

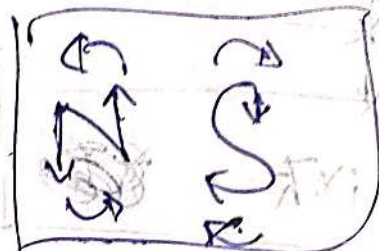
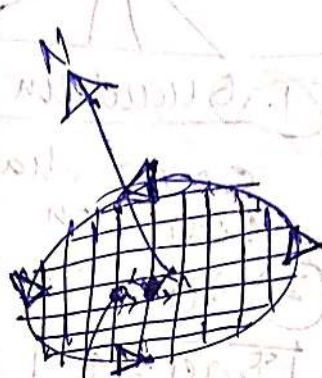
Magnetism & SC materials

In Natn.

Magnetic moment

$$\mu = \int d\mu = I \int ds$$

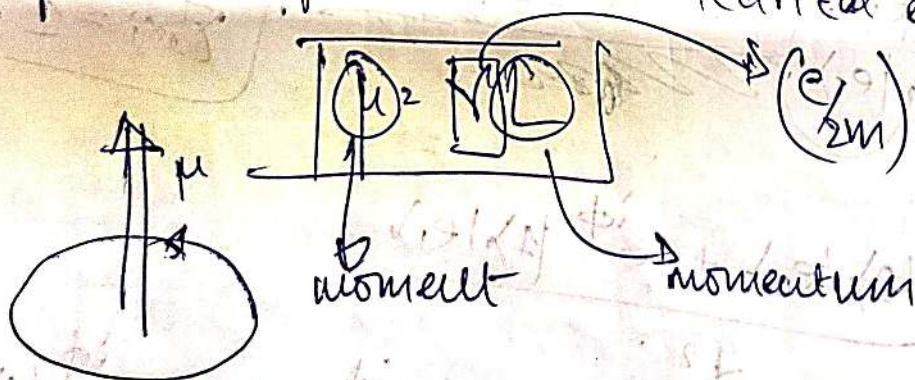
$$= I A \quad (A \cdot m^2)$$



Net current around boundary.

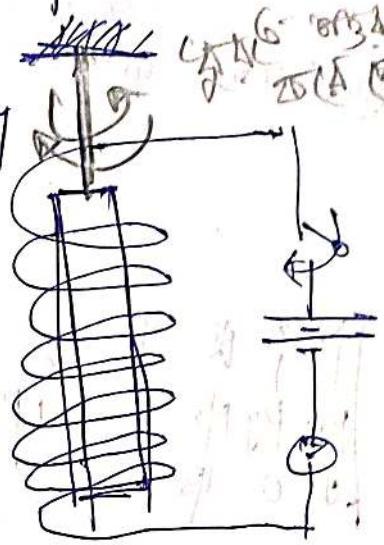
Magnetic moment & gyromagnetic ratio & momentum

A current carrying loop occurs as a consequence of motion of one or more electrical charges.



Einstein - de Haas effect.

- ① A (fm) [ferromagnetic] rod is suspended from a thin fibre
- ② A coil provides a magnetic field which magnetizes the fm vertically along its length.
- ③ to conserve net angular momentum a rotation is produced due to net magnetization.
- ④ frequency of rotation direction flipping = frequency of change in current flow direction.



Demo of relativity w/ μ & L .

Precession

Potential Energy of $\vec{\mu}$ in \vec{B}
 $E = \vec{\mu} \cdot \vec{B}$

$$\begin{aligned} \mu_z &= \mu \cos \theta \\ \mu_y &= \mu \sin \theta \\ \mu_x &= 0 \end{aligned}$$

U is min when $\mu \parallel B$.

