

Context and issues - Coupling of the acoustic and elastic wave equations

Objectives

- Accurate modelisation and simulation of seismo-acoustic waves **through heterogeneous domains with complex geometries**
- Treatment of realistic cases of interest
 - High Performance Computing (HPC)**

Issues

- Difficulty to mesh complex geometries
- High-order precision needed to accurately capture waves
 - Hybrid discontinuous methods (HDG/HHO)**

Acoustic wave equation:

$$\begin{cases} \rho_F \partial_t \mathbf{v}^F(t) + \nabla p(t) = \mathbf{0} \\ \frac{1}{\kappa} \partial_t p(t) + \nabla \cdot \mathbf{v}^F(t) = g(t) \end{cases}$$

Elastic wave equation:

$$\begin{cases} \partial_t \boldsymbol{\varepsilon}(t) - \nabla^s \mathbf{v}^S(t) = \mathbf{0} \\ \rho_S \partial_t \mathbf{v}^S(t) - \nabla \cdot (\mathbf{C} : \boldsymbol{\varepsilon}(t)) = \mathbf{f}(t) \end{cases}$$

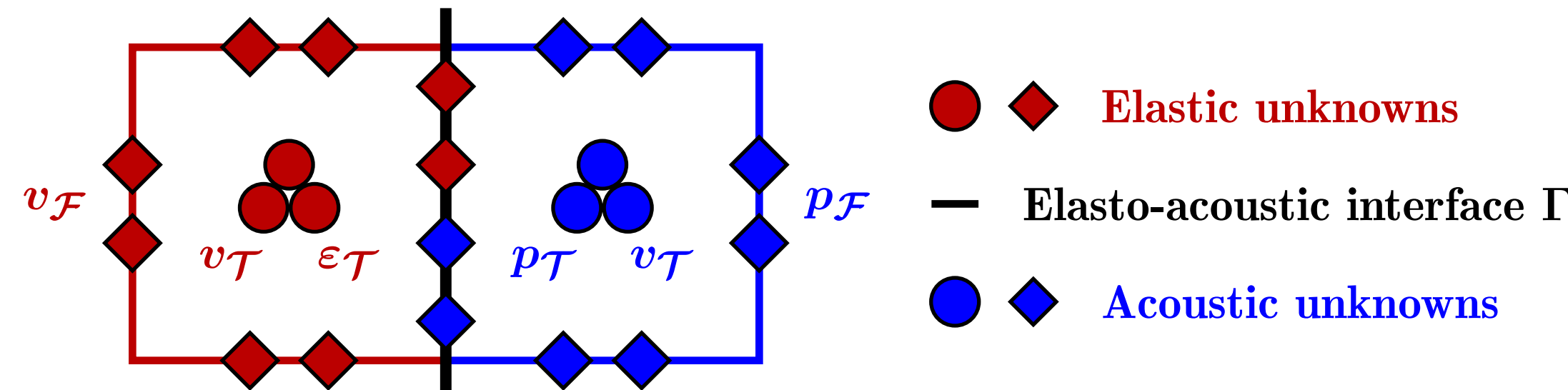
Coupling condition:

$$\begin{cases} \llbracket \mathbf{v}(t) \cdot \mathbf{n}_\Gamma \rrbracket_\Gamma = 0 \\ (\mathbf{C} : \boldsymbol{\varepsilon}(t)) \cdot \mathbf{n}_\Gamma = p(t) \mathbf{n}_\Gamma \end{cases}$$

Application of HHO method to seismo-acoustic coupling

Approximation spaces:

- Acoustic domain:** $\mathbf{V}_{\mathcal{T}_h}^k := \bigtimes_{T \in \mathcal{T}_h^F} \mathbb{P}^k(T; \mathbb{R}^d)$, $\hat{\mathbf{V}}_h^F := \bigtimes_{T \in \mathcal{T}_h^F} \mathbb{P}^{k'}(T; \mathbb{R}) \times \bigtimes_{F \in \mathcal{F}_h^F} \mathbb{P}^k(F; \mathbb{R})$
- Elastic domain:** $\mathbf{Z}_{\mathcal{T}_h}^k := \bigtimes_{T \in \mathcal{T}_h^S} \mathbb{P}^k(T; \mathbb{R}^{d \times d})$, $\hat{\mathbf{V}}_h^S := \bigtimes_{T \in \mathcal{T}_h^S} \mathbb{P}^{k'}(T; \mathbb{R}^d) \times \bigtimes_{F \in \mathcal{F}_h^S} \mathbb{P}^k(F; \mathbb{R}^d)$

Fig. 1: Elasto-acoustic unknowns with a mixed-order ($k' = k + 1 = 2$) discretization.

