

Generalized probabilistic theories in a nutshell

Rodrigo Ramos

December 10, 2023

Structure of this presentation

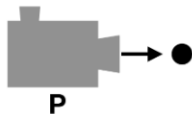
- 1 Introduction
- 2 Single Systems
- 3 Example 1
- 4 Composite Systems
- 5 Example 2
- 6 Conclusions

Operational Theories

- The main primitives of any operational theory are:

Operational Theories

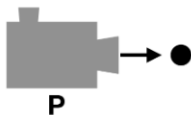
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Preparations

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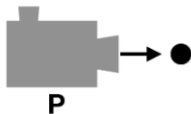
Preparations



Transformations

Operational Theories

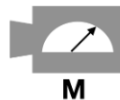
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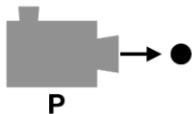
Transformations



Measurements

Operational Theories

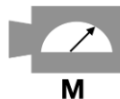
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Preparations



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Measurements

$$p(a \mid P, T, M)$$

Probabilities of outcomes

States and state space

- Let P and P' be distinct preparation procedures. If for any measurement procedure M and outcome a it is true that

$$p(a \mid P, M) = p(a \mid P', M),$$

then we say that P and P' are equivalent.

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Definition

A state is defined as an equivalence class of preparation procedures.

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Let K be a state space. We say that $x \in K$ is an extreme point of K , or equivalently that x is a *pure state*, if for all $y, z \in K$ and $\lambda \in (0, 1)$ such that $x = \lambda y + (1 - \lambda)z$, we have $x = y = z$.

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Naturally, the non-pure states are called *mixed states*.

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Theorem

Let K be a state space and denote by $\text{ext}(K)$ the set of pure states of K , then $K = \text{conv}(\text{ext}(K))^a$.

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In other words, knowing the pure states of a state space is enough to reconstruct it.

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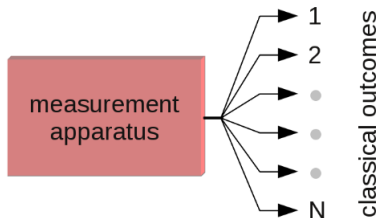
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Definition

Let K be a state space and denote by $A(K)$ the vector space of affine^a functions on K . The *effect algebra* over K is defined as

$$E(K) \doteq \{f \in A(K) : 0 \leq f \leq 1_K\},$$

where 1_K is the constant function $1_K(x) = 1 \ \forall x \in K$.

^aThis means that $f(\lambda x + (1 - \lambda)y) = \lambda f(x) + (1 - \lambda)f(y)$

Effects and effect algebra

- Besides $E(K)$, another important subset of $A(K)$ is the cone of positive functionals:

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and its dual $A(K)^{*+}$.

- One can construct the state space starting from $A(K)^{*+}$:

$$K = \{\varphi \in A(K)^{*+} : \langle \varphi | 1_K \rangle = 1\}.$$

Finite-dimensional Classical theory

- In a classical theory with $n \in \mathbb{N}$ independent outcomes, the state space is the simplex

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where s_j denotes the deterministic probability distribution associated to the j -th outcome.

- The effect algebra $E(S_n)$ is defined by the effects b_1, \dots, b_n that satisfy

$$\langle s_i | b_j \rangle = \delta_{ij}.$$

Bipartite systems

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 - BP4) The unit effect $1_{K_{AB}}$ is equivalent to Alice and Bob applying their unit effects 1_{K_A} and 1_{K_B} .
 - BP5) Tomographic locality holds, i.e local effects applied by Alice and Bob are sufficient to distinguish states in K_{AB} .

Bipartite systems

- Another valid bipartite state space is

$$K_A \otimes_{\max} K_B \doteq \{ \varphi \in K^*(A) \otimes K^*(B) : \langle \varphi | f_A \otimes f_B \rangle \geq 0, \\ \langle \varphi | 1_{K_A} \otimes 1_{K_B} \rangle = 1, \forall f_A \in E(K_A), f_B \in E(K_B) \}$$

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- For any valid bipartite system:

Theorem

$$K_A \otimes_{\min} K_B \subseteq K_{AB} \subseteq K_A \otimes_{\max} K_B.$$

Bipartite systems

- What is really interesting about entanglement in composite scenarios is that

Theorem

Let K be any state space and S_n be a classical state space. Then:

$$K_A \otimes_{\min} S_n = K \otimes_{\max} S_n^a.$$

^aSee the proof in M. Plávala. General probabilistic theories: An introduction. Physics Reports, 1033:1–64, sep 2023. doi: 10.1016/j.physrep.2023.09.001.

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- Before we proceed, we must say that in the GPT framework an *n-outcome measurement* is an affine map $m : K \rightarrow S_n$.