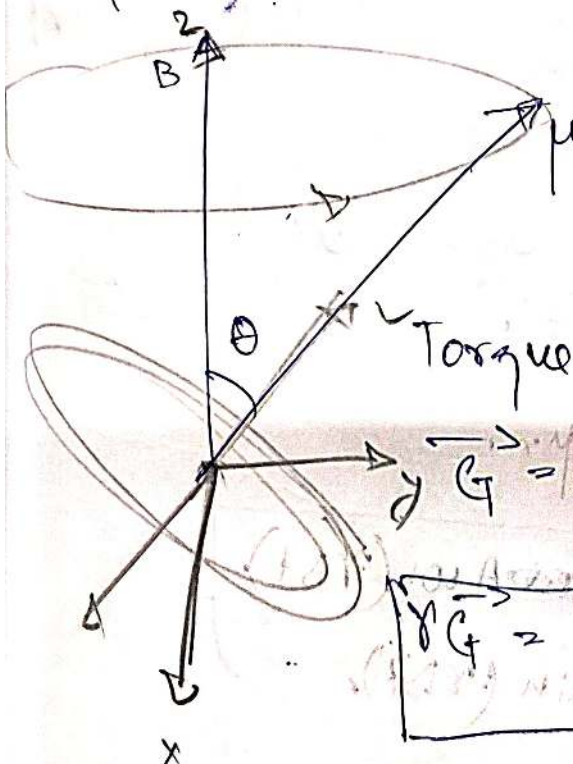
Torque in moment due to \vec{B} .

$$\vec{\tau} = \vec{\mu} \times \vec{B} = \frac{d\vec{L}}{dt}$$

$$\vec{\tau} = \frac{d\vec{\mu}}{dt}$$

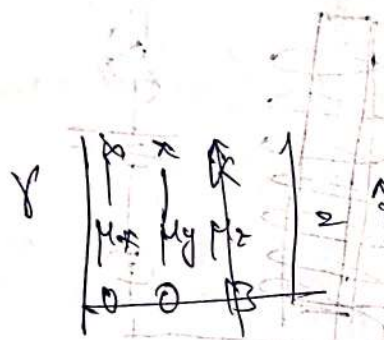
$$(\vec{\mu} \times \vec{B}) = \frac{d\vec{\mu}}{dt}$$

Rate of change of μ is \perp to μ .
 Therefore, circular motion.



The magnetic moment $\vec{\mu}$ in a magnetic field precesses around the field @ the Larmor frequency (ω_L) and a cone of semi angle θ .

consider when



$$\vec{\mu} = \mu_x \hat{i} + \mu_y \hat{j} + \mu_z \hat{k}$$

$$\gamma B (\mu_y \hat{i} - \mu_x \hat{j}) = \frac{d\mu_x}{dt} \hat{i} + \frac{d\mu_y}{dt} \hat{j} + \frac{d\mu_z}{dt} \hat{k}$$

$$\begin{cases} \dot{\mu}_x = \gamma B \mu_y \\ \dot{\mu}_y = -\gamma B \mu_x \\ \dot{\mu}_z = 0 \end{cases}$$

Solve.

$$\mu_z = \text{const}$$

$$\mu_x = A e^{i\omega t} + B e^{-i\omega t}$$

$$\omega = \gamma B \quad \text{at } t=0 \quad A+B=0$$

$$\mu_y = C e^{i\omega t} + D e^{-i\omega t}$$

$$\mu_y = \mu \sin \theta \quad \text{at } t=0$$

$$\ddot{\mu}_x = \gamma B \dot{\mu}_y$$

$$\mu_y = \mu \sin \theta \cos(\gamma B t)$$

$$\ddot{\mu}_x = -(\gamma B)^2 \mu_x$$

$$\Rightarrow \ddot{\mu}_x + (\gamma B)^2 \mu_x = 0$$

$$\frac{d}{dt} \left(\frac{d\mu_x}{dt} \right) + (\gamma B)^2 \mu_x = 0 \Rightarrow \mu_x = A \cos(\gamma B t)$$

$$\mu_x = \mu \sin \theta \sin(\gamma B t) \quad \mu_y = B \sin(\gamma B t)$$

~~10/9/25~~

$$\mu_z = |\mu| \cos \theta$$

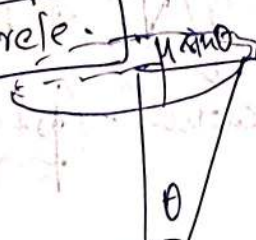
$$\mu_x = |\mu| \sin \theta \cos(\gamma B t)$$

$$\mu_y = |\mu| \sin \theta \sin(\gamma B t)$$

Dynamics of electron in B mag field & non 0 magnetic moment.

$$\mu_x^2 + \mu_y^2 = \mu^2 \sin^2 \theta$$

circle



$\omega_L = \gamma B$
Larmor frequency.

