Tópicos Especiais em Matemática Aplicada - 2025-1 UERJ

05 - Metodo dos Elementos Finitos - Caso 1D estacionário

Rodrigo Madureira rodrigo.madureira@ime.uerj.br

 $\textbf{Github}: \ https://github.com/rodrigoIrmadureira/ElementosFinitos$

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No caso unidimensional, cada subintervalo $[x_e, x_{e+1}]$, $e=1,2,\ldots,m-1$, do domínio $\Omega=[x_1,x_m]$ é um elemento finito de tamanho $h=x_{e+1}-x_e$.

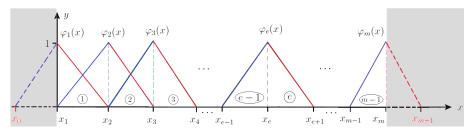


Figura: Funções de base de Lagrange linear $\varphi_i(x)$

Ao invés de "olhar" para todo o domínio, vamos "olhar" para um único elemento e.

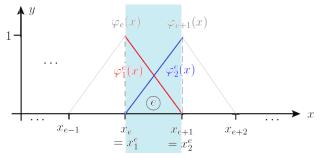


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos referenciar um nó da discretização do domínio de duas formas:

- x_i , onde i é o índice da discretização do intervalo $[x_1, x_m]$, i = 1, 2, ..., m. Neste caso, i é a **numeração global**. Assim, x_i é o **nó global**.
- x_a^e , onde α é o índice do nó no elemento finito e, $\alpha = 1, 2$. Neste caso, α é a **numeração local**. Assim, x_a^e é o **nó local**.

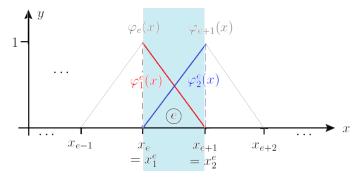


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

$$\begin{cases}
 \phi_1^e(x) = \frac{x_{e+1} - x}{x_{e+1} - x_e} = \frac{x_{e+1} - x}{x_2^e - x_1^e} = \frac{x_{e+1} - x}{h}, \\
 \phi_2^e(x) = \frac{x - x_e}{x_{e+1} - x_e} = \frac{x - x_e}{x_2^e - x_1^e} = \frac{x - x_e}{h}.
\end{cases}$$
(1)

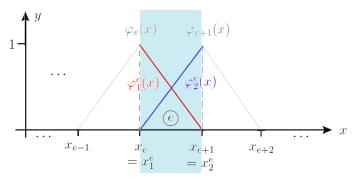


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que:

$$\begin{cases} \phi_1^e(x) = \phi_e(x), \text{ para } x \in [x_e, x_{e+1}], \\ \phi_2^e(x) = \phi_{e+1}(x), \text{ para } x \in [x_e, x_{e+1}]. \end{cases}$$
 (2)

Matriz local Ke

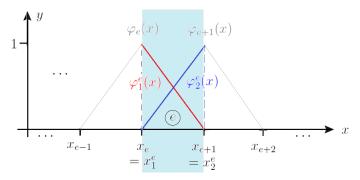


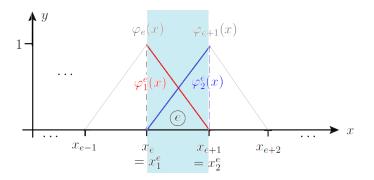
Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que restritos ao elemento finito e, obtemos o elemento da **matriz local** K^e :

$$\mathsf{K}^e_{ab} = \alpha(\phi^e_{ax},\phi^e_{bx}) + \beta(\phi^e_a,\phi^e_b) = \alpha \int_{x^e_1}^{x^e_2} \phi^e_{ax}(x) \phi^e_{bx}(x) dx + \beta \int_{x^e_1}^{x^e_2} \phi^e_a(x) \phi^e_b(x) dx,$$

para a, b = 1, 2.

