

6/2/25

Magnetic

SPDF GHI

(Term)

(2L+1)

symbol =



(2S+1)

L_J

$$(2L+1)(2S+1) = \sum_{J=L-S}^{L+S} (2J+1)$$

$$J = 3/2, 5/2, 7/2, 9/2$$

Which one is true for g_s of particular ion?

⇒ Hund's rule.

$$(2 \cdot 3/2 + 1) + (2 \cdot 5/2 + 1) + (2 \cdot 7/2 + 1) + (2 \cdot 9/2 + 1)$$

$$= 4 + 6 + 8 + 10 = 28$$

1. Maximize S
2. Minimize L
3. $J = L - S$, less than half filled
 $J = L + S$, more than half filled.

$$g_i = 1 + \frac{S(S+1) + L(L+1) - J(J+1)}{2J(J+1)}$$

$$\mu_i = \mu_{eff} / \mu_B$$

3. D transition metal ions. (Magnetic g_s for 3D ions using Hund's rule)

(Mg → d shell)

n	d shell	ion	shell	S	L	J	term
1	↓	Ti ³⁺ , V ⁴⁺	3d ¹	1/2	2	3/2	² D _{3/2}

n	d shell	ion (ex)	shell	S	L	J	Term	$\mu_i (g_i \sqrt{J(J+1)})$	μ_{exp}	$\mu_i (2S+1)$
1	↓	Ti ³⁺ , V ⁴⁺	3d ¹	1/2	2	3/2	² D _{3/2}	1.55	1.7	1.73
2	↓ ↓	V ³⁺	3d ²	1	3	2	³ F ₂	1.63	2.1	2.83
3	↓ ↓ ↓	Cr ³⁺ , V ³⁺	3d ³	3/2	3	3/2	⁴ F _{3/2}	1.73	2.85	3.87
4	↓ ↓ ↓ ↓	Mn ³⁺ , Cr ²⁺	3d ⁴	2	2	0	⁵ D ₀	0	4.82	4.9
5	↓ ↓ ↓ ↓ ↓ (Half filled)	Fe ³⁺ , Mn ²⁺	3d ⁵	5/2	0	5/2	⁶ S _{5/2}	5.92	5.82	5.92

6	$\uparrow\downarrow\downarrow\downarrow\downarrow$	Fe^{2+}	$3d^6$	2	2	4	$5D_4$	6.7	5.36	4.98
7	$\uparrow\downarrow\downarrow\downarrow\downarrow$	Co^{2+}	$3d^7$	$3/2$	3	$9/2$	$4F_{9/2}$	6.3	4.9	3.87
8	$\uparrow\uparrow\downarrow\downarrow\downarrow$	Ni^{2+}	$3d^8$	1	3	4	$3F_4$	5.9	3.12	2.83
9	$\uparrow\uparrow\uparrow\downarrow\downarrow$	Cu^{2+}	$3d^9$	$1/2$	2	$5/2$	$2D_{5/2}$	3.55	<u>1.83</u>	11.72
10	$\uparrow\uparrow\uparrow\uparrow\downarrow$	Zn^{2+}	$3d^{10}$ (perfectly non magnetic)	0	0	0	$1S_0$	0	0	0

Reason for p_1 & p_{exp} mismatch.

Due to crystal field effect there is ~~quenching~~ quenching of orbital C momentum (L). Thus, it matches p_2 but not p_1 .

No p_2 since no crystal field eff

Magnetic ground state for 4f ions using Hund's rule

4f	orientation	ion	shell	S	L	J	Term	p_1	p_{exp}	p_2 (not reqd)
1	\downarrow	Ce^{3+}	$4f^1$	$1/2$	3	$5/2$	$2F_{5/2}$	8.54	8.51	
2	$\downarrow\downarrow$	Pr^{3+}	$4f^2$	1	5	4	$3H_4$	3.58	3.56	
3	$\downarrow\downarrow\downarrow$	Nd^{3+}	$4f^3$	$3/2$	6	$9/2$	$4I_{9/2}$	3.62	3.7	
4	$\downarrow\downarrow\downarrow\downarrow$	Pm^{3+}	$4f^4$	2	6	4	$6I_4$	2.68	—	
5	$\downarrow\downarrow\downarrow\downarrow\downarrow$	Sm^{3+}	$4f^5$	$5/2$	5	$5/2$	$6H_{5/2}$	0.85	1.74	
6			$4f^6$	3	3	0	$7F_0$			
7			$4f^7$	$7/2$	0	2	$8S_{7/2}$			

crystal field eff is unimportant for lanthanide series (rare earth ions) because their partially filled 4f shell lies deep inside the ion (lanthanide contraction).

In contrast, the partially filled d shell of transition metal ions are outermost electronic shells more strongly influenced by their crystalline env. thus exhibiting strong crystal field eff i.e. quenching of orbital momentum takes place (i.e. $L \approx 0$).

