

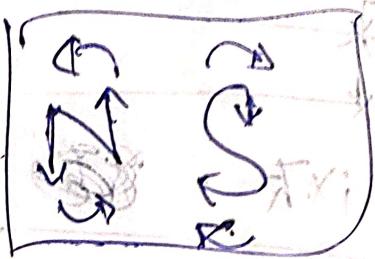
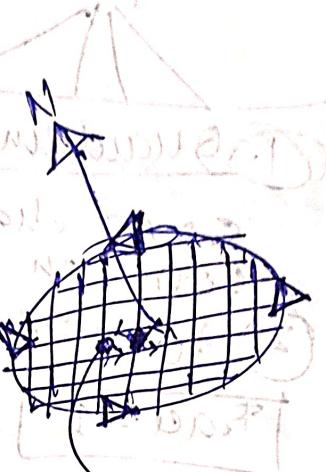
# Magnetism & SC Materials

## FMR Math.

Magnetic moment:

$$\mu = \int d\mu = I dS$$

$$= IA \quad (A-m^2)$$

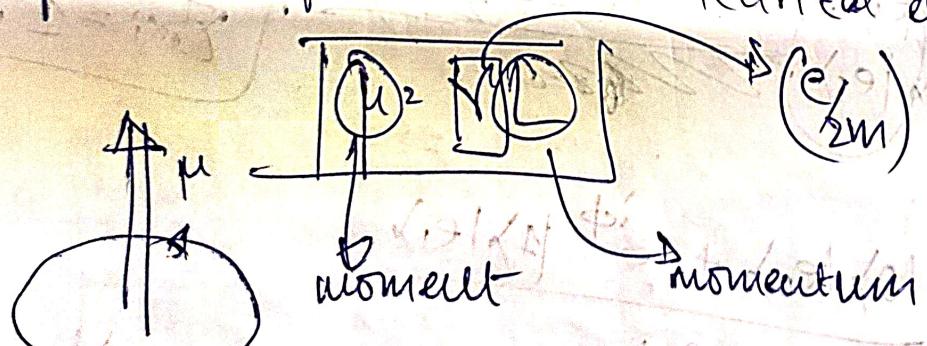


cancels out

Net current around secondary

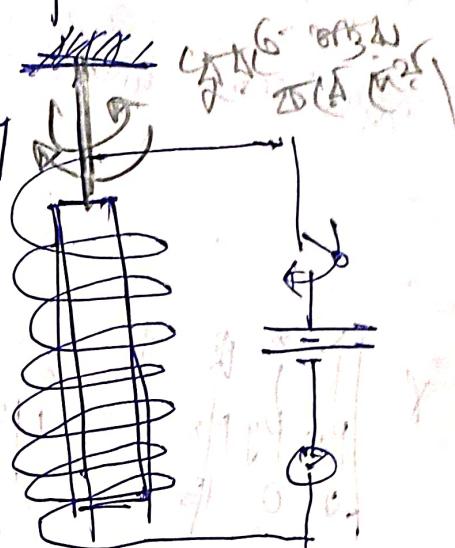
Magnetic moment & gyromagnetic ratio of momentum

A current carrying loop occurs as a consequence of motion of one or more electrical charges.



## Einstein - de Haas Effect

- ① A  $(\mu)$  [ferromagnetic] rod is suspended from a thin fibre.
  - ② A coil provides a magnetic field which magnetizes the bar vertically along its length.
  - ③ To conserve net angular momentum a rotation is produced due to net magnetization.
  - ④ Frequency of rotation direction  
flipping = frequency of change in current flow direction.
- Demo of reln b/w  $\mu$  &  $C$ .