$$|\vec{\mu}| = -\frac{ev}{2\pi r} \cdot \pi r^2 = -\frac{evr}{2}$$

Magnetic moment

$$L = m_e v r = m_e \cdot v r$$

$$v = \frac{L}{m_e r}$$

$$I = -\frac{e}{2\pi r} \cdot \frac{L}{m_e r} = \left(\frac{-eL}{2\pi r^2 m_e} \right)$$

$$|\mu| = -\frac{e}{2} \cdot \frac{L}{m_e r} = \left(\frac{-e}{2m} \right) L$$

$\mu = \gamma L$

$1 \text{ BM} = 9.274 \times 10^{-24} \text{ A m}^2$

$$L = n\hbar$$

$$\begin{aligned} \mu &= \gamma \cdot n\hbar \\ &= n(\gamma\hbar) \\ &= n \cdot \left(\frac{e\hbar}{2m} \right) \end{aligned}$$

Momentum in a magnetic field (canonical momentum)

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

Lorentz force

$$\vec{F} = m \frac{d\vec{v}}{dt}$$

Electric field = $-\vec{\nabla}V = -\frac{\partial \vec{A}}{\partial t}$

Mag field = $\vec{\nabla} \times \vec{A}$

$$m \frac{d\vec{v}}{dt} = q \left[-\vec{\nabla}V - \frac{\partial \vec{A}}{\partial t} + \vec{v} \times \vec{\nabla} \times \vec{A} \right]$$

$$\mu = n \mu_B$$

$$\mu_B = \frac{e\hbar}{2m}$$

"Bohr Magnetron"

Size of Atomic magnetic moment.

Let us consider a hydrogen atom, (e, m_e)
electron.
circulating circularly around the nucleus with orbit radius (r).
How much magnetic moment?

$$|\vec{\mu}| = I \cdot A$$

$$T = \frac{2\pi r}{v} \quad f = \frac{v}{2\pi r}$$

$$I = qf = -\frac{ev}{2\pi r}$$

$$A = \pi r^2$$

with speed (v)

$$m \frac{d\vec{v}}{dt} = q \left[-\nabla V + \nabla(\vec{v} \cdot \vec{A}) \right] - q \left[(\vec{v} \cdot \nabla) \vec{A} + \frac{\partial \vec{A}}{\partial t} \right]$$

electric potential
velocity

$$\begin{aligned} & \left(\frac{\partial}{\partial x} \frac{\partial}{\partial x} A_x \right) + \left(\frac{\partial}{\partial y} \frac{\partial}{\partial y} A_y \right) + \left(\frac{\partial}{\partial z} \frac{\partial}{\partial z} A_z \right) + \frac{\partial}{\partial t} A_t \\ & \Rightarrow \frac{\partial}{\partial t} A_x + \frac{\partial}{\partial t} A_y + \frac{\partial}{\partial t} A_z + \frac{\partial}{\partial t} A_t \\ & \Rightarrow \frac{\partial}{\partial t} [A_x + A_y + A_z + A_t] \end{aligned}$$

$$\Rightarrow \frac{d}{dt} [A(x, y, z, t)] \quad \rightarrow \text{total derivative [convective derivative]}$$

New momentum " μ "

$$\mu = p + qA$$

effective P.E.

$$PE = q(V - \vec{v} \cdot \vec{A})$$

$$\text{Momentum} = (p + qA)$$

$$\mu = -i\hbar \nabla \cdot (\nabla \psi)$$

$$\frac{d}{dt} (mv) + \frac{d}{dt} (qA) = -q \left[\nabla V - \nabla(\vec{v} \cdot \vec{A}) \right]$$

$$\frac{d}{dt} (mv + qA) = -\nabla [qV - q\vec{v} \cdot \vec{A}]$$

canonical momentum

$$KE = \frac{mv^2}{2} = \frac{p^2}{2m}$$

$$\approx \frac{1}{2m} (\mu - eA)^2$$

mechanical KE.

BOHR - Van Leeuwen theorem.

In a classical system, @ thermal equilibrium, there is NO magnetization possible.

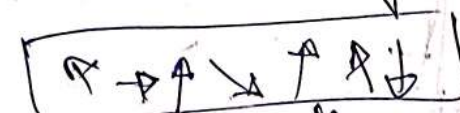
Thus the magnetization must be zero in a classical system. Magnetism is an entirely QUANTUM phenomena & cannot be ~~explained~~ to classically.

