Tópicos Especiais em Matemática Aplicada - 2025-1 UERJ

05 - Metodo dos Elementos Finitos - Caso 1D estacionário

Rodrigo Madureira rodrigo.madureira@ime.uerj.br

Github: https://github.com/rodrigolrmadureira/ElementosFinitos

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No caso unidimensional, cada subintervalo $[x_e, x_{e+1}]$, $e=1,2,\ldots,m-1$, do domínio $\Omega=[x_1,x_m]$ é um elemento finito de tamanho $h=x_{e+1}-x_e$.

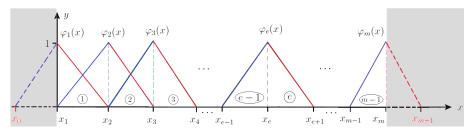


Figura: Funções de base de Lagrange linear $\varphi_i(x)$

Ao invés de "olhar" para todo o domínio, vamos "olhar" para um único elemento e.

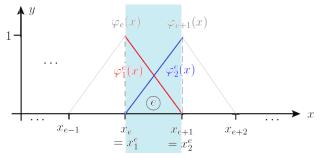


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos referenciar um nó da discretização do domínio de duas formas:

- x_i , onde i é o índice da discretização do intervalo $[x_1, x_m]$, i = 1, 2, ..., m. Neste caso, i é a **numeração global**. Assim, x_i é o **nó global**.
- x_a^e , onde α é o índice do nó no elemento finito e, $\alpha = 1, 2$. Neste caso, α é a **numeração local**. Assim, x_a^e é o **nó local**.

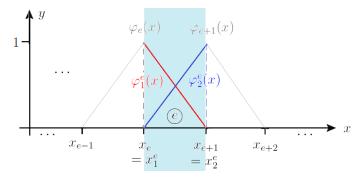


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

$$\begin{cases} \varphi_1^e(x) = \frac{x_{e+1} - x}{x_{e+1} - x_e} = \frac{x_{e+1} - x}{x_2^e - x_1^e} = \frac{x_{e+1} - x}{h}, \\ \varphi_2^e(x) = \frac{x - x_e}{x_{e+1} - x_e} = \frac{x - x_e}{x_2^e - x_1^e} = \frac{x - x_e}{h}. \end{cases}$$
(1)

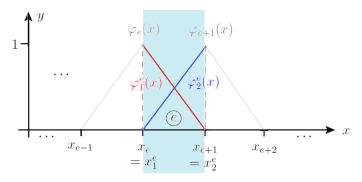


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que:

$$\begin{cases} \phi_1^e(x) = \phi_e(x), \text{ para } x \in [x_e, x_{e+1}], \\ \phi_2^e(x) = \phi_{e+1}(x), \text{ para } x \in [x_e, x_{e+1}]. \end{cases}$$
 (2)

Matriz local Ke

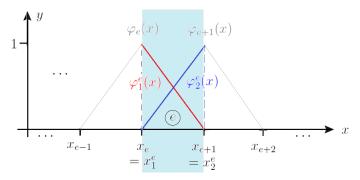


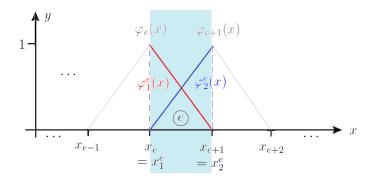
Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que restritos ao elemento finito e, obtemos o elemento da **matriz local** K^e :

$$K^e_{ab} = \alpha(\phi^e_{ax},\phi^e_{bx}) + \beta(\phi^e_a,\phi^e_b) = \alpha \int_{x^e_1}^{x^e_2} \phi^e_{ax}(x) \phi^e_{bx}(x) dx + \beta \int_{x^e_1}^{x^e_2} \phi^e_a(x) \phi^e_b(x) dx,$$

para a, b = 1, 2.

Vetor força local Fe



Analogamente, obtemos o elemento do **vetor força local** F^e :

$$\mathsf{F}_{\alpha}^{e}=(\mathsf{f},\phi_{\alpha}^{e})=\alpha\int_{x_{1}^{e}}^{x_{2}^{e}}\mathsf{f}(x)\phi_{\alpha}^{e}(x)dx,\ \mathsf{para}\ \alpha=1,2.$$



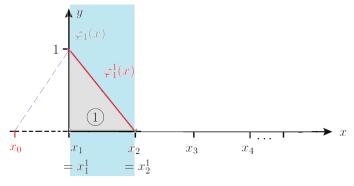
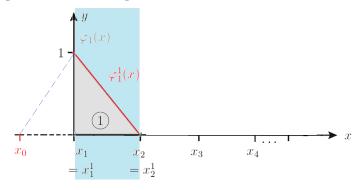


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

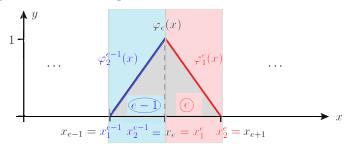
Podemos reescrever a matriz global K em termos de elementos da matriz local K^e , $e=1,2,\ldots,m$. Vimos na abordagem anterior que:

$$K_{11} = \alpha(\phi_{1x},\phi_{1x}) + \beta(\phi_{1},\phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx$$



