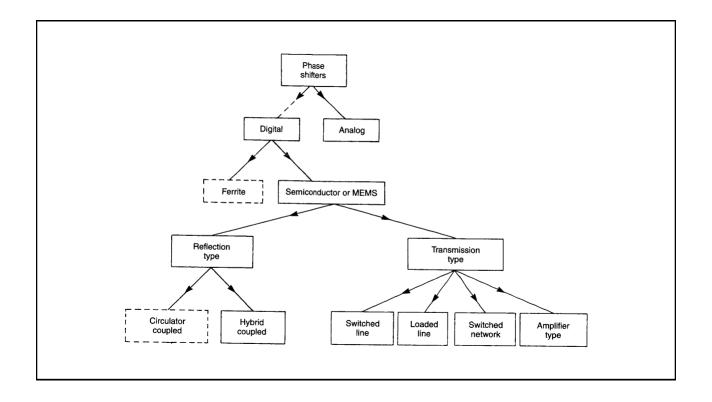
## **Phase Shifter Design**

- A phase shifter is a two-port network with a provision that the phase difference between the output and the input signals may be controlled by a control signal (DC bias).
- Phase shifters are called DIGITAL when the differential phase shift can be changed by only a few predetermined discrete values, such as 180°, 90°, 45°, 22.5°, and 11.25°. On the other hand, in ANALOG phase shifters, the differential phase shift can be varied in a continuous manner by a corresponding continuous variation of the control signal.



### **Switched-Line Phase Shifters**

Two SPDT switches are used to route the signal via one of the two alternative transmission path lengths  $l_1$  or  $l_2$ .

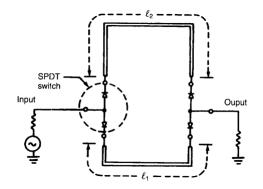
When the signal passes through the longer path, it goes through an additional phase delay given by

$$\Delta \phi = \beta (l_2 - l_1) = \frac{2\pi f}{v_p} (l_2 - l_1)$$

The differential phase shift  $\Delta \phi$  is directly proportional to frequency. Because of this feature, switched-line phase shifters are also called switchable time delay networks.

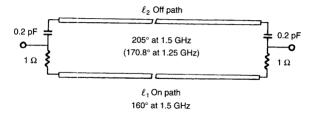
The time delay  $\tau_d$  is given by:

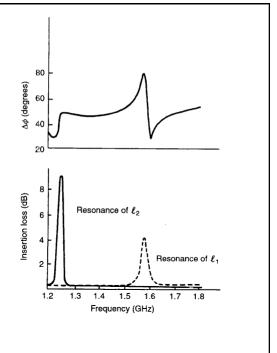
$$\tau_d = \frac{l_2 - l_1}{v_p}$$



[Single-bit switched-line phase shifter]

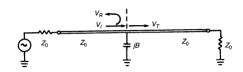
# **Problems in Switched-Line Phase Shifters** (Off-path Resonances)





## **Loaded-Line Phase Shifters**

A very common design for 45° and 22.5° phase bits is known as the loaded-line phase shifter. The mechanism of phase shift in this circuit is based on the loading of a uniform transmission line by a small reactance as shown in Fig.



It can be shown that the transmitted wave undergoes a phase shift  $\Delta \phi$  that depends upon the normalized susceptance b=B/Y. The reflection caused by b is given by

$$\Gamma = \frac{1 - (1 + jb)}{1 + (1 + jb)} = \frac{-jb}{2 + jb}$$

The voltage associated with the transmitted wave is  $V_T = V_I + V_R$ , where  $V_I$  and  $V_R$  are voltages of the incoming wave and the reflected wave, respectively.

The transmission coefficient T can therefore be written as

$$T = \frac{V_T}{V_I} = \frac{V_I + V_R}{V_I} = 1 + \Gamma = \frac{2}{2 + jb}$$

$$V_T = TV_I = V_I \frac{2}{2 + jb} = V_I \left(\frac{4}{4 + b^2}\right)^{1/2} \exp\left[-j \tan^{-1}\left(\frac{1}{2}b\right)\right]$$

The phase difference introduced (phase of  $V_I$  – phase of  $V_T$ ) may therefore be written as

$$\Delta \phi = \tan^{-1}\left(\frac{1}{2}b\right)$$

If the normalized susceptance b=0.2 (i.e., capacitive),  $\Delta\phi=\tan^{-1}(0.1)=0.1$  rad = 5.7°. If, on the other hand, we load the line with b=-0.2 (a shunt inductor), the phase delay  $\Delta\phi$  is negative, that is,  $\Delta\phi=\tan^{-1}(-0.1)\simeq -5.7^{\circ}$ .