Precission

Potential Energy of  $\vec{\mu}$  in  $\vec{B}$

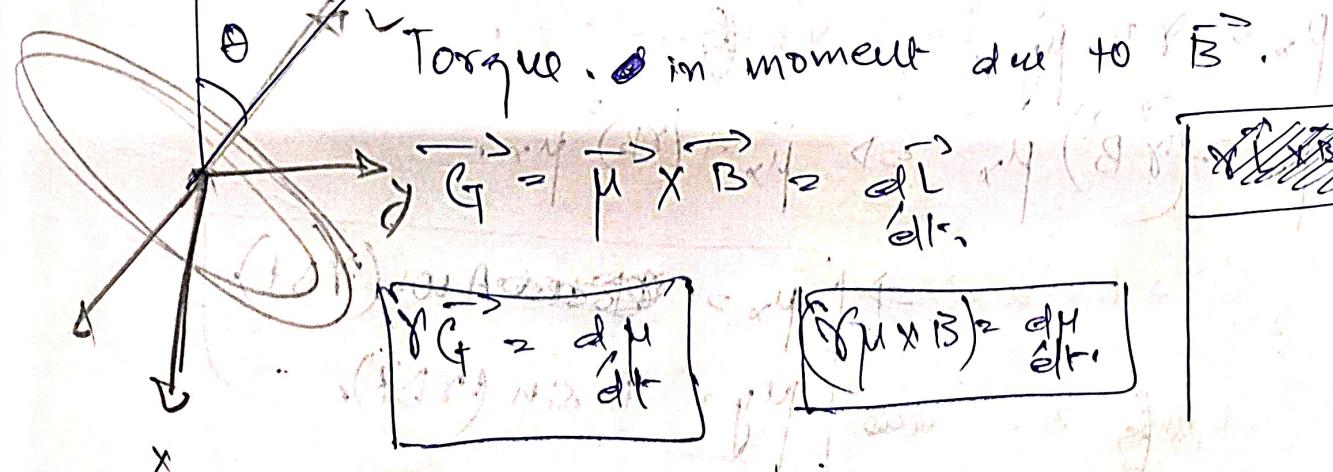
$$E = \mu \cdot B$$

$$N_z = \mu \cos \theta$$

$$\mu_x = 0$$

$$\mu_y = 0$$

$\vec{\omega}$  is given when  $\vec{\mu} \parallel \vec{B}$ .

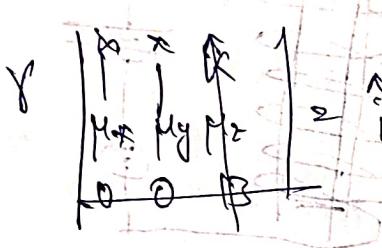


$$(\vec{\mu} \times \vec{B}) = \frac{d\vec{\mu}}{dt}$$

Rate of change if  $\mu$  is  $\perp$  to  $\mu$ .  
Therefore, circular motion.

The magnetic moment in a magnetic field precesses around the field at the Larmor frequency ( $\omega_L$ ) and a cone of semi-angle ( $\theta$ ).

Consider when the initial phase is  $\phi_0$ .



$$\nabla \cdot (\mu \hat{i} - \mu \hat{j}) = d \frac{d\mu_x}{dt} + d \frac{d\mu_y}{dt} + d \frac{d\mu_z}{dt}$$

$$\begin{cases} \ddot{\mu}_x = \gamma B \mu_y \\ \ddot{\mu}_y = -\gamma B \mu_x \\ \ddot{\mu}_z = 0 \end{cases}$$

$$\mu_x = A e^{i\omega t} + B e^{-i\omega t}, \quad \mu_y = C e^{i\omega t} + D e^{-i\omega t}, \quad \mu_z = \text{const}$$

$$\mu_y = C e^{i\omega t} + D e^{-i\omega t}, \quad \mu_y = \mu \sin(\omega t) \quad @ t=0$$

$$\ddot{\mu}_x = \gamma B \mu_y, \quad \mu_y = \mu \sin(\omega t)$$

$$\ddot{\mu}_x = (\gamma B)^2 \mu_x \quad \Rightarrow \quad \ddot{\mu}_x = (\gamma B)^2 \mu_x$$

$$\frac{d}{dt} \left( \frac{d\mu_x}{dt} \right) + (\gamma B)^2 \mu_x = 0 \quad \Rightarrow \quad \mu_x = A \cos(\gamma B t)$$

$$\mu_x = \mu \sin(\gamma B t), \quad \mu_y = B \sin(\gamma B t)$$

With initial conditions:

$\mu_x = \mu \sin(\gamma B t)$ ,  $\mu_y = B \sin(\gamma B t)$

$$10 \int g \int Q^5$$

$$\mu_z = |\mu| \cos \theta.$$

$$\mu_x = |\mu| \sin \theta (\gamma B_t).$$

$$\mu_y = |\mu| \sin \theta \sin(\gamma B_t).$$

$$\mu_x^2 + \mu_y^2 = q\mu^2 \sin^2 \theta.$$

circle -  $\omega_L = \gamma B$ .