Vetor força local Fe



Analogamente, obtemos o elemento do **vetor força local** F^e :

$$\mathsf{F}_{\mathfrak{a}}^{\mathfrak{e}}=(\mathsf{f},\phi_{\mathfrak{a}}^{\mathfrak{e}})=\alpha\int_{x_{1}^{\mathfrak{e}}}^{x_{2}^{\mathfrak{e}}}\mathsf{f}(x)\phi_{\mathfrak{a}}^{\mathfrak{e}}(x)dx, \text{ para } \mathfrak{a}=1,2.$$



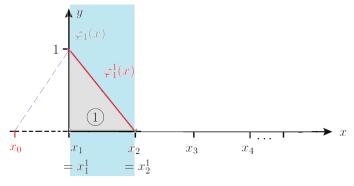
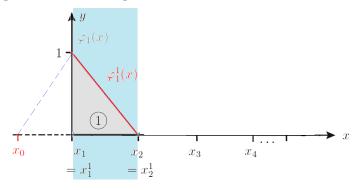


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

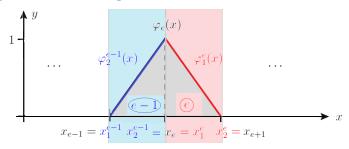
Podemos reescrever a matriz global K em termos de elementos da matriz local K^e , $e=1,2,\ldots,m$. Vimos na abordagem anterior que:

$$K_{11} = \alpha(\phi_{1x},\phi_{1x}) + \beta(\phi_{1},\phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx$$



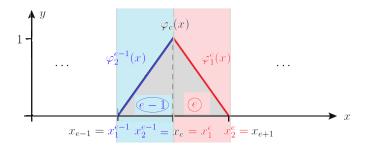
Podemos reescrever K₁₁ como:

$$\begin{split} K_{11} &= \alpha(\phi_{1x},\phi_{1x}) + \beta(\phi_{1},\phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1x}^{1}(x))^{2} dx + \beta \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1}^{1}(x))^{2} dx = \alpha(\phi_{1x}^{1},\phi_{1x}^{1}) + \beta(\phi_{1}^{1},\phi_{1}^{1}) = K_{11}^{1} \end{split}$$



Podemos reescrever K_{ii} como:

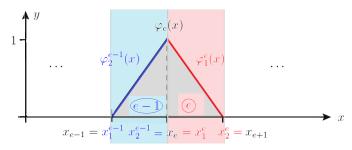
$$\begin{split} K_{ee} &= \alpha(\phi_{ex}, \phi_{ex}) + \beta(\phi_{e}, \phi_{e}) = \alpha \int_{x_{e-1}}^{x_{e+1}} (\phi_{ex}(x))^{2} dx + \beta \int_{x_{e-1}}^{x_{e+1}} (\phi_{e}(x))^{2} dx \\ &= \alpha \Big(\int_{x_{e-1}}^{x_{e}} (\phi_{ex}(x))^{2} dx + \int_{x_{e}}^{x_{e+1}} (\phi_{ex}(x))^{2} dx \Big) \\ &+ \beta \Big(\int_{x_{e-1}}^{x_{e}} (\phi_{e}(x))^{2} dx + \int_{x_{e}}^{x_{e+1}} (\phi_{e}(x))^{2} dx \Big) \end{split}$$



Podemos reescrever K_{ee} como:

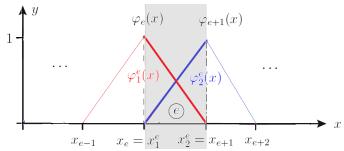
$$\begin{split} K_{ee} &= \alpha \biggl(\int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2x}^{e-1}(x))^2 dx + \int_{x_1^{e}}^{x_2^{e}} (\phi_{1x}^{e}(x))^2 dx \biggr) \\ &+ \beta \biggl(\int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2}^{e-1}(x))^2 dx + \int_{x_1^{e}}^{x_2^{e}} (\phi_{1}^{e}(x))^2 dx \biggr) \end{split}$$