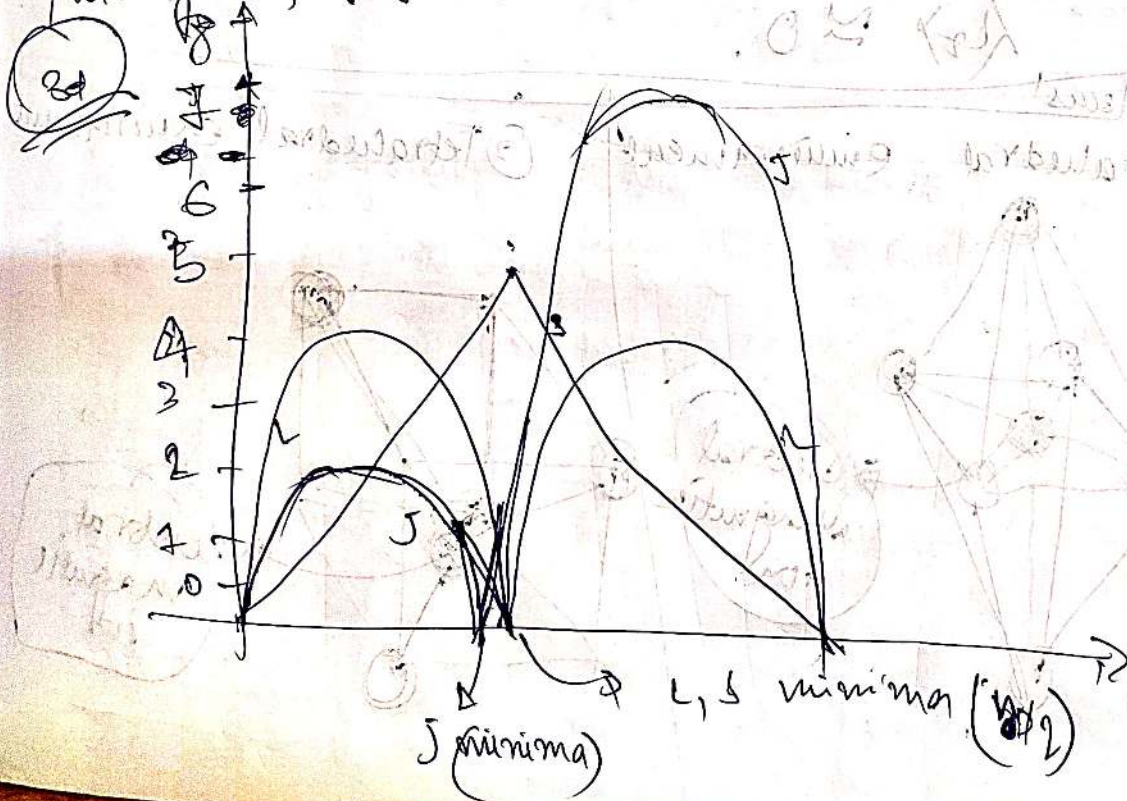
$$\therefore J = L, \quad p = p_2 \Rightarrow g_J \sqrt{J(J+1)} = 2 \sqrt{S(S+1)}$$

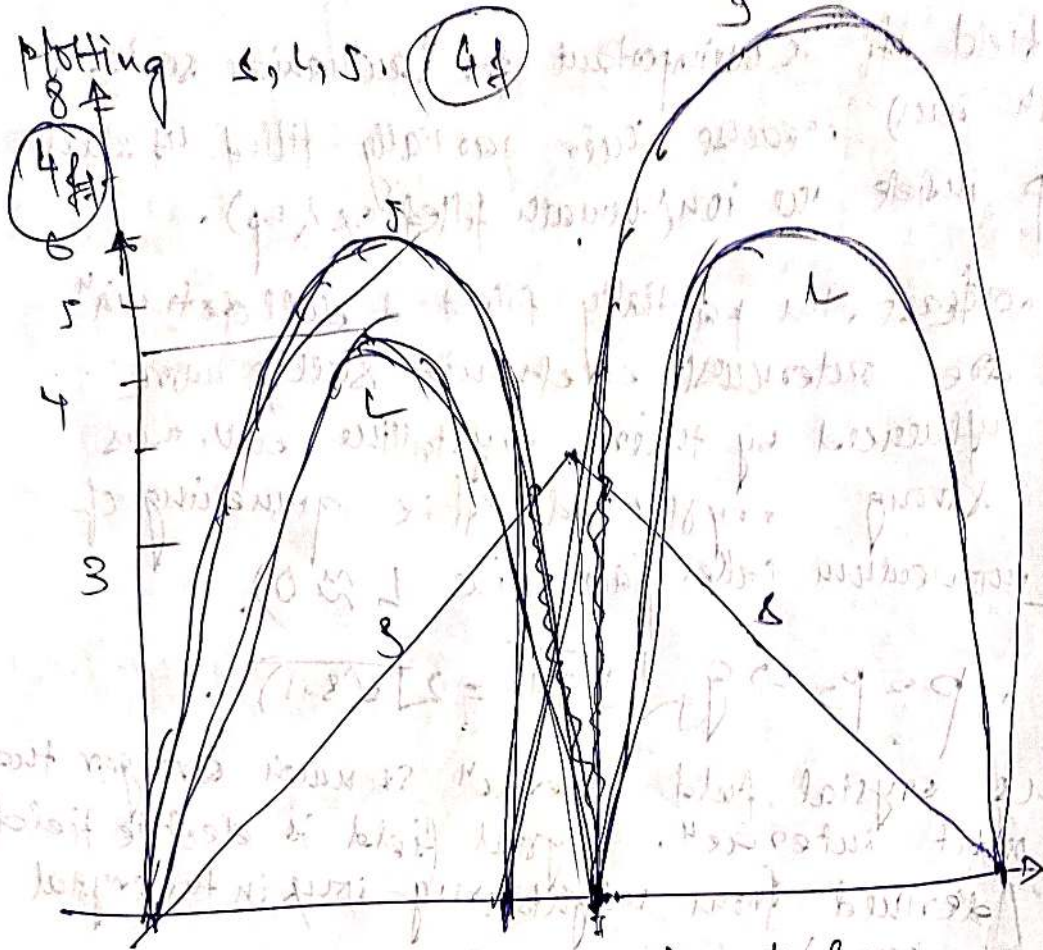
In 3d ions crystal field interaction is much stronger than the spin orbit interaction. crystal field is electric field originating/derived from neighbouring ions in the crystal (ligand field theory).

Eff of d orbitals on central ion & their overlapping with s orbitals on surrounding ions.

The size & nature of crystal field depends on symmetry & env.

Plotting Δ , L & J var.





