Bootcamp : Physics

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Contents

Ι	Brownian motion	2
1	Historical aspects:	2
2	Random-walk	2
3	Bias Random walk	5

Part I

Brownian motion

Diffusion give rise to *material transport*. It gives a quantitative understanding of many phenomenon (permeability...).

1 Historical aspects:

1834: Clapeyron proposed a phenomenological equation for the ideal gas : PV = nRT, $PV = N.K_B.T$ (N : number of particules, K_B Boltzman constant $1.38 \times 10^{-23} \frac{kg.m}{K_s^2}$)

1828: Brown: studied pollen that was moving

1905: Einstein: $\langle x^2 \rangle = 2DT$ (D: diffusion coefficient)

1908: Perrin: confirmed the Einstein guess

2 Random-walk

Simplest model that describes Brownian motion. A particle has the same probability to move to the right and to the left. After two jumps, there is one way to $2(\frac{1}{4})$ and $-2(\frac{1}{4})$, and 2 ways to $0(\frac{1}{2})$.

$$\langle a \rangle = \frac{1}{m} \sum_{i=1}^{m} a_i = \sum_{i=1}^{m} p_i a_i$$

$$\langle a^2 \rangle = \frac{1}{m} \sum_{i=1}^m a_i^2$$

 $\langle a \rangle$ the average position, with m the number of realisations, i the i-th realisation and a_i the i-th realisation. In our example : $\langle a \rangle = \frac{1}{4}(2+-2+0+0) = 0$ and $\langle a^2 \rangle = \frac{1}{4}(2^2+-2^2+0^2+0^2) = 2$

In the 3 jumps case:

- 1 way to 3
- 1 way to -3
- 3 ways to 1
- 3 ways to -1

$$\langle a \rangle = \frac{1}{8}(3 + -3 + 3 \times 1 + 3 \times -1) = 0$$
 and $\langle a^2 \rangle = \frac{1}{8}(3^2 + -3^2 + (3 \times 1)^2 + (3 \times -1)^2) = 3$

$$N - jumps = \langle a \rangle = 0$$
 and $\langle a^2 \rangle = N$

If N is large enough, distribution of position is a binomial distribution.

If we say that the boxes have a l length, then $x = a \times l$, t_0 is the time delay between successive jumps. So $\frac{t}{t_0} = N$

$$\langle x \rangle = \langle a \rangle l = 0$$

$$\langle x^2 \rangle = \langle a^2 \rangle l^2$$

$$= l^2 N$$

$$= t \frac{l^2}{t_0}$$

$$= 2t \frac{l^2}{2t_0}$$

$$So D = \frac{l^2}{2t_0}$$

Example:

- $D = 100 \mu m^2 / s$
- $l_{bacteria} = 1 \mu m$
- $l_{cell} = 20 \mu m$
- $l_{arm} = 1m$

$$\begin{split} t_{bacteria} &= \frac{l^2}{2D} = \frac{(1\mu m)^2}{2(100\mu m^2/s)} = 5ms \\ t_{cell} &= \frac{l^2}{2D} = \frac{(20\mu m)^2}{2(100\mu m^2/s)} = 2s \\ t_{arm} &= \frac{l^2}{2D} = \frac{(1m)^2}{2(100\mu m^2/s)} = 300 years \end{split}$$

Q: How long does it take to tRNA to find a ribosome?

- $N_{tRNA} = 200000$ molecules / bacteria
- the length of the bacteria is around $1\mu m$ long
- The volume of the bacteria : $V = 1\mu m^3$
- $C_{tRNA} = \frac{200000}{1}$
- $d = \sqrt[3]{\frac{1}{C_{tRNA}}} \approx 0.016 \mu m$
- So $D=100\mu m/s^2 \rightarrow t=\frac{1^2}{2D}\approx 1\mu s$
- But ribosomes synthetize at 40 nucleotides per seconds, the the synthesis is about of 2 seconds for a 80 mucleotides tRNA.