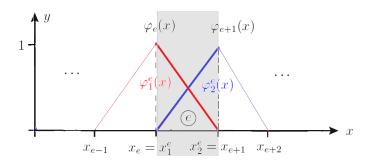
Podemos reescrever Kee como:

$$\begin{split} K_{ee} &= \alpha \int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2x}^{e-1}(x))^2 dx + \beta \int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_2^{e-1}(x))^2 dx \\ &+ \alpha \int_{x_1^{e}}^{x_2^{e}} (\phi_{1x}^{e}(x))^2 dx + \beta \int_{x_1^{e}}^{x_2^{e}} (\phi_1^{e}(x))^2 dx \\ &= \alpha (\phi_{2x}^{e-1}, \phi_{2x}^{e-1}) + \beta (\phi_2^{e-1}, \phi_2^{e-1}) + \alpha (\phi_{1x}^{e}, \phi_{1x}^{e}) + \beta (\phi_1^{e}, \phi_1^{e}) = K_{22}^{e-1} + K_{11}^{e} \end{split}$$



Podemos reescrever $K_{e,e+1}$ como:

$$\begin{split} K_{e,e+1} &= \alpha(\phi_{ex}, \phi_{(e+1)x}) + \beta(\phi_{e}, \phi_{e+1}) \\ &= \alpha \int_{x_{e}}^{x_{e+1}} \phi_{ex}(x) \phi_{(e+1)x}(x) dx + \beta \int_{x_{e}}^{x_{e+1}} \phi_{e}(x) \phi_{e+1}(x) dx \\ &= \alpha \int_{x_{1}^{e}}^{x_{2}^{e}} \phi_{1x}^{e}(x) \phi_{2x}^{e}(x) dx + \beta \int_{x_{1}^{e}}^{x_{2}^{e}} \phi_{1}^{e}(x) \phi_{2}^{e}(x) dx \\ &= \alpha(\phi_{1x}^{e}, \phi_{2x}^{e}) + \beta(\phi_{1}^{e}, \phi_{2}^{e}) = K_{12}^{e} \end{split}$$



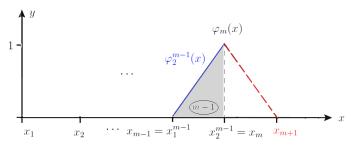
De forma análoga, obtemos:

$$K_{e+1,e} = K_{21}^e$$

Como K é simétrica $(K = K^T)$, logo:

$$K_{12}^e = K_{21}^e$$





Podemos reescrever K_{mm} como:

$$\begin{split} K_{mm} &= \alpha(\phi_{mx}, \phi_{mx}) + \beta(\phi_{m}, \phi_{m}) \\ &= \alpha \int_{x_{m-1}}^{x_{m}} (\phi_{mx}(x))^{2} dx + \beta \int_{x_{m-1}}^{x_{m}} (\phi_{m}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2x}^{m-1}(x))^{2} dx + \beta \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2}^{m-1}(x))^{2} dx \\ &= \alpha(\phi_{2x}^{m-1}, \phi_{2x}^{m-1}) + \beta(\phi_{2x}^{m-1}, \phi_{2x}^{m-1}) = K_{22}^{m-1} \end{split}$$

Montagem do vetor global F

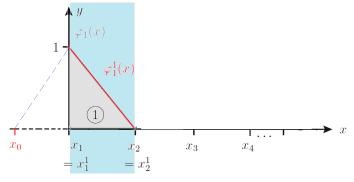
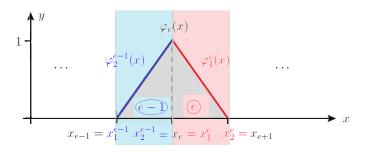


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos reescrever a **vetor global** F em termos de elementos do **vetor local** F^e , e = 1, 2, ..., m. Vimos na abordagem anterior que:

$$F_1 = (f, \phi_1) = \int_{x_1}^{x_2} f(x)\phi_1(x)dx = \int_{x_1^1}^{x_2^1} f(x)\phi_1^1(x)dx = (f, \phi_1^1) = F_1^1$$