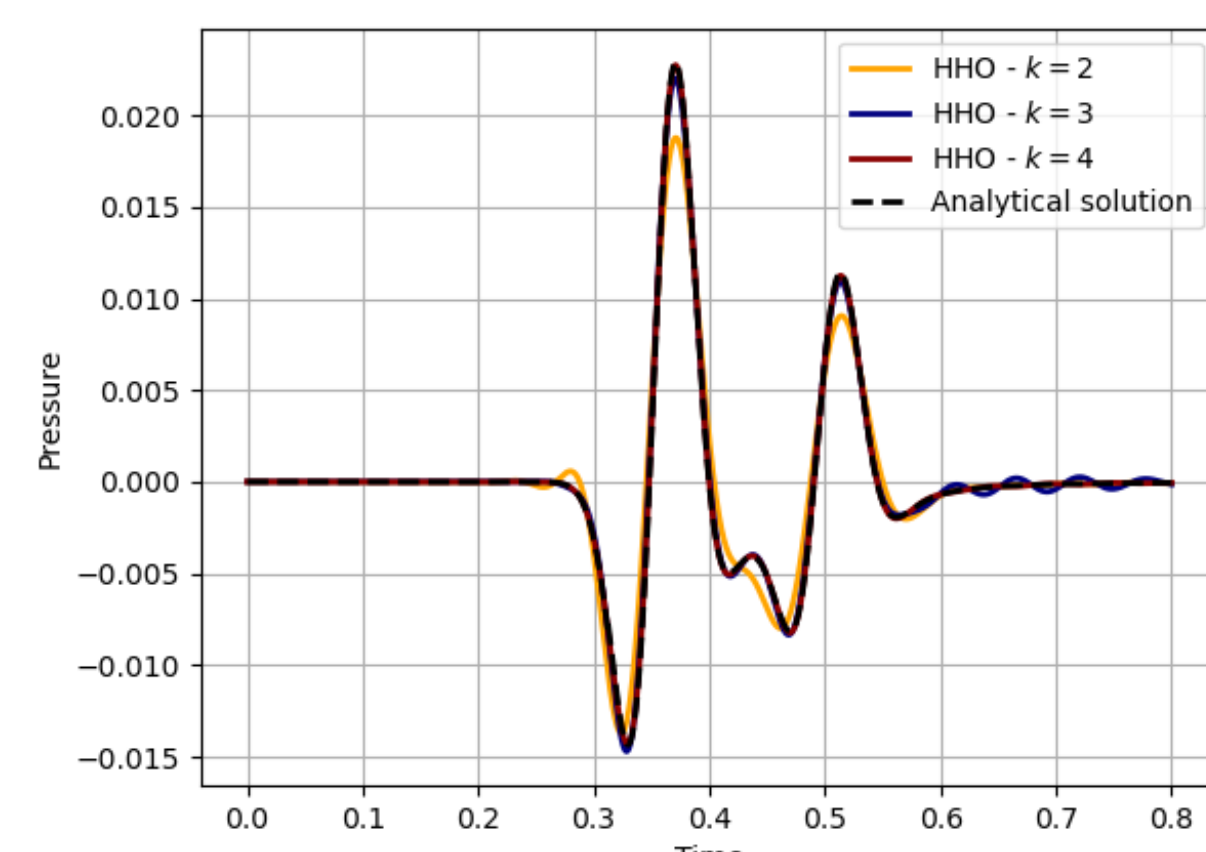
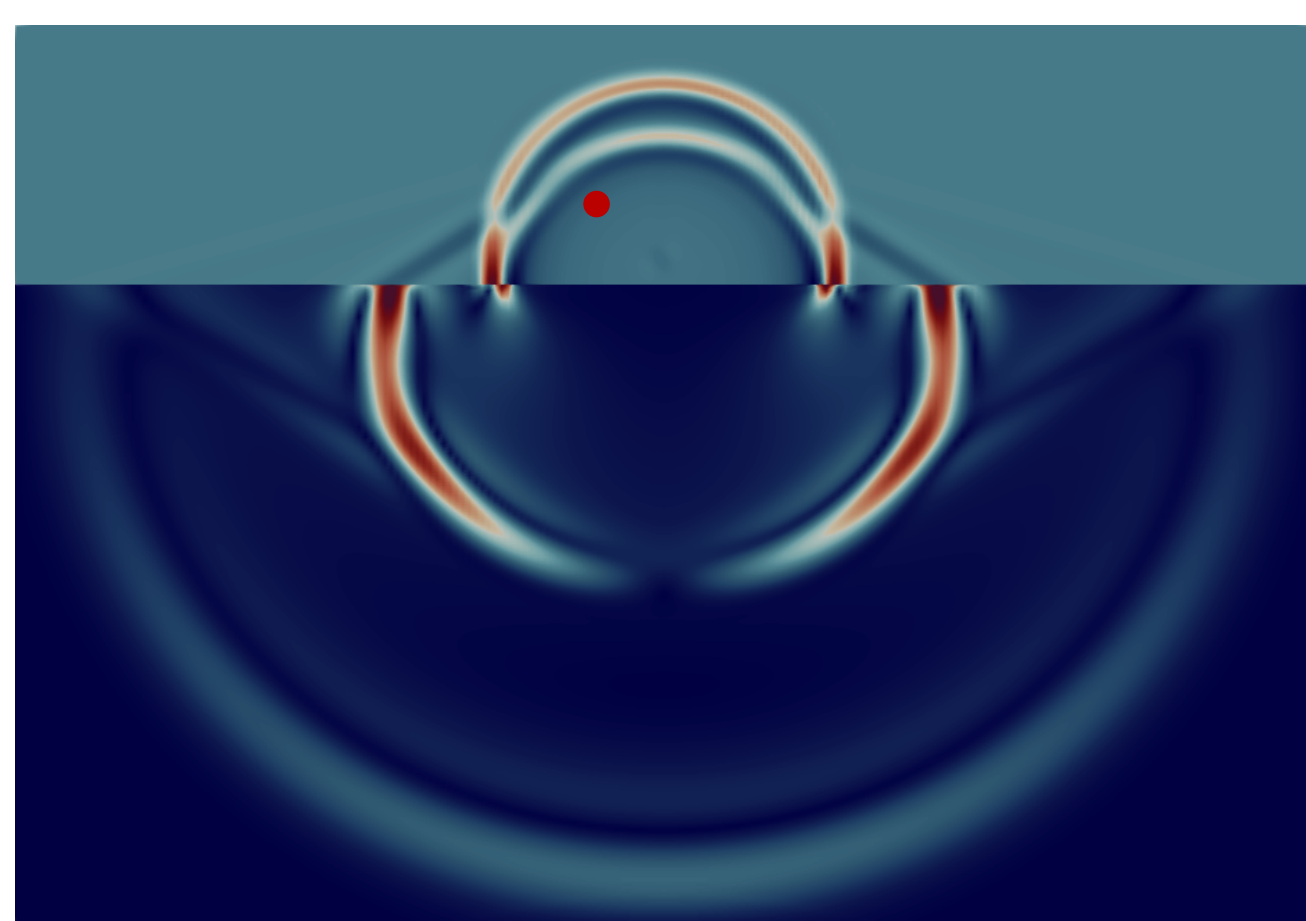
Algebraic realization:

