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## Problem. $HZ$ production at a $e^+e^-$ collider

Consider the process

$$e^-(p_1)e^+(p_2) \rightarrow h(k_1)Z(k_2) \quad (1)$$

with  $q = p_1 + p_2$ . The Feynman rule for  $Zf\bar{f}$ -vertex is

$$-i\frac{g}{4\cos\theta_W}\gamma_\mu(V - A\gamma_5), \quad V = 2I_W^3 - 4Q\sin^2\theta_W, \quad A = 2I_W^3. \quad (2)$$

being  $I_W^3 = \pm 1/2$  the isospin of the fermion and  $Q$  its electric charge. In the electron case ( $f = e^-$ ):

$$V = -1 + 2\sin^2\theta_W, \quad A = -1. \quad (3)$$

The Feynman rule for the  $hZZ$  vertex is

$$ig\frac{m_Z}{\cos\theta_W}g_{\mu\nu} = i\frac{2m_Z^2}{v}g_{\mu\nu}. \quad (4)$$

- (a) Draw the Feynman diagram for the process and write down the amplitude of the process using the couplings mentioned above. Use unitary gauge.
- (b) Write down the amplitude in terms of  $G_F$  by using

$$\frac{m_W}{m_Z} = \cos\theta_W, \quad \frac{g^2}{8m_W^2} = \frac{G_F}{\sqrt{2}} \quad (5)$$

- (c) Compute the unpolarised squared matrix-element  $|\bar{M}|^2$ , summed over final state spin and averaged over the initial polarizations, taking  $m_e \rightarrow 0$
  - (d) Compute the differential cross section of the process in the C.O.M frame.
  - (e) Compute the total cross section of the above process. Check how the total cross section behaves with variation of  $m_h$ .
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P.1. HZ production at a  $e^+e^-$  collider. (ROHAN KULKARNI)

Consider :  $e^-(p_1)e^+(p_2) \rightarrow h(k_1) Z(k_2)$  with  $q = p_1 + p_2$ .

$$Z\bar{f}f \text{ vertex} : -i \frac{g}{4 \cos \theta_w} \gamma_\mu (V - A \gamma_5), \quad V = 2 Z_w^2 - 4 Q \sin^2 \theta_w \xrightarrow{\text{Charge}}$$

$$A = 2 Z_w^3 \quad \downarrow \quad Z_w = \pm \frac{1}{2}$$

$$\text{For } (f = e^\pm) : \quad V = -1 + 2 \sin^2 \theta_w, \quad A = -1 \quad \text{isospin of } f.$$

$$hZZ \text{ vertex} : \quad i g \frac{m_Z}{\cos \theta_w} g_{\mu\nu} = i \frac{2 m_Z^2}{V} g_{\mu\nu}$$

- (a) Draw a feynman diagram for the process & write down amplitude using couplings abv. Use unitary gauge.

$$iM = \begin{array}{c} \text{Feynman Diagram: } e^+ \text{ and } e^- \text{ enter from left, } Z \text{ and } H \text{ exit to right. } \\ \text{Diagram shows } e^+ \text{ interacting with } Z \text{ to produce } H, \text{ which then interacts with } Z \text{ to produce } e^-. \end{array} = \bar{V}(p_2) \frac{-i g}{4 \cos \theta_w} \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} \left( g_{\mu\nu} - \frac{k^\mu k^\nu}{m_Z^2} \right) \frac{i g m_Z}{\cos \theta_w} g_{\mu\nu} \epsilon^\nu(k_2)$$

$$(b) \text{ Write amplitude in terms of } G_F \text{ using} : \frac{m_w}{m_Z} = \cos \theta_w ; \frac{g^2}{8 \bar{h}_w^2} = \frac{G_F}{\sqrt{2}}$$

$$\rightarrow M_w = \sqrt{\frac{1}{8 G_F}} g \Rightarrow \frac{1}{\cos \theta_w} = \sqrt{\frac{8 G_F}{\sqrt{2}}} \frac{m_Z}{g}$$

$$\begin{aligned} \therefore iM &= \bar{V}(p_2) \frac{1}{g} \frac{8 G_F}{\sqrt{2}} m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} \left( -g_{\mu\nu} + \frac{k^\mu k^\nu}{m_Z^2} \right) g_{\mu\nu} \epsilon^\nu(k_2) \\ &= \bar{V}(p_2) \sqrt{2} G_F m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} \left( -g_{\mu\nu} + \frac{k^\mu k^\nu}{m_Z^2} \right) g_{\mu\nu} \epsilon^\nu(k_2) \\ &= \bar{V}(p_2) \sqrt{2} G_F m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} \left( \frac{k^2}{m_Z^2} - 1 \right) \epsilon^\mu(k_2) \end{aligned}$$

