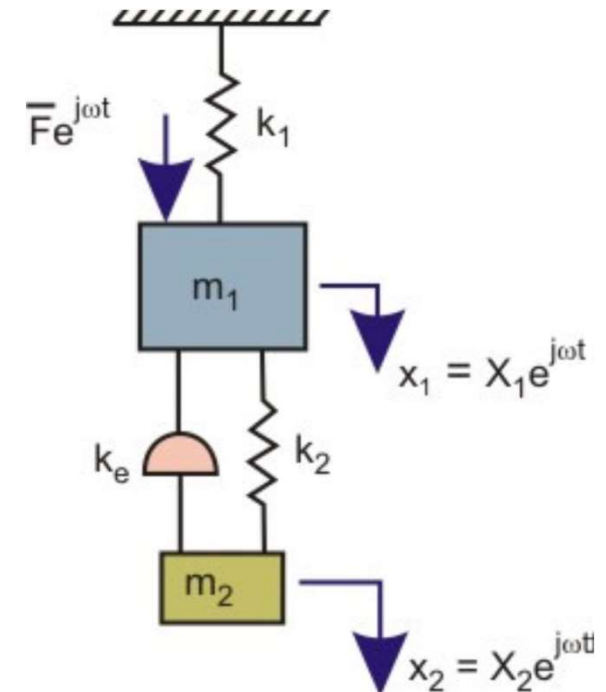


# Active Dynamic Vibration Absorber

- It may be noted that Passive Neutralizer **eliminates primary response** only at a **particular frequency**.
- Use of **active element** - for example, a hydraulic actuator would increase the advantage of tuned mass damping for a broad frequency range.

Here,  $m_1$  denotes the **primary mass** and  $k_1$  the **primary stiffness**. The **damping** of the **primary system** is **neglected**. The system is subjected to a harmonic excitation  $\bar{F}e^{j\omega t}$ .

The primary system is **attached** to a secondary system of fixed mass  $m_2$  and stiffness  $k_2$ . However, there is an additional spring element with **variable stiffness** ' $k_e$ ' representative of a **hydraulic actuator**.



The governing EOM of the two DOF system may be written as

$$m_1 \ddot{x}_1 + k_1 x_1 + k_2(x_1 - x_2) = \bar{F} e^{j\omega t} + k_e x_2 \quad (1)$$

$$m_2 \ddot{x}_2 + k_2(x_2 - x_1) = k_e x_2 \quad (2)$$

Using  $x_1 = X_1 e^{j\omega t}$  and  $x_2 = X_2 e^{j\omega t}$ , we get

From equation (2),

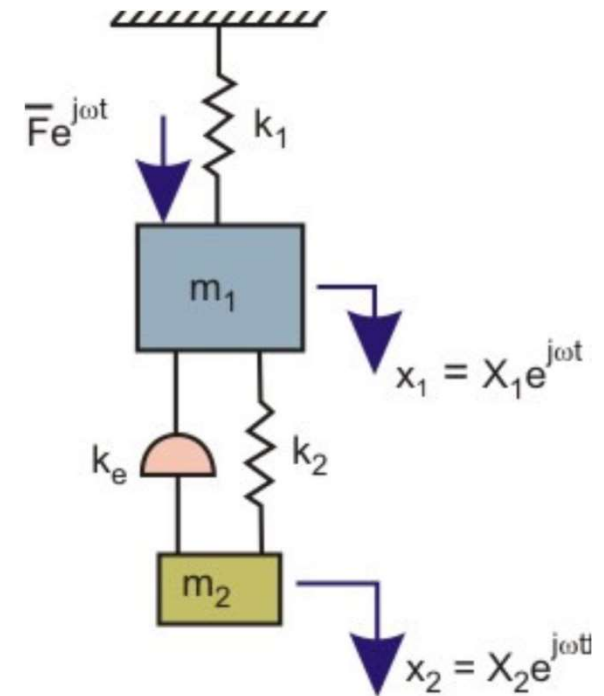
$$-\omega^2 m_2 X_2 + k_2(X_2 - X_1) = -k_e X_2$$

