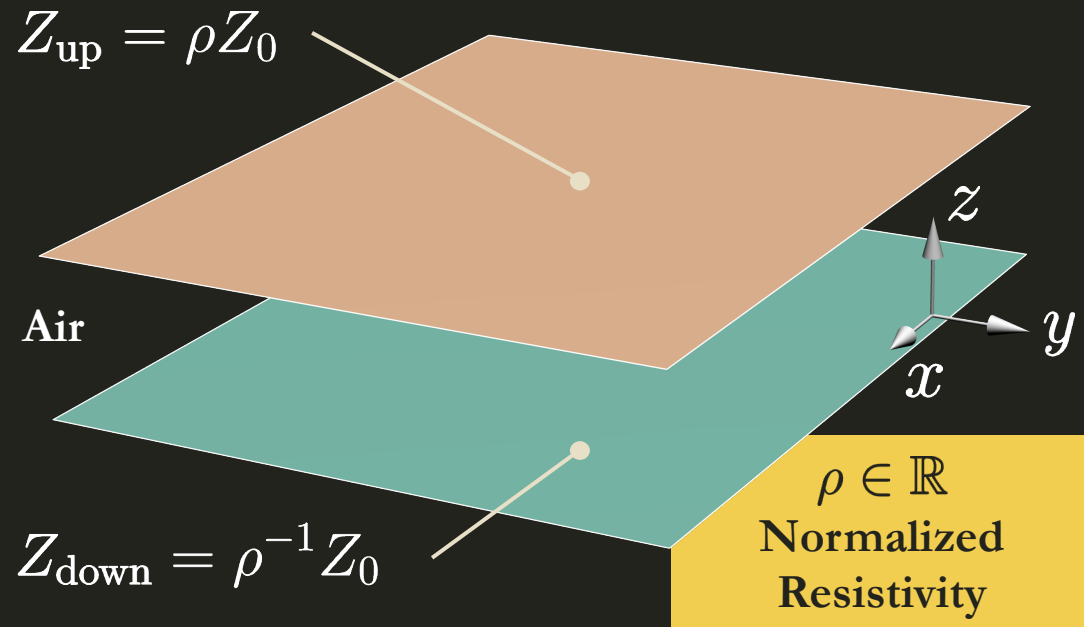
Parity Transformation \mathcal{P}
Reflection about xOy Plane
 $(x, y, z) \rightarrow (x, y, -z)$

+

Duality Transformation \mathcal{D}
Response Exchange
Permittivity $\epsilon \rightleftarrows \mu$ Permeability

Reciprocal and \mathcal{PD} -Symmetric Waveguides

1



$$\rho = 0$$

Perfect Conductors

Electric (PEC)

Magnetic (PMC)

Conservative
Hermitian
Electrodynamics

$$\rho \neq 0$$

Lossy Plates

Non-Conservative
Non-Hermitian
Electrodynamics

Reciprocal and \mathcal{PD} -Symmetric Waveguides

2

Symmetry-Induced Decoupling of Maxwell's Equations

$$\pm \frac{ic}{\epsilon(z)} \begin{pmatrix} 0 & \partial_z & \partial_y \\ -\partial_z & 0 & -\partial_x \\ \partial_y & -\partial_x & 0 \end{pmatrix} \cdot \Psi^\pm(-z) = \omega \Psi^\pm(z)$$

$$\Psi^\pm(z) = \begin{pmatrix} E_x(z) \mp Z_0 H_x(-z) \\ E_y(z) \mp Z_0 H_y(-z) \\ E_z(z) \pm Z_0 H_z(-z) \end{pmatrix}$$