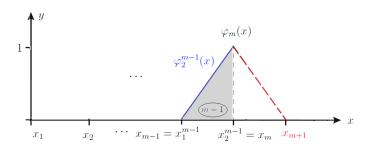
Montagem do vetor global F



Podemos reescrever F_e como:

$$\begin{split} F_e &= (f,\phi_e) = \int_{x_{e-1}}^{x_{e+1}} f(x) \phi_e(x) dx = \int_{x_{e-1}}^{x_e} f(x) \phi_e(x) dx + \int_{x_e}^{x_{e+1}} f(x) \phi_e(x) dx \\ &= \int_{x_1^{e-1}}^{x_2^{e-1}} f(x) \phi_2^{e-1}(x) dx + \int_{x_1^{e}}^{x_2^{e}} f(x) \phi_1^{e}(x) dx = (f,\phi_2^{e-1}) + (f,\phi_1^{e}) = F_2^{e-1} + F_1^{e} \end{split}$$

Montagem do vetor global F



Podemos reescrever F_{m.} como:

$$F_{m} = (f, \phi_{m}) = \int_{x_{m-1}}^{x_{m}} f(x)\phi_{m}(x)dx = \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} f(x)\phi_{2}^{m-1}(x)dx$$
$$= (f, \phi_{2}^{m-1}) = F_{2}^{m-1}$$



Resultados

Matriz de rigidez K:

$$K_{11} = K_{11}^{1};$$

$$K_{ee} = K_{22}^{e-1} + K_{11}^{e}, \text{ para } e = 2, 3, ..., m-1;$$
(3)

$$\begin{split} K_{e,e+1} = K_{12}^{\mathfrak{i}} \Rightarrow K_{e+1,e} = K_{21}^{\mathfrak{i}}, \text{ para } e = 1,2,3,\ldots,m-1; \text{ (Simetria: } K = K^{\mathsf{T}}\text{)} \\ K_{m,m} = K_{22}^{m-1}. \end{split}$$

Vetor força F:

$$\begin{aligned} F_1 &= F_1^1; \\ F_e &= F_2^{e-1} + F_1^e, \text{ para } i = 2,3,\dots,m-1; \\ F_m &= F_2^{m-1}. \end{aligned} \tag{4}$$

Resultados

A matriz de rigidez global K,

é montada a partir das matrizes locais

$$K^e = \begin{bmatrix} K_{11}^e & K_{12}^e \\ K_{21}^e & K_{22}^e \end{bmatrix},$$

para e = 1, 2, ..., m - 1.



Resultados

O vetor força global,

$$F = \begin{bmatrix} F_1^1 \\ F_2^1 + F_1^2 \\ F_2^2 + F_1^3 \\ F_2^3 + F_1^4 \\ \vdots \\ F_2^{m-2} + F_1^{m-1} \\ F_2^{m-1} \end{bmatrix}$$

é montado a partir dos vetores locais

$$F^e = \begin{bmatrix} F_1^e \\ F_2^e \end{bmatrix},$$

para e = 1, 2, ..., m - 1.



A regra da **Quadratura Gaussiana** para integração numérica consiste em aproximar

$$\int_{-1}^{1} g(\xi) d\xi \approx \sum_{i=1}^{npg} w_i \cdot g(\xi_i),$$

onde npg é o número de pontos de Gauss,

 ξ_i , $i = 1, 2, \dots, npg$, são os **pontos de Gauss**,

 w_i , i = 1, 2, ..., npg, são os **pesos de Gauss**.

Comparação com outros métodos de integração numérica:

Regra de	Número	Grau do polinômio
integração	de pontos	com integração
numérica	(nós)	exata
Trapézios	2	1
Simpson	3	2
Quadratura Gaussiana	2	3
	n	2n - 1, , , , , , , , , , , , , , , , , ,

Com n = 2, é possível mostrar que:

Pontos de Gauss: $\xi_1 = -\sqrt{3}/3$, $\xi_2 = \sqrt{3}/3$.

Pesos de Gauss: $w_1 = 1$, $w_2 = 1$.