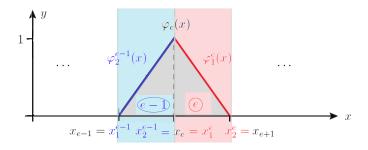
Podemos reescrever K₁₁ como:

$$\begin{split} K_{11} &= \alpha(\phi_{1x},\phi_{1x}) + \beta(\phi_{1},\phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1x}^{1}(x))^{2} dx + \beta \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1}^{1}(x))^{2} dx = \alpha(\phi_{1x}^{1},\phi_{1x}^{1}) + \beta(\phi_{1}^{1},\phi_{1}^{1}) = K_{11}^{1} \end{split}$$



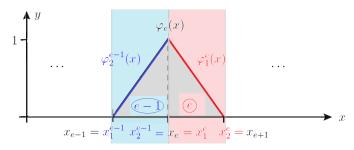
Podemos reescrever K_{ii} como:

$$\begin{split} K_{ee} &= \alpha(\phi_{ex}, \phi_{ex}) + \beta(\phi_{e}, \phi_{e}) = \alpha \int_{x_{e-1}}^{x_{e+1}} (\phi_{ex}(x))^{2} dx + \beta \int_{x_{e-1}}^{x_{e+1}} (\phi_{e}(x))^{2} dx \\ &= \alpha \Big(\int_{x_{e-1}}^{x_{e}} (\phi_{ex}(x))^{2} dx + \int_{x_{e}}^{x_{e+1}} (\phi_{ex}(x))^{2} dx \Big) \\ &+ \beta \Big(\int_{x_{e-1}}^{x_{e}} (\phi_{e}(x))^{2} dx + \int_{x_{e}}^{x_{e+1}} (\phi_{e}(x))^{2} dx \Big) \end{split}$$



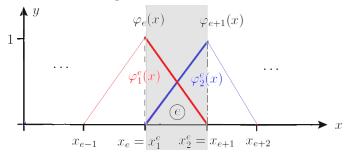
Podemos reescrever K_{ee} como:

$$\begin{split} K_{ee} &= \alpha \biggl(\int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2x}^{e-1}(x))^2 dx + \int_{x_1^{e}}^{x_2^{e}} (\phi_{1x}^{e}(x))^2 dx \biggr) \\ &+ \beta \biggl(\int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2}^{e-1}(x))^2 dx + \int_{x_1^{e}}^{x_2^{e}} (\phi_{1}^{e}(x))^2 dx \biggr) \end{split}$$



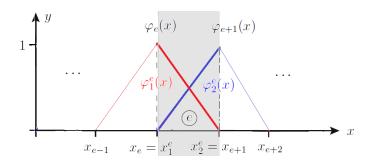
Podemos reescrever Kee como:

$$\begin{split} K_{ee} &= \alpha \int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_{2x}^{e-1}(x))^2 dx + \beta \int_{x_1^{e-1}}^{x_2^{e-1}} (\phi_2^{e-1}(x))^2 dx \\ &+ \alpha \int_{x_1^{e}}^{x_2^{e}} (\phi_{1x}^{e}(x))^2 dx + \beta \int_{x_1^{e}}^{x_2^{e}} (\phi_1^{e}(x))^2 dx \\ &= \alpha (\phi_{2x}^{e-1}, \phi_{2x}^{e-1}) + \beta (\phi_2^{e-1}, \phi_2^{e-1}) + \alpha (\phi_{1x}^{e}, \phi_{1x}^{e}) + \beta (\phi_1^{e}, \phi_1^{e}) = K_{22}^{e-1} + K_{11}^{e} \end{split}$$



Podemos reescrever $K_{e,e+1}$ como:

$$\begin{split} K_{e,e+1} &= \alpha(\phi_{ex}, \phi_{(e+1)x}) + \beta(\phi_{e}, \phi_{e+1}) \\ &= \alpha \int_{x_{e}}^{x_{e+1}} \phi_{ex}(x) \phi_{(e+1)x}(x) dx + \beta \int_{x_{e}}^{x_{e+1}} \phi_{e}(x) \phi_{e+1}(x) dx \\ &= \alpha \int_{x_{1}^{e}}^{x_{2}^{e}} \phi_{1x}^{e}(x) \phi_{2x}^{e}(x) dx + \beta \int_{x_{1}^{e}}^{x_{2}^{e}} \phi_{1}^{e}(x) \phi_{2}^{e}(x) dx \\ &= \alpha(\phi_{1x}^{e}, \phi_{2x}^{e}) + \beta(\phi_{1}^{e}, \phi_{2}^{e}) = K_{12}^{e} \end{split}$$



