

Mesh partial differential equations (PDEs): application and implementation

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Abstract

Niall: Eventually, this abstract needs to be more direct and detailed... Also, needs to highlight Algorithm 2, to emphasise the novelty

We present a novel method for generating layer-adapted meshes using a combination of h - and r -refinement, specifically using mesh PDEs (MPDEs). We show how MPDEs can be used to create Bakhvalov type meshes that are appropriate to use when solving one-dimensional singularly-perturbed problems whose solutions exhibit layers and sublayers. We demonstrate how to extend this method to generate two-dimensional, non-tensor product, layer-adapted meshes. These meshes are a suitable tool to use when solving PDEs where diffusion varies spatially. We present algorithms and Python code to implement these methods in FEniCS [1]. Numerical examples with parameter robust solutions are included to support our methods.

§1. Introduction; Keep “generic” stuff, and move that which is specific to 1D to Section 2(?). I expect, we’ll combine FEM and FEniCS.

§2. 1D

2.1 Model problem (with boundary layers)

2.2 FEM

2.3 Necessity of layer-adapted meshes (with B-mesh?)

2.4 MPDE+B-mesh

2.5 Algorithms

2.6 Examples

§3. 2D

Temp: outline of paper

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1 Introduction

^{1 2} Our goal is to present a method, to enable the reader to efficiently generate qualitatively and quantitatively accurate numerical solutions to differential equations. Often solutions to differential equations have features of interest, for example, regions where there is a rapid change in the solution. We use mesh PDEs (MPDEs) to generate non-uniform meshes that are suitably concentrated so as to resolve these features. We will focus on differential equations whose solutions exhibit layers, primarily because they are an important class of problem for which highly non-uniform meshes are used. We present algorithms to generate these meshes and Python code to implement them in FEniCS [13]. In addition, we present numerical results, demonstrating how these algorithms can be applied to both one-dimensional and two-dimensional problems. We discuss differential equations in Section 1.1, in particular, singularly perturbed differential equations. In Section 2 we present our model one-dimensional problems. We show how to obtain solutions to these problems using finite element methods (FEMs), how to implement our method in FEniCS and outline the software system in Section 2. Then, in Section 2.2, we show why layer resolving meshes are important, presenting a particularly important example, the graded *Bakhvalov mesh* [17]. MPDEs are introduced in Section 2.3 we show why they are appropriate as the basis for algorithms to drive mesh adaption. In Section 2.3 we describe how MPDEs can be used to generate one-dimensional Bakhvalov type meshes. We present two different algorithms for generating these meshes in Section 2.4. Numerical results are given in Section 2.5 for singularly-perturbed reaction-diffusion equations, including a coupled system of equations. Our Python code for solving a differential equation on a uniform mesh in FEniCS is included in Appendix A. The code for generating one-dimensional Bakhvalov meshes using an MPDE is in Appendix B. In Section 3 we describe how the

¹Niall: Need to rewrite all this later, but only after the rest of the paper is finished

²Niall: Keep an eye on when abbreviations, such as ODE, PDE and FEM are introduced and used. RH: done.

method can be extended to generate two-dimensional layer-adapted meshes. We focus on a two-dimensional reaction-diffusion equation with diffusion that varies spatially. We present an algorithm for generating these meshes in Section 3.2, and numerical results in Section 3.3. Python code for generating two-dimensional layer-adapted meshes using an MPDE is in Appendix C.³ We use the notation presented in [4][§2.1] and it's Glossary of symbols: Notation for specific vector spaces.

1.1 Differential equations

Differential equations arise in many aspects of science and engineering. They model how phenomena change over time or space, such as how a cancerous tumour grows, how a string vibrates, or how a pollutant disperses. So their solution is of fundamental importance to science and engineering. Analytical solutions to differential equations are rarely available so we want to calculate reliable approximate solutions to these problems using numerical schemes. Often there is a subsection of the solution that is of particular interest and we want to resolve these small scales effects without excessive computational cost.

Singularly perturbed problems are differential equations whose leading term is multiplied by a small positive constant, usually denoted ε . They are singularly perturbed in the sense that they are well posed for all non-zero values of ε , but ill-posed if one formally sets $\varepsilon = 0$. For a more formal definition see, e.g., [17, p.2] or [11, p.2-3]. Solutions to singularly perturbed problems typically exhibit boundary layers and may also contain interior layers. The elliptic singularly perturbed differential reaction-diffusion equations that we focus on in this paper are

$$-\varepsilon^2 \Delta u + ru = f \quad \text{in } \Omega \subseteq \mathbb{R}^d, \quad \text{with } u|_{\partial\Omega} = 0, \quad (1)$$

where $d = 1, 2$. Here r and f are given functions, $\varepsilon > 0$ and $r \geq \beta^2$ on $\bar{\Omega}$, where β is some positive constant.

1.2 Numerical methods and mesh adaptivity

⁴ To ensure the layer regions of singularly-perturbed problems are resolved when solving them numerically, the local mesh width needs to be very small, in a way that depends on ε . Often $\varepsilon \ll 1$ and consequentially, uniform meshes are unsuitable.

Since, uniform meshes are usually not appropriate, we use adapted meshes. There are three primary strategies for generating adapted meshes: h -refinement, p -refinement and r -refinement. The most extensively used are the h -refinement methods which adapt an initial mesh by adding mesh points to create a finer mesh in the layer regions. In p -refinement methods the order of the elements is increased in areas of interest without changing the number of mesh points. However, r -refinement methods, where our interest lies, relocates mesh points to control the error in the solution while fixing the mesh topology and the number of mesh points. Our objective is to compute “parameter-robust” solutions, on meshes generated using an r -refinement method. That is, solutions that have errors of similar magnitude and order of convergence independent of the value of ε .

³Niall: Suggest reading through some books, and picking a good introductory one, and then follow its notation. SIAM, if possible? RH: Is this sufficient -Section 1

⁴NM: Ignore for now, and maybe rewrite later

2 One-dimensional problems

2.1 A scalar problem, and its FEM solution

Our first boundary layer problem is the following one-dimensional singularly perturbed reaction-diffusion equation

$$-\varepsilon^2 u''(x) + ru(x) = f, \quad \text{for } x \in \Omega := (a, b), \quad (2a)$$

with boundary conditions,

$$u(a) = u(b) = 0, \quad (2b)$$

where $0 < \varepsilon \ll 1$, r and f are given functions and $r \geq \beta^2 > 0$ where β is some positive constant. In spite of its simplicity, this is a very frequently studied problem. With reasonable assumptions on f it features boundary layers at one or both boundaries. Indeed, it is one of the problems solved by the first boundary-layer adapted mesh of [3]. That paper applied a finite difference scheme, but here we use FEMs, which we now describe, along with a simple FEniCS implementation.

The bilinear form, associated with (2) is

$$a(u, v) := \varepsilon^2 (u', v') + (ru, v), \quad (3)$$

⁵ where (\cdot, \cdot) is the usual L^2 inner product. Then the weak form of (2) is: find $u \in H_0^1(\Omega)$, such that⁶

$$a(u, v) = (f, v), \quad \text{for all } v \in H_0^1(\Omega). \quad (4)$$

We obtain a finite element method by restricting the choice of u and v to a finite dimensional subspace, V_h , of $H_0^1(\Omega)$. That is, we find $u_h \in V_h$, such that

$$a(u_h, v_h) = (f, v_h), \quad \text{for all } v_h \in V_h. \quad (5)$$

Specifically, we compute the coefficients of the u_h , with respect to some basis, that solves (5).

To be more precise, we will focus on the simplest finite element method: the Galerkin method with piecewise linear basis functions (\mathcal{P}_1 -FEM). On an arbitrary mesh $a = x_0 < x_1 < \dots < x_N = b$ define the usual “hat functions” $\{\psi_i\}_{i=1}^{N-1}$ as

$$\psi_i(x) = \begin{cases} (x - x_{i-1})/(x_i - x_{i-1}) & \text{if } x_{i-1} \leq x < x_i, \\ (x_{i+1} - x)/(x_{i+1} - x_i) & \text{if } x_i \leq x < x_{i+1}, \\ 0 & \text{otherwise.} \end{cases}$$

V_h is the space spanned by these. Then, u_h can be written as

$$u_h = \sum_{i=1}^{N-1} \lambda_i \psi_i.$$

⁷ Taking

$$v_h = \psi_i \quad \text{for } i = 1, \dots, N-1,$$

in (5) yields $N-1$ equations, which can be solved to give the λ_i .

