

# Disentangling Canted Phases and Phase Separation Regions with Spin Waves in Doped Manganites

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**The system:** Doped manganites present a rich phase structure at zero temperature: antiferromagnetic, canted and ferromagnetic phases, together with phase separation regions.

**The problem:** Upon increasing the doping the system goes from an AF-insulator to a F-conductor. It is not clear if in its way passes through a canted phase or a phase separation region (F-AF).

**A solution:** The interaction of the spin waves with the charge carriers lead to a doping dependence of the dispersion relation of the spin waves. This may allow to experimentally differentiate canted phases from phase separation regions in doped manganites.

# Doped Manganites

The manganites are a generic group of alloys which present a rich magnetic structure in their ground state, from antiferromagnetism to ferromagnetism going through canted phases (i.e., a non-collinear arrangement of the magnetic moments with  $0 < \theta < \pi$ ) or phase separation regions [1, 2, 3, 4, 5].

## General Formula:

$$\begin{aligned} & \textcolor{teal}{A} = Ca^{2+}, Sr^{2+}, Ba^{2+} \\ La_{1-\textcolor{violet}{x}}\textcolor{teal}{A}_{\textcolor{violet}{x}}MnO_3 & \quad (0 \leq \textcolor{violet}{x} \leq 1) \\ & La^{3+}, Mn^{\textcolor{red}{3}+}, \text{ or } Mn^{\textcolor{red}{4}+} \end{aligned}$$

Each substitution  $\textcolor{green}{La}^{3+} \rightarrow \textcolor{green}{A}^{2+}$  corresponds in the manganese atom to  $\textcolor{green}{Mn}^{3+} \rightarrow \textcolor{green}{Mn}^{4+}$ .

## Electronic Structure:

$$\textcolor{blue}{3d} \text{ (5-fold)} \longrightarrow \begin{cases} \textcolor{blue}{t}_{2g} & (3\text{-fold}) \\ \textcolor{blue}{e}_g & (2\text{-fold}) \end{cases}$$

## Magnetic Structure:

$$\begin{aligned} LaMnO_3 \text{ and } AMnO_3 & \longrightarrow \text{Ground State: } \textcolor{teal}{AF}\text{-insulator.} \\ La_{1-\textcolor{violet}{x}}\textcolor{violet}{A}_{\textcolor{violet}{x}}MnO_3 \text{ } (0.2 \leq \textcolor{violet}{x} \leq 0.4) & \longrightarrow \text{Ground State: } \textcolor{teal}{F}\text{-conductor.} \end{aligned}$$

# Continuum Double Exchange Model

Double Exchange Model by Zener [2] suggested a relation between the electronic and magnetic properties through the exchange of electrons among the  $Mn^{3+}$  and  $Mn^{4+}$  ions.

Low energy and momentum properties of the system can be studied from a continuum model, Continuum Double Exchange Model, where the fields representing the excitations are slowly varying in space and time:

Background Magnetizations are generated by the magnetic moments in the  $t_{2g}$ -band:  $\mathbf{M}_1(x)$  and  $\mathbf{M}_2(x)$ .

Charge Carriers correspond to the holes generated in the  $e_g$ -band as the doping increases:  $\psi_1(x)$  and  $\psi_2(x)$ , which interact locally (Hund coupling) with the corresponding magnetization.

$$\begin{aligned}\mathcal{L}(x) = & \psi_1^\dagger(x) \left[ (1 + i\epsilon)i\partial_0 + \frac{\partial_i^2}{2m} + \mu + J_H \frac{\sigma}{2} \mathbf{M}_1(x) \right] \psi_1(x) \\ & + \psi_2^\dagger(x) \left[ (1 + i\epsilon)i\partial_0 + \frac{\partial_i^2}{2m} + \mu + J_H \frac{\sigma}{2} \mathbf{M}_2(x) \right] \psi_2(x) \\ & + t \left( \psi_1^\dagger(x) \psi_2(x) + \psi_2^\dagger(x) \psi_1(x) \right) - J_{AF} \mathbf{M}_1(x) \mathbf{M}_2(x).\end{aligned}$$

$$2m \sim \frac{z}{a^2 t}, \quad (a - \text{lattice spacing, } z = 6 - \text{coordination number})$$

