Preface

This textbook is intended for use by students of physics, physical chemistry, and theoretical chemistry. The reader is presumed to have a basic knowledge of atomic and quantum physics at the level provided, for example, by the first few chapters in our book *The Physics of Atoms and Quanta*. The student of physics will find here material which should be included in the basic education of every physicist. This book should furthermore allow students to acquire an appreciation of the breadth and variety within the field of molecular physics and its future as a fascinating area of research.

For the student of chemistry, the concepts introduced in this book will provide a theoretical framework for that entire field of study. With the help of these concepts, it is at least in principle possible to reduce the enormous body of empirical chemical knowledge to a few basic principles: those of quantum mechanics. In addition, modern physical methods whose fundamentals are introduced here are becoming increasingly important in chemistry and now represent indispensable tools for the chemist. As examples, we might mention the structural analysis of complex organic compounds, spectroscopic investigation of very rapid reaction processes or, as a practical application, the remote detection of pollutants in the air.

April 1995

Walter Olthoff Program Chair ECOOP'95

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Table of Contents

Hamiltonian Mechanics unter besonderer Berücksichtigung der höhreren Lehranstalten

Ivar Ekeland¹, Roger Temam² Jeffrey Dean, David Grove, Craig Chambers, Kim B. Bruce, and Elsa Bertino

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$$x(0) = x(T)$$

with $H(t,\cdot)$ a convex function of x, going to $+\infty$ when $||x|| \to \infty$.

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$$\gamma := \text{smallest eigenvalue of } B_{\infty} - A_{\infty}$$
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Theorem 1 tells us that if $\lambda + \gamma < 0$, the boundary-value problem:

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has at least one solution \overline{x} , which is found by minimizing the dual action functional:

$$\psi(u) = \int_{o}^{T} \left[\frac{1}{2} \left(\Lambda_{o}^{-1} u, u \right) + N^{*}(-u) \right] dt \tag{4}$$

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Proposition 1. Assume H'(0) = 0 and H(0) = 0. Set:

$$\delta := \liminf_{x \to 0} 2N(x) \|x\|^{-2} . \tag{7}$$

If $\gamma < -\lambda < \delta$, the solution \overline{u} is non-zero:

$$\overline{x}(t) \neq 0 \quad \forall t \ .$$
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Proof. Condition (7) means that, for every $\delta' > \delta$, there is some $\varepsilon > 0$ such that

$$||x|| \le \varepsilon \Rightarrow N(x) \le \frac{\delta'}{2} ||x||^2$$
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It is an exercise in convex analysis, into which we shall not go, to show that this implies that there is an $\eta > 0$ such that

$$f \|x\| \le \eta \Rightarrow N^*(y) \le \frac{1}{2\delta'} \|y\|^2 . \tag{10}$$

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Since u_1 is a smooth function, we will have $||hu_1||_{\infty} \leq \eta$ for h small enough, and inequality (10) will hold, yielding thereby:

$$\psi(hu_1) \le \frac{h^2}{2} \frac{1}{\lambda} \|u_1\|_2^2 + \frac{h^2}{2} \frac{1}{\delta'} \|u_1\|^2 . \tag{11}$$

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Corollary 1. Assume H is C^2 and (a_{∞}, b_{∞}) -subquadratic at infinity. Let ξ_1, \ldots, ξ_N be the equilibria, that is, the solutions of $H'(\xi) = 0$. Denote by ω_k the smallest eigenvalue of $H''(\xi_k)$, and set:

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If:

$$\frac{T}{2\pi}b_{\infty} < -E\left[-\frac{T}{2\pi}a_{\infty}\right] < \frac{T}{2\pi}\omega\tag{14}$$

then minimization of ψ yields a non-constant T-periodic solution \overline{x} .

We recall once more that by the integer part $E[\alpha]$ of $\alpha \in \mathbb{R}$, we mean the $a \in \mathbb{Z}$ such that $a < \alpha \le a+1$. For instance, if we take $a_{\infty} = 0$, Corollary 2 tells us that \overline{x} exists and is non-constant provided that:

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Proof. The spectrum of Λ is $\frac{2\pi}{T}Z + a_{\infty}$. The largest negative eigenvalue λ is given by $\frac{2\pi}{T}k_o + a_{\infty}$, where

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Hence:

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The condition $\gamma < -\lambda < \delta$ now becomes:

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which is precisely condition (14).

Lemma 1. Assume that H is C^2 on $\mathbb{R}^{2n}\setminus\{0\}$ and that H''(x) is non-degenerate for any $x \neq 0$. Then any local minimizer \widetilde{x} of ψ has minimal period T.

Proof. We know that \widetilde{x} , or $\widetilde{x} + \xi$ for some constant $\xi \in \mathbb{R}^{2n}$, is a T-periodic solution of the Hamiltonian system:

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This would contradict (21), and thus cannot happen.

Notes and Comments. The results in this section are a refined version of [1]; the minimality result of Proposition 14 was the first of its kind.