\mathcal{PTD} Symmetry-Protected Scattering Anomaly in Optics

3



M. G. Silveirinha, Physical Review B 95 (2017).

\mathcal{PD} -Symmetry

+ Invariance under Time Reversal \mathcal{T}

Reflectionless Light Channels

System 1 System 2



N Modes Towards
Junction

Antisymmetric
Scattering Matrix

$$\mathbf{S} = -\mathbf{S}^\top$$

$$\mathbf{S} = \begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix}$$

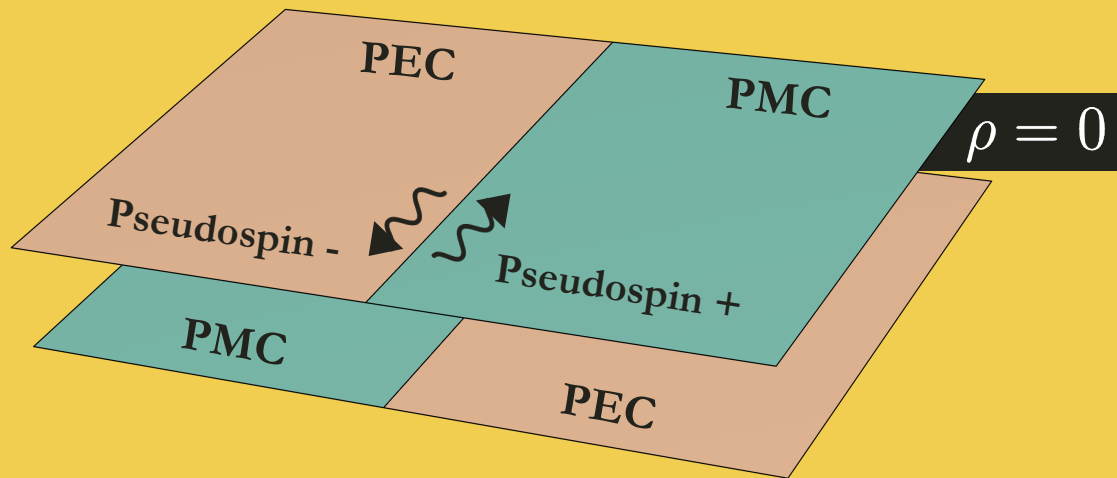
$$\det \mathbf{S}_{11} = (-1)^N \det \mathbf{S}_{11} = 0 \text{ (odd } N\text{)}$$

\mathcal{PTD} Symmetry-Protected Scattering Anomaly in Optics

4



W.-J. Chen, S.-J. Jiang, X.-D. Chen, B. Zhu, L. Zhou, J.-W. Dong, and C. T. Chan, Nature Communications 5 (2014).

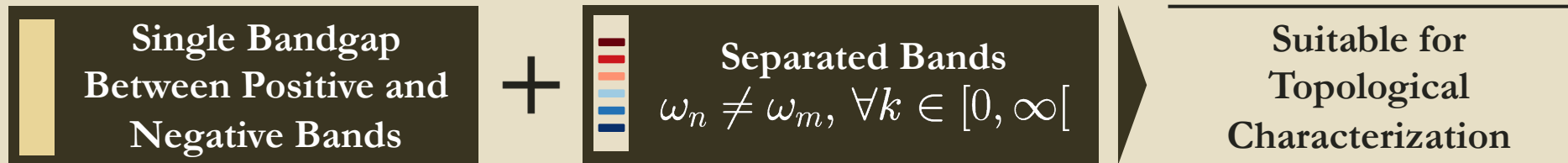
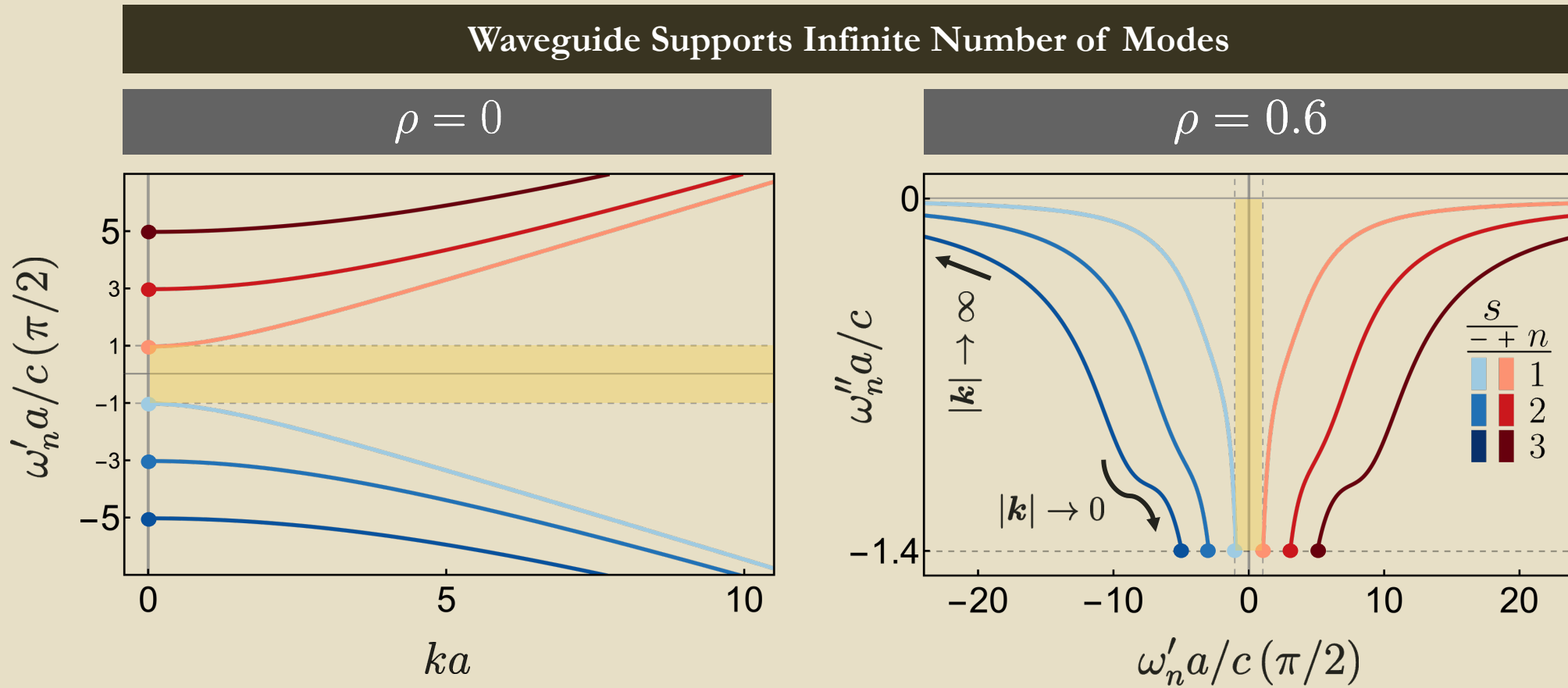