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Theorem

If $m : K \rightarrow S_n$ is an *n-outcome measurement*, then m belongs to

$$A(K)^+ \otimes_{\min} A(S_n)^{*+} = \text{conv}(\text{cone}\{f \otimes s : f \in E(K), s \in E(S_n)\}).$$

Bipartite systems

Corollary

An n -outcome measurement $m : K \rightarrow S_n$ is uniquely defined by effects $\{f_1, \dots, f_n\} \subset E(K)$ such that

$$\sum_{j=1}^n f_j = 1_K.$$

Finite-dimensional quantum theory

Let \mathcal{H} be a finite-dimensional Hilbert space and denote the set of its self-adjoint operators by $\mathcal{B}_H(\mathcal{H})$. As we know, the state space in quantum theory is

$$\mathcal{D}(\mathcal{H}) \doteq \{\rho \in \mathcal{B}_H^+(\mathcal{H}) : \text{Tr}(\rho) = 1\},$$

where

$\mathcal{B}_H^+(\mathcal{H}) \equiv$ set of all positive semi-definite operators on \mathcal{H} .

Finite-dimensional quantum theory

- The space of affine functions is $A(\mathcal{D}(\mathcal{H})) = \mathcal{B}_H(\mathcal{H})$, since the map $X \in \mathcal{B}_H(\mathcal{H}) \mapsto f_X \in A(\mathcal{D}(\mathcal{H}))$,

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- As consequence:

$$A^+(\mathcal{D}(\mathcal{H})) = \mathcal{B}_H^+(\mathcal{H}).$$

Finite-dimensional quantum theory

- Recall that $(\mathcal{B}_H(\mathcal{H}), \langle \cdot | \cdot \rangle_{H-\mathcal{S}})$ is a Hilbert space. As Hilbert spaces are self-dual, due to the previous isomorphism it also holds that

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The equality above tells us that quantum theory obeys what is called **strong self-duality**.

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$$\sum_{k=1}^n E_j = \mathbb{1}. \quad (POVMS!)$$

Where do GPTs show up?

PRX QUANTUM **2**, 010331 (2021)

Characterization of Noncontextuality in the Framework of Generalized Probabilistic Theories

David Schmid^{1,*}, John H. Selby^{1,2,†}, Elie Wolfe¹, Ravi Kunjwal^{1,3} and Robert W. Spekkens¹

¹*Perimeter Institute for Theoretical Physics, 31 Caroline Street North, Waterloo, Ontario N2L 2Y5 Canada*

²*ICTQT, University of Gdańsk, Wita Stwosza 63, Gdańsk 80-308, Poland*

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Where do GPTs show up?

ARTICLE OPEN



Post-quantum steering is a stronger-than-quantum resource for information processing

Paulo J. Cavalcanti¹, John H. Selby¹, Jamie Sikora^{2✉}, Thomas D. Galley³ and Ana Belén Sainz¹

We present the first instance where post-quantum steering is a stronger-than-quantum resource for information processing – remote state preparation. In addition, we show that the phenomenon of post-quantum steering is not just a mere mathematical curiosity allowed by the no-signalling principle, but it may arise within compositional theories beyond quantum theory, hence making its study fundamentally relevant. We show these results by formulating a new compositional general probabilistic theory – which we call Witworld – with strong post-quantum features, which proves to be a intuitive and useful tool for exploring steering and its applications beyond the quantum realm.

npj Quantum Information (2022)8:76; <https://doi.org/10.1038/s41534-022-00574-8>

Where do GPTs show up?

Quantum Darwinism and the spreading of classical information in non-classical theories

Roberto D. Baldijão^{*,1,2}, Marius Krumm^{*,2,3}, Andrew J. P. Garner², and Markus P. Müller^{2,4,5}

¹Instituto de Física Gleb Wataghin, Universidade Estadual de Campinas, Campinas, SP 13083-859, Brazil

²Institute for Quantum Optics and Quantum Information, Austrian Academy of Sciences, Boltzmanngasse 3, A-1090 Vienna, Austria

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⁴Vienna Center for Quantum Science and Technology (VCQ), Faculty of Physics, University of Vienna, Vienna, Austria

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Where do GPTs show up?

A no-go theorem on the nature of the gravitational field beyond quantum theory

Thomas D. Galley^{§1}, Flaminia Giacomini^{§1}, and John H. Selby^{§2}

¹Perimeter Institute for Theoretical Physics, 31 Caroline St. N, Waterloo, Ontario, N2L 2Y5, Canada

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03-08-2022

Take-home message

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- GPTs help us perceive quantum theory from a bird's eye-view (no restriction hypothesis, local tomography, strong self-duality).
- They serve as a tool to tinker around nonclassical theories beyond quantum.