$$\begin{bmatrix} \mathbf{M}_{\mathcal{T}\mathcal{T}}^V & 0 & 0 & 0 & 0 & 0 \\ 0 & \mathbf{M}_{\mathcal{T}\mathcal{T}}^F & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \mathbf{M}_{\mathcal{T}\mathcal{T}}^S & 0 & 0 \\ 0 & 0 & 0 & 0 & \mathbf{M}_{\mathcal{T}\mathcal{T}}^S & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \partial_t \mathbf{V}_{\mathcal{T}}^F \\ \partial_t \mathbf{P}_{\mathcal{T}} \\ \partial_t \mathbf{P}_{\mathcal{F}} \\ \partial_t \mathbf{S}_{\mathcal{T}} \\ \partial_t \mathbf{V}_{\mathcal{T}} \\ \partial_t \mathbf{V}_{\mathcal{F}} \end{bmatrix} + \begin{bmatrix} 0 & -\mathbf{G}_{\mathcal{T}} & -\mathbf{G}_{\mathcal{F}} & 0 & 0 & 0 \\ \mathbf{G}_{\mathcal{T}}^\dagger & \Sigma_{\mathcal{T}\mathcal{T}}^F & \Sigma_{\mathcal{T}\mathcal{F}}^F & 0 & 0 & 0 \\ \mathbf{G}_{\mathcal{F}}^\dagger & \Sigma_{\mathcal{F}\mathcal{T}}^F & \Sigma_{\mathcal{F}\mathcal{F}}^F & 0 & 0 & \mathbf{C}_{\mathcal{T}} \\ 0 & 0 & 0 & 0 & -\mathbf{E}_{\mathcal{T}} & -\mathbf{E}_{\mathcal{F}} \\ 0 & 0 & 0 & \mathbf{E}_{\mathcal{T}}^\dagger & \Sigma_{\mathcal{T}\mathcal{T}}^S & \Sigma_{\mathcal{T}\mathcal{F}}^S \\ 0 & 0 & -\mathbf{C}_{\mathcal{T}}^\dagger & \mathbf{E}_{\mathcal{F}}^\dagger & \Sigma_{\mathcal{F}\mathcal{T}}^S & \Sigma_{\mathcal{F}\mathcal{F}}^S \end{bmatrix} \begin{bmatrix} \mathbf{V}_{\mathcal{T}}^F \\ \mathbf{P}_{\mathcal{T}} \\ \mathbf{P}_{\mathcal{F}} \\ \mathbf{S}_{\mathcal{T}} \\ \mathbf{V}_{\mathcal{T}}^S \\ \mathbf{V}_{\mathcal{F}}^S \end{bmatrix} = \begin{bmatrix} 0 \\ \mathbf{G}_{\mathcal{T}} \\ 0 \\ 0 \\ \mathbf{F}_{\mathcal{T}} \\ 0 \end{bmatrix}$$