2 Chiral Edge Waves

Transport Protected by Pseudospin-Decoupling

Dispersion Diagram of Bulk Waves

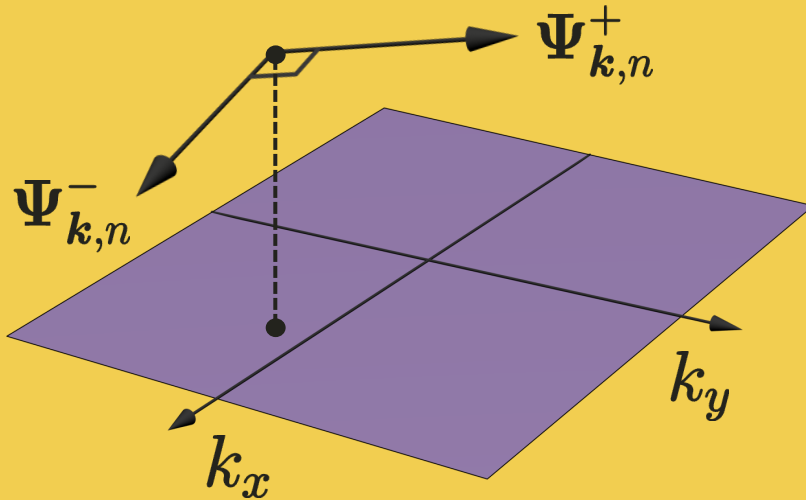
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Topological Characterization

6

Decoupled Pseudospinors $\Psi_{k,n}^{\pm} \in \mathbb{C}^3$



Wavevector Domain

Euclidean Plane

Base Manifold

$$\mathbb{R}^2$$

Typical Fiber

$$U(1) \oplus U(1)$$



Pair of Band
Chern Numbers

$$(\mathcal{C}_n^+, \mathcal{C}_n^-) \in \mathbb{Z} \times \mathbb{Z}$$

Electromagnetic
Reciprocity

$$\mathcal{C}_n^- = -\mathcal{C}_n^+$$



\mathbb{Z} -Valued
Topological
Invariant

$$\mathcal{C}_n^+$$

Topological Characterization

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H. Shen, B. Zhen, and L. Fu,
Physical Review Letters
120 (2018).



M. G. Silveirinha,
Physical Review B
92 (2015).

Topological Band Theory of
Non-Hermitian Systems

Berry Potential

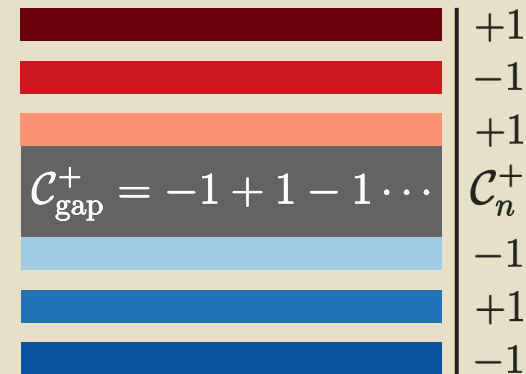
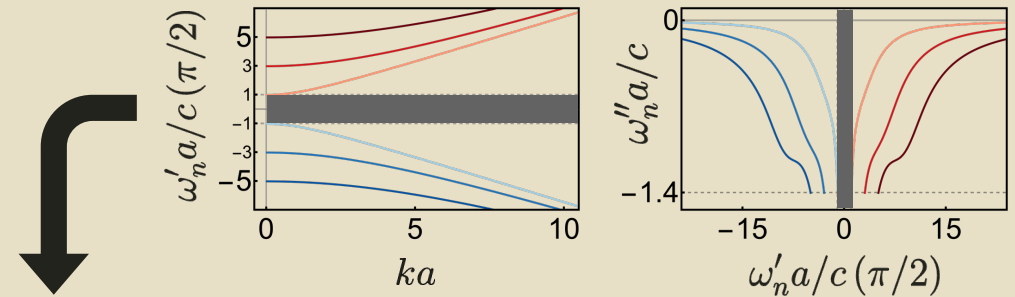
$$\mathcal{A}_{\mathbf{k},n}^{\pm} = i \left\langle \Psi_{\mathbf{k},n}^{\pm}(-\rho) \left| \nabla_{\mathbf{k}} \Psi_{\mathbf{k},n}^{\pm}(\rho) \right. \right\rangle$$