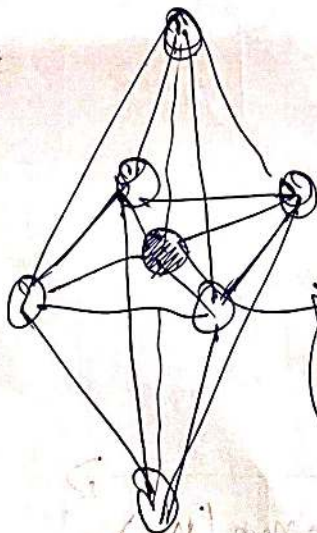
J (minima) diff from L & S minima.

A semiclassical interpretation of orbital L momentum quenching is that the orbital L momentum precesses in the crystal field so that its magnitude is unquenched but its components all average to zero.

$$\langle L_x \rangle \approx 0$$

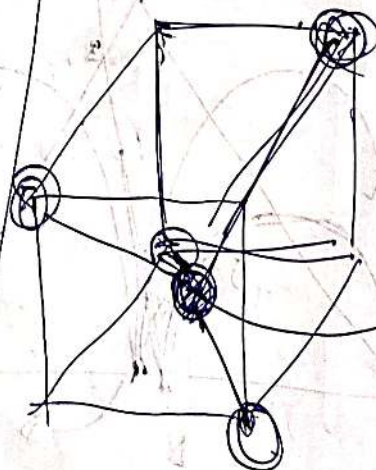
3d systems

① Octahedral environment



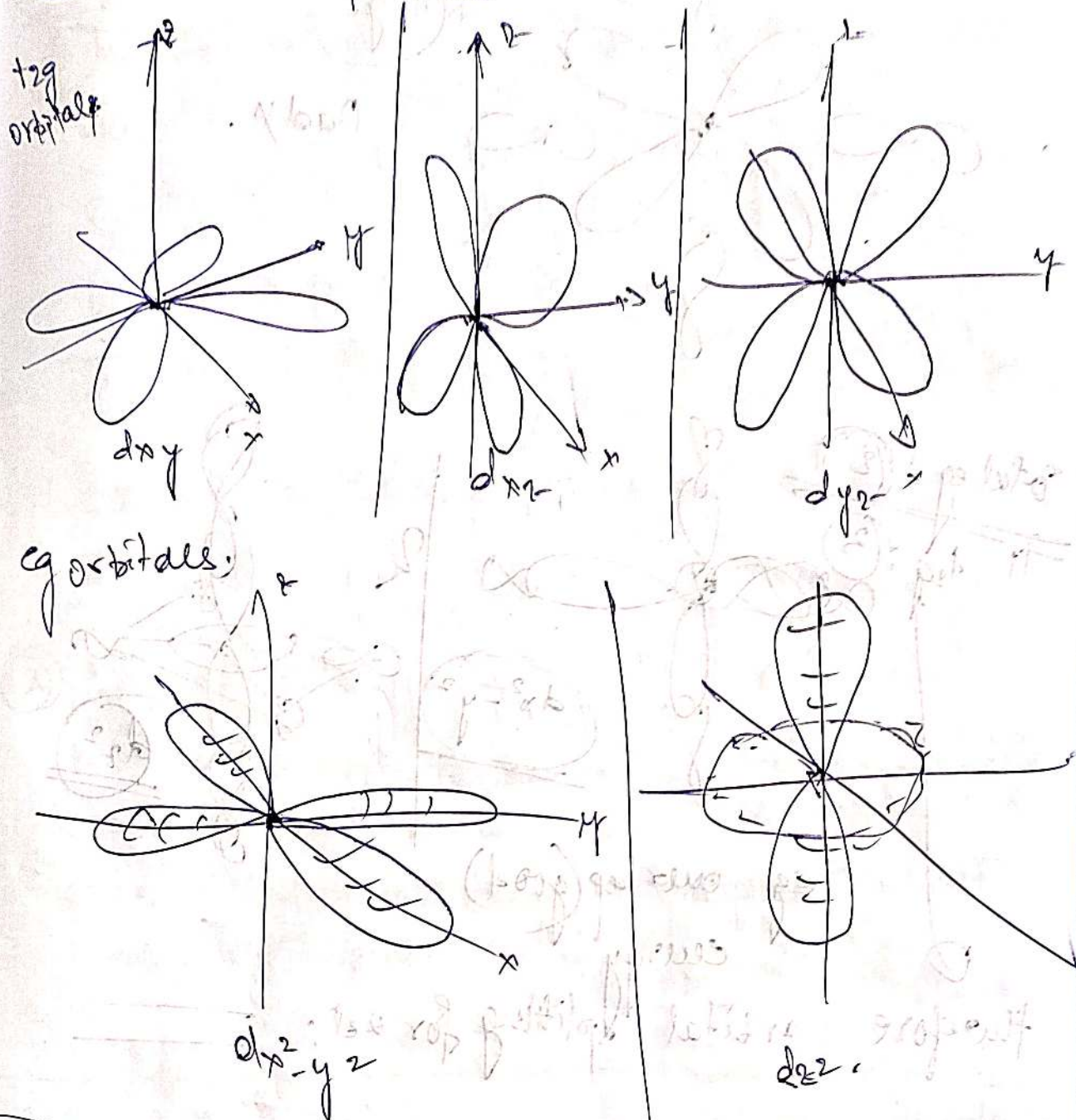
Central magnetic ion

② Tetrahedral environment



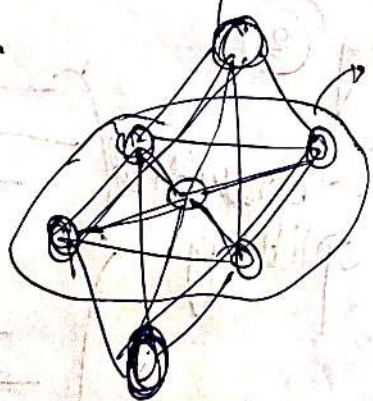
Central magnetic ion

3d orbital shapes:



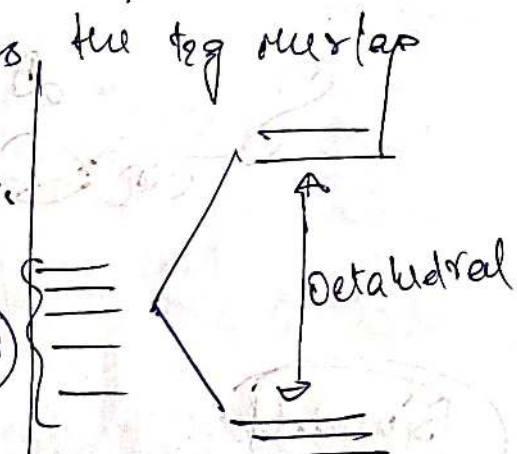
In 3d systems (oct & tet), there is an influence of orbital overlap & thus field strength.