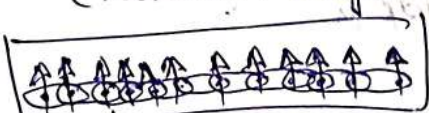
several types of magnetic order.

- (a) diamagnetism.
- (b) paramagnetism.
- (c) ferrimagnetism.
- (c') ferromagnetism.
- (d') antiferromagnetism.
- (d) helimagnetic.

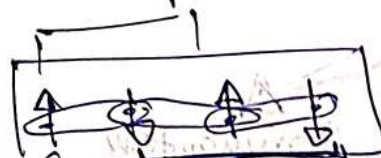
BC is a perfect diamagnet.

Dia \Rightarrow All materials (rarely) diamagnet.

Para:  $\chi = \frac{M}{H}$ (dimensionless No unit).

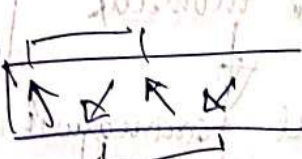
Ferro:  "collective phenomena" $\chi \gg 1$.

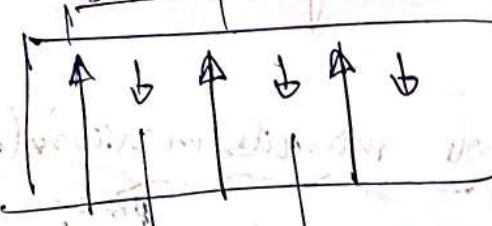
$H=0$.
interacting (in absence of field)
aligned correctly. "spontaneous magnetization"
 $M_r = 1 + \chi$
 $\langle M \rangle \neq 0$

AF \rightarrow 
"collective" pheno
"interacting" sub lattice

$H=0$. $M=0$

$M_A - M_B \approx 0$

(canted) 

Ferri \rightarrow 
sub lattice.

$M_A - M_B \gg 0$

$M_A - M_B \neq 0$

For any phase transition, there should be some order parameter.