Dynamics of electron in B mag field & non II magnetic moment.

$$\omega_L = \gamma B.$$

Larmor frequency.

$$L = \frac{mv}{2\pi r} \cdot \pi r^2 = \frac{\pi evr}{2}$$

Magnetic moment.

Size of ATOMIC magnetic moment.

Let us consider a hydrogen atom, ( $e, m_e$ )

electron circulating circularly around the nucleus with orbit radius ( $r$ ).

How much Magnetic moment?

$$\mu = I \cdot A.$$

$$T = \frac{2\pi r}{\nu}, f = \frac{\nu}{2\pi},$$

$$I = qf = -ev$$

$$A = \pi r^2.$$

$$L = m_e v r = m_e v r, I = -e \frac{1}{2\pi r \cdot m_e} \left( \frac{-eL}{2\pi r^2 m_e} \right).$$

$$\nu = \frac{L}{m_e r}.$$

$$(\mu) = -\frac{e\pi r}{2} \cdot \frac{1}{m_e r} \cdot f \left( -\frac{e}{2m} \right) L$$

$$\mu = \gamma L$$

$$1 \text{ BM} = 9.274 \times 10^{-24} \text{ A m}^2$$

$$L = n \hbar, \mu = \gamma n \hbar,$$

$$= n (\gamma \hbar),$$

$$= n \left( \frac{e\hbar}{2m} \right)$$

Momentum in a magnetic field [canonical moment].

$$F = q [\vec{E} + \vec{v} \times \vec{B}]$$

$$F = q \frac{d\vec{p}}{dt}$$

$$\text{Electric field} = -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t},$$

$$\text{Mag field} = \vec{\nabla} \times \vec{A}.$$

lorentz force,

$$\mu = \mu_B \mu_B$$

$$\mu_B = \frac{e\hbar}{2m}$$

"Bohr Magnet"

$$\frac{d\vec{p}}{dt} = q \left[ -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t} + \vec{v} \times \vec{\nabla} \times \vec{A} \right].$$

$$m \frac{d\vec{v}}{dt} = q \left[ -\nabla V + \vec{A} \cdot \frac{d\vec{A}}{dt} \right]$$

velocity

electric potential

$$\left( \frac{\partial}{\partial t} A_x \right) \hat{i} + \left( \frac{\partial}{\partial y} A_y \right) \hat{j} + \left( \frac{\partial}{\partial z} A_z \right) \hat{k}$$

$$\begin{aligned} &= \frac{\partial}{\partial t} A_x + \frac{\partial}{\partial x} A_y + \frac{\partial}{\partial y} A_z + \frac{\partial}{\partial t} A_t \\ &= \frac{\partial}{\partial t} [A_x + A_y + A_z + A_t]. \end{aligned}$$

$$\frac{\partial}{\partial t} [A(x, y, z, t)]$$

total derivative  
convection derivative

New momentum " $\vec{p}'$ "

$$\vec{p}' = \vec{p} + q\vec{A}.$$

Effective P.E.

$$PE = q(V - \vec{V} \cdot \vec{A}).$$

$$\text{Momentum} = (\vec{p} + q\vec{A})$$

$$p_L = -i\hbar \nabla \cdot (\vec{q}\vec{A}).$$

~~mass~~

$$\frac{d}{dt}(mv) + \frac{d}{dt}(qA) = -q \left[ \nabla V - \nabla(\vec{A} \cdot \vec{A}) \right].$$

$$\frac{d}{dt}(mv + qA) = -\nabla [qV - q\vec{A} \cdot \vec{A}].$$

canonical  
momentum

$$KE = \frac{mv^2}{2} = \frac{p^2}{2m}$$

$$= \frac{1}{2m} (\vec{p} - q\vec{A})^2$$

mechanical  
KE.

BOHR- van Leeuwen theorem.

In a classical system, @ thermal equilibrium, there is NO magnetization possible.

Thus the magnetization must be zero in a classical system.  
Magnetism is explained entirely QUANTUM phenomena & cannot be ~~explained~~ to classically.

# several types of magnetic orders.

- (a) diamagnetism:
- (b) paramagnetism:
- (c) ferrimagnetism:

(e) ferromagnetism -

(f) anti-ferromagnetism -

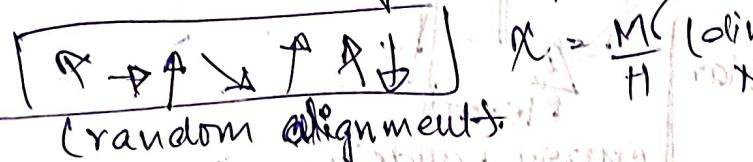
(d') canted AF.

