1. Counterpropagating waves, energy transport

Consider a superposition of two y-polarized plane harmonic electromagnetic waves with angular frequency  $\omega$  and vacuum wavenumber k. Assume that these waves propagate in opposite directions along the z-axis, through a homogeneous, nonmagnetic medium with refractive index n. This relates to an electric field of the form

$$E_y(z,t) = a_f e^{\mathbf{i}(\omega t - knz)} + a_b e^{\mathbf{i}(\omega t + knz)},$$

with complex amplitudes  $a_f$  and  $a_b$ . Evaluate the time averaged energy density and the time averaged power flux density S. Compare with the absolute squares  $|E|^2$ ,  $|H|^2$  of the electric and magnetic field.

## 2. Reflection of a plane wave at a surface of a potentially attenuating medium

Assume that the x-y-plane spans the interface between two regions (1), (2) filled with linear, nonmagnetic, and uncharged media. While region (1), z < 0, is filled with a lossless dielectric medium with (real) permittivity  $\epsilon_1 = n_1^2$  and refractive index  $n_1$ , for the half-space (2), z > 0, we assume a potentially complex permittivity  $\epsilon_2$ . There are no free charges or currents on the interface.

This is a 2-D problem; let the coordinates be oriented such that all electromagnetic fields are constant in the y-direction. We restrict things to s-polarized waves. Their propagation is governed by the scalar 2-D TE Helmholtz equation

$$(\partial_x^2 + \partial_z^2 + k^2 \epsilon) E_y = 0 \tag{1}$$

for the principal electric field component  $E_y(x,z)$ . Continuity of this field and of its normal derivative is required across all interfaces.  $k=2\pi/\lambda$  is the vacuum wavenumber, related to the vacuum wavelength  $\lambda$ . All fields oscillate in time  $\sim \exp(\mathrm{i}\omega t)$  with angular frequency  $\omega=kc$ .

Assume that a plane with complex amplitude  $E_{\rm I}$  propagates in region 1 towards the interface at an angle of incidence  $\theta$ . The interface causes a reflected wave in region (1) with amplitude  $E_{\rm R}$ . For region (2) we start with an ansatz of a separable field and an amplitude  $E_0$ , such that the electric field can be stated in the form

$$E_{y}(x,z) = \begin{cases} E_{I} e^{-ikn_{1}(z\cos\theta + x\sin\theta)} + E_{R} e^{-ikn_{1}(-z\cos\theta + x\sin\theta)} & \text{for } z < 0, \\ E_{0} X(x) Z(z) & \text{for } z > 0. \end{cases}$$
(2)

Here the functions X and Z are yet to be determined. Without loss of generality we choose amplitudes X(0) = 1 and Z(0) = 1.

- (a) Draw a sketch to clarify the geometry. Verify that the ansatz (2) satisfies Eq. (1) in region (1).
- (b) Use the continuity of  $E_y$  across the interface to determine X, and to relate  $E_0$  to  $E_I$  and  $E_R$ .
- (c)  $E_y$  needs to satisfy Eq. (1) in region (2). Show that Z is of the form

$$Z(z) = e^{-ik\kappa z}$$
, with  $\kappa^2 = \epsilon_2 - n_1^2 \sin^2 \theta$ , (3)

where  $\kappa$  is a potentially complex constant.

(d) Further the normal derivative  $\partial_z E_y$  needs to be continuous at the interface. Use this requirement to derive the equations

$$E_{R} = \frac{n_{1}\cos\theta - \kappa}{n_{1}\cos\theta + \kappa}E_{I}, \qquad E_{0} = \frac{2n_{1}\cos\theta}{n_{1}\cos\theta + \kappa}E_{I}$$
(4)

that relate the amplitudes of the reflected and "transmitted" waves to the amplitude of the incoming wave. Compare with the Fresnel-equations for s-polarized waves.

- (2., continued)
- (e) First consider the case of a lossless medium in region (2), with  $\epsilon_2 \in \mathbb{R}$ . Introduce the refractive index  $n_2$  for region (2). Distinguish between two cases with  $n_2^2 > n_1^2 \sin^2 \theta$  and  $n_2^2 < n_1^2 \sin^2 \theta$ .
  - i. For  $n_2^2 > n_1^2 \sin^2 \theta$ , select a sign for  $\kappa$  such that the field in region (2) is an outgoing travelling wave propagating at an angle  $\theta_2$  that is related to the incoming wave by Snell's law. State the field in region (2), and verify that  $|E_{\rm R}|^2 < |E_{\rm I}|^2$ , due to the power carried by the transmitted wave
  - ii. For  $n_2^2 < n_1^2 \sin^2 \theta$ , select a sign for  $\kappa$  such that the field in region (2) is damped in the +z-direction. State the field in region (2), and verify that  $|E_{\rm R}|^2 = |E_{\rm I}|^2$ , i.e. that there is total reflection at the interface.
- (f) Now specialize to a lossy medium in region (2), with  $\epsilon_2 \notin \mathbb{R}$ . Assume a permittivity  $\epsilon_2 = \epsilon' i\epsilon''$  with  $\epsilon' < 0$  and  $\epsilon'' > 0$ , as is the case for many metals.
  - i. Evaluate the wavenumber  $\kappa$ , choosing signs of the roots such that  $\text{Re}\kappa > 0$ , and  $\text{Im}\kappa < 0$ . Verify that the field in region (2) is then an outwards (+z) propagating damped wave.
  - ii. Show that  $|E_R|^2/|E_I|^2 < 1$ , i.e. that the reflection at the metallic surface is always accompanied by losses.
  - iii. Evaluate the reflectance  $R=|E_{\rm R}|^2/|E_{\rm I}|^2$  (Why is this correct here?) for the reflection of a plane wave coming in from air  $(n_1=1)$  towards a metal surface, as a function of the angle of incidence  $\theta$ . Use the material properties [1] of silver  $\epsilon_{\rm Ag}=-14.5-i1.2$  and copper  $\epsilon_{\rm Cu}=-11.7-i2.1$  at the wavelength  $\lambda=0.633\,\mu{\rm m}$  of a He-Ne-laser.
  - [1] M. N. Polyanskiy. "Refractive index database", http://refractiveindex.info (accessed May 05, 2016).
- 3. Commercial software dedicated to simulations in photonics / integrated optics is being offered by a number of companies. Examples are (state of early 2022, a list without any claim to completeness):
  - JCMwave,
  - Optiwave,
  - Synopsys (RSoft, PhoeniX),
  - · Photon Design,
  - CST,
  - · Ansys/Lumerical,
  - VPIphotonics.

In view of the classes of typical simulation tasks that we discussed during the lecture: Try to get an overview of what types of solvers are offered, and try to get an idea about the numerical schemes that the programs rely on.