or

$$(k_2 - m_2 \omega^2 + k_e) X_2 = k_2 X_1$$

or

$$X_2 = \frac{k_2}{k_2 - m_2 \omega^2 + k_e} X_1 \quad (3)$$



Similarly, from Equation 1, we get

$$m_1 \ddot{x}_1 + k_1 x_1 + k_2 (x_1 - x_2) = \bar{F} e^{j\omega t} + k_e x_2$$

$$(k_1 - \omega^2 m_1) X_1 + k_2 (X_1 - X_2) = \bar{F} + k_e X_2$$

or

$$(k_1 + k_2 - \omega^2 m_1) X_1 - (k_2 + k_e) X_2 = \bar{F}$$

or

$$(k_1 + k_2 - \omega^2 m_1) X_1 - (k_2 + k_e) X_2 = \bar{F}$$

or

$$\left[ (k_1 + k_2 - \omega^2 m_1) - \frac{k_2^2 + k_e k_2}{k_2 - \omega^2 m_2 + k_e} \right] X_1 = \bar{F}$$



$$\left[ (k_1 + k_2 - \omega^2 m_1) - \frac{k_2^2 + k_e k_2}{k_2 - \omega^2 m_2 + k_e} \right] X_1 = \bar{F}$$

Thus, when the hydraulic actuator is switched on the active displacement of the primary mass  $X_{1a}$  may be written as:

$$X_{1a} = \bar{F} \frac{(k_2 - \omega^2 m_2 + k_e)}{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2 + k_e) - k_2^2 - k_e k_2}$$

When the hydraulic system is switched off, the passive displacement of the primary mass  $X_{1p}$  may be written as:

$$X_{1p} = \bar{F} \frac{(k_2 - \omega^2 m_2)}{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2) - k_2^2}$$

The ratio of active and passive displacement of the primary mass brings out the efficiency of the new system. Therefore,

$$\frac{X_{1a}}{X_{1p}} = \frac{(k_2 - \omega^2 m_2 + k_e)}{(k_2 - \omega^2 m_2)} \times \frac{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2) - k_2^2}{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2 + k_e) - k_2^2 - k_e k_2}$$



$$\frac{X_{1a}}{X_{1p}} = \frac{(k_2 - \omega^2 m_2 + k_e)}{(k_2 - \omega^2 m_2)} \times \frac{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2) - k_2^2}{(k_1 + k_2 - \omega^2 m_1)(k_2 - \omega^2 m_2 + k_e) - k_2^2 - k_e k_2}$$

For a simple case, use  $k_1 = k_2 = k$ ,  $m_1 = m_2 = m$ ,  $\Omega^2 = \frac{\omega^2}{(k/m)}$

$$\frac{X_{1a}}{X_{1p}} = \frac{\left(1 + \frac{k_e}{k} - \Omega^2\right)}{1 - \Omega^2} \times \frac{(2 - \Omega^2)(1 - \Omega^2) - 1}{(2 - \Omega^2)\left(1 + \frac{k_e}{k} - \Omega^2\right) - 1 - \frac{k_e}{k}}$$

As a test case,

$$\text{for } \frac{k_e}{k} = -2, \quad \frac{X_{1a}}{X_{1p}} = \left| \frac{\Omega^6 - 2\Omega^4 - 2\Omega^2 + 1}{\Omega^6 - 2\Omega^4 + 1} \right|$$