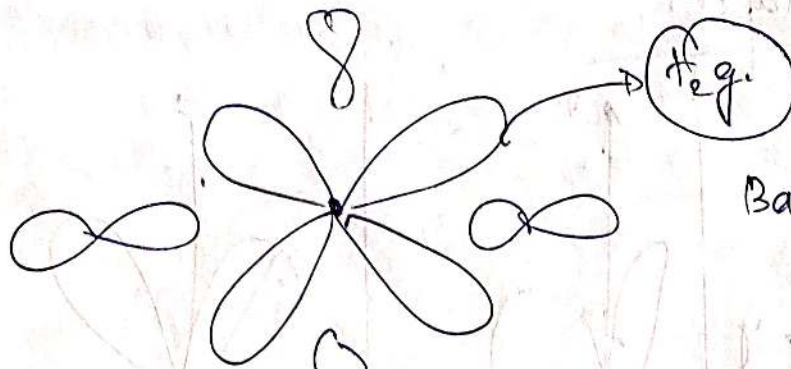
eg orbital overlap & thus field strength.



close overlap.

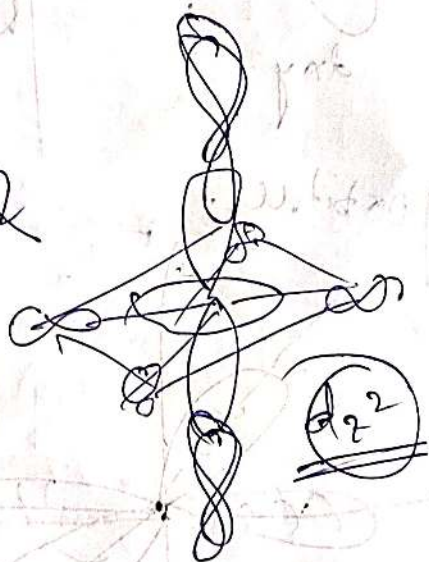
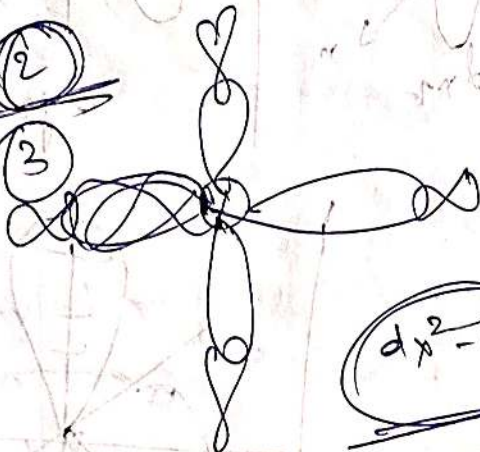
3d





Badly.

Total eg: (2)
 11 d_{xy} (3)



eg overlap (good)
 energy

Therefore orbital splitting for oct: _____

For tetrahedral case:

Lg is diagonal interaction

Lg (3)

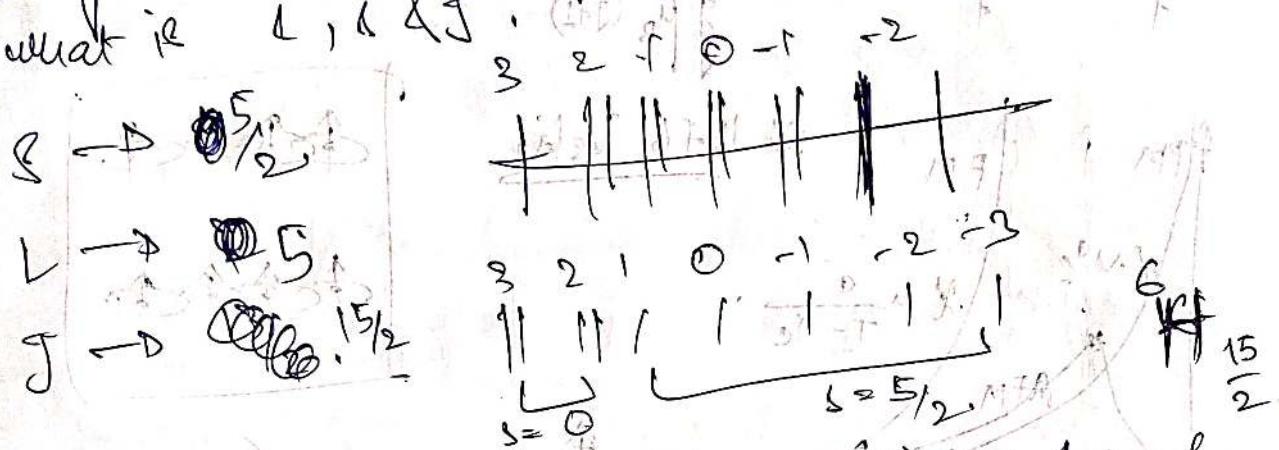
\therefore tetrahedral splitting

Blundell

Problem.