# Effective Potential: Phase Diagram

The integration of the fermionic fields in the path-integral yields an **effective potential** for doped manganites ( $t \ll J_H$ ) in terms of  $y = \cos(\theta/2)$  ( $\theta$  is the angle between  $\mathbf{M}_1$  and  $\mathbf{M}_2$ ) in doped manganites.

$$V_{eff} = V_0 \left[ (2y^2 - 1) - A \left( (y_0 + y)^{5/2} \theta(y_0 + y) + (y_0 - y)^{5/2} \theta(y_0 - y) \right) \right]$$

$$V_0 = J_{AF} M^2 \quad A = \frac{z^{3/2}}{15\pi^2} \frac{t}{(J_{AF} a^3 M^2)}$$

$$\text{Chemical potential} \sim y_0 \longrightarrow \text{Doping: } x = -\frac{a^3}{t} \frac{\partial V_{eff}}{\partial y_0}$$

Upon minimization with respect to  $y$  the following phases arise:

**Antiferromagnetic:**  $y = 0$ , *AFI*, *AFC2*

**Canted:**  $0 < y < 1$ , *CC1*, *CC2*

**Ferromagnetic:**  $y = 1$ , *FC1*

- For typical values of the coupling constants  $A \sim 1 - 2$ .
- In this range either, **canted phases** or **phase separation regions** can arise from the phase diagram shown in **fig. 1**.