Propagation of an acoustic (water) pulse into an elastic medium (granit)

Simulation parameters:

- Computational domain:**
 - Water on the upper side
 - Granit on the lower side
- Homogeneous Dirichlet conditions**
- Initial condition:** velocity Ricker wave in the acoustic medium:
- Time integration scheme:** SDIRK(3,4)
- Time step:** $\text{dt} = 0, 1 \times 2^{-9}$
- Mixed-order discretization:** $k' = k + 1 = 3$
- $\mathbf{v}_0^F(x, y) := \theta \exp\left(-\pi^2 \frac{r^2}{\lambda^2}\right) (x - x_c, y - y_c)^T$

Fig. 5: Left panel: Two-dimensional distribution of the acoustic pressure and the elastic velocity norm, at $t = 0, 4375$ s. Right panel: Pressure as a function of time at the sensor in the water (coarse mesh).

General principles of the HHO method

HHO is a finite element method similar to Hybrid Discontinuous Galerkin method (HDG)

Design of the HHO method

- Degrees of freedom:**
 - Polynomial unknowns located in the cells (degree k') and on the faces (degree k): $\hat{\mathbf{u}}_h := (\mathbf{u}_{\mathcal{T}}, \mathbf{u}_{\mathcal{F}})$
 - Equal-order discretization: $k' = k$
 - Mixed-order discretization: $k' = k + 1$
- Operators:**
 - Gradient reconstruction operator: $\nabla \mathbf{u} \rightarrow \mathbf{G}(\hat{\mathbf{u}}_h)$,
 - Stabilization operator: $\mathbf{s}(\hat{\mathbf{u}}_h, \hat{\mathbf{w}}_h)$
 - Penalization at the element level to ensure stability while preserving the approximation properties of the reconstruction.

Advantages over classical finite element methods

- Mesh flexibility:**
 - Complex geometries
 - Unstructured and polyhedral meshes
 - Local mesh refinement
- Local conservativity**
- Optimal error estimates for smooth solutions**
- Attractive computational costs:**
 - Global problem couples only face dofs
 - Cell dofs recovered by local post-processing