S P D F G H I
0 1 2 3 4 5

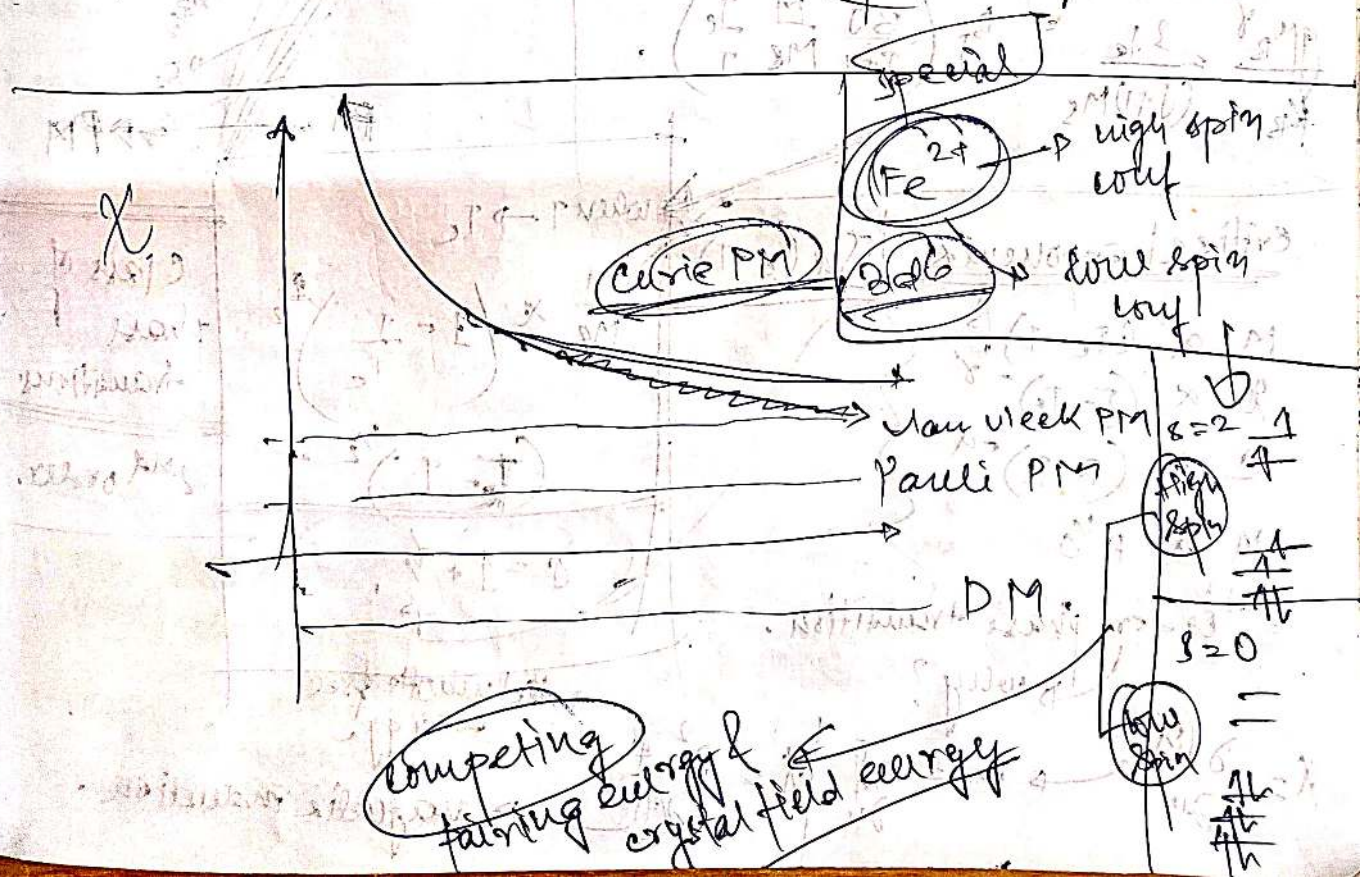
Magnetic moment of Py^{3+} ion has 9e in 4f shell
what is L, S & J.



calculate susceptibility containing 1 gram at 4K.

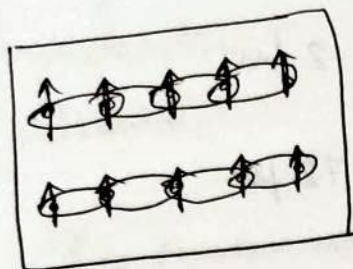
2 Paramag of $S=1$ system

find the magnetization as a funcⁿ of magnetic field & temp of a system with $S=1$, moment M & $\chi_{\text{one}} = \text{N atoms/unit volume}$. Show that the limit $\mu_B \gg k_B T$ leads to $M \approx (2 \mu_B N m^2 H) / 3 k_B T$.



7/2/25 / Friday

FERROMAGNETISM

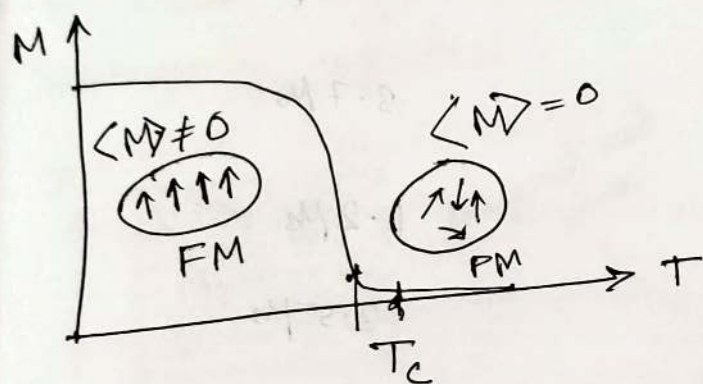


$$H=0$$

$$\langle M_{sp} \rangle \neq 0$$

Collective Cooperative phenomenon

Example \rightarrow Fe, Ni, Co, $\text{La}_{0.7}\text{Sr}_{0.3}\text{MnO}_3$, SrRuO_3



$\langle M_{sp} \rangle =$ order parameter

FM \rightarrow PM \rightarrow 2nd order phase transition

Order parameter changes drastically across the phase transition.