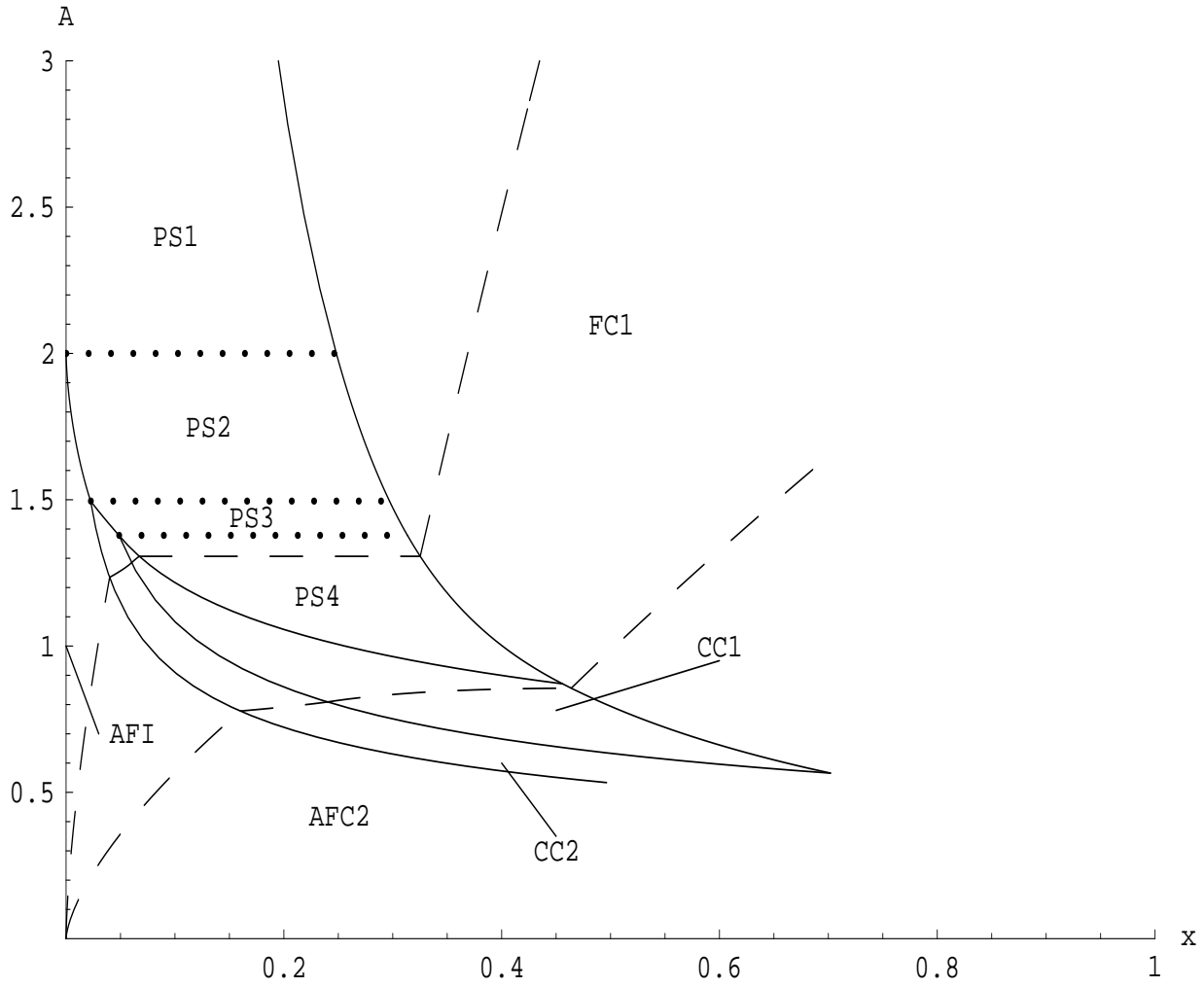


Figure 1: Phase diagram in the  $(x, A)$  plane.

- AFI**: Antiferromagnetic Insulator ( $x = 0$ )
- AFC2**: Antiferromagnetic Conductor (2-bands)
- CC2**: Canted Conductor (2-bands)
- CC1**: Canted Conductor (1-band)
- FC1**: Ferromagnetic Conductor (1-band)
- PS $i$** : Phase Separation regions ( $i = 1, 2, 3, 4$ )

# Effective Lagrangian for Spin Waves

In magnetically ordered systems the spontaneous symmetry breaking of the  $SU(2)$  group and the crystallographic group allows to write an effective lagrangian for the lowest lying excitation: spin waves [6].

$$\left. \begin{array}{c} \text{charge carriers} \\ + \\ \text{background magnetization} \end{array} \right\} \longrightarrow \left\{ \begin{array}{c} \text{doping dependence} \\ \text{of the} \\ \text{dispersion relation} \end{array} \right.$$

**Canted Spin Waves:**  $(0 < \theta < \pi)$

$$\begin{aligned} \mathcal{L}(x) &= \pi^- i \partial_t \pi^+ - \frac{1}{2m'} \partial_i \pi^- \partial_i \pi^+ && \longleftarrow \text{One branch} \\ &+ \frac{1}{2} \partial_t \pi^3 \partial_t \pi^3 - \frac{v^2}{2} \partial_i \pi^3 \partial_i \pi^3 && \longleftarrow \text{One branch} \end{aligned}$$

$$\frac{1}{2m'} \sim \frac{1}{(2M + x)y} \sqrt{2 - 4y^2 + \frac{5A}{2} \left( \frac{6\pi^2 x}{z^{3/2}} \right)} \times \sqrt{\frac{5A}{2} \left( \frac{6\pi^2 x}{z^{3/2}} \right) - 2y^2 - 2 \Pi_{+-} (1 - y^2)}$$

$$v^2 \sim \frac{1}{M^2} \left[ 2 + \frac{5A}{2} \left( \frac{6\pi^2 x}{z^{3/2}} \right) \right] (1 - y^2)$$

$y$  and  $\Pi_{+-}$  depend on the doping  $x$ .

## Ferromagnetic Spin Waves: $(\theta = 0)$

$$\mathcal{L}(x) = \pi^- i \partial_t \pi^+ - \frac{1}{2m'} \partial_i \pi^- \partial_i \pi^+ \quad \longleftarrow \quad \text{One branch}$$

$$\frac{1}{2m'} \sim \frac{1}{(2M + x)} \left[ -2 + \frac{5A}{2} \left( \frac{6\pi^2 x}{z^{3/2}} \right) \right]$$

## Antiferromagnetic Spin Waves: $(\theta = \pi)$

$$\mathcal{L}(x) = \partial_t \pi^- \partial_t \pi^+ - v^2 \partial_i \pi^- \partial_i \pi^+ \quad \longleftarrow \quad \text{Two branches}$$

$$v^2 \sim \frac{1}{M^2} \left[ 2 + \frac{5A}{2} \left( \frac{6\pi^2 x}{z^{3/2}} \right) \right]$$