- (c) Compute  $|M|^2$  : Summed over final spin state & Avg over init. polariz.

$$m_c \rightarrow 0$$

$$+ |M|^2 = \frac{G_F^2}{2(s - m_Z^2)^2} m_Z^6 \sum_{S, S'} [\bar{V}_S \gamma_\mu (V - A \gamma_5) u_S] \left( \sum_p \epsilon_p^\mu \epsilon_p^\nu \right) \left( \frac{k^2}{m_Z^2} - 1 \right) [ \bar{U}_{S'} \gamma_\mu (V - A \gamma_5) \bar{V}_{S'} ]$$

$$\begin{aligned}
&= \frac{G_F^2}{2(s-m_Z^2)} m_Z^6 \left( \frac{k^2}{m_Z^2} - \zeta \right) \sum_{S,S'} \left\{ \overline{[U_S \gamma_\mu (V - A \gamma_5) U_S]} \right. \\
&\quad \left. \left( \frac{k_2^\mu k_2^\nu}{m_Z^2} - g^{\mu\nu} \right) \overline{[U_{S'} \gamma_\nu (V - A \gamma_5) U_{S'}]} \right\} \\
&= \frac{G_F^2}{2(s-m_Z^2)} m_Z^6 \left( \frac{k^2}{m_Z^2} - \zeta \right) \times \left[ \text{Tr} [\gamma_2 \gamma_\mu (V - A \gamma_5) \gamma_1 \gamma_\nu (V - A \gamma_5)] \left( \frac{k_2^\mu k_2^\nu}{m_Z^2} - g^{\mu\nu} \right) \right]
\end{aligned}$$

$$\text{Tr} [\gamma_2 \gamma_\mu (V - A \gamma_5) \gamma_1 \gamma_\nu] = p_2^\lambda p_1^\sigma \text{Tr} [\gamma_\lambda \gamma_\mu \gamma_\sigma \gamma_\nu (V - A \gamma_5) \gamma_\alpha \gamma_\nu (V - A \gamma_5)]$$

$$\begin{aligned}
&= p_2^\lambda p_1^\sigma \left\{ -V^2 \text{Tr} [\gamma_\lambda \gamma_\mu \gamma_\sigma \gamma_\nu] - VA \text{Tr} [\gamma_\lambda \gamma_\mu \gamma_\sigma \gamma_\nu \gamma_\rho] \right. \\
&\quad \left. - VA \text{Tr} [\gamma_\lambda \gamma_\mu \gamma_\nu \gamma_\sigma \gamma_\rho] + A^2 \text{Tr} [\gamma_\lambda \gamma_\mu \gamma_\nu \gamma_\sigma \gamma_\rho] \right\}
\end{aligned}$$

$$\begin{aligned}
&= p_2^\lambda p_1^\sigma \left\{ (V^2 + A^2) (g_{\lambda\mu} g_{\sigma\nu} - g_{\lambda\nu} g_{\mu\nu} + g_{\lambda\nu} g_{\mu\sigma}) \right. \\
&\quad \left. - 8VA \epsilon_{\lambda\mu\nu\sigma\rho} \right\}
\end{aligned}$$

$$\Rightarrow g^{\mu\nu} \text{Tr} [\gamma_2 \gamma_\mu (V - A \gamma_5) \gamma_1 \gamma_\nu]$$

$$\begin{aligned}
&= \left\{ (V^2 + A^2) [(\bar{p}_2 \cdot p_1) - \zeta (\bar{p}_2 \cdot p_1) + (\bar{p}_2 \cdot p_2)] \right\} \\
&= -2(V^2 + A^2) (\bar{p}_2 \cdot p_1)
\end{aligned}$$

$$k^\mu k^\nu \text{Tr} [\gamma_2 \gamma_\mu (V - A \gamma_5) \gamma_1 \gamma_\nu] = \left\{ (V^2 + A^2) [\bar{2}(\bar{p}_1 k_2) (\bar{p}_2 k_2) - (\bar{p}_1 \bar{p}_2) k_2^2] \right\}$$