⁵RH: If I use $a(\cdot, \cdot)$ to denote the bilinear form is it okay to also use the interval (a, b) ?

⁶Niall: Check $H_0^1(\Omega)$ and other places domain is specified. RH: done.

⁷Do we need to do something similar in 2D. Niall: Suggest find a good standard reference, for both 1D and 2D

The Galerkin FEM is very powerful, and yet can be quite complicated to implement. It is powerful in the sense that it generalises to many classes of PDEs, domains of various dimensions and shapes, and incorporates many spaces of approximating functions, all within the same theoretical framework. Due to the “power of abstraction” [6] the mathematical formulation of the method is essentially the same regardless of these possible complications.

However, FEMs can be quite complicated to implement. One must discretise the domain and compute transformations of quantities onto a reference element ([7]); almost always quadrature is required to compute the associated integrals. Contributions from the elements are assembled to generate a system of linearly independent algebraic equations [18]. Finally, this system of equations is solved; careful selection of suitable solvers for larger problems usually necessary. Normally, code to visualise the solution and calculate *a posteriori* error estimates is also needed.

The code needed differs depending on the order of the basis functions, the shape and dimension of the domain and it’s elements. The power of abstraction seen in the mathematical notation does not apply to the code required to find the numerical solution, unless we use purpose built software. which automate many of the tasks needed to implement FEMs. Notable examples include FEniCS [1], deal.II [2], MFEM [14], FreeFEM [8], and Firedrake [15].

Our focus is on the use of FEniCS, a free, open-source software system, that automates many aspects of implementing FEMs while still allowing the user access to lower-level features. Using either the Python or C++ interfaces, the user specifies the domain and how to discretise it (i.e., define a mesh), and their choice of basis functions to use. Then, they define the weak form of the problem using a syntax that is very similar to the mathematical formulation. Other components of the suite then take care of assembling and solving the related system of algebraic equations.

To demonstrate a simple use of FEniCS we consider the following specific example of (2):

$$-\varepsilon^2 u''(x) + u(x) = 1 - x, \text{ for } x \in \Omega := (0, 1), \quad \text{with } u(0) = u(1) = 0, \quad (6)$$

where $0 < \varepsilon \ll 1$; its solution exhibits a layer near $x = 0$.

In Listing 1, we show a snippet of FEniCS code, the syntax is similar to the mathematical formulation of (4).

Listing 1: The weak form of (6) for FEniCS.

```
f = Expression ('1-x[0]', degree = 2)
a = epsilon*epsilon*inner(grad(u), grad(v))*dx + inner(u,v)*dx
L = f*v*dx
solve(a==L, uN, bc) # Solve, applying the boundary conditions strongly
```

In Appendix A⁸ we present the complete Python code for solving the one-dimensional singularly perturbed reaction-diffusion equation (6) in FEniCS on a uniform mesh.

2.2 The necessity of layer resolving meshes

A useful numerical solution to (6) should be both qualitatively and quantitatively representative of the exact solution. That is, it should accurately determine both the layer location and width. However, consider the numerical solutions to (6), shown in Figure 1, which were generated using a uniform mesh. One can see that the solutions change rapidly near $x = 0$, but outside this layer region it closely resembles the solution

⁸Niall: Review code and comments; include author, and file name; RH: done. time to plan how/where to publish the code

to (6) with $\varepsilon = 0$ and the boundary conditions neglected, viz, $1 - x$. In Figure 1(a), $\varepsilon = 10^{-2}$, and one can see that a mesh with 16 intervals is not adequate to resolve the layer region, but on a mesh with $N = 128$ intervals the layer is resolved, but only because N is $\mathcal{O}(\varepsilon^{-1})$. This is clear from Figure 1(b) where one can observe that, when $\varepsilon = 10^{-4}$, the solution with $N = 128$ is unstable, and fails to resolve the layer. In general, a satisfactory solution is obtained with a uniform mesh only when N is $\mathcal{O}(\varepsilon^{-1})$, which is infeasible since we wish to accurately solve this problem for any positive ε , irrespective of how small it is.⁹

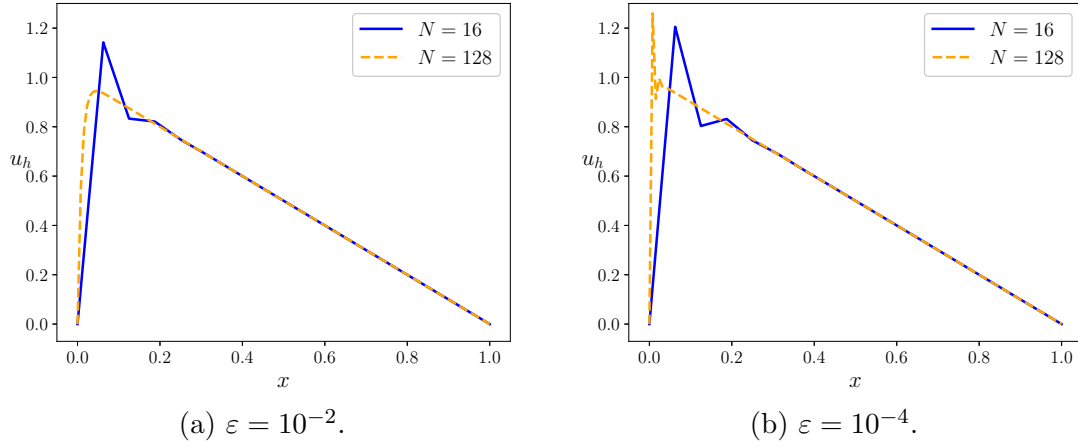


Figure 1: FEM solutions to (6) on uniform meshes.

Obviously, for a problem such as (6) we need a mesh that is very fine in the layer region. One strategy for generating such a mesh is to use asymptotic information concerning the solution (and its derivatives) and, specifically the location and width of the layer. Where this information is available, one can construct *a priori* fitted layer-resolving meshes. Among the class of such meshes that have been successfully applied to solving singularly perturbed problems, the piecewise uniform mesh of Shishkin [19] and the graded mesh of Bakhvalov [3] are the most widely studied. Both these methods yield parameter-robust solutions.¹⁰

Recall that a mesh is a partition of the domain into intervals in one-dimension, and triangles or quadrilaterals in two-dimensions. We denote a generic one-dimensional mesh, with N intervals on $[a, b]$, as $\omega : a = x_0 < x_1 < \dots < x_{N-1} < x_N = b$. One can describe a mesh by explicitly listing the coordinates of each x_i . Here, we will describe a mesh in terms of a “mesh generating function”.

Definition 2.1 (Mesh generating function). A mesh generating function is a strictly monotonic bijective function $\varphi : \bar{\Omega}^{[c]} := [0, 1] \rightarrow \bar{\Omega} := [a, b]$ that maps a uniform mesh $\xi_i = i/N$, for $i = 0, 1, \dots, N$, to a potentially non-uniform mesh $x_i = \varphi(i/N)$, for $i = 0, 1, \dots, N$, with $\varphi(0) = a$ and $\varphi(1) = b$.

For example, if $\varphi(\xi) = \xi$, then the resulting mesh will be a uniform mesh on the interval $\bar{\Omega}$ as shown in Figure 2. If $\varphi(\xi) = \xi^4$, then the resulting mesh is graded on the interval $\bar{\Omega}$, with mesh points closest together when ξ is close to 0, as shown in Figure 3.

⁹Niall: Humour me, and add, just for comparison, a version of Figure 1(a) but with $\varepsilon = 10^{-4}$. RH: Done, I think I prefer it.

¹⁰Niall: Have we explained what parameter-robust means? RH: Yes, see last line of Section 1.2

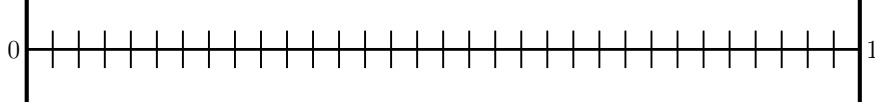


Figure 2: Mesh generated when $\varphi(\xi) = \xi$.

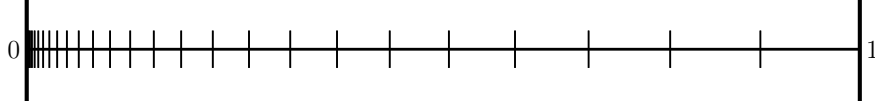


Figure 3: Mesh generated when $\varphi(\xi) = \xi^4$.