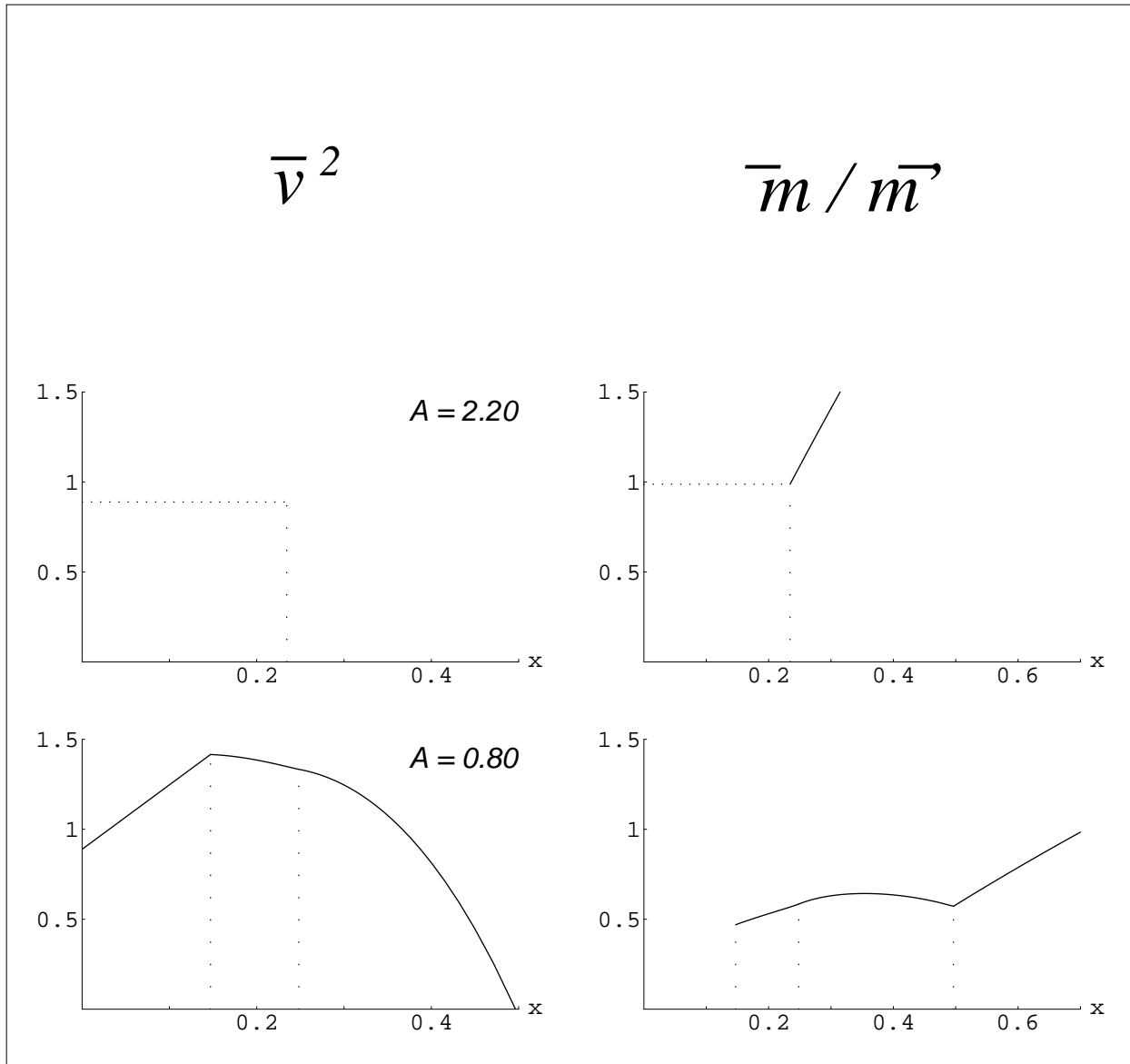


Figure 2: **Doping dependence** of the velocities and the masses.

$A = 2.20 \longrightarrow F - AF$  phase separation region

$A = 0.80 \longrightarrow$  canted phase

Horizontal dotted lines correspond to the phase separation regions, and the vertical dotted lines correspond to the phase transitions.



# Canted Phases vs. Phase Separation

## Phase Separation Region guess:

- Two macroscopic  $F$  and  $AF$  domains.
- The interphase does not modify qualitatively the properties.

Under such circumstances:

PS: One  $F$  and two  $AF$  spin wave branches

C: One  $F$  and one  $AF$  spin wave branches (see [7])

PS: Splitting of the  $AF$  branches by a magnetic field

C: No modification of the  $AF$  branch by a magnetic field

Different behavior of the mass and velocity with the doping

PS: Phase Separation: in fig. 2  $A = 2.20$

C: Canted Phase: in fig. 2  $A = 0.80$

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