① Two distinct characteristics of FM \rightarrow

- ① The FM has spontaneous magnetization. $\langle M_{sp} \rangle$
- ② the existence of magnetic ordering temp. (T_c)

Material	$T_c(K)$	Spontaneous Magnetization (μ_B/atom)
Fe	1043	2.2 μ_B
Co	1394	1.72 μ_B
Ni	631	0.61 μ_B
Gd	289	7.5 μ_B
La Sr Mn O_3 0.7 0.3	360	3.7 μ_B
La Sr Ru O_3	160	1.2 μ_B
	587	3.5 μ_B
Mn Sb		

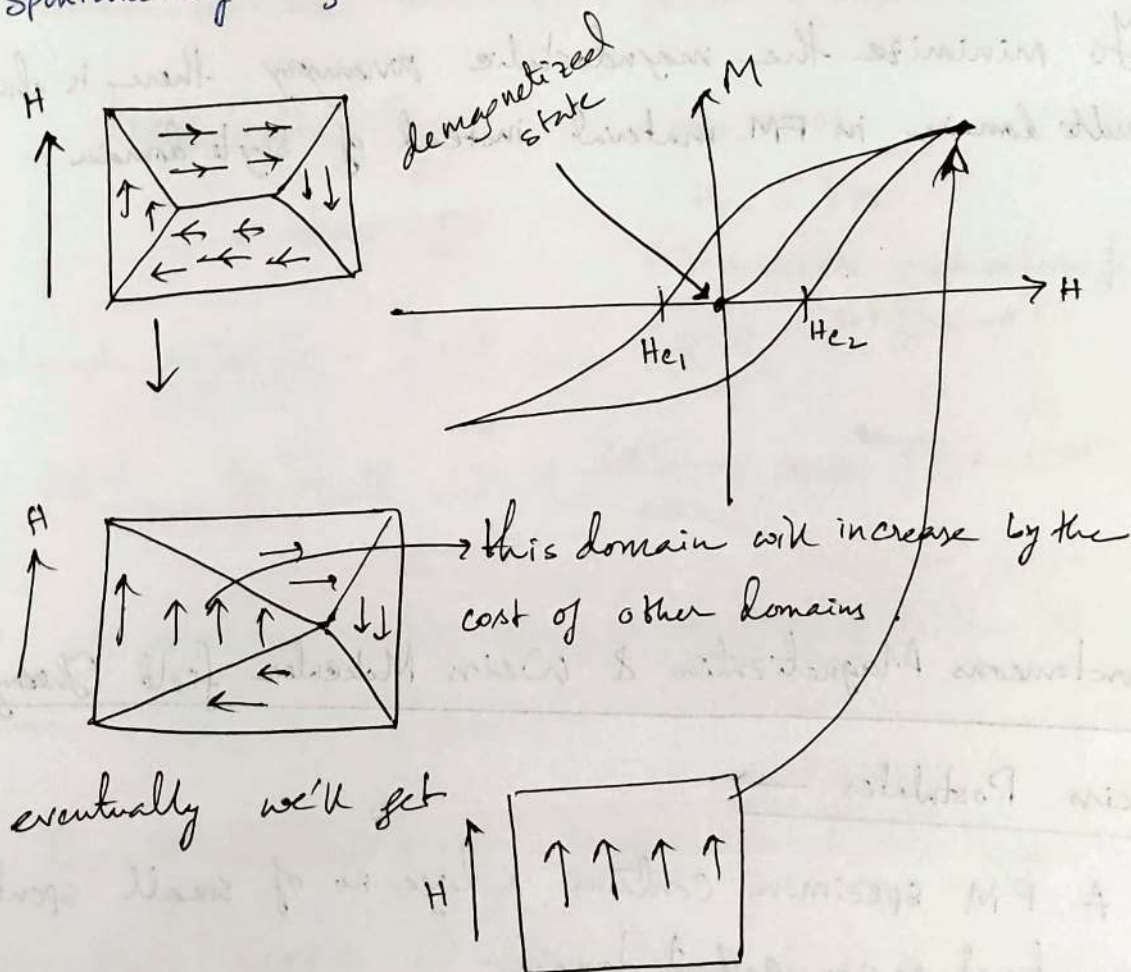
* $\mu_B H_m \approx k_B T_c$

$H_m \approx$ internal molecular field that keeps the spin aligned
each other in Ferromagnetism

$$\begin{aligned}
 H_m &\approx \frac{k_B T_c}{\mu_B} = \frac{0.026 \times 2 \times 1.6 \times 10^{-19} \times 10^7 \text{ eV}}{9.21 \times 10^{-21} \text{ eV/Oe}} \\
 &= 10^7 \text{ Oe} \\
 &= 500 \text{ T}
 \end{aligned}$$

② Properties :-

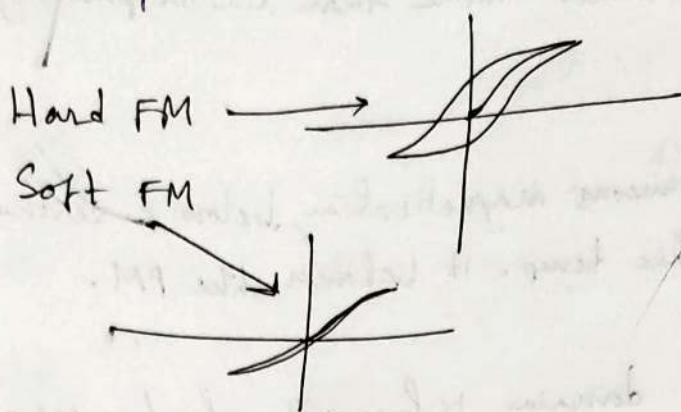
- ① They're very strongly magnetic for a given field. The magnetization is 10^6 times more than dia or para-magnetic substance.
- ② They possess spontaneous magnetization below a certain temp. (T_c). Above this temp. it behaves like PM.
- ③ In a bulk material ^{FM} domain is formed which are spontaneously magnetized.