Condition Met for
Integer \mathcal{C}_n^+

$$|\mathbf{k}|^2 \left[\nabla_{\mathbf{k}} \times \mathcal{A}_{\mathbf{k},n}^+ \right]_z \rightarrow 0$$

when $|\mathbf{k}|^2 \rightarrow \infty$

Ordered Scheme of Bands in the Dispersion Diagram



Band Chern Numbers

Alternate

Gap Chern Number

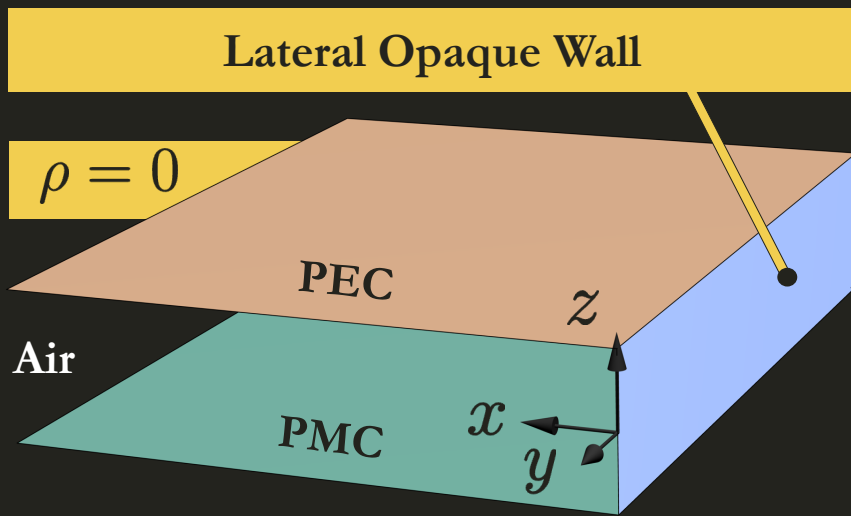
Non-Convergent Series

Bulk-Edge Correspondence holds for

$$\mathcal{C}_{\text{gap}}^+ = -1 + 1 - 1 \dots ?$$

Chiral Transport via Infinitely Many Edge Waves

9



Leontovich Boundary Condition

$$\mathbf{Z} \cdot (\hat{\mathbf{x}} \times \mathbf{H}_{\text{tan}}) = \mathbf{E}_{\text{tan}}$$

Reactive Wall with $Z_y Z_z = Z_0^2$

(as in Corrugated Conducting Plates)

Anisotropic Surface Impedance

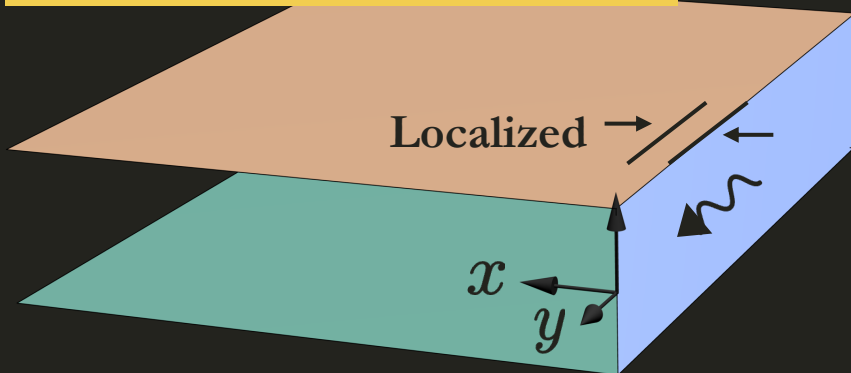
$$\mathbf{Z} = Z_y \hat{\mathbf{y}} \otimes \hat{\mathbf{y}} + Z_z \hat{\mathbf{z}} \otimes \hat{\mathbf{z}}$$

Dispersion of Edge Waves

Chiral Transport via Infinitely Many Edge Waves

9

$$\Psi_n^\pm = \Psi_{\mathbf{k},n}^\pm e^{-i\omega_n t} e^{ik_y y} e^{-\alpha_x x}$$



Leontovich Boundary Condition

$$\mathbf{Z} \cdot (\hat{\mathbf{x}} \times \mathbf{H}_{\text{tan}}) = \mathbf{E}_{\text{tan}}$$

Reactive Wall with $Z_y Z_z = Z_0^2$

$$\frac{\omega_n^2/c^2 - \kappa_n^2}{\kappa_n} = i \frac{Z_z}{Z_0} \left(\pm (-1)^n k_y - \frac{\omega_n}{c \kappa_n} s \times \sqrt{k_y^2 + \kappa_n^2 - \omega_n^2/c^2} \right)$$

(as in Corrugated Conducting Plates)