Demonstração: Se com npg=2, o resultado da integração numérica é exato para um polinômio de grau 3,

$$P_3(x) = a_0 + a_1x + a_2x^2 + a_3x^3$$
,

então,

$$\int_{-1}^{1} P_3(x) dx = \sum_{i=1}^{2} w_i \cdot P_3(\xi_i) = w_1 \cdot P_3(\xi_1) + w_2 \cdot P_3(\xi_2)$$

$$\Rightarrow \int_{-1}^{1} (a_0 + a_1 x + a_2 x^2 + a_3 x^3) dx = w_1 \cdot (a_0 + a_1 \xi_1 + a_2 \xi_1^2 + a_3 \xi_1^3) + w_2 \cdot (a_0 + a_1 \xi_2 + a_2 \xi_2^2 + a_3 \xi_2^3)$$

$$\begin{split} & \Rightarrow a_0 \int_{-1}^1 dx + a_1 \int_{-1}^1 x dx + a_2 \int_{-1}^1 x^2 dx + a_3 \int_{-1}^1 x^3 dx \\ & = (a_0 w_1 + a_1 w_1 \xi_1 + a_2 w_1 \xi_1^2 + a_3 w_1 \xi_1^3) \\ & + (a_0 w_2 + a_1 w_2 \xi_2 + a_2 w_2 \xi_2^2 + a_3 w_2 \xi_2^3) \\ & = a_0 (w_1 + w_2) + a_1 (w_1 \xi_1 + w_2 \xi_2) + a_2 (w_1 \xi_1^2 + w_2 \xi_2^2) + a_3 (w_1 \xi_1^3 + w_2 \xi_2^3) \end{split}$$

Igualando os coeficientes de a_0 , a_1 , a_2 , a_3 , obtemos:

$$\int_{-1}^{1} dx = w_1 + w_2 \Rightarrow w_1 + w_2 = 2;$$

$$\int_{-1}^{1} x dx = w_1 \xi_1 + w_2 \xi_2 \Rightarrow w_1 \xi_1 + w_2 \xi_2 = 0;$$

$$\int_{-1}^{1} x^2 dx = w_1 \xi_1^2 + w_2 \xi_2^2 \Rightarrow w_1 \xi_1^2 + w_2 \xi_2^2 = 2/3;$$

$$\int_{-1}^{1} x^3 dx = w_1 \xi_1^3 + w_2 \xi_2^3 \Rightarrow w_1 \xi_1^3 + w_2 \xi_2^3 = 0.$$

Resolvendo este sistema, obtemos: $w_1=1$, $w_2=1$, $\xi_1=\sqrt{3}/3$, $\xi_2=\sqrt{3}/3$, $\xi_3=\sqrt{3}/3$

npg	Pontos	Pesos
1	0	2
2	$\pm\sqrt{3}/3 \approx \pm 0.5773502692$	1
3	$0, \pm \sqrt{3/5} \approx \pm 0.7745966692$	0.8888888889, 0.555555556
4	$\pm 0.8611363116, \pm 0.3399810436$	0.3478548451, 0.6521451549
5	± 0.9061798459 , ± 0.5384693101 ,	0.2369268851, 0.4786286705,
	0	0.5688888889
6	$\pm 0.9324695142, \pm 0.6612093865,$	0.1713244924, 0.3607615730,
	± 0.2386191861	0.4679139346
7	± 0.9491079123 , ± 0.7415311856 ,	0.1294849662, 0.2797053915,
	$\pm 0.4058451514, 0$	0.3818300505, 0.4179591837
8	± 0.9602898565 , ± 0.7966664774 ,	0.1012285363, 0.2223810345,
	± 0.5255324099 , ± 0.1834346425	0.3137066459, 0.3626837834

Tabela: Tabela para Quadratura Gaussiana

Os elementos da **matriz local** K^e são dados por:

$$\begin{split} \mathsf{K}^e_{ab} &= \alpha(\phi^e_a,\phi^e_b) = \alpha(\phi^e_{ax},\phi^e_{bx}) + \beta(\phi^e_a,\phi^e_b) \\ &= \alpha \int_{x^e_1}^{x^e_2} \phi^e_{ax}(x) \phi^e_{bx}(x) \ dx + \beta \int_{x^e_1}^{x^e_2} \phi^e_a(x) \phi^e_b(x) \ dx, \ \mathsf{para} \ a,b \in \{1,2\}. \end{split}$$