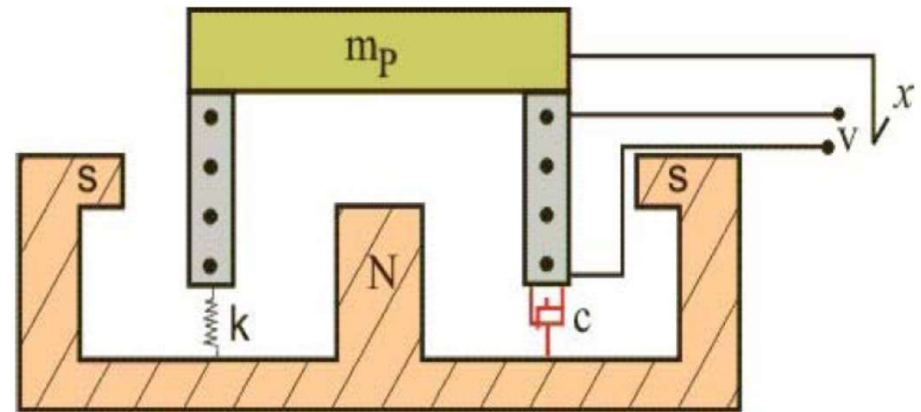
$$\text{for } \frac{k_e}{k} = +2, \quad \frac{X_{1a}}{X_{1p}} = \left| \frac{\Omega^6 - 6\Omega^4 + 10\Omega^2 - 3}{\Omega^6 - 6\Omega^4 + 8\Omega^2 - 3} \right|$$

From these expressions, one can check that the negative feedback system with  $k_e/k = -2$  works better for a wider frequency range.



# Active DVA

## Another Active DVA – Proof Mass Actuator



- A **Proof mass**  $m_p$  is connected to a **magnetic base** with **spring**  $k$  and **damper**  $c$ .
- The proof mass is placed over a solenoid in which magnetic field could be generated by passing current through coils.
- The **resistance** of the coil is  $R$ , **inductance**  $L$ , **current** passing through the coil is  $i$  and the **proportionality constant** corresponding to **back EMF** is  $k_b$ .



# What is **back EMF**?

- When device like a refrigerator or an air conditioner (anything with a motor) first turns on in our house, the lights often dim momentarily.
- Because, a motor has coils turning inside magnetic field, and a coil turning inside a magnetic field induces a **back emf**, which acts against the applied voltage and reduces the current flowing through the coils of the motor.

If the applied voltage is  $\Delta V$ , then the initial current flowing through a motor with coils of resistance  $R$  is:

$$I = \frac{\Delta V}{R}, \text{ for example } I = 120/6 = 20A$$

- A device **drawing that much current reduces the voltage and current** provided to other electrical equipment in your house, causing **lights to dim**.
- When the motor is spinning and generating a back emf  $E_b$  the current is reduced to:

$$I = \frac{\Delta V - E_b}{R}, \text{ for example } I = \frac{120-108}{6} = 2A$$

$$E_b = \frac{PZ \phi N}{A 60} = k_b \frac{\phi N}{60}$$

Back EMF regulates the flow of armature current i.e., it automatically changes the armature current to meet the load requirement.

$P$  – Number of poles of the machine  
 $\phi$  – Flux per pole in Weber.  
 $Z$  – Total number of armature conductors.  
 $N$  – Speed of armature in r.p.m.  
 $A$  – Number of parallel paths in the armature winding  
 $k_b$  - Back emf constant



The EOM are:

$$Ri + L \frac{di}{dt} = V - k_b \dot{x} \quad (1)$$

$$m_p \ddot{x} + k_p x + c_p \dot{x} = k_a i \quad (2)$$

$k_a$  = Current constant  
 $k_p$  = Proof mass stiffness  
 $c_p$  = Proof mass damping  
 $m_p$  = Proof mass  
 $L$  = inductance  
 $k_b$  = back emf constant

Converting equation (1) into frequency domain using Laplace transform,

$$RI + LsI = \bar{V} - k_b sX$$

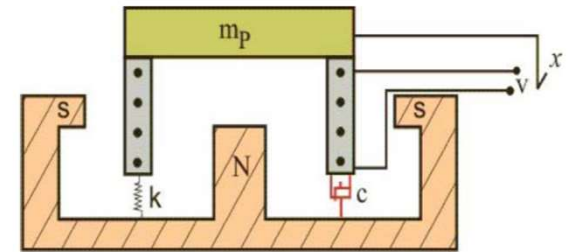
$$I = \frac{1}{R+Ls} (\bar{V} - k_b sX) \quad (3)$$