Anisotropic Surface Impedance

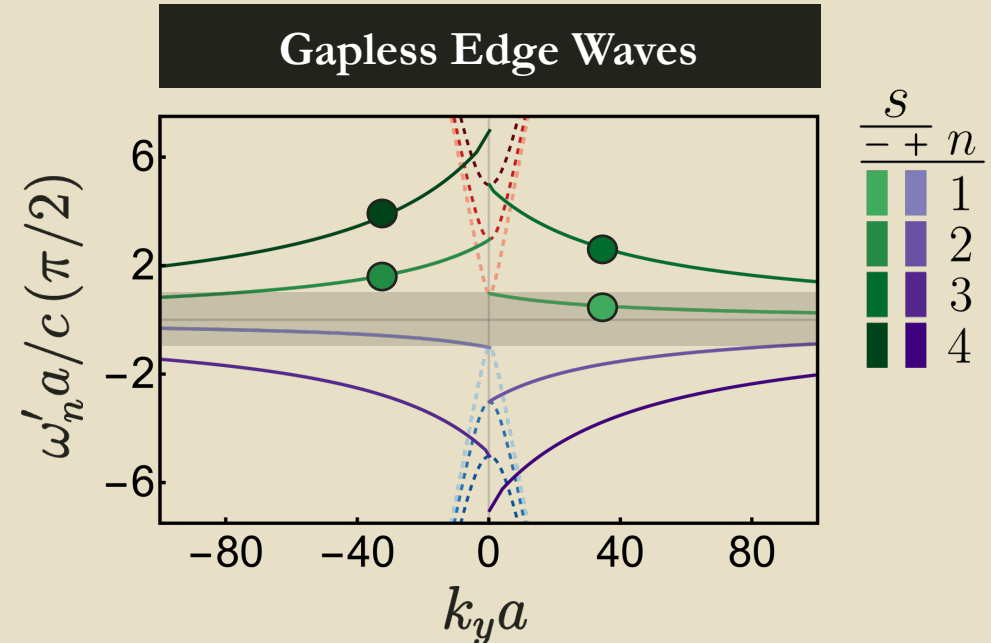
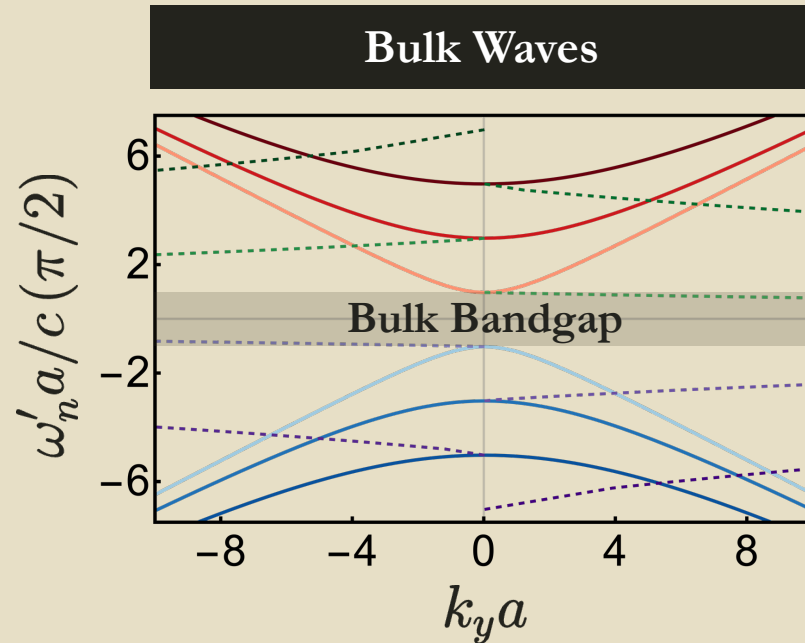
$$\mathbf{Z} = Z_y \hat{\mathbf{y}} \otimes \hat{\mathbf{y}} + Z_z \hat{\mathbf{z}} \otimes \hat{\mathbf{z}}$$

Dispersion of Edge Waves

$$\begin{array}{ccc} + & \rightarrow & - \\ k_y & \rightarrow & -k_y \end{array} \quad \blacklozenge \quad \begin{array}{l} \text{Reciprocity} \\ \mathcal{C}_n^- = -\mathcal{C}_n^+ \end{array}$$

Chiral Transport via Infinitely Many Edge Waves

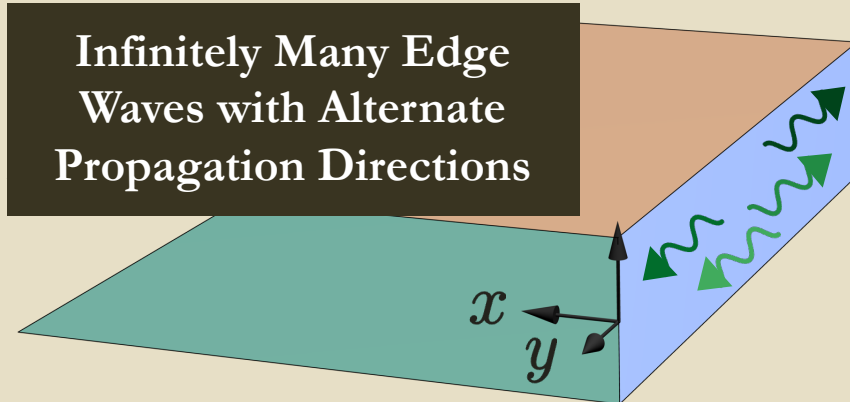
10



$$C_{\text{gap}}^+ = \sum_{n \in \mathbb{N}}^{s=-} C_n^+ = -1 + 1 - 1 + \dots \iff$$

