Tópicos Especiais em Matemática Aplicada - 2025-1 UERJ

05 - Metodo dos Elementos Finitos - Caso unidimensional

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Github: https://github.com/rodrigolrmadureira/ElementosFinitos

Sumário

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No caso unidimensional, cada subintervalo $[x_e, x_{e+1}]$, $e=1,2,\ldots,m-1$, do domínio $\Omega=[x_1,x_m]$ é um elemento finito de tamanho $h=x_{e+1}-x_e$.

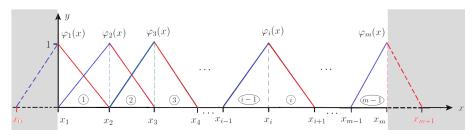


Figura: Funções de base de Lagrange linear $\varphi_i(x)$

Ao invés de "olhar" para todo o domínio, vamos "olhar" para um único elemento e.

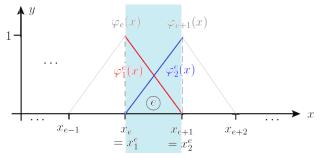


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos referenciar um nó da discretização do domínio de duas formas:

- x_i , onde i é o índice da discretização do intervalo $[x_1, x_m]$, i = 1, 2, ..., m. Neste caso, i é a **numeração global**. Assim, x_i é o **nó global**.
- x_a^e , onde α é o índice do nó no elemento finito e, $\alpha = 1, 2$. Neste caso, α é a **numeração local**. Assim, x_a^e é o **nó local**.

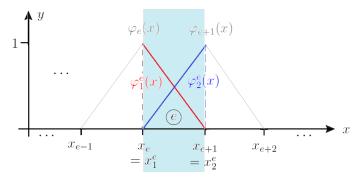


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

$$\begin{cases}
\varphi_1^e(x) = \frac{x_{e+1} - x}{x_{e+1} - x_e} = \frac{x_{e+1} - x}{x_2^e - x_1^e} = \frac{x_{e+1} - x}{h}, \\
\varphi_2^e(x) = \frac{x - x_e}{x_{e+1} - x_e} = \frac{x - x_e}{x_2^e - x_1^e} = \frac{x - x_e}{h}.
\end{cases} (1)$$

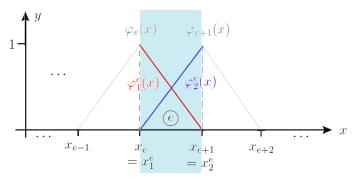


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que:

$$\begin{cases} \phi_1^e(x) = \phi_e(x), \text{ para } x \in [x_e, x_{e+1}], \\ \phi_2^e(x) = \phi_{e+1}(x), \text{ para } x \in [x_e, x_{e+1}]. \end{cases}$$
 (2)

Matriz local Ke

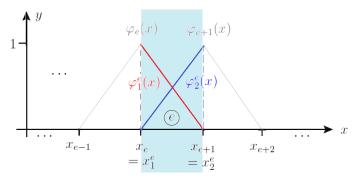


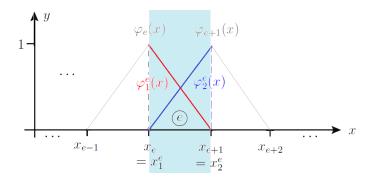
Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Note que restritos ao elemento finito e, obtemos o elemento da **matriz local** K^e :

$$K^e_{ab} = \alpha(\phi^e_{ax},\phi^e_{bx}) + \beta(\phi^e_a,\phi^e_b) = \alpha \int_{x^e_1}^{x^e_2} \phi^e_{ax}(x) \phi^e_{bx}(x) dx + \beta \int_{x^e_1}^{x^e_2} \phi^e_a(x) \phi^e_b(x) dx,$$

para a, b = 1, 2.

Vetor força local Fe



Analogamente, obtemos o elemento do **vetor força local** F^e :

$$\mathsf{F}_{\alpha}^{e}=(\mathsf{f},\phi_{\alpha}^{e})=\alpha\int_{x_{1}^{e}}^{x_{2}^{e}}\mathsf{f}(x)\phi_{\alpha}^{e}(x)dx,\ \mathsf{para}\ \alpha=1,2.$$



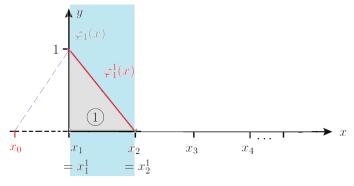
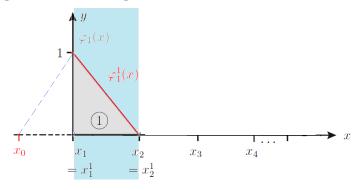


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

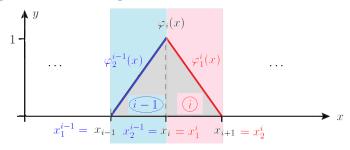
Podemos reescrever a matriz global K em termos de elementos da matriz local K^e , $e=1,2,\ldots,m$. Vimos na abordagem anterior que:

$$K_{11} = \alpha(\phi_{1x},\phi_{1x}) + \beta(\phi_{1},\phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx$$