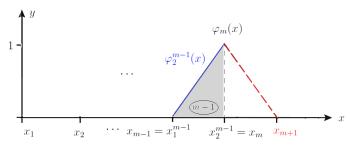
De forma análoga, obtemos:

$$K_{e+1,e} = K_{21}^e$$

Como K é simétrica ($K = K^T$), logo:

$$K_{12}^e=K_{21}^e$$





Podemos reescrever K_{mm} como:

$$\begin{split} K_{mm} &= \alpha(\phi_{mx}, \phi_{mx}) + \beta(\phi_{m}, \phi_{m}) \\ &= \alpha \int_{x_{m-1}}^{x_{m}} (\phi_{mx}(x))^{2} dx + \beta \int_{x_{m-1}}^{x_{m}} (\phi_{m}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2x}^{m-1}(x))^{2} dx + \beta \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2}^{m-1}(x))^{2} dx \\ &= \alpha(\phi_{2x}^{m-1}, \phi_{2x}^{m-1}) + \beta(\phi_{2x}^{m-1}, \phi_{2x}^{m-1}) = K_{22}^{m-1} \end{split}$$

Montagem do vetor global F

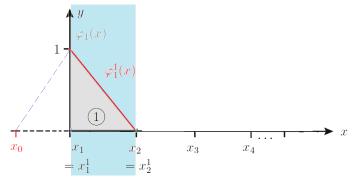
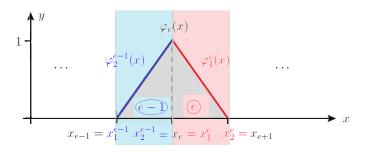


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos reescrever a **vetor global** F em termos de elementos do **vetor local** F^e , e = 1, 2, ..., m. Vimos na abordagem anterior que:

$$F_1 = (f, \phi_1) = \int_{x_1}^{x_2} f(x)\phi_1(x)dx = \int_{x_1^1}^{x_2^1} f(x)\phi_1^1(x)dx = (f, \phi_1^1) = F_1^1$$

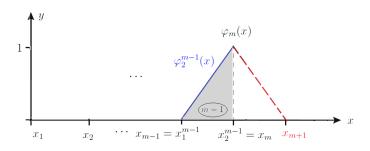
Montagem do vetor global F



Podemos reescrever F_e como:

$$\begin{split} F_e &= (f,\phi_e) = \int_{x_{e-1}}^{x_{e+1}} f(x) \phi_e(x) dx = \int_{x_{e-1}}^{x_e} f(x) \phi_e(x) dx + \int_{x_e}^{x_{e+1}} f(x) \phi_e(x) dx \\ &= \int_{x_1^{e-1}}^{x_2^{e-1}} f(x) \phi_2^{e-1}(x) dx + \int_{x_1^e}^{x_2^e} f(x) \phi_1^e(x) dx = (f,\phi_2^{e-1}) + (f,\phi_1^e) = F_2^{e-1} + F_1^e \end{split}$$

Montagem do vetor global F



Podemos reescrever F_{m.} como:

$$F_{m} = (f, \phi_{m}) = \int_{x_{m-1}}^{x_{m}} f(x)\phi_{m}(x)dx = \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} f(x)\phi_{2}^{m-1}(x)dx$$
$$= (f, \phi_{2}^{m-1}) = F_{2}^{m-1}$$



Resultados

Matriz de rigidez K:

$$K_{11} = K_{11}^{1};$$

$$K_{ee} = K_{22}^{e-1} + K_{11}^{e}, \text{ para } e = 2, 3, ..., m-1;$$
(3)

$$\begin{split} K_{e,e+1} = K_{12}^{\mathfrak{i}} \Rightarrow K_{e+1,e} = K_{21}^{\mathfrak{i}}, \text{ para } e = 1,2,3,\ldots,m-1; \text{ (Simetria: } K = K^{\mathsf{T}}\text{)} \\ K_{m,m} = K_{22}^{m-1}. \end{split}$$

Vetor força F:

$$\begin{aligned} F_1 &= F_1^1; \\ F_e &= F_2^{e-1} + F_1^e, \text{ para } i = 2,3,\dots,m-1; \\ F_m &= F_2^{m-1}. \end{aligned} \tag{4}$$

Resultados

A matriz de rigidez global K,

é montada a partir das matrizes locais

$$K^{e} = \begin{bmatrix} K_{11}^{e} & K_{12}^{e} \\ K_{21}^{e} & K_{22}^{e} \end{bmatrix},$$

para e = 1, 2, ..., m - 1.



Resultados

O vetor força global,

$$F = \begin{bmatrix} F_1^1 \\ F_2^1 + F_1^2 \\ F_2^2 + F_1^3 \\ F_2^3 + F_1^4 \\ \vdots \\ F_2^{m-2} + F_1^{m-1} \\ F_2^{m-1} \end{bmatrix}$$

é montado a partir dos vetores locais

$$F^e = \begin{bmatrix} F_1^e \\ F_2^e \end{bmatrix},$$

para e = 1, 2, ..., m - 1.



Referências I



Liu, I.S.; Rincon, M.A.. Introdução ao Método de Elementos Finitos, Análise e Aplicação. IM/UFRJ, 2003.