The first layer adapted mesh for singularly perturbed problems was proposed in [3]. It is graded, and very fine, in layer regions, and uniform elsewhere. We will now outline its construction based on the presentation in [11]. To motivate the mesh we will consider the reaction-diffusion problem (6), whose solution has a single boundary layer which is located near $x = 0$. Let $\exp(-\beta x/\varepsilon)$ be a function that represents this boundary layer. Then the mesh points x_i near $x = 0$ are chosen to satisfy

$$\exp\left(\frac{-\beta x_i}{\sigma\varepsilon}\right) = 1 - \frac{i}{qN} \quad \text{for } i = 0, 1, \dots, \quad (7)$$

where $q \in (0, 1)$ is roughly the portion of mesh points used to resolve the layer, $\sigma > 0$ is a constant that depends on the underlying method, typically chosen to be the formal order of the underlying method, and N is the number of mesh intervals.

The mesh is expressed in terms of the mesh generating function,

$$\varphi(t) = \begin{cases} \chi(t) := -\frac{\sigma\varepsilon}{\beta} \ln\left(1 - \frac{t}{q}\right), & \text{for } t \in [0, \tau], \\ \pi(t) := \chi(\tau) + (t - \tau)\varphi'(\tau), & \text{elsewhere.} \end{cases} \quad (8)$$

The mesh generated has mesh points $x_i = \varphi(i/N)$ for $i = 0, 1, \dots, N$.

We measure the error in our numerical solution in the energy norm,

$$\|u\|_{E,\Omega} = \sqrt{a(u, u)},$$

with

$$e_h(x) = |u_{2h}(x) - u_h(x)|,$$

where u_h is the \mathcal{P}_1 -FEM solution and u_{2h} is the \mathcal{P}_2 -FEM solution on the same mesh, that is, the Galerkin method with quadratic basis functions. A proof that the error in the \mathcal{P}_1 -FEM solutions of singularly-perturbed problems, calculated on Bakhvalov meshes, measured in the energy norm converges linearly is given in [16].

¹¹RH: Does this image add anything?

¹²Niall: Need to think if error estimates are appropriate here? What is a good source? RH: Still not quite right.!!!

2.3 MPDEs

We are interested in developing parameter-robust methods, where the quality of the solution is independent of the value of the perturbation parameter, ε . One way to achieve this is to use layer adaptive methods which generate meshes that concentrate mesh points in regions where large variations in the solution occur. There are various ways one can express a formula for such a non-uniform mesh. In this article we define meshes in terms of a “*mesh generating function*”.

A mesh PDE, is presented in [9], as a way of performing r -refinement. First, a PDE whose solution is a mesh generating function is posed. The PDE features a coefficient that controls the concentration of points in the resulting mesh. As we will see it may be defined explicitly in a way that generates *a priori* layer adaptive meshes, or in terms of error estimates.

If one wants to find the solution to (2) on a fixed number of mesh points, how should these points be selected? One approach is to choose a *mesh density function*, $\rho : \Omega \rightarrow \mathbb{R}_{>0}$ and then construct a mesh so that the integral of ρ is the same on each mesh interval. This is known as the “*equidistribution principle*” and was introduced by [5]. A mesh $\omega : a = x_0 < x_1 < \dots < x_N = b$, with N intervals, equidistributes ρ when

$$\int_{x_i}^{x_{i+1}} \rho(x) dx = \frac{1}{N} \int_a^b \rho(x) dx, \quad \text{for all } i = 0, 1, \dots, N. \quad (9a)$$

or, equivalently,

$$\int_a^{x_i} \rho(x) dx = \frac{i}{N} \int_a^b \rho(x) dx \quad \text{for } i = 0, 1, \dots, N, \quad (9b)$$

If we choose $\rho = C$, a constant, then $x(\xi) = \xi$ is a solution to (9a) and the mesh generated is uniform, see Figure 2. More typically, ρ is not constant and the resulting mesh is finer where ρ is large. For example, if $\rho = 1/\xi^3$, then $x(\xi) = \xi^4$ is a solution to (9a) and the mesh generated is as shown in Figure 3.

Since the equidistribution problem (9b) is nonlinear, typically one must solve it using an iterative process until the resulting mesh equidistributes ρ . One way to think of layer-adapted meshes is that if the numerical solution, u_h , is calculated on the adapted mesh, then the mesh will be finer in the layer region than elsewhere on the domain. We solve the MPDE on a computational domain $\Omega^{[c]}$ using a uniform mesh and the PDE on a physical domain Ω using the layer-adapted mesh generated from the solution to the MDPE. If $u_h(x)$ is transformed from Ω onto $\Omega^{[c]}$, the layer region will be stretched and no sharp layer exists in $\bar{u}_h(x(\xi))$, as can be seen in Figure 4.

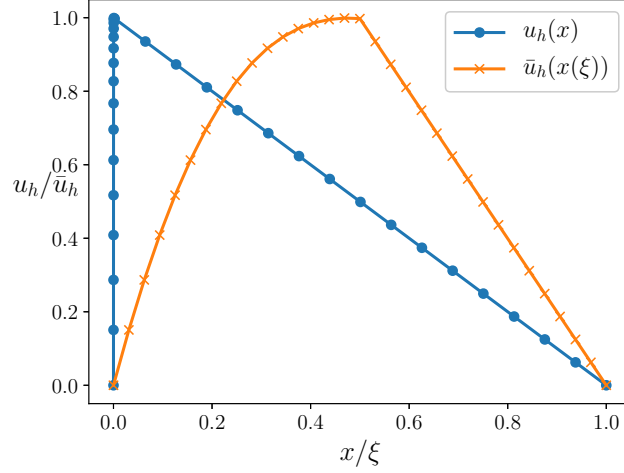


Figure 4: Solution to (6) with $\varepsilon = 10^{-4}$ and $N = 32$ generated on an adapted mesh $\{x\}$, and the solution transformed onto a uniform mesh $\{\xi\}$.

We can derive an MPDE by considering the equidistribution principle as a mapping $x(\xi) : \bar{\Omega}^{[c]} := [0, 1] \rightarrow \bar{\Omega} := [a, b]$ from the computational coordinate ξ to a physical coordinate x , which satisfies

$$\int_a^{x(\xi)} \rho(x) dx = \xi \int_a^b \rho(x) dx. \quad (10)$$

First, we differentiate (10) with respect to ξ . We can see that $x(\xi)$ satisfies

$$\rho(x) \frac{dx}{d\xi} = \int_a^b \rho(x) dx, \quad (11)$$

where the right-hand side is independent of ξ . Differentiating again with respect to ξ gives the MPDE,

$$(\rho(x)x'(\xi))' = 0, \quad \text{for } \xi \in \Omega^{[c]}, \quad (12a)$$

whose solution is a mesh generating function, see Definition 2.1. We impose the boundary conditions,

$$x(0) = a \quad \text{and} \quad x(1) = b, \quad (12b)$$

where a and b are the end points of the resulting mesh. The weak form of (12) is: find, $x(\xi) \in H^1(\Omega)$, such that

$$\int_0^1 \rho(x)x'(\xi)v'd\xi = 0, \quad \text{for all } v \in H_0^1(\Omega), \quad \text{with } x(0) = a \quad \text{and} \quad x(1) = b. \quad (13)$$

Since (13) is a nonlinear problem, we use a fixed point iterative method to find a solution. At each iteration, k , we solve

$$\sum_{i=1}^N \int_{\xi_{i-1}}^{\xi_i} \rho(x_h^{[k-1]})(x_h^{[k]})' \varphi_i' d\xi = 0 \quad (14)$$

The process is repeated until the mesh generated $x_h^{[k]}$ equidistributes $\rho(x_h^{[k]})$; that is, ρ calculated on that mesh.

The definition of a Bakhvalov mesh given in Section 2.2 is essentially that as originally proposed in [3]. Alternatively, this Bakhvalov mesh can be generated by equidistributing the mesh density function

$$\rho(x) = \max \left\{ 1, K \frac{\beta}{\varepsilon} \exp \left(-\frac{\beta x}{\sigma \varepsilon} \right) \right\}, \quad (15)$$

where K is a positive constant that determines the proportion of mesh points that resolve the layer [11]. When,

$$K \frac{\beta}{\varepsilon} \exp\left(-\frac{\beta\varphi(\tau)}{\sigma\varepsilon}\right) = 1, \quad (16)$$

the transition point, $\varphi(\tau)$, of the mesh generated is the same as in (8). This motivates our choice of K . We insert the calculation for $\chi(\tau)$ in (8) into (16) to get

$$K = \frac{\varepsilon q}{\beta(q - \tau)}.$$

One observes in practice that,

$$\tau \approx q - \frac{\sigma\varepsilon(1 - q)}{\beta},$$

so we set

$$K = \frac{q}{\sigma(1 - q)}.$$

We can use (16) to generate a Bakhvalov type mesh using an MPDE. When ρ is as defined in (15), a Bakhvalov mesh is a solution to (12), making it a good choice of MPDE for this problem. In Section 2.4 we include algorithms to generate a Bakhvalov type mesh using an MPDE.