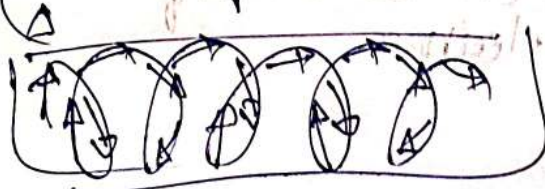
Ferri \rightarrow Ferro

second order

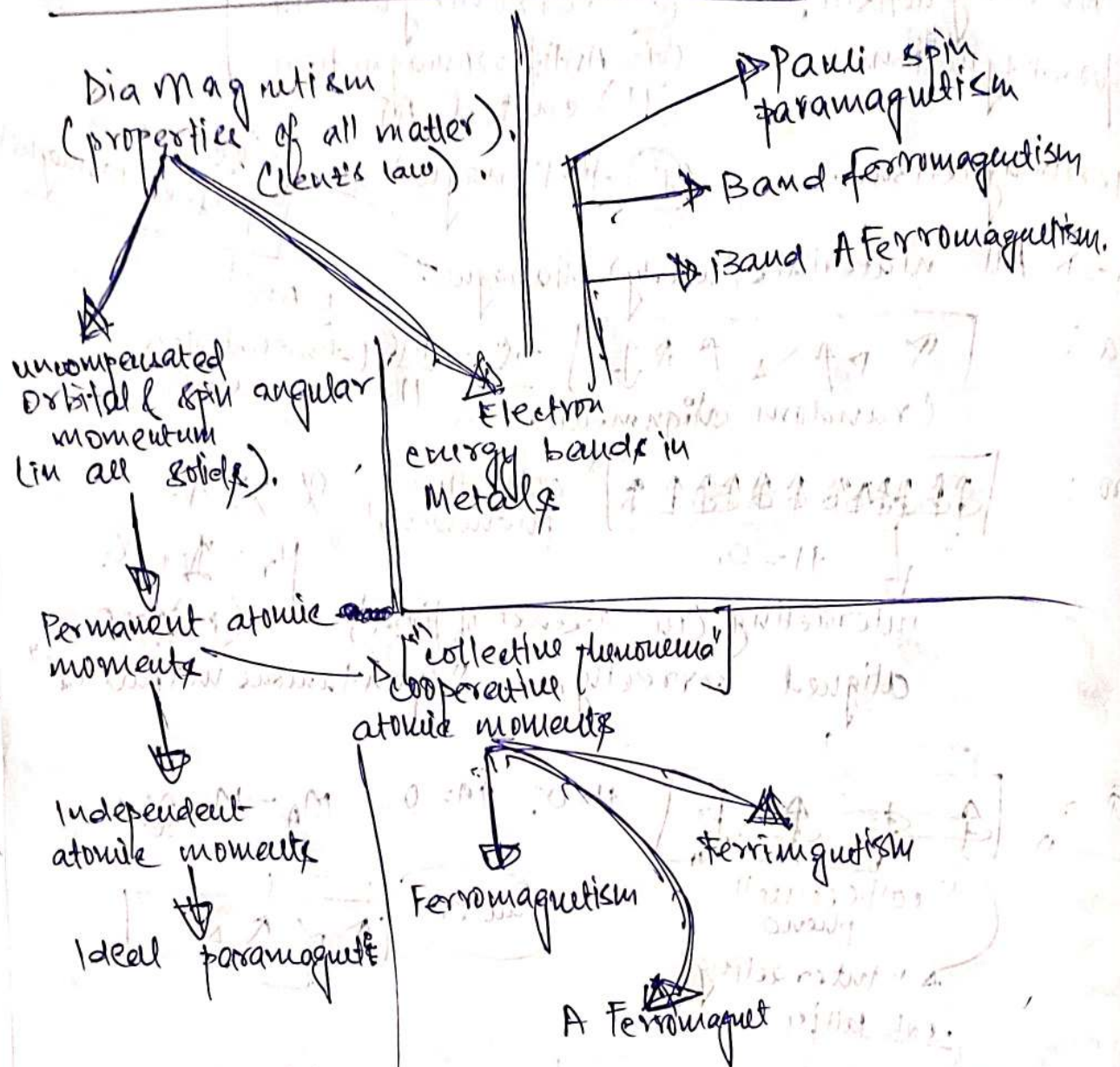
PT (phase transition)

even if ext mag field = 0.

Hel: spins vectors align themselves according to a helix, "helical pattern".



FAMILY TREE OF MAGNETISM.



ORIGIN of permanent dipole moments in materials.

An atom is an elementary magnet.

(*) In atom we find the following 3 contribution:

(i) Orbital motion of an electron contributes OAM. $\mu = \gamma L$

$$L = \vec{r} \times \vec{p} = m \vec{r} \times \vec{v}$$

$$L^2 |\psi\rangle = \hbar^2 l(l+1) |\psi\rangle$$

$$L_z |\psi\rangle = \hbar m_l |\psi\rangle$$

Dipole moment due to the orbital motion of the electron.



Dipole moment $\Rightarrow \mu = \gamma L$. \rightarrow Bohr magneton.

$\Rightarrow \mu_L = \left(\frac{-e}{2m} \hbar \right) \sqrt{l(l+1)}$. (Orbital Mag moment)

2) Dipole moment due to spin of electron.
 $\Rightarrow \mu_S \rightarrow$ similar treatment

$\mu_S = \left(\frac{e\hbar}{2m} \right) \sqrt{s(s+1)}$.

$s^2 |\chi_s\rangle = \hbar^2 s(s+1) |\chi_s\rangle$

$S_z |\chi_s\rangle = \hbar s |\chi_s\rangle$.

$m_s = \pm \frac{1}{2}$ (spin half).

$S_z |\chi_s\rangle = \pm \frac{\hbar}{2} |\chi_s\rangle$.

Bohr magneton

Landé g factor

$\frac{m_H}{m_e} = 1836$

3) Dipole moment due to spin of nucleus.

$\mu_N = \frac{e\hbar}{2m_N} \sqrt{I(I+1)}$.

Big mass of nucleus.

(neglected)

Prove that (Lie algebra). $[H, \omega]$.

(a) $[L^2, L_i] = 0$ (b) $[L_i, L_j] = i\hbar \epsilon_{ijk} L_k$.

(c) $L_x = -i\hbar \frac{\partial}{\partial y} = i\hbar \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$.

(d) $S^2 = S_x^2 + S_y^2 + S_z^2$
 $= \frac{1}{2} (S_+ S_- + S_- S_+) + S_z^2$.

(f) $[S_+, S_-] = 2 S_z$.

(e) $[S^2, S_i] = 0$.