That is, the transmitted wave advances in phase as compared with the transmitted wave for b=0.

The insertion loss is given by:

$$\frac{V_T}{V_I} = \left(\frac{4}{4+b^2}\right)^{1/2}$$

As |b| is equal in two states (only its sign is changed), insertion loss has the same value in two states. For b=0.2,  $V_T/V_I=0.995$ , which corresponds to a loss of 0.04 dB. Although, in the present case, the value of insertion loss is not significant, the presence of this loss is an unfavorable feature of the circuit

This drawback of reflection from the susceptance producing the phase shift can be overcome by using two identical susceptances separated by a quarter-wave line length. Such an arrangement is shown in Fig.

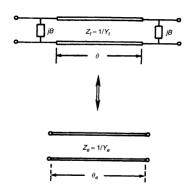
If we consider partial reflections caused by the two susceptances, these reflected waves are almost equal in magnitude and out-of-phase (180° phase difference) when looking into the input terminals. Thus these undesirable reflections cancel each other.

A quantitative assessment of the situation may be obtained by writing a uniform transmission line equivalent circuit as shown in Fig.

Equivalent  $Y_e$  and  $\theta_e$  may be obtained by comparing ABCD matrices of the two networks.

An ABCD matrix of the transmission line of length  $\theta$ , shunt-loaded by susceptances at either end, may be obtained by multiplying ABCD matrices of the shunt susceptance, of the line section, and of the second susceptance at the other end. This yields

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ jB & 1 \end{bmatrix} \begin{bmatrix} \cos \theta & jZ_t \sin \theta \\ jY_t \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} 1 & 0 \\ jB & 1 \end{bmatrix}$$
$$= \begin{bmatrix} (\cos \theta - BZ_t \sin \theta), & j(Z_t \sin \theta) \\ j(2B \cos \theta + Y_t \sin \theta - B^2 Z_t \sin \theta), & (\cos \theta - BZ_t \sin \theta) \end{bmatrix}$$



Circuit configuration for loaded-line phase shifter and equivalent representation.

ABCD of the equivalent transmission line may be written as:

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} = \begin{bmatrix} \cos \theta_e & jZ_e \sin \theta_e \\ jY_e \sin \theta_e & \cos \theta_e \end{bmatrix}$$

By comparing the *A* terms of the above two equations, we can find the length of the equivalent uniform line, as

$$\cos \theta_e = \cos \theta - BZ_t \sin \theta$$

The equivalent admittance  $Y_e$  is obtained by comparing the ratios C/B, as

$$Y_e = Y_t[1 - (BZ_t)^2 + 2BZ_t \cot \theta]^{1/2}$$

When  $\theta = 90^{\circ}$ , the above two relations may be written as:

$$\cos \theta_e = -BZ_t$$

or

$$\theta_e = \frac{1}{2}\pi + BZ_t + \frac{1}{6}(BZ_t)^3$$

$$Y_e = Y_t[1 - (BZ_t)^2]^{1/2}$$

 $Y_e < Y_t$  and its magnitude is independent of the sign of B.

Thus, if only the sign of B is changed (and its magnitude kept unchanged), the circuit can remain matched in both states of the phase shifter. For example, if  $b_1 = 0.2$  and  $b_2 = -0.2$ , we have  $\theta_{e1} = 101.54^{\circ}$  and  $\theta_{e2} = 78.46^{\circ}$ , resulting in a differential phase shift  $\Delta \phi = \theta_{e2} - \theta_{e1} = 23.08^{\circ} \simeq 0.4$  rad.

## Various Configurations for Loaded-Line Phase Shifters

Various designs for loaded-line phase shifters differ to the extent that the susceptances  $B_1$  and  $B_2$  for two states of the phase shifter are realized by different circuit configurations.

The basic design equations are obtained by combining

$$\cos \theta_e = \cos \theta - BZ_t \sin \theta \qquad Y_e = Y_t [1 - (BZ_t)^2 + 2BZ_t \cot \theta]^{1/2}$$

with

$$\Delta \phi = \tan^{-1}\left(\frac{1}{2}b\right)$$

And may be written as:

$$Y_T = Y_0 \sec(\frac{1}{2}\Delta\phi) \sin \theta$$
  

$$B_{1,2} = Y_0 \left[\sec(\frac{1}{2}\Delta\phi) \cos \theta \pm \tan(\frac{1}{2}\Delta\phi)\right]$$

#### 1. Main-Line Mounted Circuit:

In this case, the switching devices are mounted directly across the main line. The idea is to use the high-impedance state capacitance  $C_j$  and the low-impedance state inductance  $L_s$  directly for  $B_1$  and  $B_2$ , respectively. However, it becomes necessary to add an external inductance  $L_e$  in series with the device as shown in Fig.



$$B_1 = \frac{\omega C_j}{1 - \omega^2 L C_j} \qquad B_2 = \frac{-1}{\omega L}$$

where  $L = L_s + L_e$ .