$$\begin{aligned}
\Rightarrow |M|^2 &= \frac{G_F^2}{2(s-m_Z^2)} m_Z^6 \left( \frac{k^2}{m_Z^2} - \zeta \right) (V^2 + A^2) \left[ \frac{2(\bar{p}_1 k_2) (\bar{p}_2 k_2)}{m_Z^2} - (\bar{p}_1 \bar{p}_2) \right. \\
&\quad \left. + 3(\bar{p}_1 \bar{p}_2) \right]
\end{aligned}$$

$$= \frac{G_F^2}{2(s-m_Z^2)} m_Z^6 \left[ \frac{k^2}{m_Z^2} - \zeta \right] 2(V^2 + A^2) \left\{ (\bar{p}_1 \bar{p}_2) + \frac{(\bar{p}_1 k_2) (\bar{p}_2 k_2)}{m_Z^2} \right\}$$

(d) Cross Section in Com frame.

$$\frac{d\sigma}{d\Omega} = \frac{1}{6\pi T^2 S} \frac{p_F}{p_i} \times \frac{G_F^2}{2(s-m_Z^2)} m_Z^6 \left( \frac{k^2}{m_Z^2} - \zeta \right) 2(V^2 + A^2) \left\{ (\bar{p}_1 \bar{p}_2) + \frac{(\bar{p}_1 k_2) (\bar{p}_2 k_2)}{m_Z^2} \right\}$$

(3) Width of the  $W$  boson.

- (a) Write all the allowed 2-body decays of the  $W$  boson in the SM. Neglect quark mixing ( $V_{CKM} = 1$ )

→ Leptons :  $W^- \rightarrow e^- \bar{\nu}_e$

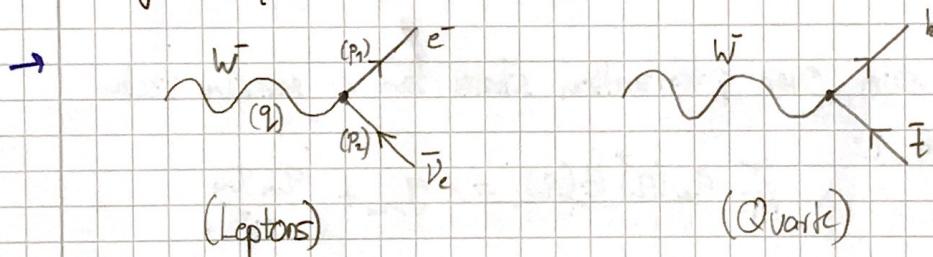
$$W^- \rightarrow \tau^- \bar{\nu}_\tau$$

$$W^- \rightarrow \mu^- \bar{\nu}_\mu$$

Quarks :

- $W^- \rightarrow \bar{u}s$
- $W^- \rightarrow \bar{d}u$
- $W^- \rightarrow \bar{c}s$
- $W^- \rightarrow \bar{s}c$
- $W^- \rightarrow \bar{u}d$
- $W^- \rightarrow \bar{d}u$
- $W^- \rightarrow \bar{t}b$
- $W^- \rightarrow \bar{b}t$
- $W^- \rightarrow \bar{e}\bar{\nu}_e$
- $W^- \rightarrow \bar{\nu}_e e^-$

- (b) Draw the Feynman diagrams for a decay to Leptons & for a decay to quarks.



- (c) Neglecting quark mixing, the Feynman rule for  $W\bar{f}f'$  interaction is the same for quarks & leptons :  $i \frac{g}{\sqrt{2}} \gamma_\mu P_L = i \frac{g}{2\sqrt{2}} \gamma_\mu (1 - \gamma_5)$

Write down the amplitude for  $W^-(q) \rightarrow e^-(p_1) \bar{\nu}_e(p_2)$

→ Rules :

- Incoming  $W$ -boson :  $E_\mu(q)$
- Electron :  $\bar{u}(p_1)$
- Anti-neutrino :  $\nu(p_2)$
- Vertex factor :  $-i \frac{g}{\sqrt{2}} \gamma^\mu P_L = i \frac{g}{2\sqrt{2}} \gamma^\mu (1 - \gamma_5)$

$$-i M = E_\mu(q) \bar{u}(p_1) \left( -i \frac{g}{\sqrt{2}} \gamma^\mu (1 - \gamma_5) \right) \nu(p_2)$$

$$M = \frac{g}{\sqrt{2}} E_\mu(q) \bar{u}(p_1) \frac{g}{2} \gamma^\mu (1 - \gamma_5) \nu(p_2)$$

If we take weak-charged current as :

$$j^\mu = \bar{u}(p_1) \gamma^\mu \frac{1}{2} (1 - \gamma^5) v(p_2)$$

$$M = \frac{g}{\sqrt{2}} \epsilon_\mu(q) j^\mu$$

(d) Take the Fermion masses to be zero, & compute the squared matrix element  $|M|^2$  for  $W^- \rightarrow e^- \bar{\nu}_e$  by

→ Averaging over the polarizations of the initial state

→ Summing over spin of final state particles.