(g) helimagnetism:

→  $\vec{B}$  is a perfect diamagnet.

Dia  $\Rightarrow$  All materials (nearly) diamagnet.

Para :



$$\chi_r = \frac{M}{H} \text{ (dimensionless)}$$

Ferro :



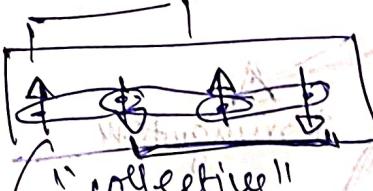
$$\chi \gg 1.$$

$$H=0.$$

interacting (in absence of field)  
aligned correctly, "spontaneous magnetization"

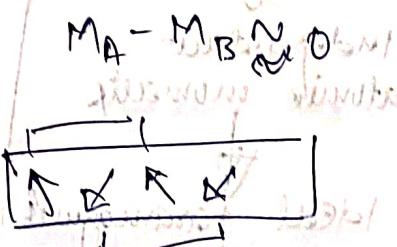
$$\mu_r = 1 + \chi.$$

~~AF~~ →  
AF

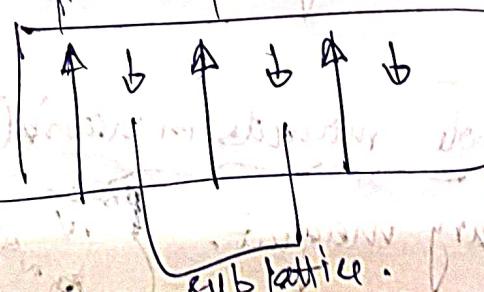


$$H=0, M=0$$

cauted



Ferrimagnetism →



$$M_A - M_B \gg 0.$$

$$M_A - M_B \neq 0$$

For any phase transition, there should be some order parameter.

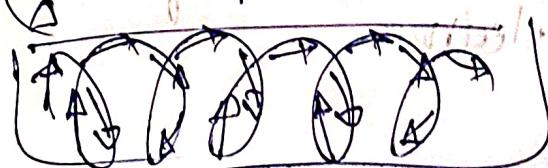
Ferro  $\rightarrow$  Ferro  
(second order)

PT (phase transition)