Similarly for equation (2),

$$s^2 m_p X + k_p X + c_p sX = k_a I$$

$$(s^2 m_p + k_p + c_p s)X = k_a I \quad (4)$$

Now, using equation (3) in (4), we get





$$(s^2 m_p + k_p + c_p s)X = k_a \frac{1}{R + Ls} (\bar{V} - k_b sX)$$

$$\text{Denoting, } \frac{1}{R+Ls} = \frac{1}{Z(s)} = \tilde{A}(s)$$

$$(s^2 m_p + k_p + c_p s + s k_a k_b \tilde{A}(s))X = k_a \tilde{A}(s) \bar{V}$$

$$X = \frac{k_a \tilde{A}(s)}{(s^2 m_p + k_p + s(c_p + k_a k_b \tilde{A}(s)))} \bar{V} \quad (5)$$

$k_a$  = Current constant  
 $k_p$  = Proof mass stiffness  
 $c_p$  = Proof mass damping  
 $m_p$  = Proof mass  
 $L$  = inductance  
 $k_b$  = back emf constant  
 $Z(s)$  = Impedance  
 $\tilde{A}(s)$  = Admittance

Also, force exerted by the proof mass actuator on the base is  $F(t) = -m_p \ddot{x}$

Using Laplace transform, thus,  $F(s) = -s^2 m_p X$

Using equation (5), we get

$$F(s) = \frac{-s^2 m_p k_a \tilde{A}(s)}{(s^2 m_p + k_p + s(c_p + k_a k_b \tilde{A}(s)))} \bar{V} \quad (6)$$

This relationship tells us how **force F** will be **generated** by the **proof mass actuator** upon **application of voltage  $\bar{V}$** .



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## Application of active DVA

Consider a SDOF system (undamped, mass  $m_1$ ) subjected to base excitation.

$$m_1 \ddot{x}_1 + k_1(x_1 - y) = F(t)$$

Converting above equation into frequency domain,

$$s^2 m_1 X_1 + k_1 X_1 - k_1 Y = F(s)$$

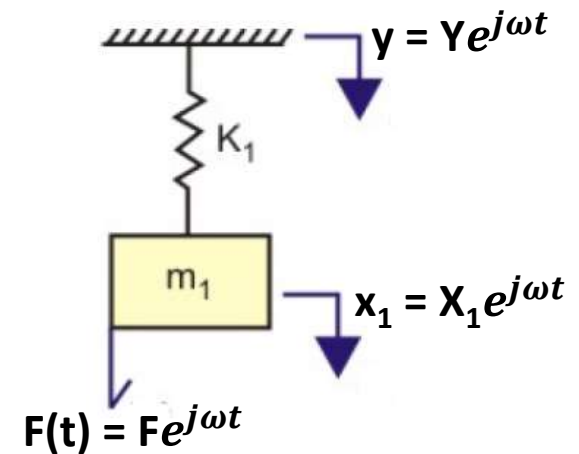
$$(s^2 m_1 + k_1) X_1 = F(s) + k_1 Y$$

Therefore,

$$X_1 = \frac{F(s) + k_1 Y}{s^2 m_1 + k_1}$$

Now, we already derived this equation

$$F(s) = \frac{-s^2 m_p k_a \tilde{A}(s)}{(s^2 m_p + k_p + s(c_p + k_a k_b \tilde{A}(s)))} \bar{V} \quad (6)$$