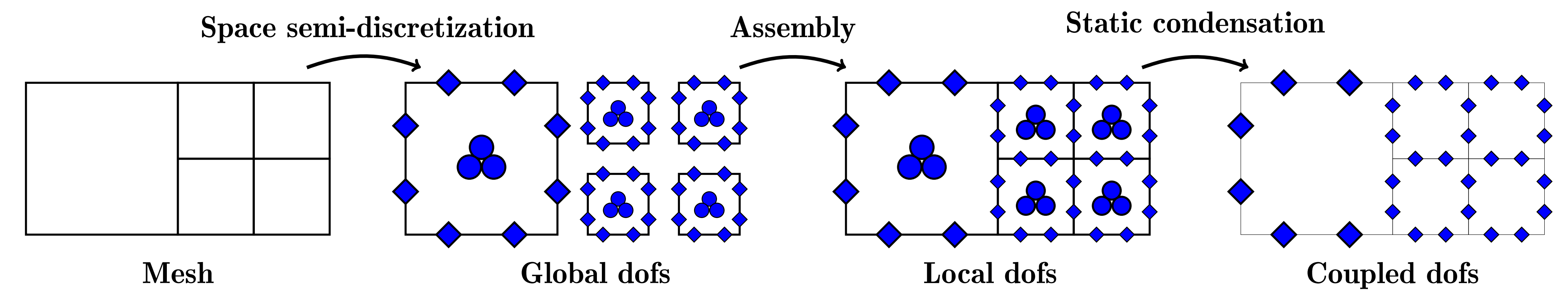


Fig. 2: Static condensation procedure.

Validation on academic test cases

Verification of convergence rates on sinusoidal analytical solutions:

- $\mathcal{O}(h^{k+1})$ in H^1 -norm
- $\mathcal{O}(h^{k+2})$ in L^2 -norm (superconvergence)

Verification of the energy conservation of the scheme with a no contrast test case:

- $\rho_S = \rho_F = 1$, $c_S^P = c_F^P = \sqrt{3}$, $c_S^S = 1$
- Initial condition:** velocity Ricker wave in the acoustic medium:

$$\mathbf{v}_0^F(x, y) := \theta \exp\left(-\pi^2 \frac{r^2}{\lambda^2}\right) (x - x_c, y - y_c)^T$$