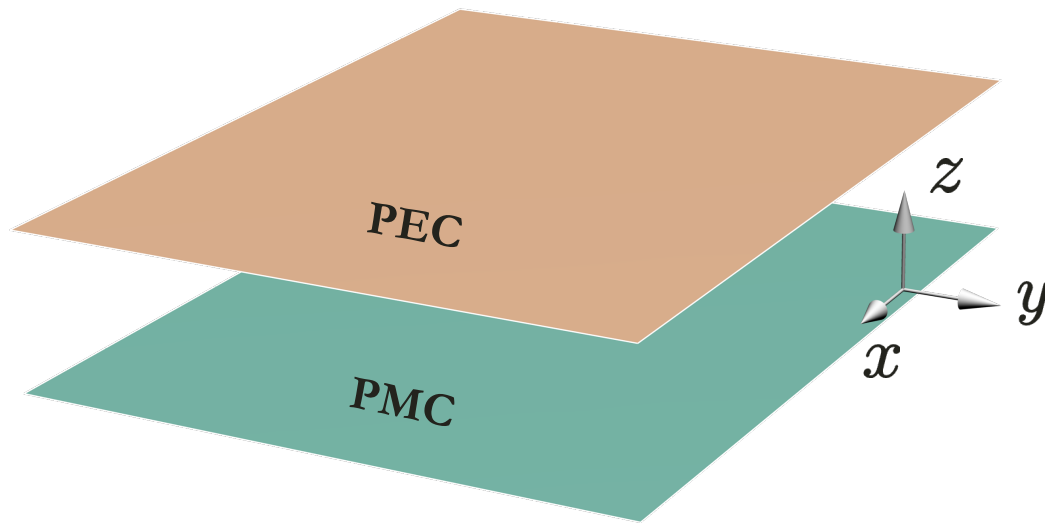
$$C_{\text{gap}}^+ = \sum_{n \in \mathbb{N}}^{s=-} C_n^+ = -1 + 1 - 1 + \dots \quad \text{Infinite Alternate}$$