Another method to generate a mesh which is similar to a Bakhvalov mesh, and which has a computation advantage over (15), is to equidistribute the mesh density function¹³

$$\rho(x) = 1 + K \frac{\beta}{\varepsilon} \exp\left(-\frac{\beta x}{\sigma\varepsilon}\right). \quad (17)$$

In the layer region the mesh is very similar to a Bakhvalov mesh since the addition of 1 to the value of ρ is insignificant. Outside the layer region the value of the second term is small so the resulting mesh is almost uniform.

2.4 Algorithms

In Algorithm 1 and Algorithm 2 we show how to generate a Bakhvalov type mesh using an MPDE. This is a suitable mesh for the singularly-perturbed problem (6) whose solution has a single boundary layer, near $x = 0$, when the perturbation parameter $\varepsilon \ll 1$. Let $\exp(-\beta s/\varepsilon)$ represent the boundary layer. In Algorithm 3 we show how to generate a two-dimensional layer-adapted mesh using a combination of an MPDE and h -refinement. This mesh is suitable for the two-dimensional singularly-perturbed problem (25) which has layers near $x = 1$ and $y = 1$ which vary in width spatially.

In this section the mesh is defined by the mesh density function,

$$\rho(s) := \max\left\{1, K \frac{\beta}{\varepsilon} \exp\left(-\frac{\beta s}{\sigma\varepsilon}\right)\right\}, \quad \text{for } s \in \Omega := (0, 1). \quad (18)$$

To ensure the mesh generated approximately equidistributes ρ , we calculate the residual of (13) as

$$\text{res}^{[k]} := \sum_{i=1}^N \int_{\xi_{i-1}}^{\xi_i} \rho(x_h^{[k]})(x_h^{[k]})' \varphi_i' d\xi, \quad (19)$$

at each iteration k . We stop iterating when $\|\text{res}^{[k]}\|_{0,\Omega} \leq TOL$, an appropriate tolerance level.

¹³RH: I didn't find a reference for this yet.

Niall: Check margins on algorithm boxes. RH: done, I needed to put them in minipages

Input: N , the number of intervals in the mesh.
Input: ρ , a mesh density function.
Input: TOL .

- 1 Set $\omega^{[c]} := \{\xi_0, \xi_1, \dots, \xi_N\}$ to be a uniform mesh, with N intervals, discretising $\overline{\Omega}^{[c]}$;
- 2 Set $x(\xi) = \xi$ for $\xi \in \overline{\Omega}^{[c]}$;
- 3 $s \leftarrow x$;
- 4 $k \leftarrow 0$;
- 5 **do**
- 6 set x to be the \mathcal{P}_1 -FEM solution, on $\omega^{[c]}$, to

$$(\rho(s)x'(\xi))' = 0, \quad \text{for } \xi \in \Omega^{[c]}, \quad \text{with } x(0) = 0, \quad x(1) = 1, \quad (20)$$
- $s \leftarrow x$;
- 7 calculate $\|\text{res}^{[k]}\|_{0,\Omega^a}$;
- 8 $k \leftarrow k + 1$;
- 9 **while** $\|\text{res}^{[k]}\|_{0,\Omega} > TOL$;
- 10 Set $\omega^{[p]} := x(\omega^{[c]})$ to be the adapted mesh on $\overline{\Omega}$;

Algorithm 1: Generate a layer resolving mesh using an MPDE

^aNiall: use $\|\text{res}^{[j]}\|_{0,\Omega}$ instead of r . RH: done.

Table 1: Number of iterations to create a Bakhvalov type mesh using Algorithm 1.

$\varepsilon \backslash N$	32	64	128	256	512	1024
1	1	1	1	1	1	1
10^{-1}	4	3	2	1	1	1
10^{-2}	15	14	15	15	14	13
10^{-3}	21	37	56	68	68	70
10^{-4}	21	37	69	132	256	343
10^{-5}	21	37	69	133	260	517
10^{-6}	21	37	69	133	261	517
10^{-7}	21	37	69	133	261	517
10^{-8}	21	37	69	133	261	517

In practice we see that, when using Algorithm 1 to create a Bakhvalov type mesh, the number of iterations required to generate the mesh is $\mathcal{O}(qN)$, where q is the portion of mesh points used to resolve the layer region and N is the number of mesh intervals, see Table 1. This occurs as only one mesh point is added to the layer region at each iteration.

We amend the algorithm by alternating between r - and h -refinement. Initially, we use a uniform mesh with 4 intervals, we refine the mesh three times using r -refinement then double the number of mesh intervals using a uniform h -refinement. In Algorithm 2 we show how to generate a Bakhvalov type mesh using a combination of r - and h -refinement. The number of iterations required on the final mesh are shown in Table 2.

Input: N , the number of intervals in the final mesh.

Input: ρ , a mesh density function.

Input: TOL .

```

1 Set  $\omega^{[c;0]} := \{\xi_0, \xi_1, \dots, \xi_4\}$  to be the uniform mesh on  $\overline{\Omega}^{[c]}$  with 4 intervals;
2 Set  $x(\xi) = \xi$  for  $\xi \in \overline{\Omega}^{[c]}$ ;
3  $s \leftarrow x$ ;
4  $i, k \leftarrow 0$ ;
5 while  $2^{i+2} < N$  do
6   for  $j$  in 1:3 do
7     set  $x$  to be the  $\mathcal{P}_1$ -FEM solution, on  $\omega^{[c;i]}$ , to
      
$$(\rho(s)x'(\xi))' = 0, \quad \text{for } \xi \in \Omega^{[c]}, \quad x(0) = 0, \quad x(1) = 1; \quad (21)$$

       $s \leftarrow x$ ;
8   end
9    $w^{[c;i+1]} \leftarrow$  uniform  $h$ -refinement of  $w^{[c;i]}$ ;
10   $s \leftarrow s$  interpolated onto  $w^{[c;i+1]}$ ;
11   $i \leftarrow i + 1$ ;
12 end
13 do
14   set  $x$  to be the  $\mathcal{P}_1$ -FEM solution, on  $\omega^{[c;i]}$ , to
      
$$(\rho(s)x'(\xi))' = 0, \quad \text{for } \xi \in \Omega^{[c]}, \quad x(0) = 0, \quad x(1) = 1; \quad (22)$$

       $s \leftarrow x$ ;
15   calculate  $\|\text{res}^{[k]}\|_{0,\Omega}$ ;
16    $k \leftarrow k + 1$ ;
17 while  $\|\text{res}^{[k]}\|_{0,\Omega} > TOL$ ;
18 Set  $\omega^{[p]} := x(\omega^{[c;i]})$  to be the adapted mesh on  $\overline{\Omega}$ ;
```

Algorithm 2: Generate a Bakhvalov type mesh using an MPDE, and h -refinement.

An implementation of this algorithm in FEniCS is in Appendix B.

Table 2: Number of iterations to create a Bakhvalov type mesh using Algorithm 2 on final computational domain.

$\varepsilon \backslash N$	32	64	128	256	512	1024
1	1	1	1	1	1	1
10^{-1}	1	1	1	1	1	1
10^{-2}	1	1	1	1	1	1
10^{-3}	1	1	1	1	1	1
10^{-4}	2	1	1	1	1	1
10^{-5}	3	2	2	1	1	1
10^{-6}	3	3	2	2	1	1
10^{-7}	3	3	3	3	2	1
10^{-8}	3	3	3	3	2	2

From Table 2, we can see that using Algorithm 2 reduces the number of iterations required to generate a Bakhvalov type mesh, including the refinements on the coarser

meshes, to $\mathcal{O}(\log(N))$.

2.5 Numerical results

In this section we present numerical results for the one-dimensional reaction-diffusion equation (6) and a one-dimensional coupled set of reaction-diffusion equations (23).

2.5.1 A scalar reaction-diffusion equation

We present results for (6) solved on meshes generated using Algorithm 2. When generating the meshes we choose $\sigma = 2.5$, $\beta = 0.99$ and $K = 0.4$ as the values of the variables in our mesh generating function, (18). We set the tolerance level for the L_2 -norm of the residual (19) as 0.4.