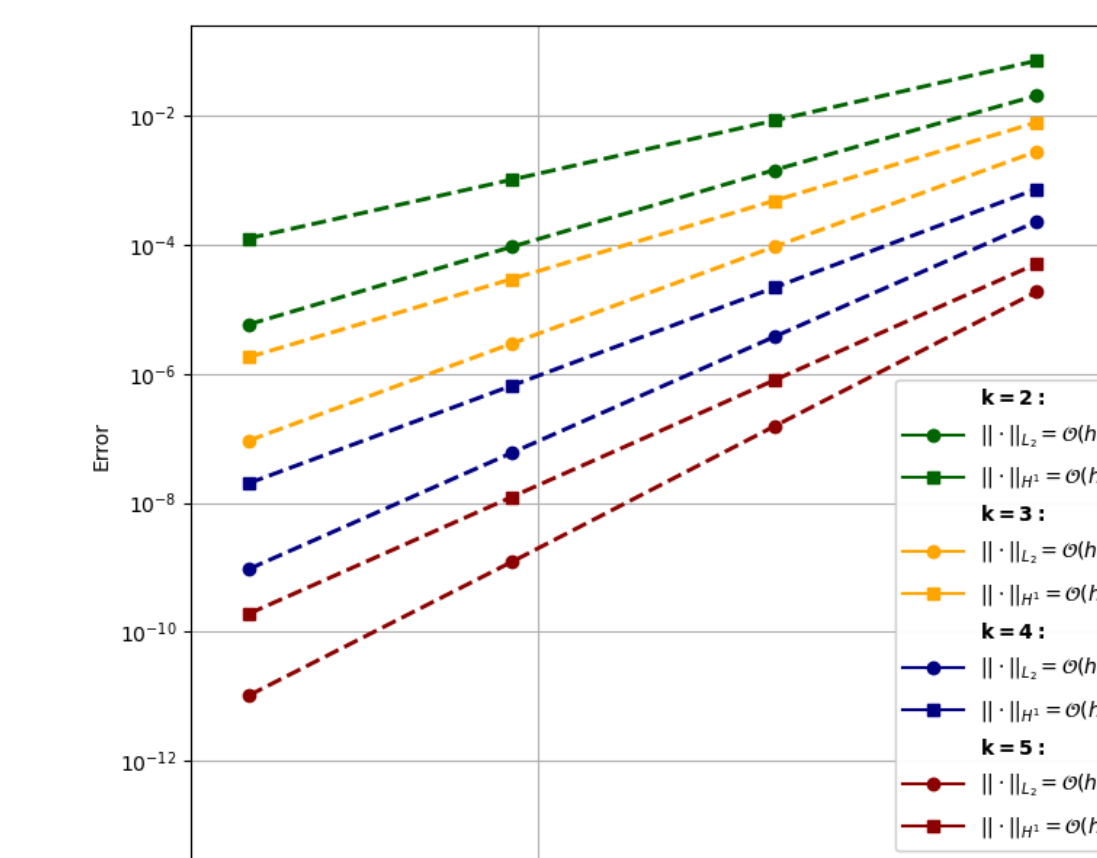
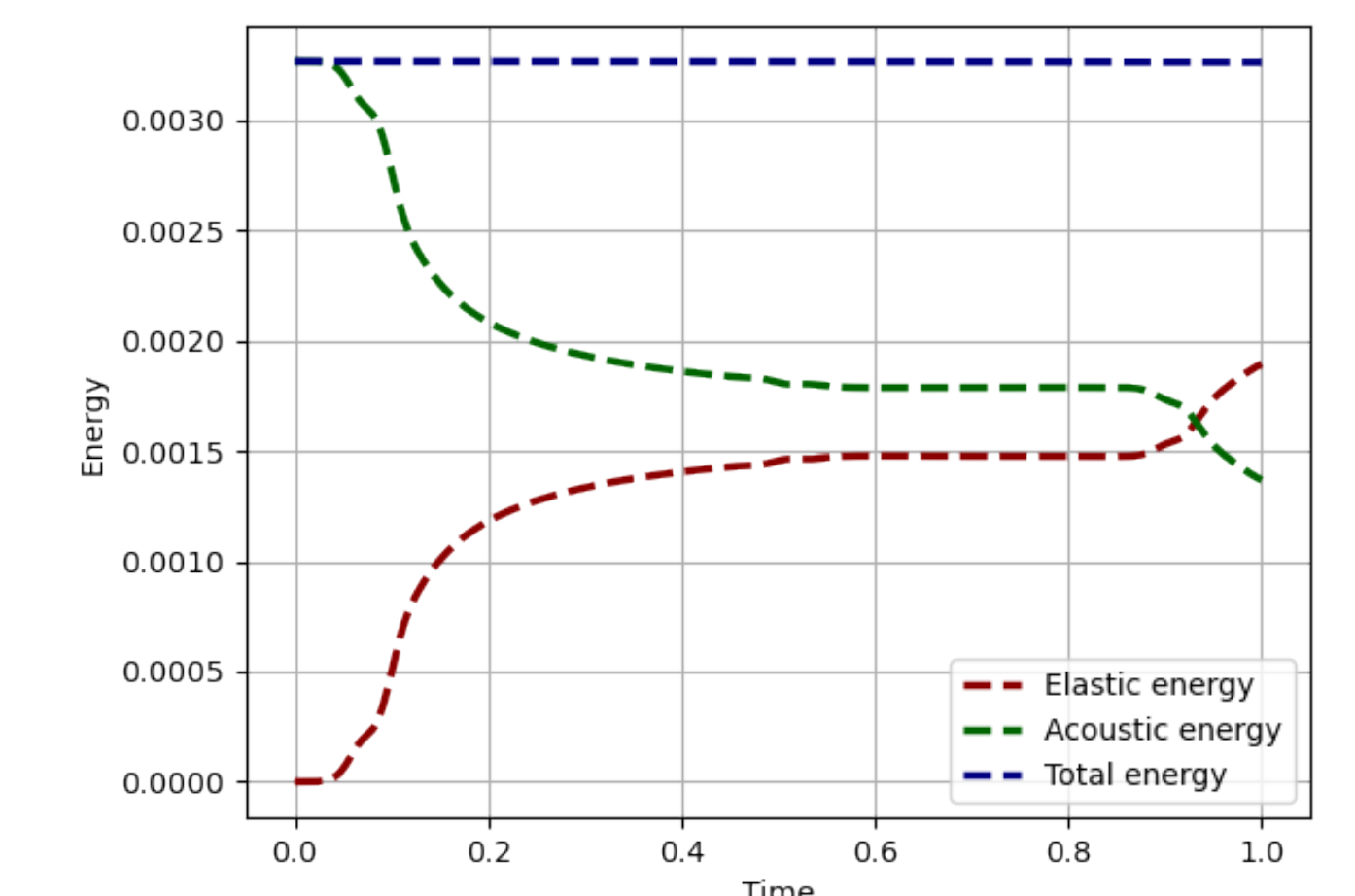
Fig. 3: Errors as a function of the mesh size with $\Delta t = 0.1 \times 2^{-5}$.

Fig. 4: Repartition of the elastic and acoustic energy after propagation of the wave.