Infinitely Many Edge Waves with Alternate Propagation Directions



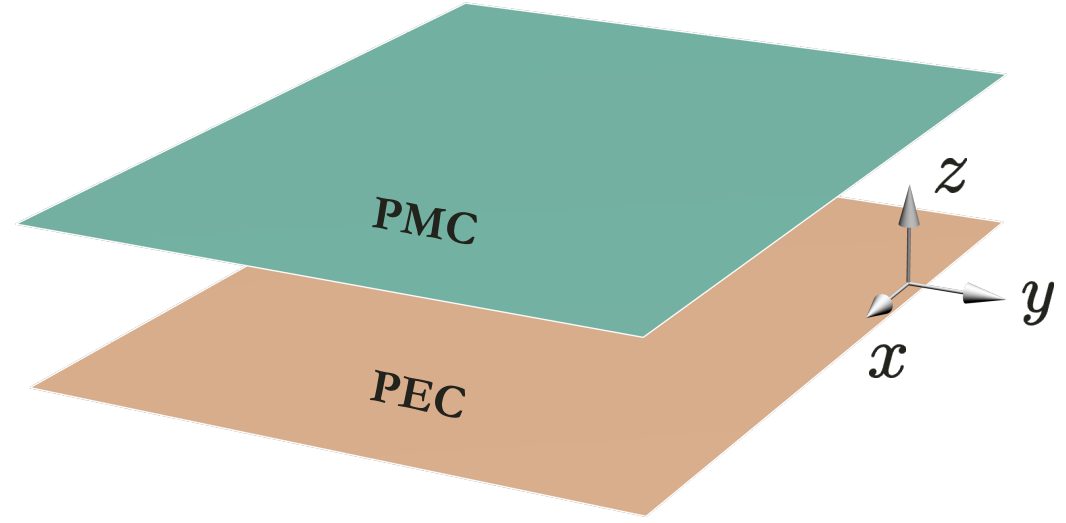
Topological Phase Transition

11



$$\rho = 0$$

$$\mathcal{C}_{\text{gap}}^+ = -1 + 1 - 1 + \dots$$



$$\rho = \infty$$

$$\mathcal{C}_{\text{gap}}^+ = +1 - 1 + 1 - \dots$$

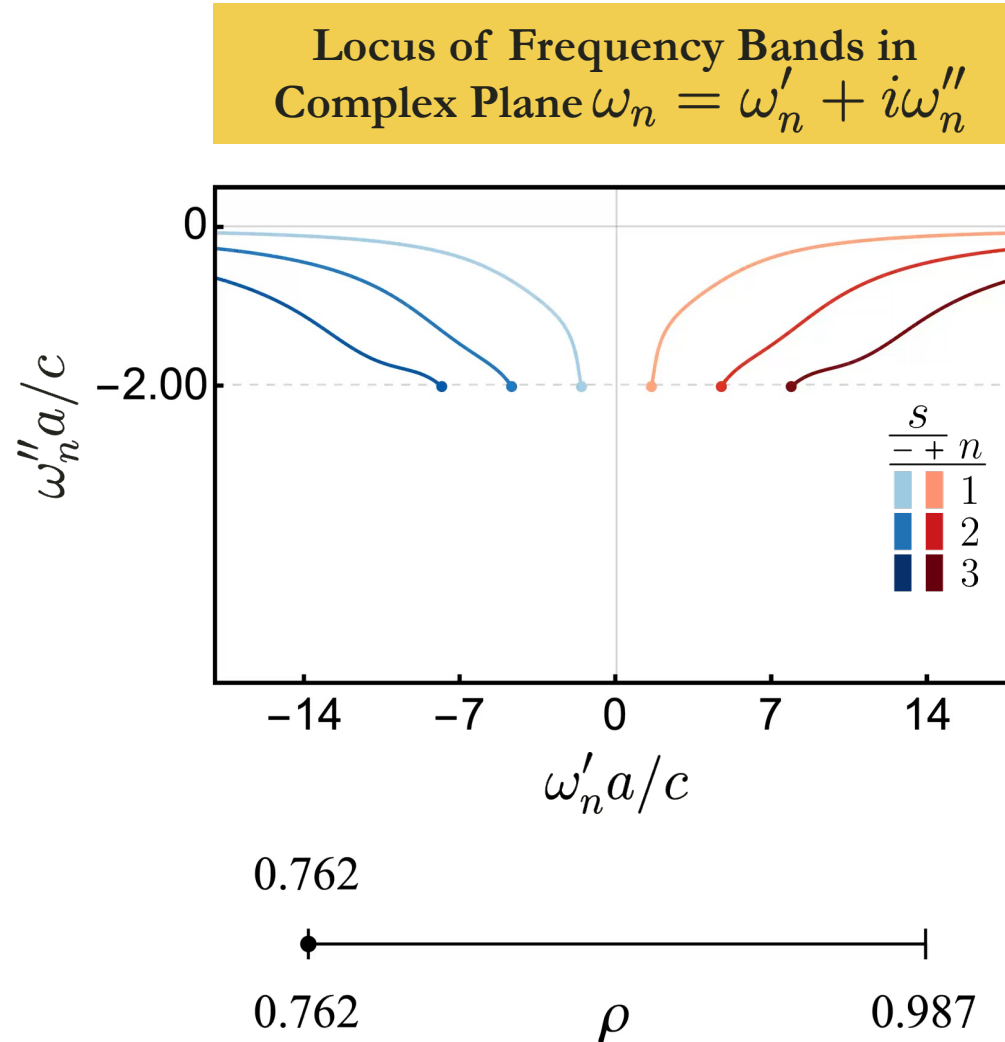


Topological Numbers
Switch Sign



Band Structure near $\rho = 1$

12



Bandgap Closes

$$0.94 < \rho < 0.94^{-1}$$

First-Order Bands Coalesce

$$\omega_{n=1}^{s=+} = \omega_{n=1}^{s=-}$$

+

(Electromagnetic Fields Match)