Write it in terms of  $G_F$  &  $m_W$

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8m_W^2}$$

$$\frac{m_W}{m_Z} = \cos \theta_W$$

& work out the dependence on the particles' momenta,  
working in  $W$  rest frame.

Hint: The sum over polarization states for a massive vector boson is

$$\sum_{\text{pol}} \epsilon_\mu(q)^* \epsilon_\nu(q) = -g_{\mu\nu} + \frac{q_\mu q_\nu}{m^2}$$

$$\text{Also, } \gamma_\mu \gamma^\nu \gamma^\lambda = -2 \gamma_\nu$$

→ Start with Lepton current :  $j^\mu = \bar{u}(p_1) \gamma^\mu \frac{1}{2} (1 - \gamma^5) v(p_2)$

- Now use the rest frame of the  $W^-$  boson & also  $m_W \gg m_e$  to neglect  $m_{\text{lepton}}$ .

$$q = (m_W, \vec{0})$$

$$p_1 = ($$

$$|M|^2 = M M^* = \sum \frac{g^2}{2} \epsilon_\mu^{(q)} j_\mu^{(q)} (\epsilon_\nu^{(q)*})^* (j_\nu^{(q)})^*$$

$$= \frac{g^2}{2} \sum \sum \epsilon_\mu^{(q)} (\epsilon_\nu^{(q)})^* \sum j^\mu$$

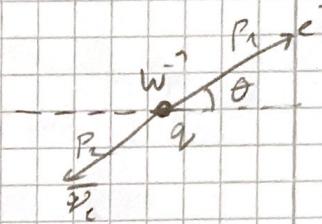
$$\begin{aligned}
\rightarrow |M|^2 &= MM^\dagger = \underbrace{\frac{1}{3} \sum_{\sigma_1} \sum_{\sigma_2}}_{\text{Avg over pol.}} \frac{g^2}{2} E_m^{(\alpha)}(q) E_n^{(\alpha)}(q)^* J_{(\sigma_1)}^\mu J_{(\sigma_2)}^\nu \\
&= \frac{g^2}{6} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] \sum_{\sigma_1} \overline{U}_{(\sigma_1)}(p_1) \gamma^\mu \frac{1}{2} (1-\gamma^5) V_{(\sigma_1)}(p_2) \overline{V}_{(\sigma_2)}(p_2) \frac{1}{2} (1-\gamma^5) U_{(\sigma_2)}(p_1) \\
&\quad \xrightarrow{\text{Sum over } \sigma} \text{Dirac sum} \\
&= \frac{g^2}{6} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] \sum_{\sigma_1} \overline{U}_{(\sigma_1)}(p_1) \gamma^\mu \frac{1}{2} (1-\gamma^5) (p_1 + m_{\tilde{e}_1}) \frac{\gamma^\nu}{2} (1-\gamma^5) U_{(\sigma_1)}(p_1) \\
&\quad \xrightarrow{\text{Sum over } \sigma} \text{Dirac sum} \\
&= \frac{g^2}{6 \cdot 5} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] \gamma^\mu (1-\gamma^5) (p_1 + m_{\tilde{e}_1}) \frac{\gamma^\nu}{2} (1-\gamma^5) \\
&= \frac{g^2}{25} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] \gamma^\mu (1-\gamma^5) \gamma^\beta(p_1)_\beta \gamma^\nu(p_2)_\nu \gamma^\lambda (1-\gamma^5) \\
&= \frac{g^2}{25} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] (p_1)_\mu (p_2)_\nu \text{Tr} [\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\lambda (1-\gamma^5)] \\
&= \dots \text{Tr} [\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\lambda (1-\gamma^5)] \\
&= \dots \text{Tr} [(\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\lambda) - (\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\lambda)] \\
&= \dots \left[ \frac{1}{4} (g^{\mu\alpha} g^{\nu\lambda} - g^{\mu\nu} g^{\beta\lambda} + g^{\mu\beta} g^{\nu\lambda}) \right. \\
&\quad \left. - 4 i E^{\mu\beta\nu\lambda} \right] \xrightarrow{\text{Doesn't contribute.}} \\
&= \frac{g^2}{5} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] (p_1)_\mu (p_2)_\nu [g^{\mu\beta} g^{\nu\lambda} - g^{\mu\nu} g^{\beta\lambda} + g^{\mu\lambda} g^{\beta\nu}] \\
&= \frac{g^2}{5} \left[ -g_{\mu\nu} + \frac{g_\mu g_\nu}{m_w^2} \right] ((p_1)^\mu (p_2)^\nu - (p_1)^\nu (p_2)^\mu + (p_1)^\mu \cdot (p_2)^\nu \\
&= \frac{g^2}{5} \left[ \cancel{(p_1 \cdot p_2)} + (p_1 \cdot p_2)^* - \cancel{(p_1 \cdot p_2)} + \frac{1}{m_w^2} \left[ \cancel{(p_1 \cdot q)} \cancel{(p_2 \cdot q)} - \cancel{(p_1 \cdot q)} \cancel{(p_2 \cdot q)} \right. \right. \\
&\quad \left. \left. + (q \cdot q) (p_1 \cdot p_2) \right] \right] \\
&\equiv \frac{g^2}{5} \left[ -\cancel{(p_1 \cdot p_2)} + \frac{(q \cdot q)(p_1 \cdot p_2)}{m_w^2} \right]
\end{aligned}$$