Coercive field = H_c the amount of -ve field required to bring back the magnetization to zero.

$$H_{c1} \neq H_{c2} \text{ always} \quad H_c = \frac{|H_{c1}| + |H_{c2}|}{2}$$

we get the hysteresis loop and as well as we can take snapshots of the domains by using Lorentz microscope.



To minimize the magnetostatic energy there is a multi domain in FM material instead of single domain.

Spontaneous Magnetization & Weiss Molecular field Theory

Weiss Postulates →

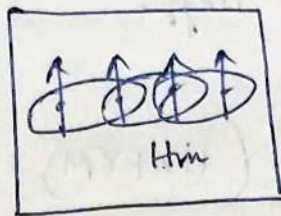
- ① A FM specimen contains a large no. of small spontaneously magnetized regions, called domains.
- ② Within each domain the spontaneous magnetization is due to the existence of an internal molecular field which tends to produce a parallel alignment of atomic dipoles.

$$H_{\text{eff}} = H_a + H_m$$

↓
applied
field

↓
internal molecular
field.

H_a



- ③ The strong interaction which tends to align atomic dipoles parallel in FM materials may be considered as equivalent to some internal molecular field. (H_m)

$$H_{\text{eff}} = H_a + H_m$$

$$= H_a + \gamma M$$

$$H_m \propto M$$

$$H_m = \gamma M$$

(γ = Weiss molecular field coefficient)

Internal molecular field $H_m = \gamma M$

$$\mu_B H_m = k_B T_c \Rightarrow H_m = \frac{k_B T_c}{\mu_B} \approx 10^7 \text{ Oe} \approx 1500 \text{ T}.$$

Let us consider a FM material containing N atoms per unit volume
 $M = N J g \mu_B B_J(x)$

$$B_J(x) = \text{Brillouine fn where } x = \frac{g \mu_B B}{k_B T} H_{\text{eff}}$$

H_m = keeping all the dipoles along H_a .

$$\chi = \frac{g \mu_B B J}{k_B T} H_{eff}$$

$$= \frac{g \mu_B B J}{k_B T} (H_a + \gamma M)$$

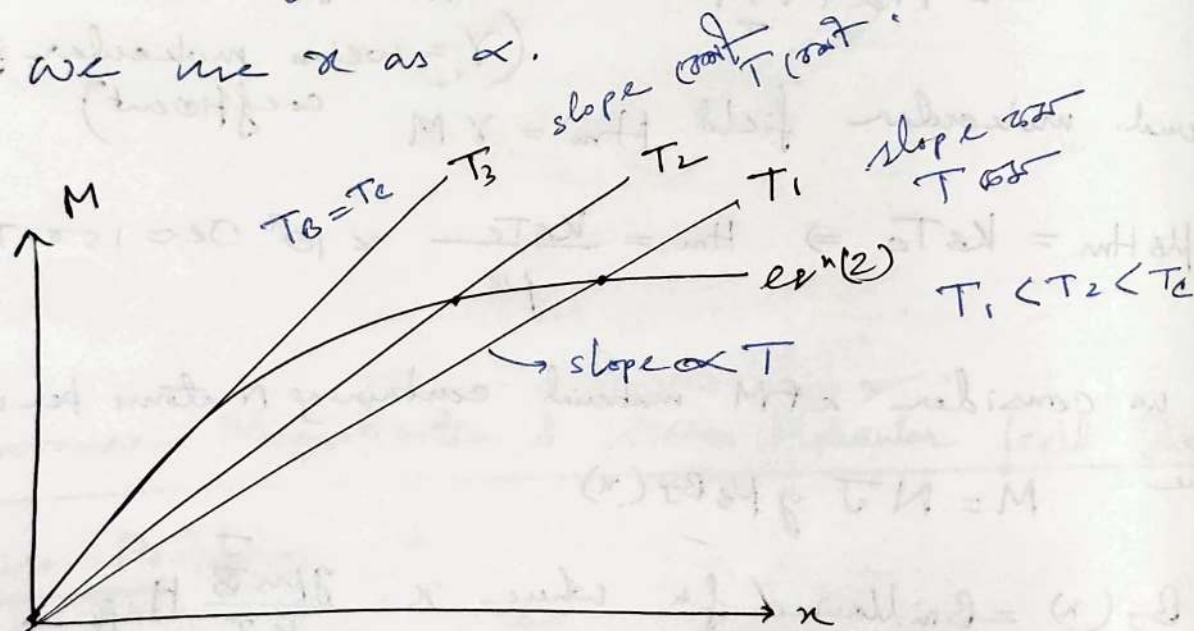
For spontaneous magnetization ($H_a = 0$)

$$\chi = \frac{g \mu_B J}{k_B T} \gamma M$$

$$M = \frac{\chi k_B T}{g \mu_B J \gamma} \quad \text{--- (1)}$$

$$M = N J g \mu_B B_J(\chi) \quad \text{--- (2)}$$

We use χ as x .



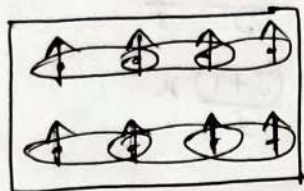
① if $T < T_c$ we get non zero value of M . so for $T < T_c$, There is a spontaneous magnetization.

② If $T = T_c$ the curve ① just the tangent to curve ②. At the origin & spontaneous magnetization just vanishes.

③ $T > T_c$, no solⁿ of spontaneous magnetization.

$T = T_c$ = Curie Temp. of Ferromagnetism.