Para usar a **quadratura gaussiana** nessas integrais, devemos usar mudança de variáveis do domínio $[x_1^e, x_2^e]$ para o domínio [-1, 1]:



As distâncias são proporcionais nas duas retas. Logo,

$$\frac{x-x_1^e}{x_2^e-x_1^e}=\frac{\xi-(-1)}{1-(-1)}\Rightarrow\frac{x-x_1^e}{h}=\frac{\xi+1}{2}\Rightarrow x(\xi)=x_1^e+h\Big(\frac{\xi+1}{2}\Big)$$

Logo,
$$\frac{dx}{d\xi} = \frac{h}{2} \Rightarrow dx = \frac{h}{2}d\xi$$
.

Fazendo a mudança de variável nas funções locais $\phi^e_a(x)$ para a=1,2, obtemos:

$$\begin{split} \phi_1^e(x) &= \frac{x_2^e - x}{x_2^e - x_1^e} = \frac{x_2^e - x}{h} \\ \Rightarrow \phi_1^e(\xi) &= \phi_1^e(x(\xi)) = \frac{x_2^e - x(\xi)}{x_2^e - x_1^e} = \frac{x_2^e - x(\xi)}{h} = \frac{x_2^e - \left(x_1^e + h\left(\frac{\xi + 1}{2}\right)\right)}{h} \\ \Rightarrow \phi_1^e(\xi) &= \frac{\left(x_2^e - x_1^e\right) - h\left(\frac{\xi + 1}{2}\right)}{h} = \frac{h - h\left(\frac{\xi + 1}{2}\right)}{h} = 1 - \left(\frac{\xi + 1}{2}\right) = \frac{1 - \xi}{2} \end{split}$$

De forma análoga, obtemos:

$$\phi_2^e(x) = \frac{x-x_1^e}{x_2^e-x_1^e} = \frac{x-x_1^e}{h} \Rightarrow \phi_2^e(\xi) = \frac{1+\xi}{2}$$



Portanto, as funções da base de Lagrange linear locais no domínio $\xi \in [-1,1]$ são:

$$\begin{cases}
\varphi_1^e(\xi) = \frac{1-\xi}{2}, \\
\varphi_2^e(\xi) = \frac{1+\xi}{2},
\end{cases} (5)$$

e as suas derivadas são:

$$\begin{cases} \varphi_{1\xi}^{e}(\xi) = \frac{d\varphi_{1}^{e}}{d\xi}(\xi) = -\frac{1}{2}, \\ \varphi_{2\xi}^{e}(\xi) = \frac{d\varphi_{1}^{e}}{d\xi}(\xi) = \frac{1}{2}. \end{cases}$$
 (6)

Com a mudança de variável, obtemos os elementos de K^e no domínio [-1,1]. Lembrando que ao usar a regra da cadeia, para $a,b\in\{1,2\}$,

$$\begin{split} \phi_{\alpha x}(x) &= \phi_{\alpha x}(\xi(x)) = \frac{d\phi_{\alpha}(\xi)}{d\xi} \frac{d\xi}{dx} = \phi^{e}_{\alpha \xi}(\xi) \frac{d\xi}{dx} = \phi^{e}_{\alpha \xi}(\xi) \frac{2}{h}. \\ \text{Logo,} \end{split}$$

$$\begin{split} K^e_{ab} &= \alpha(\phi^e_a,\phi^e_b) = \alpha(\phi^e_{ax},\phi^e_{bx}) + \beta(\phi^e_a,\phi^e_b) \\ &= \alpha \int_{x^e_1}^{x^e_2} \phi^e_{ax}(x) \phi^e_{bx}(x) \ dx + \beta \int_{x^e_1}^{x^e_2} \phi^e_a(x) \phi^e_b(x) \ dx \\ &= \alpha \int_{-1}^1 \phi^e_{a\xi}(\xi) \frac{2}{h} \phi^e_{b\xi}(\xi) \frac{2}{h} \frac{h}{2} d\xi + \beta \int_{-1}^1 \phi^e_a(\xi) \phi^e_b(\xi) \ \frac{h}{2} d\xi \\ &= \frac{2\alpha}{h} \int_{-1}^1 \phi^e_{a\xi}(\xi) \phi^e_{b\xi}(\xi) \ d\xi + \frac{\beta h}{2} \int_{-1}^1 \phi^e_a(\xi) \phi^e_b(\xi) \ d\xi \end{split}$$