To understand the nontriviality conditions, such as the one in formula (16), one may think of a one-parameter family x_T , $T \in (2\pi\omega^{-1}, 2\pi b_{\infty}^{-1})$ of periodic solutions, $x_T(0) = x_T(T)$, with x_T going away to infinity when $T \to 2\pi\omega^{-1}$, which is the period of the linearized system at 0.

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$$\forall t$$
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$$\exists c \in \mathbb{R} : H(t,x) \le \frac{1}{2} (B_{\infty}(t)x, x) + c \quad \forall x .$$
 (35)

If $A_{\infty}(t) = a_{\infty}I$ and $B_{\infty}(t) = b_{\infty}I$, with $a_{\infty} \leq b_{\infty} \in \mathbb{R}$, we shall say that H is (a_{∞}, b_{∞}) -subquadratic at infinity. As an example, the function $\|x\|^{\alpha}$, with $1 \leq \alpha < 2$, is $(0, \varepsilon)$ -subquadratic at infinity for every $\varepsilon > 0$. Similarly, the Hamiltonian

$$H(t,x) = \frac{1}{2}k \|k\|^2 + \|x\|^{\alpha}$$
(36)

is $(k, k + \varepsilon)$ -subquadratic for every $\varepsilon > 0$. Note that, if k < 0, it is not convex.

Notes and Comments. The first results on subharmonics were obtained by Rabinowitz in [5], who showed the existence of infinitely many subharmonics both in the subquadratic and superquadratic case, with suitable growth conditions on H'. Again the duality approach enabled Clarke and Ekeland in [2] to treat the same problem in the convex-subquadratic case, with growth conditions on H only.

Recently, Michalek and Tarantello (see [3] and [4]) have obtained lower bound on the number of subharmonics of period kT, based on symmetry considerations and on pinching estimates, as in Sect. 5.2 of this article.

References

- Clarke, F., Ekeland, I.: Nonlinear oscillations and boundary-value problems for Hamiltonian systems. Arch. Rat. Mech. Anal. 78, 315–333 (1982)
- 2. Clarke, F., Ekeland, I.: Solutions périodiques, du période donnée, des équations hamiltoniennes. Note CRAS Paris 287, 1013–1015 (1978)
- 3. Michalek, R., Tarantello, G.: Subharmonic solutions with prescribed minimal period for nonautonomous Hamiltonian systems. J. Diff. Eq. 72, 28–55 (1988)
- 4. Tarantello, G.: Subharmonic solutions for Hamiltonian systems via a \mathbb{Z}_p pseudoindex theory. Annali di Matematica Pura (to appear)
- 5. Rabinowitz, P.: On subharmonic solutions of a Hamiltonian system. Comm. Pure Appl. Math. 33, 609–633 (1980)

Hamiltonian Mechanics2

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$$N(t,x) \ge n(\|x\|)$$
 with $n(s)s^{-1} \to +\infty$ as $s \to +\infty$ (34)

$$\exists c \in \mathbb{R} : H(t,x) \le \frac{1}{2} (B_{\infty}(t)x, x) + c \quad \forall x .$$
 (35)

If $A_{\infty}(t) = a_{\infty}I$ and $B_{\infty}(t) = b_{\infty}I$, with $a_{\infty} \leq b_{\infty} \in \mathbb{R}$, we shall say that H is (a_{∞}, b_{∞}) -subquadratic at infinity. As an example, the function $\|x\|^{\alpha}$, with $1 \leq \alpha < 2$, is $(0, \varepsilon)$ -subquadratic at infinity for every $\varepsilon > 0$. Similarly, the Hamiltonian

$$H(t,x) = \frac{1}{2}k \|k\|^2 + \|x\|^{\alpha}$$
(36)

is $(k, k + \varepsilon)$ -subquadratic for every $\varepsilon > 0$. Note that, if k < 0, it is not convex.

Notes and Comments. The first results on subharmonics were obtained by Rabinowitz in 1985, who showed the existence of infinitely many subharmonics both in the subquadratic and superquadratic case, with suitable growth conditions on H'. Again the duality approach enabled Clarke and Ekeland in 1981 to treat the same problem in the convex-subquadratic case, with growth conditions on H only.

Recently, Michalek and Tarantello (see Michalek, R., Tarantello, G. 1982 and Tarantello, G. 1983) have obtained lower bound on the number of subharmonics of period kT, based on symmetry considerations and on pinching estimates, as in Sect. 5.2 of this article.

References

- Clarke, F., Ekeland, I.: Nonlinear oscillations and boundary-value problems for Hamiltonian systems. Arch. Rat. Mech. Anal. 78, 315–333 (1982)
- Clarke, F., Ekeland, I.: Solutions périodiques, du période donnée, des équations hamiltoniennes. Note CRAS Paris 287, 1013–1015 (1978)
- Michalek, R., Tarantello, G.: Subharmonic solutions with prescribed minimal period for nonautonomous Hamiltonian systems. J. Diff. Eq. 72, 28–55 (1988)
- Tarantello, G.: Subharmonic solutions for Hamiltonian systems via a \mathbb{Z}_p pseudoindex theory. Annali di Matematica Pura (to appear)
- Rabinowitz, P.: On subharmonic solutions of a Hamiltonian system. Comm. Pure Appl. Math. 33, 609–633 (1980)

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