② Relation Between γ and $T_c \rightarrow$



$$H_m = \gamma M = k_B T_c$$

The expected relation is that T_c increases with increasing γ .

$$\gamma \cdot B_J(x) = \frac{J+1}{J} \frac{x}{3} \quad (x \ll 1)$$

$$\text{eqn } ② \quad M = N g \mu_B J B_J(x)$$

$$= N g \mu_B J \frac{J+1}{J} \frac{x}{3} \quad (x \ll 1)$$

$$M = \frac{N g \mu_B (J+1)}{3} x$$

$$\text{eqn } ① \quad M = \frac{x k_B T}{g \mu_B J \gamma}$$

when $x \ll 1$ both the slopes should be equal from the previous figure.

$$\mu_{\text{eff}}^2 = g^2 \mu_B^2 J(J+1)$$

$$\boxed{\frac{N \mu_{\text{eff}}^2}{3 k_B} = C}$$

$$\frac{k_B T}{g \mu_B J \gamma} = \frac{N g \mu_B (J+1)}{3}$$

$$T_c = \frac{N g^2 \mu_B^2 J(J+1) \gamma}{3 k_B} = C \gamma$$

$$\boxed{T_c = C \gamma}$$

for a strong ferromagnetic material γ = internal field coefficient as well as T_c becomes high.

② Curie Weiss Law:-

$$M = N g \mu_B J B_J(x) = N g \mu_B J \frac{J+1}{J} \frac{x}{3}$$

$$= N g \mu_B \frac{(J+1)}{3} x$$

$$x = \frac{g \mu_B J}{kT} (H + \gamma M)$$

$$M = N g \frac{\mu_B (J+1)}{3} \frac{g \mu_B J}{kT} (H + \gamma M)$$

$$= \frac{N g^2 \mu_B^2 J (J+1)}{3 kT} (H + \gamma M)$$

$$M = \frac{C}{T} (H + \gamma M)$$

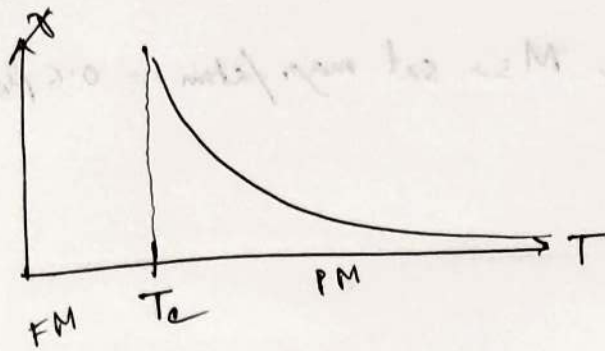
$$M \left(1 - \frac{\gamma C}{T} \right) = \frac{C}{T} H$$

$$\Rightarrow \frac{M}{H} = \chi = \frac{\frac{C}{T}}{1 - \frac{\gamma C}{T}} = \frac{\frac{C}{T}}{1 - \frac{T_c}{T}} = \frac{\frac{C}{T}}{\frac{T - T_c}{T}}$$

$$\Rightarrow \boxed{\chi = \frac{C}{T - T_c}}$$

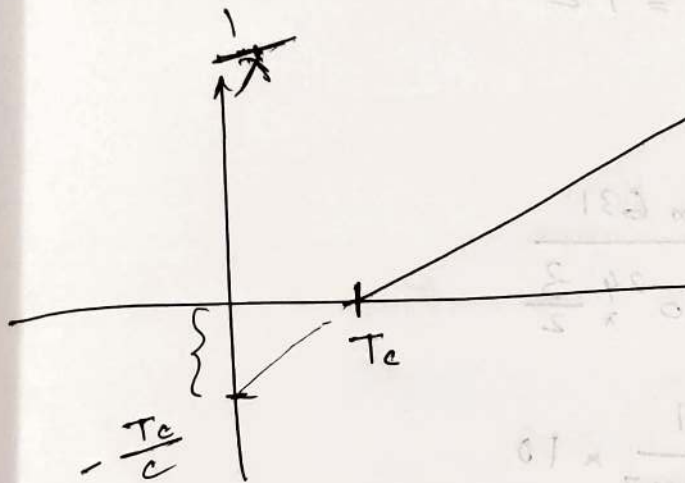
Curie-Weiss Law.

$$\Rightarrow \frac{1}{\chi} = \frac{T}{C} - \frac{T_c}{C}$$



$$X = \frac{c}{T - T_c} \quad \text{at } T = T_c \quad X = \infty$$

$$\frac{1}{X} = \frac{T}{c} - \frac{T_c}{c}$$



$$T_c = c \gamma = \frac{N g^2 J (J+1) \mu_B^2}{3 k_B \gamma}$$

$$= \frac{N g J \mu_B \gamma g (J+1) \mu_B}{3 k_B}$$

$$M = N g \mu_B J B_J(x)$$

$$x \gg 1$$

$$B_J(x) = 1$$

$$M = N g \mu_B J = M_S$$

$$T_c = \frac{M_S \gamma g (J+1) \mu_B}{3 k_B} = M_S \gamma \frac{g (J+1) \mu_B}{3 k_B}$$

$$A_m = B_{mf} M_S = \frac{3 k_B T_c}{g \mu_B (J+1)}$$

★★
(problem)

Problem → Calculate the molecular field of Ni

$$T_c = 631 \text{ K}, J = \frac{1}{2}, M_s = \text{sat mag./atom} = 0.6 \mu_B$$

$$H_m = \frac{3k_B T_c}{g \mu_B (J+1)}$$

$$M_s = N g \mu_B J = 0.6 \mu_B / \text{atom} \times N = 0.6 N \mu_B$$

$$\rightarrow N g \mu_B J = 0.6 N \mu_B$$

$$\Rightarrow g = \frac{0.6}{\frac{1}{2}} = 1.2$$

$$H_m = \frac{3 \times 1.38 \times 10^{-23} \times 631}{1.2 \times 9.27 \times 10^{-24} \times \frac{3}{2}}$$

$$= \frac{2 \times 2 \times 1.38 \times 631}{1.2 \times 9.27 \times 3} \times 10$$

$$= 156.5 \times 10 = 1565 \text{ T}$$