Então,

$$\begin{split} K_{11}^e &= \frac{2\alpha}{h} \int_{-1}^1 (\phi_{1\xi}^e(\xi))^2 \ d\xi + \frac{\beta h}{2} \int_{-1}^1 (\phi_{1}^e(\xi))^2 d\xi \\ &= \frac{2\alpha}{h} \int_{-1}^1 \left(-\frac{1}{2} \right)^2 d\xi + \frac{\beta h}{2} \int_{-1}^1 \left(\frac{1}{2} \right)^2 d\xi \\ &= \frac{\alpha}{2h} \int_{-1}^1 \ d\xi + \frac{\beta h}{8} \int_{-1}^1 (1-\xi)^2 \ d\xi \end{split}$$

Usando **Quadratura Gaussiana** com npg=2, temos $w_1=w_2=1$, $\xi_1=-\sqrt{3}/3$, $\xi_2=\sqrt{3}/3$. Assim, obtemos:

$$\int_{-1}^{1} d\xi = \sum_{i=1}^{2} w_i \cdot \xi_i = w_1 \cdot g(\xi_1) + w_2 \cdot g(\xi_2)$$
$$= 1 \cdot g(-\sqrt{3}/3) + 1 \cdot g(\sqrt{3}/3) = 1 \cdot 1 + 1 \cdot 1 = 2$$

$$\begin{split} \int_{-1}^{1} (1 - \xi)^2 d\xi &= \sum_{i=1}^{2} w_i \cdot \xi_i = w_1 \cdot g(\xi_1) + w_2 \cdot g(\xi_2) \\ &= 1 \cdot g(-\sqrt{3}/3) + 1 \cdot g(\sqrt{3}/3) \\ &= 1 \cdot \left(1 - \left(-\frac{\sqrt{3}}{3}\right)\right)^2 + 1 \cdot \left(1 - \left(\frac{\sqrt{3}}{3}\right)\right)^2 \\ &= \left(1 + \frac{\sqrt{3}}{3}\right)^2 + \left(1 - \frac{\sqrt{3}}{3}\right)^2 = \frac{8}{3} \end{split}$$

Portanto, os elementos de Ke são:

$$K_{11}^e = \frac{\alpha}{2h} \cdot 2 + \frac{\beta h}{8} \cdot \frac{8}{3} = \frac{\alpha}{h} + \frac{\beta h}{3}$$

Analogamente, verifique que:

$$\begin{split} K_{12}^e &= K_{21}^e = -\frac{\alpha}{h} + \frac{\beta h}{6}, \text{ pois K \'e sim\'etrica}, \\ K_{22}^e &= \frac{\alpha}{h} + \frac{\beta h}{3} \end{split}$$



Assim, temos a matriz de rigidez local Ke também simétrica:

$$K^{e} = \begin{bmatrix} K_{11}^{e} & K_{12}^{e} \\ K_{21}^{e} & K_{22}^{e} \end{bmatrix} = \begin{bmatrix} \frac{\alpha}{h} + \frac{\beta h}{3} & -\frac{\alpha}{h} + \frac{\beta h}{6} \\ -\frac{\alpha}{h} + \frac{\beta h}{6} & \frac{\alpha}{h} + \frac{\beta h}{3} \end{bmatrix}$$