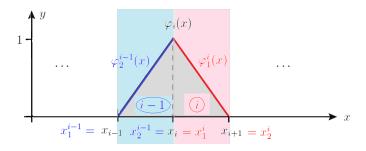
Podemos reescrever K₁₁ como:

$$\begin{split} K_{11} &= \alpha(\phi_{1x}, \phi_{1x}) + \beta(\phi_{1}, \phi_{1}) = \alpha \int_{x_{1}}^{x_{2}} (\phi_{1x}(x))^{2} dx + \beta \int_{x_{1}}^{x_{2}} (\phi_{1}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1x}^{1}(x))^{2} dx + \beta \int_{x_{1}^{1}}^{x_{2}^{1}} (\phi_{1}^{1}(x))^{2} dx = \alpha(\phi_{1x}^{1}, \phi_{1x}^{1}) + \beta(\phi_{1}^{1}, \phi_{1}^{1}) = K_{11}^{1} \end{split}$$



Podemos reescrever K_{ii} como:

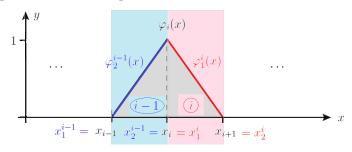
$$\begin{split} K_{ii} &= \alpha(\phi_{ix},\phi_{ix}) + \beta(\phi_{i},\phi_{i}) = \alpha \int_{x_{i-1}}^{x_{i+1}} (\phi_{ix}(x))^{2} dx + \beta \int_{x_{i-1}}^{x_{i+1}} (\phi_{i}(x))^{2} dx \\ &= \alpha \bigg(\int_{x_{i-1}}^{x_{i}} (\phi_{ix}(x))^{2} dx + \int_{x_{i}}^{x_{i+1}} (\phi_{ix}(x))^{2} dx \bigg) \\ &+ \beta \bigg(\int_{x_{i-1}}^{x_{i}} (\phi_{i}(x))^{2} dx + \int_{x_{i}}^{x_{i+1}} (\phi_{i}(x))^{2} dx \bigg) \end{split}$$



Podemos reescrever K_{ii} como:

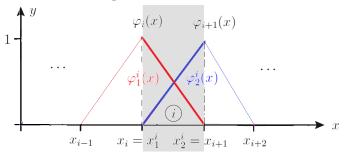
$$\begin{split} K_{ii} &= \alpha \biggl(\int_{x_1^{i-1}}^{x_2^{i-1}} (\phi_{2x}^{i-1}(x))^2 dx + \int_{x_1^{i}}^{x_2^{i}} (\phi_{1x}^{i}(x))^2 dx \biggr) \\ &+ \beta \biggl(\int_{x_1^{i-1}}^{x_2^{i-1}} (\phi_{2}^{i-1}(x))^2 dx + \int_{x_1^{i}}^{x_2^{i}} (\phi_{1}^{i}(x))^2 dx \biggr) \end{split}$$





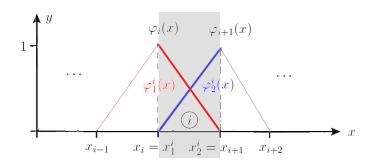
Podemos reescrever K_{ii} como:

$$\begin{split} K_{ii} &= \alpha \int_{x_1^{i-1}}^{x_2^{i-1}} (\phi_{2x}^{i-1}(x))^2 dx + \beta \int_{x_1^{i-1}}^{x_2^{i-1}} (\phi_2^{i-1}(x))^2 dx \\ &+ \alpha \int_{x_1^{i}}^{x_2^{i}} (\phi_{1x}^{i}(x))^2 dx + \beta \int_{x_1^{i}}^{x_2^{i}} (\phi_1^{i}(x))^2 dx \\ &= \alpha (\phi_{2x}^{i-1}, \phi_{2x}^{i-1}) + \beta (\phi_2^{i-1}, \phi_2^{i-1}) + \alpha (\phi_{1x}^{i}, \phi_{1x}^{i}) + \beta (\phi_1^{i}, \phi_1^{i}) = K_{22}^{i-1} + K_{11}^{i} \end{split}$$



Podemos reescrever $K_{i,i+1}$ como:

$$\begin{split} K_{i,i+1} &= \alpha(\phi_{ix}, \phi_{(i+1)x}) + \beta(\phi_{i}, \phi_{i+1}) \\ &= \alpha \int_{x_{i}}^{x_{i+1}} \phi_{ix}(x) \phi_{(i+1)x}(x) dx + \beta \int_{x_{i}}^{x_{i+1}} \phi_{i}(x) \phi_{i+1}(x) dx \\ &= \alpha \int_{x_{1}^{i}}^{x_{2}^{i}} \phi_{1x}^{i}(x) \phi_{2x}^{i}(x) dx + \beta \int_{x_{1}^{i}}^{x_{2}^{i}} \phi_{1}^{i}(x) \phi_{2}^{i}(x) dx \\ &= \alpha (\phi_{1x}^{i}, \phi_{2x}^{i}) + \beta(\phi_{1}^{i}, \phi_{2}^{i}) = K_{12}^{i} \end{split}$$