Plugging the output of the Active DVA into the system and using Equation (6), we get the new displacement of  $m_1$  as

$$X_1 = \frac{k_1 Y}{s^2 m_1 + k_1} - \frac{s^2 m_p k_a \tilde{A}(s)}{(s^2 m_p + k_p + s(c_p + k_a k_b \tilde{A}(s)))} \bar{V} \quad (7)$$

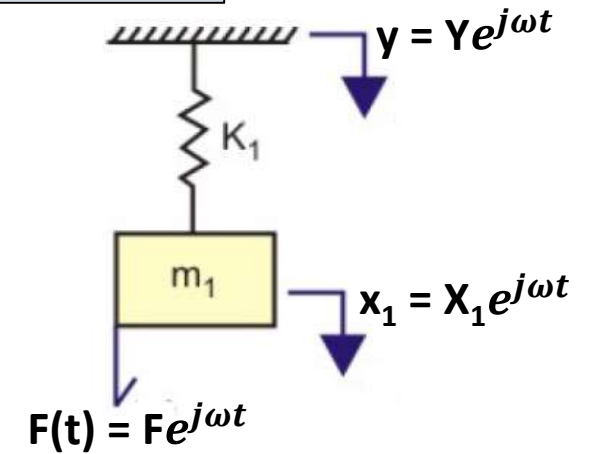
**A Special case:**

When  $\tilde{A}(s)$  is constant,  $\omega_1^2 = \frac{k_1}{m_1}$ ,  $\omega_p^2 = \frac{k_p}{m_p}$ ,  $\bar{V} = kX_1$

$$X_1 = \frac{\omega_1^2}{s^2 + \omega_1^2} Y - \frac{\bar{\alpha} s^2}{(s^2 + \bar{\beta} s + \omega_p^2)} X_1$$

where,

$$\bar{\alpha} = \frac{\tilde{A} k_a}{k} \quad \bar{\beta} = \frac{c_p + \tilde{A} k_a k_b}{m_p}$$



$$\text{Therefore, } \left(1 + \frac{\bar{\alpha} s^2}{(s^2 + \bar{\beta} s + \omega_p^2)}\right) X_1 = \frac{\omega_1^2}{s^2 + \omega_1^2} Y$$