Para calcular os elementos da matriz local F^e usando **Quadratura Gaussiana** com npg=2, temos:

$$\begin{split} F_{\alpha}^e &= (f,\phi_{\alpha}^e) = \int_{x_1^e}^{x_2^e} f(x)\phi_{\alpha}^e(x)dx = \int_{-1}^1 f(\xi)\phi_{\alpha}^e(\xi)\frac{h}{2}\ d\xi \\ &= \frac{h}{2}\int_{-1}^1 f(\xi)\phi_{\alpha}^e(\xi)\ d\xi \approx \frac{h}{2}\sum_{i=1}^2 w_i \cdot f(\xi_i)\phi_{\alpha}^e(\xi_i) \\ &\approx \frac{h}{2}\Big[1 \cdot f\Big(-\frac{\sqrt{3}}{3}\Big)\phi_{\alpha}^e\Big(-\frac{\sqrt{3}}{3}\Big) + 1 \cdot f\Big(\frac{\sqrt{3}}{3}\Big)\phi_{\alpha}^e\Big(\frac{\sqrt{3}}{3}\Big)\Big] \\ &\approx \frac{h}{2}\Big[f\Big(-\frac{\sqrt{3}}{3}\Big)\phi_{\alpha}^e\Big(-\frac{\sqrt{3}}{3}\Big) + f\Big(\frac{\sqrt{3}}{3}\Big)\phi_{\alpha}^e\Big(\frac{\sqrt{3}}{3}\Big)\Big] \end{split}$$

Então,

$$\begin{split} F_1^e &= (f,\phi_1^e) \approx \frac{h}{2} \Big[f\Big(-\frac{\sqrt{3}}{3}\Big) \Big(\frac{1-(-\sqrt{3}/3)}{2}\Big) + f\Big(\frac{\sqrt{3}}{3}\Big) \Big(\frac{1-(\sqrt{3}/3)}{2}\Big) \Big] \\ &\approx \frac{h}{4} \Big[f\Big(-\frac{\sqrt{3}}{3}\Big) \Big(1+\frac{\sqrt{3}}{3}\Big) + f\Big(\frac{\sqrt{3}}{3}\Big) \Big(1-\frac{\sqrt{3}}{3}\Big) \Big] \end{split}$$

Analogamente,

$$\begin{split} F_2^e &\approx \frac{h}{2} \bigg[f \bigg(-\frac{\sqrt{3}}{3} \bigg) \phi_2^e \bigg(-\frac{\sqrt{3}}{3} \bigg) + f \bigg(\frac{\sqrt{3}}{3} \bigg) \phi_2^e \bigg(\frac{\sqrt{3}}{3} \bigg) \bigg] \\ &\approx \frac{h}{2} \bigg[f \bigg(-\frac{\sqrt{3}}{3} \bigg) \bigg(\frac{1 + (-\sqrt{3}/3)}{2} \bigg) + f \bigg(\frac{\sqrt{3}}{3} \bigg) \bigg(\frac{1 + (\sqrt{3}/3)}{2} \bigg) \bigg] \\ &\approx \frac{h}{4} \bigg[f \bigg(-\frac{\sqrt{3}}{3} \bigg) \bigg(1 - \frac{\sqrt{3}}{3} \bigg) + f \bigg(\frac{\sqrt{3}}{3} \bigg) \bigg(1 + \frac{\sqrt{3}}{3} \bigg) \bigg] \end{split}$$

Os resultados vão depender do cálculo de f(x).



Assim, temos o vetor força local Fe:

$$F^{e} = \begin{bmatrix} F_{1}^{e} \\ F_{2}^{e} \end{bmatrix} = \begin{bmatrix} \frac{h}{4} \left[f \left(-\frac{\sqrt{3}}{3} \right) \left(1 + \frac{\sqrt{3}}{3} \right) + f \left(\frac{\sqrt{3}}{3} \right) \left(1 - \frac{\sqrt{3}}{3} \right) \right] \\ \frac{h}{4} \left[f \left(-\frac{\sqrt{3}}{3} \right) \left(1 - \frac{\sqrt{3}}{3} \right) + f \left(\frac{\sqrt{3}}{3} \right) \left(1 + \frac{\sqrt{3}}{3} \right) \right] \end{bmatrix}$$

Referências I



Liu, I.S.; Rincon, M.A.. Introdução ao Método de Elementos Finitos, Análise e Aplicação. IM/UFRJ, 2003.