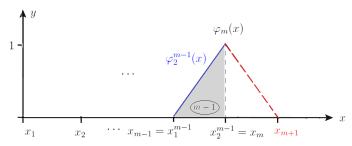
De forma análoga, obtemos:

$$K_{\mathfrak{i}+1,\mathfrak{i}}=K_{21}^{\mathfrak{i}}$$

Como K é simétrica $(K = K^T)$, logo:

$$K_{12}^{\mathfrak{i}}=K_{21}^{\mathfrak{i}}$$





Podemos reescrever K_{mm} como:

$$\begin{split} K_{mm} &= \alpha(\phi_{mx},\phi_{mx}) + \beta(\phi_{m},\phi_{m}) \\ &= \alpha \int_{x_{m-1}}^{x_{m}} (\phi_{mx}(x))^{2} dx + \beta \int_{x_{m-1}}^{x_{m}} (\phi_{m}(x))^{2} dx \\ &= \alpha \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2x}^{m-1}(x))^{2} dx + \beta \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} (\phi_{2}^{m-1}(x))^{2} dx \\ &= \alpha (\phi_{2x}^{m-1},\phi_{2x}^{m-1}) + \beta (\phi_{2}^{m-1},\phi_{2x}^{m-1}) = K_{22}^{m-1} \end{split}$$

Montagem do vetor global F

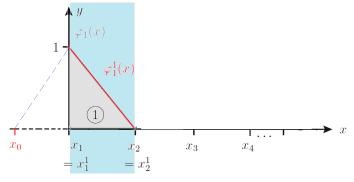
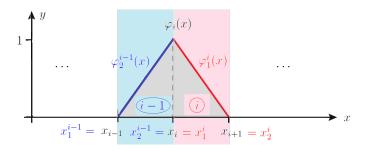


Figura: Funções do elemento e: $\varphi_1^e(x)$ e $\varphi_2^e(x)$

Podemos reescrever a **vetor global** F em termos de elementos do **vetor local** F^e , e = 1, 2, ..., m. Vimos na abordagem anterior que:

$$F_1 = (f, \phi_1) = \int_{x_1}^{x_2} f(x)\phi_1(x)dx = \int_{x_1^1}^{x_2^1} f(x)\phi_1^1(x)dx = (f, \phi_1^1) = F_1^1$$

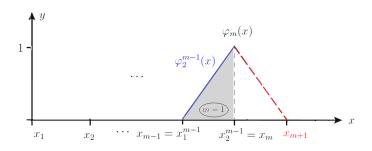
Montagem do vetor global F



Podemos reescrever F_i como:

$$\begin{split} F_i &= (f,\phi_i) = \int_{x_{i-1}}^{x_{i+1}} f(x) \phi_i(x) dx = \int_{x_{i-1}}^{x_i} f(x) \phi_i(x) dx + \int_{x_i}^{x_{i+1}} f(x) \phi_i(x) dx \\ &= \int_{x_1^{i-1}}^{x_2^{i-1}} f(x) \phi_2^{i-1}(x) dx + \int_{x_1^{i}}^{x_2^{i}} f(x) \phi_1^{i}(x) dx = (f,\phi_2^{i-1}) + (f,\phi_1^{i}) = F_2^{i-1} + F_1^{i} \end{split}$$

Montagem do vetor global F



Podemos reescrever F_m como:

$$F_{m} = (f, \phi_{m}) = \int_{x_{m-1}}^{x_{m}} f(x)\phi_{m}(x)dx = \int_{x_{1}^{m-1}}^{x_{2}^{m-1}} f(x)\phi_{2}^{m-1}(x)dx$$
$$= (f, \phi_{2}^{m-1}) = F_{2}^{m-1}$$



Resultados

Matriz de rigidez K:

$$K_{11} = K_{11}^{1};$$

$$K_{ii} = K_{22}^{i-1} + K_{11}^{i}, \text{ para } i = 2, 3, ..., m-1;$$
 (3)

$$\begin{split} K_{i,i+1} = K_{12}^i \Rightarrow K_{i+1,i} = K_{21}^i, \text{ para } i = 1,2,3,\ldots,m-1; \text{ (Simetria: } K = K^T\text{)} \\ K_{mm} = K_{22}^{m-1}. \end{split}$$

Vetor força F:

$$\begin{split} F_1 &= F_1^1; \\ F_i &= F_2^{i-1} + F_1^i, \text{ para } i = 2,3,\dots,m-1; \\ F_m &= F_2^{m-1}. \end{split} \tag{4}$$

Referências I



Liu, I.S.; Rincon, M.A.. Introdução ao Método de Elementos Finitos, Análise e Aplicação. IM/UFRJ, 2003.