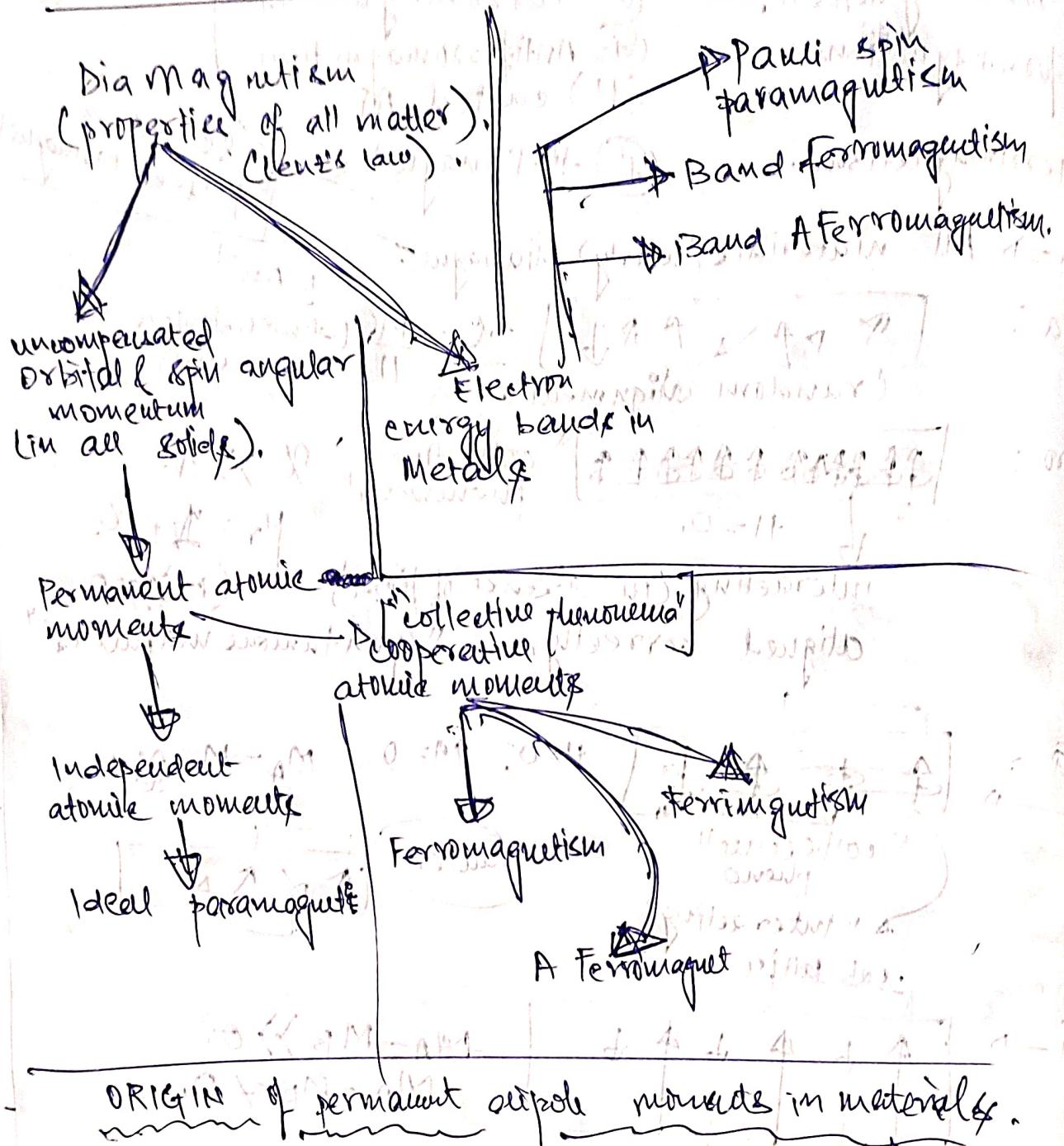
Heli : spins vectors align themselves

even if ext mag field = 0.

according to a helix, "helical pattern"



# FAMILY TREE OF MAGNETISM.



ORIGIN of permanent dipole moments in materials.

An atom is an elementary magnet.

In atom we find the following 3 contributions:

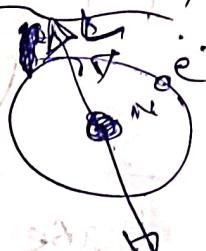
(1) Orbital motion of an electron contributes OAM.  

$$\vec{L} = \vec{r} \times \vec{p} = -ie(\vec{r} \times \vec{v})$$

$$L^2 |\Psi\rangle = \hbar^2 l(l+1) |\Psi\rangle,$$

$$L |\Psi\rangle = \text{mag} |\Psi\rangle.$$

Dipole moment due to the orbital motion of the electron.



Dipole moment  $\Rightarrow \mu = \gamma L$ .  $\rightarrow$  Bohr magneton.

$$\Rightarrow \mu_L = -\frac{e}{2m} \hbar l(l+1). \quad (\text{Orbital Mag moment})$$

(2) Dipole moment due to spin of electron.  
 $M_S$   $\rightarrow$  similar treatment

$$\mu_S = -\frac{g_s e \hbar}{2m} \cdot J \cdot f(S+1).$$

Bohr magneton

laundering factor

$$S^2 = \hbar^2 \cdot g(S+1)$$

$$S_z |N_S\rangle = \pm \text{mag} |N_S\rangle.$$

$$S_z |N_S\rangle = \pm \hbar |N_S\rangle.$$

$$\frac{M_N}{M_e} = 1836$$

(3) Dipole moment due to spin of nucleus.

$$\mu_N = +g_N \left( \frac{e \hbar}{2m_N} \right) \cdot J \cdot f_N (I_N + 1).$$

(neglected)

Big mass  
of nucleus.

Progressive Proof that (Lie algebra). [Also].

$$(a) [L_i^2, L_j^2] = 0 \quad (b) [L_i, L_j] = i \epsilon_{ijk} L_k.$$

$$(c) L_x = -i\hbar \frac{\partial}{\partial y} = -i\hbar \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right).$$

$$(d) S^2 = S_x^2 + S_y^2 + S_z^2, \\ = \frac{1}{2} (S_+ S_- + S_- S_+) + S_z^2.$$

$$(e) [S_x, S_y] = i S_z.$$

$$(f) [S_x, S_z] = i S_y.$$