We can see that the errors in the solutions to (6), measured in the energy norm achieved on the Bakhvalov type mesh, generated using an MPDE, shown in Table 3, converge quadratically, and the bound on the errors is,

$$\|e_h\|_{E,\Omega} \leq C\varepsilon^{1/2}N^{-1},$$

where C is a constant independent of ε and N . These errors are comparable to those generated on a Bakhvalov mesh, using (8), which are shown in Table 4.

Table 3: Errors measured in the energy norm for (6) using Algorithm 2

$\varepsilon \backslash N$	32	64	128	256	512	1024
1	1.10e-05	2.76e-06	6.90e-07	1.73e-07	4.31e-08	1.08e-08
10^{-1}	3.46e-04	8.65e-05	2.16e-05	5.41e-06	1.35e-06	3.38e-07
10^{-2}	2.33e-04	5.81e-05	1.46e-05	3.64e-06	9.09e-07	2.27e-07
10^{-3}	1.21e-04	2.68e-05	5.22e-06	1.29e-06	3.23e-07	8.09e-08
10^{-4}	3.50e-05	8.16e-06	1.91e-06	4.50e-07	1.06e-07	2.59e-08
10^{-5}	1.03e-05	2.54e-06	5.99e-07	1.43e-07	3.42e-08	8.32e-09
10^{-6}	3.27e-06	7.63e-07	1.87e-07	4.46e-08	1.09e-08	2.64e-09
10^{-7}	1.03e-06	2.41e-07	5.77e-08	1.40e-08	3.45e-09	8.51e-10
10^{-8}	3.26e-07	7.62e-08	1.82e-08	4.41e-09	1.09e-09	2.68e-10

Table 4: Errors measured in the energy norm for (6) on a Bakhvalov mesh.

$\varepsilon \backslash N$	32	64	128	256	512	1024
1	1.10e-05	2.76e-06	6.90e-07	1.73e-07	4.31e-08	1.08e-08
10^{-1}	5.73e-04	1.44e-04	3.59e-05	8.97e-06	2.24e-06	5.61e-07
10^{-2}	2.62e-04	6.55e-05	1.64e-05	4.10e-06	1.02e-06	2.56e-07
10^{-3}	9.07e-05	2.16e-05	5.29e-06	1.32e-06	3.29e-07	8.22e-08
10^{-4}	3.12e-05	7.13e-06	1.71e-06	4.21e-07	1.04e-07	2.60e-08
10^{-5}	1.06e-05	2.36e-06	5.54e-07	1.35e-07	3.32e-08	8.26e-09
10^{-6}	3.58e-06	7.76e-07	1.79e-07	4.30e-08	1.06e-08	2.62e-09
10^{-7}	1.20e-06	2.55e-07	5.79e-08	1.38e-08	3.36e-09	8.30e-10
10^{-8}	4.00e-07	8.35e-08	1.87e-08	4.40e-09	1.07e-09	2.63e-10

2.5.2 Coupled system of reaction-diffusion equations

Our second example shows that adapted meshes generated using MPDEs are appropriate for solving problems with solutions that contain sublayers. The chosen problem is a coupled system of two singularly-perturbed reaction-diffusion equations with variable coefficients taken from [12],

$$-\begin{pmatrix} \varepsilon_1 & 0 \\ 0 & \varepsilon_2 \end{pmatrix}^2 \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}'' + \begin{pmatrix} 2(x+1)^2 & -(1+x^3) \\ -2\cos(\pi x/4) & (1+\sqrt{2})\exp(1-x) \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \begin{pmatrix} 2x^2 \\ 10x+1 \end{pmatrix}, \quad \text{for } x \in \Omega := (0, 1), \quad (23a)$$

with boundary conditions,

$$u_1(0) = u_1(1) = 0, \quad u_2(0) = u_2(1) = 0. \quad (23b)$$

While both solutions to (23) exhibit layers near $x = 0$ and $x = 1$, only the solution to u_1 has sublayers that are dependant on both ε_1 and ε_2 as can be seen in Figure 5.

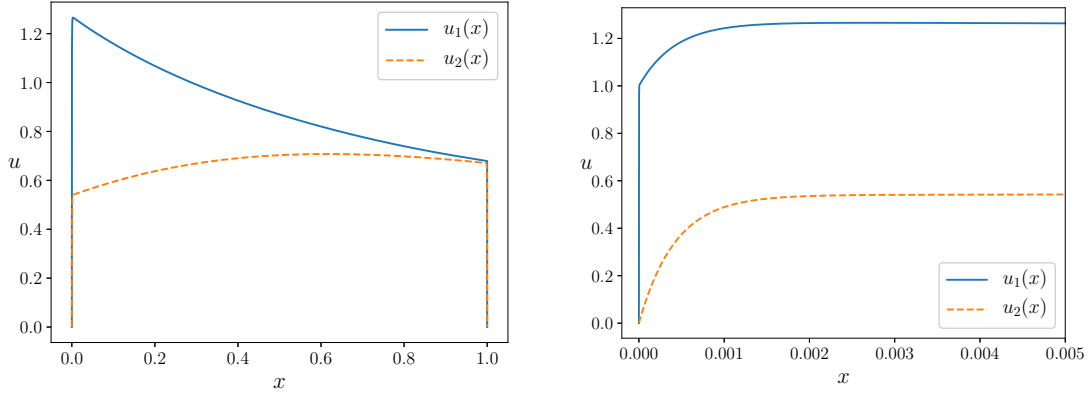


Figure 5: Solutions to (23) with $\varepsilon_1 = 10^{-6}$ and $\varepsilon_2 = 10^{-3}$.

To solve (23) we generate meshes that are graded near $x = 0$ and $x = 1$, in a way that is dependant on both ε_1 and ε_2 . To achieve this we solve MPDE (12), using Algorithm 2, with

$$\rho(x) = 1 + K \frac{\beta}{\varepsilon_1} \exp\left(-\frac{\beta x}{\sigma \varepsilon_1}\right) + K \frac{\beta}{\varepsilon_2} \exp\left(-\frac{\beta x}{\sigma \varepsilon_2}\right) + K \frac{\beta}{\varepsilon_1} \exp\left(-\frac{\beta(1-x)}{\sigma \varepsilon_1}\right) + K \frac{\beta}{\varepsilon_2} \exp\left(-\frac{\beta(1-x)}{\sigma \varepsilon_2}\right), \quad (24)$$

with $K = 0.1$, $\beta = 1.99$ and $\sigma = 2.5$. The errors measured in the energy norm $\|E_h\|_{E,\Omega}$ are shown in Table 5, the errors converge linearly, and the bound on the error is

$$\|e_h\|_E \leq C \left(\varepsilon_1^{\frac{1}{2}} + \varepsilon_2^{\frac{1}{2}} \right) N^{-1},$$

where C is a constant independent of ε_1 , ε_2 and N .

Table 5: Errors in the energy norm for (23) with $\varepsilon_1 = 10^{-6}$.

$\varepsilon_2 \backslash N$	32	64	128	256	512	1024
1	5.29e-02	2.53e-02	1.27e-02	6.34e-03	3.17e-03	1.58e-03
10^{-2}	3.46e-02	1.78e-02	8.96e-03	4.49e-03	2.25e-03	1.12e-03
10^{-4}	2.80e-03	1.45e-03	7.54e-04	3.80e-04	1.91e-04	9.58e-05
10^{-6}	2.95e-04	1.20e-04	5.59e-05	2.76e-05	1.37e-05	6.83e-06
10^{-8}	3.06e-04	1.31e-04	6.29e-05	3.10e-05	1.54e-05	7.71e-06
10^{-10}	3.29e-04	1.43e-04	6.81e-05	3.37e-05	1.67e-05	8.39e-06
10^{-12}	3.30e-04	1.43e-04	6.84e-05	3.39e-05	1.70e-05	8.48e-06

3 Two-dimensional problem

¹⁴ As described in [10][§4.1], a two-dimensional layer-adapted mesh is generated as a uniform mesh in a metric space in $\Omega \subset \mathbb{R}^2$. A matrix function $M = M(\mathbf{x})$, the mesh density function, is defined on Ω and the layer-adapted mesh is known as an *M-uniform mesh*. An advantage of a layer-adaptive mesh for singularly-perturbed problems is that the error bound on the layer adapted mesh is dependent on a smaller power of ε and the bound is smaller than if the solution had been generated on a uniform mesh as proven in [10][§2.4]. ¹⁵¹⁶

3.1 Layer-adapted mesh

We used the knowledge gained from generating one-dimensional Bakhvalov meshes using MPDEs to develop algorithms to generate two-dimensional layer-adapted meshes for reaction-diffusion equations using our MPDE approach. We studied problems where diffusion varies on the domain, so the boundary layers vary in width, spatially. Consequently, even when posed on a regular domain (such as a unit square) optimal meshes are not tensor product in nature. Whereas standard Bakhvalov meshes are tensor product, our MPDE approach has no such restriction.