Propagation of an elastic pulse into a sedimentary basin

- Composition of the sedimentary basin:**
 - Acoustic region: air
 - Sedimentary region: $\rho = 1200 \text{ kg.m}^3$, $c_P = 3400 \text{ m.s}^{-1}$, $c_S = 1400 \text{ m.s}^{-1}$
 - Elastic region: $\rho = 5350 \text{ kg.m}^3$, $c_P = 3090 \text{ m.s}^{-1}$, $c_S = 2570 \text{ m.s}^{-1}$

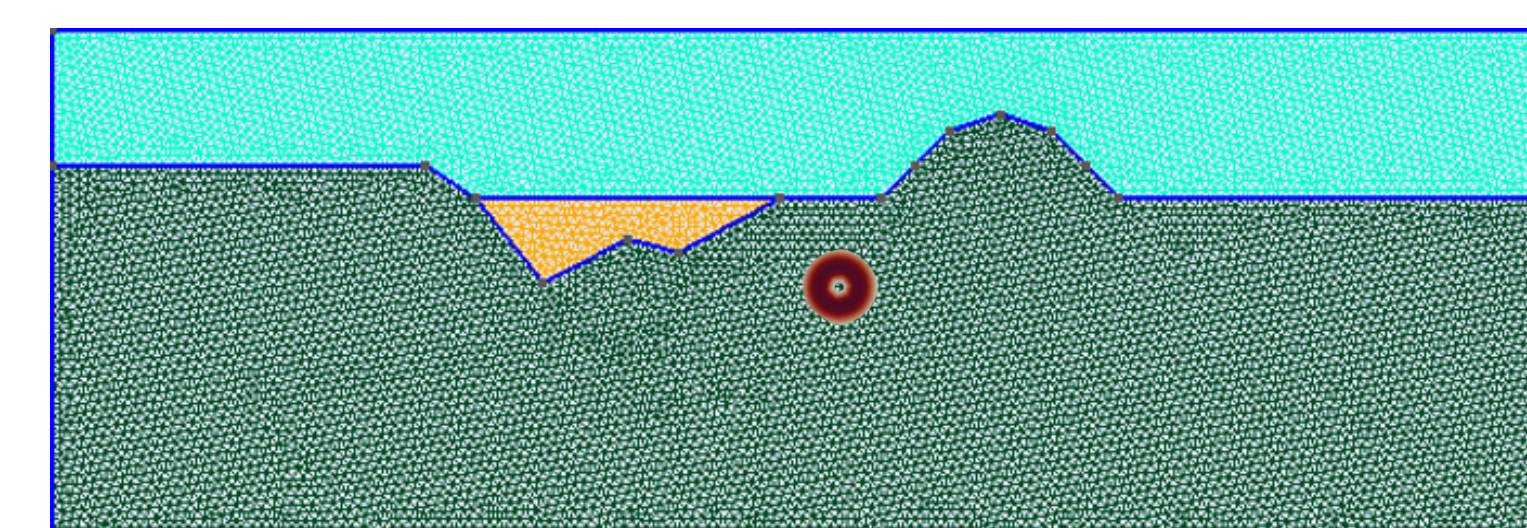


Fig. 6: Mesh of a sedimentary basin and location of the initial pulse.