Now take the rest frame for  $W^-$  in  $W^- \rightarrow e^- \bar{\nu}_e$

$$q = (m_W, \vec{0}) , \quad p_1 = (E, E \sin \theta, 0, E \cos \theta)$$

$$p_2 = (E, -E \sin \theta, 0, -E \cos \theta)$$

$$\Rightarrow (p_1 \cdot p_2) = 2E^2 \quad , \quad (q \cdot q) = m_W^2$$



Using

$$|M|^2 = \frac{g^2}{4} \left[ -\cancel{q} \cdot 2E^2 + m_W^2 \frac{2E^2}{m_W^2} \right]$$

$$= -\frac{g^2}{4} \cancel{q} \cdot 2E^2 + \frac{g^2}{4} m_W^2 \frac{2E^2}{m_W^2}$$

$$= -g^2 2E^2 + g^2 \frac{E^2}{2}$$

$$= -\frac{3}{2} E^2 g^2$$

$$= \left( -\frac{3}{2} \right) g^2 m_W^2$$

Some mistake,  
I should get

$$\frac{1}{3} g^2 m_W^2$$

- One can also find that

$$|M_-|^2 = g_{\phi}^2 m_W^2 \frac{1}{4} (1 + \cos \theta)^2$$

$$|M_L|^2 = g^2 m_W^2 \frac{1}{2} \sin^2 \theta$$

$$|M_+|^2 = g_{\phi}^2 m_W^2 \frac{1}{4} (1 - \cos \theta)^2$$

Three possible helical polarization modes

- (e) Using this result, compute the total width of write down  $|M|^2$  for  $W^- \rightarrow \bar{u}d$ . What are the changes compared to the leptonic one.

(lowest order)  $\rightarrow$

$$\Gamma(W^- \rightarrow e^- \bar{\nu}_e)$$

$$\Gamma(W^- \rightarrow \bar{u}d) = 6 \Gamma(W^- \rightarrow e^- \bar{\nu}_e)$$

I am assuming due to the heavier weight of quarks & more generation.

$$(f) \quad d\Gamma = \frac{1}{64\pi^2 m} |\bar{M}|^2 d\Omega \quad d\Omega = d\phi \, d\cos\theta$$

• Compare to scattering formula. ( $2 \rightarrow 2$ )

→ Pretty much the same

$$d\sigma = \frac{1}{64\pi^2 s} |\bar{m}|^2 d\Omega$$

$\rightarrow$  We have 2 momenta coming in.

• Compute the partial widths  $W^- \rightarrow e^- \bar{\nu}_e$

$$W^- \rightarrow \bar{u} d$$

$$W^- \rightarrow e^- \bar{\nu}_e$$

$$\int d\Gamma = \int \frac{1}{64\pi^2 m_W} |\bar{M}|^2 d\Omega = \frac{1}{64\pi^2 m_W} \int \int \frac{1}{3} g^2 \Omega^2 d\Omega$$

$$\Gamma_{e\nu} = \frac{4\pi g^2}{64\pi} = = \frac{g^2}{16}$$

$$W^- \rightarrow \bar{u} d$$

$$\Gamma_{ud} = 6 \Gamma_{e\nu} = \frac{3}{8} g^2$$