Our chosen two-dimensional singularly-perturbed reaction-diffusion equation, with varying diffusion, is

$$-\nabla \cdot \left(\begin{pmatrix} \varepsilon(1+y)^2 & 0 \\ 0 & \varepsilon(2-x)^2 \end{pmatrix} \nabla u(x, y) \right) + u(x, y) = (e^x - 1)(e^y - 1),$$

for $(x, y) \in \Omega = (0, 1) \times (0, 1)$, (25a)

with boundary conditions

$$u = 0 \quad \text{on } \partial\Omega. \quad (25b)$$

The solution to (25) exhibits layers near $x = 1$ and $y = 1$, and a corner layer near $(x, y) = (1, 1)$, as can be seen in Figure 6.

¹⁴Niall: Section title?

¹⁵RH: Is this sufficient to describe a 2d mesh? NM: Maybe – though a more precise reference is needed. Which section? RH: done.

¹⁶Niall: Also, compare with [10, §2.4.5]? RH: Done.

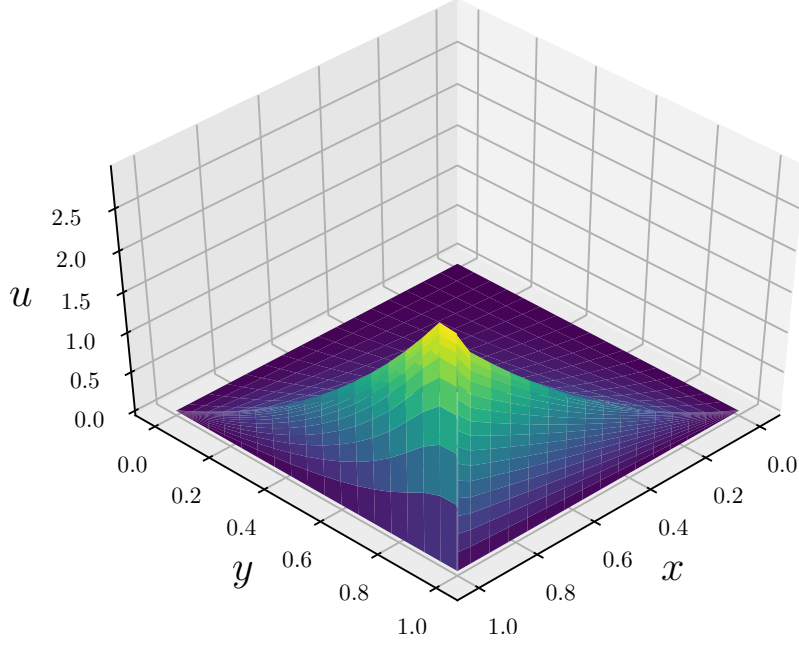


Figure 6: Solution to (25) with $\varepsilon = 10^{-4}$ and $N = 32$.

To generate a layer adapted mesh, that is suitable for (25), using an MPDE we solve the two-dimensional Poisson equation for $\mathbf{x}(\xi_1, \xi_2) = (x, y)$, on the domain $\Omega^{[c]} = (0, 1) \times (0, 1)$,

$$-\nabla \cdot (M(\mathbf{x}(\xi_1, \xi_2)) \nabla \mathbf{x}(\xi_1, \xi_2)) = \mathbf{f}(\xi_1, \xi_2), \quad \text{for } (\xi_1, \xi_2) \in \Omega^{[c]}, \quad (26a)$$

with $\mathbf{f} = (0, 0)^T$, and boundary conditions,

$$\begin{cases} \mathbf{x}(0, \xi_2) &= \left(0, \frac{\partial \mathbf{x}}{\partial \xi_2} = 0\right), \\ \mathbf{x}(1, \xi_2) &= \left(1, \frac{\partial \mathbf{x}}{\partial \xi_2} = 0\right), \\ \mathbf{x}(\xi_1, 0) &= \left(\frac{\partial \mathbf{x}}{\partial \xi_1} = 0, 0\right), \\ \mathbf{x}(\xi_1, 1) &= \left(\frac{\partial \mathbf{x}}{\partial \xi_1} = 0, 1\right). \end{cases} \quad (26b)$$

Here

$$M(\mathbf{x}) = \begin{pmatrix} \max \left\{ 1, K_1 \frac{\beta}{\varepsilon_1} \exp \left(-\frac{\beta(1-x)}{\sigma \varepsilon_1} \right) \right\} & 0 \\ 0 & \max \left\{ 1, K_2 \frac{\beta}{\varepsilon_2} \exp \left(-\frac{\beta(1-y)}{\sigma \varepsilon_2} \right) \right\} \end{pmatrix}, \quad (27)$$

with $\varepsilon_1(x, y) = \varepsilon(1+y)^2$, $\varepsilon_2(x, y) = \varepsilon(2+x)^2$, $\beta = 0.99$, $\sigma = 2.5$ and $K_1 = K_2 = 0.4$. An example of the resulting non tensor-product mesh is shown in Figure 7.

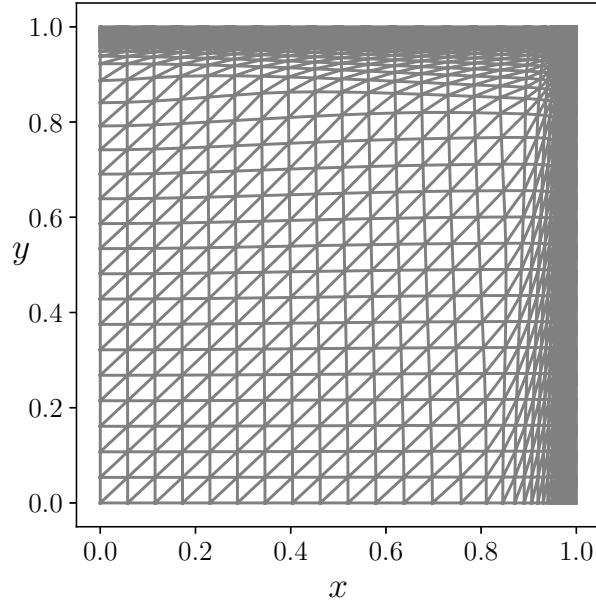


Figure 7: A 32×32 mesh generated using (26) for (25) with $\varepsilon = 10^{-2}$.

3.2 Algorithm

Input: N , where $N + 1$ is the number of mesh points in both directions on $\Omega^{[c]}$.

Input: M , a mesh density function.

- 1 Set $\omega^{[c;0]} := \{(\xi_1, \xi_2)_0, (\xi_1, \xi_2)_1, \dots, (\xi_1, \xi_2)_4\}$ to be the uniform mesh on $\overline{\Omega}^{[c]}$ with 5 mesh points in each direction on $\overline{\Omega}^{[c]}$;
- 2 Set $\mathbf{x}(\xi_1, \xi_2) = (\xi_1, \xi_2)$ for $(\xi_1, \xi_2) \in \overline{\Omega}^{[c]}$;
- 3 $\mathbf{s} \leftarrow \mathbf{x}$;
- 4 $i \leftarrow 0$;
- 5 **while** $2^{i+2} < N$ **do**
- 6 **for** j in $1 : 3$ **do**
- 7 set \mathbf{x} to be the \mathcal{P}_1 -FEM solution, on $\omega^{[c;i]}$, to

$$\nabla \cdot (M(\mathbf{s}) \nabla \mathbf{x}(\xi_1, \xi_2)) = (0, 0)^T, \quad \text{for } (\xi_1, \xi_2) \in \Omega^{[c]}, \quad (28)$$
 with boundary conditions as defined in (26b);
- 8 $\mathbf{s} \leftarrow \mathbf{x}$;
- 9 **end**
- 10 $w^{[c;i+1]} \leftarrow \text{refine } w^{[c;i]}$;
- 11 $\mathbf{s} \leftarrow \mathbf{s}$ interpolated onto $w^{[c;i+1]}$;
- 12 $i \leftarrow i + 1$;
- 13 **end**
- 14 **for** k in $1 : 4$ **do**
- 15 set \mathbf{x} to be the \mathcal{P}_1 -FEM solution, on $\omega^{[c;i]}$, to

$$\nabla \cdot (M(\mathbf{s}) \nabla \mathbf{x}(\xi_1, \xi_2)) = (0, 0)^T, \quad \text{for } (\xi_1, \xi_2) \in \Omega^{[c]}, \quad (29)$$
 with boundary conditions as defined in (26b);
- 16 $\mathbf{s} \leftarrow \mathbf{x}$;
- 17 **end**
- 18 Set $\omega^{[p]} := \mathbf{x}(\omega^{[c;i]})$ to be the adapted mesh on $\overline{\Omega}$;

Algorithm 3: Generate a two-dimensional layer-adapted mesh using an MPDE, and h -refinement.