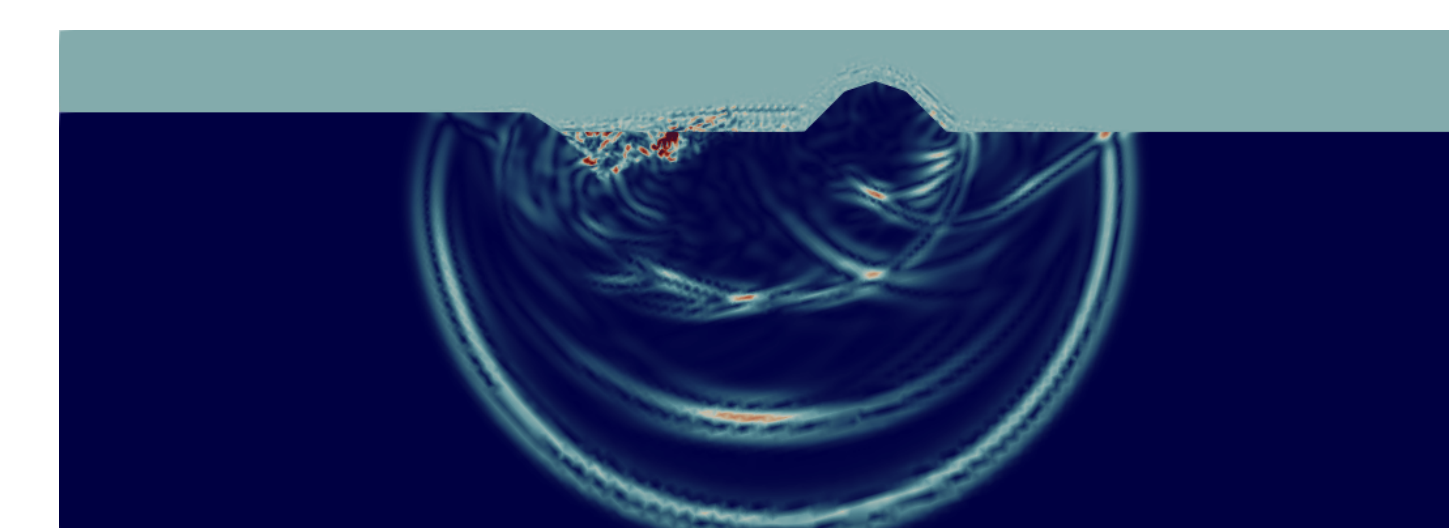


Fig. 7: Propagation through the sedimentary basin.

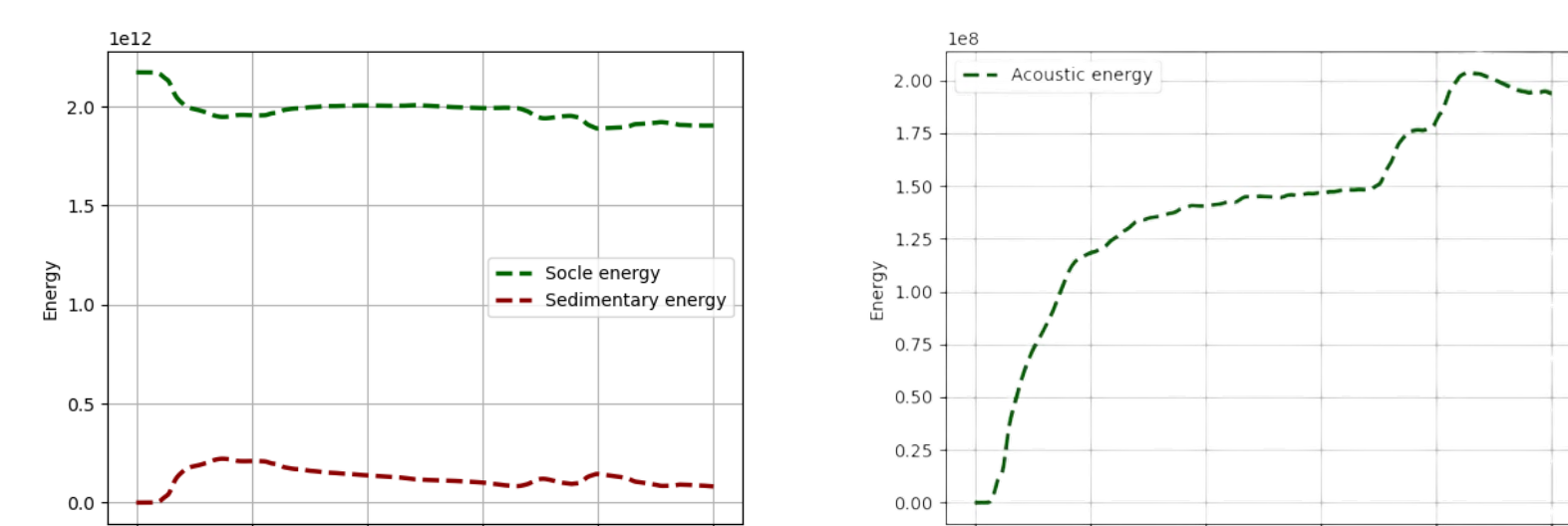


Fig. 8: Repartition of the elastic (left) and the acoustic (right) energy.

Some references

- Burman, Duran, and Ern. "Hybrid high-order methods for the acoustic wave equation in the time domain". In: *Comm. App. Math. Comp. Sci.* 4.2 (2022), pp. 597–633.
- Di Pietro and Ern. "A hybrid high-order locking-free method for linear elasticity on general meshes". In: *Comput. Meth. Appl. Mech. Engrg.* 283 (2015), pp. 1–21.
- Terrana, Vilotte, and Guillot. "A spectral hybridizable discontinuous Galerkin method for elastic-acoustic wave propagation". In: *Geophys. J. Int.* 213.1 (2017), pp. 574–602.