Pseudospinors Align

$$\Psi_{n=1}^{\pm, s=+} \parallel \Psi_{n=1}^{\pm, s=-}$$

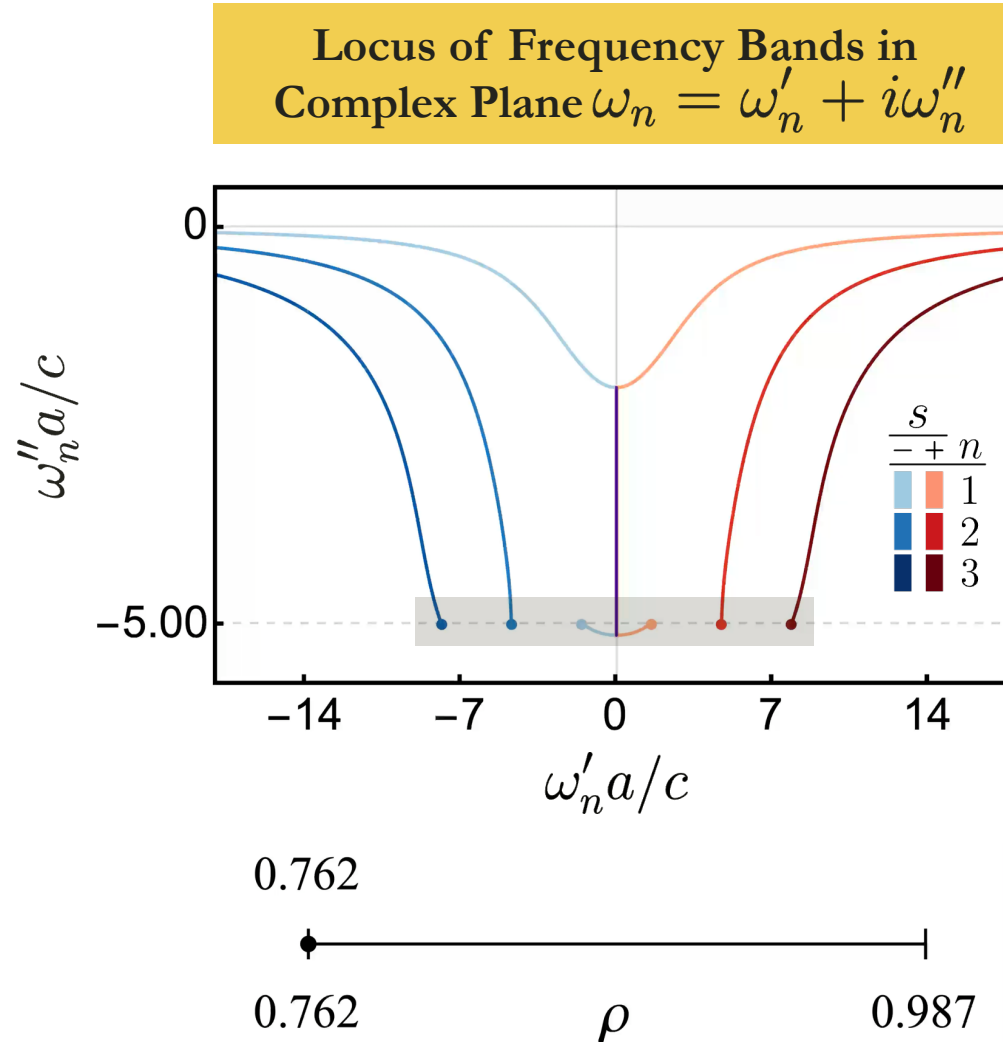
Exceptional Point Degeneracies



M.-A. Miri and A. Alù,
Science **363** (2019).

Band Structure near $\rho = 1$

12



Impedance of top plate matches
that of air when $\rho = 1$

Resonant Absorption at $k = 0$

$$\omega''_n \approx ca^{-1} \ln(\rho - 1)$$

Conclusions

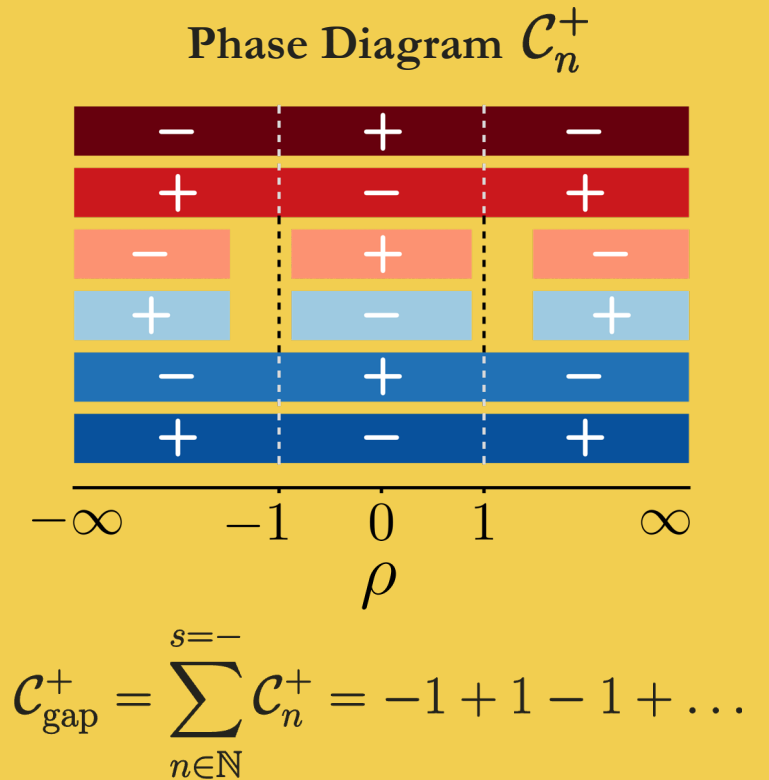
Scattering anomaly in PTD-symmetric platforms has a topological origin that is robust against dissipative effects.

Gap topological charge of PD-symmetric waveguides is a non-convergent sum of integers.

Bulk-edge correspondence can hold despite the ill-defined topology. Infinite number of decoupled chiral modes propagate along the edge of the guide.

A topological phase transition occurs with the formation of exceptional points.

At the exceptional points, there is resonant absorption of energy. Lifetime of guided modes approaches zero.





Thank you!

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