3.3 Numerical results

To solve (25) we generated meshes that are graded near $x = 1$ and $y = 1$, in a way that depends on the varying diffusion coefficient. To achieve this we solve (26) using Algorithm 3 with $M(\mathbf{x})$ as shown in (27).

We can see that the errors in the solutions to (25), measured in the energy norm achieved on the layer-adapted mesh, generated using an MPDE, shown in Table 6, converge linearly, and the bound on the errors is,

$$\|e_h\|_{E,\Omega} \leq C\varepsilon^{\frac{1}{2}}N^{-1},$$

where C is a constant independent of ε and N .

Table 6: Errors measured in the energy norm for (25) on a layer-adapted mesh generated using an MPDE.

ε/N	16	32	64	128	256	512
1	2.809e-03	1.457e-03	7.404e-04	3.725e-04	1.866e-04	9.338e-05
10^{-1}	1.663e-02	8.390e-03	4.210e-03	2.108e-03	1.054e-03	5.272e-04
10^{-2}	1.689e-02	8.332e-03	4.153e-03	2.075e-03	1.037e-03	5.185e-04
10^{-3}	8.221e-03	3.790e-03	1.867e-03	9.285e-04	4.640e-04	2.320e-04
10^{-4}	4.728e-03	1.602e-03	6.783e-04	3.221e-04	1.588e-04	7.917e-05
10^{-5}	4.180e-03	1.105e-03	3.301e-04	1.200e-04	5.304e-05	2.562e-05
10^{-6}	4.134e-03	1.044e-03	2.707e-04	7.340e-05	2.295e-05	9.010e-06

4 Conclusion

We have demonstrated how MPDEs can be used to generate *a priori* layer-adapted meshes suitable for one- and two-dimensional singularly-perturbed problems whose solution exhibit layers. These layers include sublayers and layers that vary in width spatially. We combined *h*- and *r*-refinement methods to improve the efficiency of the methods. Solutions calculated on these meshes are parameter robust. We've presented algorithms and Python code to enable the reader to generate these meshes in FEniCS.

Appendices

A Python Code to solve (6) on a uniform mesh

17

Listing 2: Code to solve (6) on a uniform mesh.

```

432  ## 1DRD_FEM_uniform_mesh.py
433  """
434
435  Solve the singularly perturbed problem
436  -epsilon^2 u'' + u = f = 1-x, for x in (0,1)
437  with u(0) = u(1) = 0
438  on a uniform mesh
439
440  This code is part of
441  Mesh partial differential equations (PDEs): application and implementation,
442  Wwritten by Roisin Hill and Niall Madden
443  Contact: Roisin Hill
444  For: <cite paper>
445  """
446
447  from fenics import *
448
449  # Define problem parameters
450  epsilon = 1E-2          # perturbation factor
451  N = 128                 # mesh intervals
452
453  # Create mesh and function space with linear element on that mesh
454  mesh = UnitIntervalMesh(N)
455  V = FunctionSpace(mesh, 'P', 1)
456
457  u = TrialFunction(V) # Trial functions on V

```

¹⁷RH: Should I include email addresses in the code?

```

458 v = TestFunction(V) # Test functions on V
459 uN = Function(V)    # Numerical solution in V
460
461 # Define the Dirichlet boundary conditions
462 g = Expression('x[0]*(1-x[0])', degree = 2)
463 bc = DirichletBC(V, g, 'on_boundary')
464
465 # Define right hand side of the SPP, and the weak form
466 f = Expression('1-x[0]', degree = 2)
467 a = epsilon*epsilon*inner(grad(u), grad(v))*dx + inner(u,v)*dx
468 L = f*v*dx
469 solve(a==L, uN, bc) # Solve, applying the boundary conditions strongly
470

```

B Python Code to generate a Bakhvalov type mesh for an SPP whose solution has one layer region near $x = 0$ using a combination of r - and h -refinement.

Listing 3: Code to generate a Bakhvalov mesh using a combination of r - and h -refinement.

```

475 ## 1D_layer_adapted_mesh_rh_refinement.py
476 """
477
478 Generate 1 sided Bakhvalov type mesh, for a singularly
479 perturbed ODE whose solution has a layer near x = 0,
480 by solving the MPDE -(rho(x) x'(xi))' = 0 for xi in (0,1) with
481 x(0) = 0 and x(1) = 1 and using uniform h-refinements.
482
483 Gauss-Lobatto quadrature rule is used:
484 i.e. replace "return GaussJacobiQuadratureLineRule(ref_el, m)"
485 with "return GaussLobattoLegendreQuadratureLineRule(ref_el, m)"
486 in the file FIAT/quadrature.py
487
488 This code is part of
489 Mesh partial differential equations (PDEs): application and implementation,
490 written by Roisin Hill and Niall Madden
491 Contact: Roisin Hill
492 For: <cite paper>
493 """
494
495 from fenics import *
496 import math
497
498 # Set degree for the Gauss-Lobatto quadrature rule
499 parameters["form_compiler"]["quadrature_degree"] = 4
500
501 # Problem parameters
502 epsilon = 1E-6 # perturbation factor of the physical PDE
503 N = 32 # mesh intervals in final mesh
504
505 # parameters for rho(x)
506 sigma = 2.5 # related to degree of elements
507 b = 0.99 # lower bound on reaction coefficient
508 q = 0.5 # proportion of mesh points in layer
509 K = q/(sigma*(1-q)) # proportion of points in the layer region
510
511 # Parameters for h-refinement steps
512 N_start = 4 # initial number of mesh intervals
513 N_step = N_start # current number of mesh intervals
514 step_index = 0 # initial step number
515
516 # Define right side of PDE and initial solution
517 f = Expression('0.0', degree = 2)
518 x_0 = Expression('x[0]', degree = 2)
519

```

```

520 # Define stopping criteria tolerance and initial value
521 residual_stop_TOL = 0.2
522 residual_stop = residual_stop_TOL + 1
523
524 # Lists for mesh and function space names
525 mesh_list = ["mesh0", "mesh1", "mesh2", "mesh3", "mesh4", "mesh5", "mesh6", "mesh7",
526             "mesh8", "mesh9"]
527 V_list = ["V0", "V1", "V2", "V3", "V4", "V5", "V6", "V7", "V8", "V9"]
528
529 # Computational function space parameters
530 def comp_space(N):
531     meshc = UnitIntervalMesh(N) # Uniform unit interval
532     mesh =
533     V = FunctionSpace(meshc, 'P', 1) # Function space V with linear
534     elements
535     v = TestFunction(V) # Test function
536     on V
537     x = TrialFunction(V) # Trial function on V
538     return meshc, V, v, x
539
540 # Define weak form of MPDE
541 def weak_form(rho, xN, x, v, f, meshz, V):
542     bc = DirichletBC(V, x_0, "on_boundary") # Define Dirichlet boundary
543     condition
544     a = rho(xN)*inner(grad(x), grad(v))*dx(meshz) # Left side of weak_form
545     L = f*v*dx # Right side of weak_form
546     xN = Function(V) # Numerical solution in V
547     solve(a==L, xN, bc) # Solve, applying the boundary
548     conditions strongly
549     return xN
550
551 # Define rho(x) for MPDE
552 def rho(xN):
553     return conditional((K*(b/epsilon)*exp(-b*xN/(sigma*epsilon)))>1.0, K*(b/
554     epsilon)*exp(-b*xN/(sigma*epsilon)), 1.0)
555
556 # Calculate the norm of the residual to use as stopping criterion
557 def calc_residual_stop(rho, V, meshN):
558     residual = assemble((rho(xN)*xN.dx(0)*v.dx(0)-f*v)*dx) # residual of xN
559     bcr = DirichletBC(V, '0.0', "on_boundary")
560     bcr.apply(residual) # apply
561     Dirichlet boundary conditions
562     residual_func = Function(V)
563     residual_func.vector()[:] = residual.get_local() # Residual function on
564     V
565     residual_stop = norm(residual_func, 'L2', meshN) # L2-norm of residual
566     return residual_stop
567
568 # Initial function space parameters
569 mesh_list[0], V_list[0], v, x = comp_space(N_step) #
570
571 # Initial value for xN
572 xN = interpolate(x_0, V_list[0])
573
574 # Iterate through uniform h-refinements
575 while N_step < N:
576
577     # Calculate solution twice on each mesh size
578     for i in range(0,2):
579         xN = weak_form(rho, xN, x, v, f, mesh_list[step_index], V_list[
580         step_index])
581
582     # Double the number of mesh intervals
583     N_step = N_step*2
584     step_index = round(math.log(N_step/N_start,2))
585
586     # Generate computational function space on the finer mesh
587     mesh_list[step_index], V_list[step_index], v, x = comp_space(N_step)
588
589     # Interpolate the solution onto the new computational function space
590     xN = interpolate(xN, V_list[step_index])
591
592 # Calculate solution and L2 norm of the residual on final function space