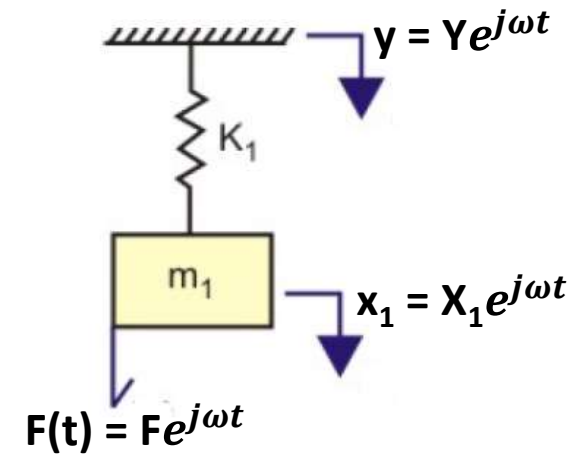


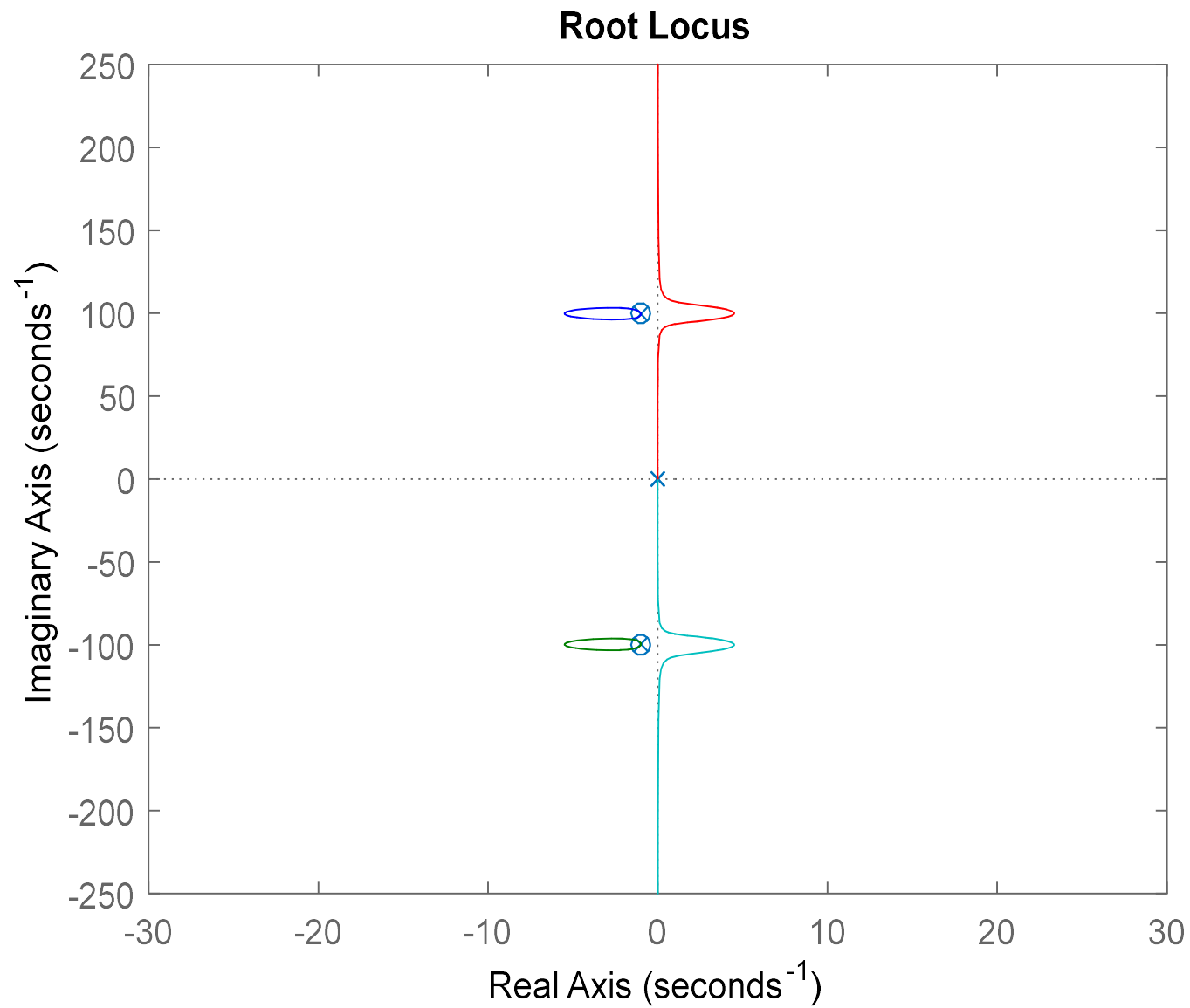
One can show a wide band amplitude reduction of the primary mass ( $X_1$ ) by suitably choosing  $k_a$  and  $k_b$  in equation (7)

$$\frac{X_1}{Y} = \frac{\omega_1^2 (s^2 + \bar{\beta}s + \omega_p^2)}{(s^2 + \omega_1^2)[(s^2(1 + \bar{\alpha}) + \bar{\beta}s + \omega_p^2)]}$$

$$= \frac{\omega_1^2}{1 + \bar{\alpha}} \frac{s^2 + \bar{\beta}s + \omega_p^2}{(s^2 + \omega_1^2)[s^2 + \frac{\bar{\beta}}{1 + \bar{\alpha}}s + \frac{\omega_p^2}{1 + \bar{\alpha}}]}$$

$$\bar{\alpha} = \frac{\tilde{A}k_a}{k} \quad \bar{\beta} = \frac{c_p + \tilde{A}k_a k_b}{m_p}$$





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