Using the above equations, the following relationship can be derived:

$$Y_0 = \omega C_j \frac{\sin \Delta \phi}{\sin^2(\Delta \phi/2) - \cos^2 \theta}$$

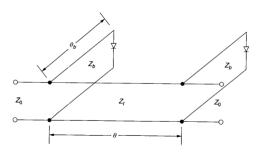
Usually one would like to select  $\theta$  so as to obtain the maximum bandwidth from the phase shifter circuit, which occurs at  $\theta=90^\circ$  . For this case

$$Y_0 = 2\omega C_j \cot\left(\frac{1}{2}\Delta\phi\right)$$

This puts restrictions on the input-output line impedance unless  $C_j$  can be chosen arbitrarily, which becomes difficult because of the limited availability of switching device capacitances. Because of these difficulties, this design configuration has not become very popular.

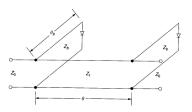
#### 2. Stub-Mounted-Type Circuit:

In the circuit arrangement shown in Figure, switching devices are mounted at the ends of two shunt-connected stubs separated by a line length  $\theta$ . This allows greater design flexibility.



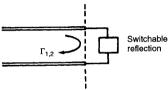
Stub impedance  $Z_b$  and length  $\theta_b$  are determined from device reactances  $X_f$  and  $X_r$  (in low and high impedance states, respectively) as follows:

$$Z_b = \left(\frac{X_f - X_r - X_f X_r (B_1 - B_2)}{B_1 - B_2 - B_1 B_2 (X_f - X_r)}\right)^{1/2}$$
$$\tan \theta_b = \frac{Z_b (1 - X_f B_1)}{X_f - B_1 Z_b^2}$$



where  $Z_t$  and  $\theta$  may be selected for wide bandwidth. It has been reported [12] that the best bandwidths are obtained for  $\theta \sim 90^{\circ}$ . These design methods make use of device reactances only and assume the device resistances to be negligible.

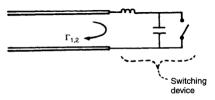
## **Reflection-Type Phase Shifters**



[Basic concept of reflection-type phase shifter]

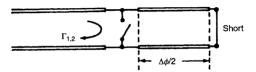
the reflection coefficient is switched from  $\Gamma_1 = |\Gamma_1| \angle \phi_1$  to  $\Gamma_2 = |\Gamma_2| \angle \phi_2$ , the reflected signal undergoes a differential phase shift  $\Delta \phi = \phi_1 - \phi_2$ .

the ratio of the reflected power to the incident power is given by  $|\Gamma|^2$ . Ideally  $|\Gamma|$  should be unity so that there is no loss associated with the phase shifting operation.



[Switchable-reactance-type reflection phase shifter]

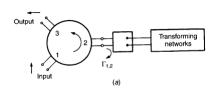
Subnetworks providing the switchable reflection coefficients may be of two different types. In the first group, the reactance terminating the line is changed (say from inductive to capacitive) as shown in Fig. Phase shifters using these subnetworks are called switchable reactance types and constitute the more commonly used variety of reflection phase shifters.

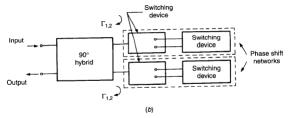


[Switchable-length-type reflection phase shifter]

In the second group of these phase shifter circuits, an additional line length is added at the reflection plane by using an SPST switch. The concept is similar to that used in switchedline phase shifters. Here, as shown in Fig. the signal travels a longer path when the switch is open. The differential phase shift, in this case, is twice the electrical length of the shorted line switched in by the SPST switch.

#### Transformation of a reflection phase bit into a two-port network:





[(a) Circulator-coupled reflection-type phase shifter, (b) Hybrid-coupled reflection-type phase shifter]

It may be noted that hybrid-coupled arrangements require two identical phase bits (and hence twice the number of devices). However, this circuit is preferred to the circulator-coupled phase shifters because of the following two features: (1) 90° hybrids are a lot more integrable in MICs and MMICs than the circulators, and (2) the use of two switching devices in hybrid-coupled phase shifters increases the power-handling capability by a factor of 2, which is needed in several systems.