```

```

593 while residual_stop > residual_stop_TOL:
594
595     xN = weak_form(rho, xN, x, v, f, mesh_list[step_index], V_list[
596         step_index])
597     residual_stop = calc_residual_stop(rho, V_list[step_index], mesh_list[
598         step_index])
599
600 # Generate the physical mesh
601 meshp = UnitIntervalMesh(N)
602 meshp.coordinates()[0] = xN.compute_vertex_values()[0]
603

```

C Python Code to generate a two-dimensional layer-adapted mesh for an SPP whose diffusion varies spatially and has layer regions near $x = 1$ and $y = 1$ using a combination of r - and h -refinement.

Listing 4: Code to generate a two-dimensional layer-adapted mesh using a combination of r - and h -refinement.

```

609 ## 2D_layer_adapted_mesh_rh_refinement.py
610 """
611 Generate a non-tensor product layer adapted mesh, for a singularly
612 perturbed PDE whose solution has layers that vary spatially near x = 1 and y =
613 1,
614 and a corner layer near (x,y) = (1,1)
615 by solving the MPDE -grad(M(x) grad(x(xi1, xi2))) = 0 for (xi1,xi2) in (0,1)^2
616 with
617 x(0,xi2) = (0,x_xi2=0),
618 x(1,xi2) = (1,x_xi2=0),
619 x(xi1,0) = (x_xi1=0,0),
620 x(xi1,1) = (x_xi1=0,1),
621 and using uniform h-refinements.
622
623 This code is part of
624 Mesh partial differential equations (PDEs): application and implementation,
625 written by Roisin Hill and Niall Madden
626 Contact: Roisin Hill
627 For: <cite paper>
628 """
629
630
631 from fenics import *
632 import math
633
634 # Problem parameters
635 epsilon = 1E-2 # perturbation factor of the physical PDE
636 N = 32 # N+1 mesh points in each direction in
637     final mesh
638
639 # Parameters for h refinement steps
640 N_start = 4 # initial mesh size
641 N_step = N_start # current mesh size
642 step_index = 0 # initial step index
643
644 # Parameters for M(x)
645 sigma = 2.5 # dependant on method, often degree of
646     elements +1 or +1.5
647 b = 0.99 # minimum coefficient of reaction or
648     convectio term as appropriate
649 q = 0.5 # proportion of mesh points in layer
650 K = q/(sigma*(1-q)) # indicates the proportion of points in the
651     layer region
652
653 # define right side of MPDE

```

```

654 f = Expression(('0.0','0.0'), degree = 2)
655
656 # Lists for mesh and function space names
657 mesh_list = ["mesh0", "mesh1", "mesh2", "mesh3", "mesh4", "mesh5", "mesh6", "mesh7",
658             "mesh8", "mesh9"]
659 V_list = ["V0", "V1", "V2", "V3", "V4", "V5", "V6", "V7", "V8", "V9"]
660
661 # generate computational function space
662 def comp_space(N):
663     mesh = UnitSquareMesh(N, N, diagonal = "crossed") # uniform mesh on
664     (0,1)^2
665     V = VectorFunctionSpace(mesh, 'P', 1) # Function space V with linear
666     elements
667     v = TrialFunction(V) # Test function
668     on V
669     x = TestFunction(V) # Trial
670     function on V
671     return mesh, V, v, x
672
673 # Define boundaries
674 TOL = 1E-14
675 def left_boundary( x, on_boundary):
676     return abs(x[0]) < TOL
677 def right_boundary(x, on_boundary):
678     return abs(1-x[0]) < TOL
679 def bottom_boundary( x, on_boundary):
680     return abs(x[1]) < TOL
681 def top_boundary( x, on_boundary):
682     return abs(1-x[1]) < TOL
683
684 # Define Dirichlet boundary conditions
685 def boundary_conditions(V):
686     bcl = DirichletBC(V.sub(0), '0.0', left_boundary)
687     bcr = DirichletBC(V.sub(0), '1.0', right_boundary)
688     bcb = DirichletBC(V.sub(1), '0.0', bottom_boundary)
689     bct = DirichletBC(V.sub(1), '1.0', top_boundary)
690     bcs = [bcl, bcr, bcb, bct]
691     return bcs
692
693 # Define the matrix M(x) for MPDE
694 def M(xN):
695     epsilon_y = Expression('epsilon*(1+xi1)*(1+xi1)', epsilon = epsilon, xi1
696     = xN.sub(0), degree = 2)
697     epsilon_x = Expression('epsilon*(2-xi2)*(2-xi2)', epsilon = epsilon, xi2
698     = xN.sub(1), degree = 2)
699     M = Expression((( 'K*(b/epsilon_x)*exp(-b*(1-xi1)/(sigma*epsilon_x))>1?_
700     K*(b/epsilon_x)*exp(-b*(1-xi1)/(sigma*epsilon_x))_<1', '0'), \
701     ('0', 'K*(b/epsilon_y)*exp(-b*(1-xi2)/(sigma*epsilon_y))>1?_K*(b/
702     epsilon_y)*exp(-b*(1-xi2)/(sigma*epsilon_y))_<1')), \
703     K=K, b=b, sigma=sigma, epsilon_x = epsilon_x, epsilon_y = epsilon_y, xi1
704     = xN.sub(0), xi2 = xN.sub(1), degree = 4)
705     return M
706
707 # Define the weak form
708 def weak_form(M, xN, x, v, f, bcs, meshz, V):
709     a = inner(M(xN)*grad(x), grad(v))*dx(meshz) # Left side of weak
710     weak_form
711     L = dot(f,v)*dx
712     # Right side of weak weak_form
713     xN = Function(V)
714     # Numerical solution in V
715     solve(a==L, xN, bcs) # Solve,
716     applying the DirichletBC strongly
717     return xN
718
719 # Initial function space parameters
720 mesh_list[0], V_list[0], v, x = comp_space(N_step)
721
722 # Set intial value for xN: xN(xi1, xi2) = (xi1, xi2)
723 a = inner(grad(x), grad(v))*dx
724 L = dot(f,v)*dx
725 xN = Function(V_list[0])
726 solve(a==L, xN, boundary_conditions(V_list[0]))

```

```

727
728 # Iterate through uniform h-refinements
729 while N_step < N:
730
731     # Generate solution four times on each mesh size
732     for i in range (0,4):
733         xN = weak_form(M, xN, x, v, f, boundary_conditions(V_list[
734             step_index]), mesh_list[step_index], V_list[step_index])
735
736     # Double the number of mesh intervals in each direction
737     N_step = N_step*2
738     step_index = round(math.log(N_step/N_start,2))
739
740     # Generate computational function space on the finer mesh
741     mesh_list[step_index], V_list[step_index], v, x = comp_space(N_step)
742
743     # Interpolate solution onto the new computational function space
744     xN = interpolate(xN,V_list[step_index])
745
746 # Calculate solution in final function space
747 iterations = 0
748 while iterations < 5:
749
750     xN = weak_form(M, xN, x, v, f, boundary_conditions(V_list[step_index]),
751         mesh_list[step_index], V_list[step_index])
752     iterations +=1
753
754 # Generate the physical mesh
755 xiX, xiY = xN.split(True)
756 meshp = UnitSquareMesh(N,N)
757 meshp.coordinates()[:,0] = xiX.compute_vertex_values()[0:(N+1)*(N+1)]
758 meshp.coordinates()[:,1] = xiY.compute_vertex_values()[0:(N+1)*(N+1)]
759

```

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