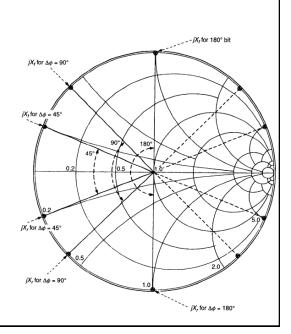
#### Design of reflection-type phase shifters:

Various designs for hybrid-coupled reflection phase shifters use different configurations for the two transforming networks.

In order to obtain a phase shift  $\Delta\phi$ , the reflection coefficient values (as seen at two coupled ports of the hybrid) should differ in phase by  $\Delta\phi$ . Ideally their magnitudes should be unity (no loss). Thus for a 180° bit, values of the reflection coefficients in two states would lie on the outer circle ( $|\Gamma| = 1$ ) of the Smith chart and their locations would be diametrically opposite.

Although an infinite number of such combinations are possible, it can be shown that for a reflection phase shifter using a two-branch branch-line hybrid, the maximum phase shifter bandwidth is obtained when the two values are located at points corresponding to normalized reactances equal to +j and -j on the Smith chart (i.e., when located symmetrically with respect to the X=0 axis).

Similarly for other values of  $\Delta \phi$ ,  $\Gamma_1$  and  $\Gamma_2$  in two states should also be located symmetrically with respect to the X=0 axis.

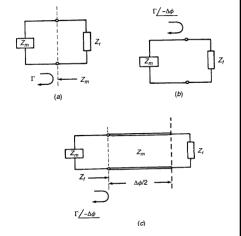


Based on the preceding discussion, we can divide the function of the transforming network into two parts.

- 1. It arranges  $\Gamma_f$  and  $\Gamma_r$  so that the desired phase shift  $\Delta \phi$  is obtained.
- 2. It locates  $\Gamma_f$  and  $\Gamma_r$  symmetrically about the X=0 axis on the Smith chart so that the condition for maximum bandwidth is satisfied.

It may be noted that the second part of the function just mentioned is implemented by connecting a suitable line length between the transforming networks designed for part 1 and the ports of the hybrid. This line length will rotate  $\Gamma_f$  and  $\Gamma_r$  points around the Smith chart without disturbing the phase relationship between them.

Let a network with the equivalent Thévenin impedance  $Z_m$  terminated by an impedance  $Z_r$  (of the high-impedance state of the switching device) produce a reflection coefficient  $\Gamma$ , as shown in Fig. a. When the same network is terminated with an impedance  $Z_f$  (of the low-impedance state of the device), we want a reflection coefficient of  $\Gamma/-\Delta\phi$  (Fig. b). A reflection coefficient of  $\Gamma/-\Delta\phi$  will also be produced when a transmission-line section of length  $\Delta\phi/2$  and impedance  $Z_m$  is connected between the network and  $Z_r$ , as shown in Fig. c. Thus networks shown in Fig. b and c should be equivalent.



For this equivalence to hold good, the following relationship should be satisfied:

$$Z_f = Z_m \frac{Z_r + jZ_m \tan(\Delta\phi/2)}{Z_m + jZ_r \tan(\Delta\phi/2)}$$

For low-loss switching devices,  $Z_f$  and  $Z_r$  may be approximated by  $jX_f$  and  $jX_r$ , respectively. In this case, ( ) may be rewritten to express  $Z_m$  explicitly

in terms of  $X_f$ ,  $X_r$ , and  $\Delta \phi$ . That is, given device reactances, the desired phase shift  $\Delta \phi$  may be obtained by choosing  $Z_m$  given by

$$Z_m = \frac{X_f - X_r}{2 \tan(\Delta \phi / 2)} \pm \sqrt{\left(\frac{X_f - X_r}{2 \tan(\Delta \phi / 2)}\right)^2 - X_f X_r}$$

For a 180° bit, Z<sub>m</sub> becomes

$$Z_m = \sqrt{-X_f X_r}$$

and for a 90° bit, we should have

$$Z_m = \frac{X_f - X_r}{2} \pm \sqrt{\left(\frac{X_f - X_r}{2}\right)^2 - X_f X_r}$$

It should be remembered that the reactance X, has a negative value (it is capacitive). This design approach holds good even when the switching device is lossy.

#### Phase-shifter using a $\lambda/4$ transforming network:

The design is simple. A quarter-wave line of impedance  $Z_t$  is used to transform the hybrid impedance  $Z_0$  to a value  $Z_m$ , which the device should look into in order to provide the desired phase shift  $\Delta \phi$ . So

$$Z_t = \sqrt{Z_0 Z_m}$$

The line length  $\theta_t$  is selected so that the impedances  $Z_b$  (in two bias states) are located symmetrically with respect to the X=0 axis on the Smith chart. For a 180° phase bit,  $Z_b$  in the low-impedance state should be  $-jZ_0$  (since a 90° line length has already been added). This yields

$$-jZ_0 = Z_b = jZ_0 \frac{-Z_t^2/X_f + Z_0 \tan \theta_t}{Z_0 + (Z_t^2/X_f) \tan \theta_t}$$

which gives

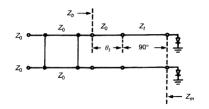
$$\theta_t|_{\Delta\phi=180^\circ} = \tan^{-1}\left(\frac{Z_t^2 - Z_0 X_f}{Z_t^2 + Z_0 X_f}\right)$$

Similarly, for a 90° bit, we should have

$$Z_b = -j2.4142Z_0$$

which yields

$$\theta_t|_{\Delta\phi=90^\circ} = \tan^{-1} \left( \frac{Z_t^2 - 2.4142X_f Z_0}{Z_0 X_f + 2.4142Z_t^2} \right)$$



## **Switched-Network Phase Shifters**

When the input signal, originally passing through network 1, is switched to pass through network 2, we get a differential phase shift  $(\phi_2-\phi_1)$ . The switched-line phase shifter is a special case of a switched-network phase shifter .

The main advantage of generalizing a switched-line phase shifter into a switched-network configuration is that one can design the variations of  $\phi_1$  and  $\phi_2$  with frequency appropriately and obtain a wider bandwidth or a desired frequency response of the phase shifter.

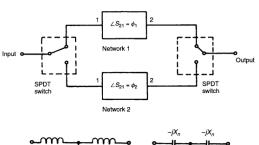
The most commonly used networks in switched-network phase shifters are the low-pass and high-pass filter configurations shown in Fig.

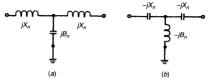
The normalized ABCD matrix of the network in Fig. a may be written as

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix}_n = \begin{bmatrix} A & B/Z_0 \\ CZ_0 & D \end{bmatrix} = \begin{bmatrix} 1 & jX_n \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ jB_n & 1 \end{bmatrix} \begin{bmatrix} 1 & jX_n \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1 - B_nX_n & j(2X_n - B_nX_n^2) \\ jB_n & 1 - B_nX_n \end{bmatrix}$$

where  $X_n$  and  $B_n$  are the reactance and susceptance shown in Fig. normalized with respect to transmission-line impedance  $Z_0$  and admittance  $Y_0$ , respectively.





[(a) Low-pass and (b) high-pass networks used in switched network phase shifters]

The transmission coefficient  $S_{21}$  is given, in terms of the normalized ABCD matrix, by

$$S_{21} = \frac{2}{A + B + C + D} = \frac{2}{2(1 - B_n X_n) + j(B_n + 2X_n - B_n X_n^2)}$$

The transmission phase  $\phi$  is given by

$$\phi = \tan^{-1} \left( -\frac{B_n + 2X_n - B_n X_n^2}{2(1 - B_n X_n)} \right)$$

When both  $B_n$  and  $X_n$  change signs (as shown in Fig. b), the phase  $\phi$  retains the same magnitude but changes sign. Amplitude of  $S_{21}$  does not change. Thus the phase shift  $\Delta \phi$  caused by switching between low-pass and high-pass networks is given by

$$\int_{X_n} \int_{|B_n|} \int_{|B$$

$$\Delta \phi = 2 \tan^{-1} \left( -\frac{B_n + 2X_n - B_n X_n^2}{2(1 - B_n X_n)} \right)$$

For the phase shifter to be matched we need

$$|S_{11}| = 0$$

Since we are considering a lossless case,

$$|S_{11}| = \sqrt{1 - |S_{21}|^2}$$

This leads to the following relationship between  $B_n$  and  $X_n$ :

$$B_n = \frac{2X_n}{X_n^2 + 1}$$

Thus the phase shift  $\Delta \phi$  can be expressed in terms of  $X_n$  alone as

$$\Delta \phi = 2 \tan^{-1} \left( \frac{2X_n}{X_n^2 - 1} \right)$$

which yields  $X_n$  in terms of  $\Delta \phi$  as

$$X_n = \tan(\frac{1}{4}\Delta\phi)$$

Substituting  $X_n$  in ( ) yields

For this case

$$B_n = \sin(\frac{1}{2}\Delta\phi)$$

The  $\pi$ -section filter shown in Fig. , may also be used in place of the T configuration .

figuration .

$$B_n = \tan\left(\frac{1}{4}\Delta\phi\right)$$

$$X_n = \sin(\frac{1}{2}\Delta\phi)$$