So the ribosomes binds quickly to the tRNA, and takes longer to synthesize.

What is the concentration ? $C = \frac{Number\ of\ particles}{Volume} = 200000molec/\mu m$. Although, the concentration is not uniform is the bacteria.

What is flux (j)? Flux the number of particules that pass through an area A per unit of time. $j \propto 2-1=1$

$$\frac{\textit{\# of particles}}{\text{time step} \times \text{area of the plane}}$$

Example:

- $N_{H_2O} = 10^{10} molecules$
- T = 3000sec
- $A = 6\mu m^2$

$$j = \frac{10^{10}}{3\times10^3\times6} = 10^6 \frac{molecules}{s.\mu m^2}$$

Fick's Law

$$j = -D\frac{\partial c}{\partial x}$$

Darcy's Law

$$v = -d \frac{\partial P}{\partial x} d$$
: mobility coeff

Ohm's Law

$$\overrightarrow{j_E} = \sigma \overrightarrow{E} = -\sigma \frac{\partial V}{\partial x}$$
 sigma : electric conductivity, V : electric potential

Fourier's Law

$$j_T = -K \frac{\partial T}{\partial x}$$
 T: temperature, K: thermal conductivity

Random walk:

$$\begin{split} j &= \frac{particles(\rightarrow) - particles(\leftarrow)}{t_0 A} \\ particles(\rightarrow) &= N(x) \times \frac{1}{2} = C(x) V_{Box} \frac{1}{2} = C(x) A l \frac{1}{2} \\ particles(\leftarrow) &= N(x+l) \frac{1}{2} = C(x+l) A l \frac{1}{2} \\ j &= \frac{C(x) \frac{l}{2} - C(x+l) \frac{l}{2}}{t_0} \\ &= \frac{l^2}{2t_0} \left(\frac{C(x) - C(x+l)}{l} \right) \\ &= -D \frac{\partial C}{\partial x} \end{split}$$

Law of mass conservation : $\#part@_{t+t_0} = \#part@_t + \#part_{entering} - \#part_{leaving}$ $\#part_{entering} = j(x)At_0$, $part_{leaving} = j(x+l)At_0$ $\#part@_{t+t_0} = C(x,t+t_0)lA$

$$(...) = \left(\frac{C(x, t + t_0) - C(x, t)}{t_0}\right) lAt_0$$
$$= \frac{\partial C}{\partial t} lAT_0$$
$$\frac{\partial C}{\partial t} = -\frac{\partial j}{\partial x}$$
$$\frac{\partial C}{\partial t} = D\frac{\partial^2 C}{\partial x^2}$$

$$\frac{\partial C}{\partial t}=0, \frac{\partial^2 C}{\partial x^2}=0, \frac{\partial C}{\partial t}=0$$

If there is a concentration gradient:

C = -Kx with K the slope of the conentration gradient

This is a non-equilibrium state

$$\begin{aligned} \frac{\partial C}{\partial x} &= -K \\ \frac{\partial C}{\partial t} &= 0 \\ \frac{\partial^2 C}{\partial x^2} &= 0 \end{aligned}$$

Example 3: There is a film of water, and a drop of ink on it. We just take a plan to consider the diffusion. d is the size of the drop.

$$C(x,t) = \frac{N}{\sqrt{4\pi Dt}} e^{\frac{x^2}{4Dt}}$$

Example 4: There are centain molecules inside a cell. There is much more molecules outside the cell. There is a flux a the interface of the membrane.

$$\begin{split} j_M &= -P_m \Delta C \text{ with : } \Delta C = C_{ext} - C_{int} \\ R &= 10 \mu m; P_m = 20 \mu m/s \\ \frac{\partial C_i}{\partial t} &= \frac{j_m A}{V} = \frac{P_m A}{V} (C_e - C_i) \\ \frac{\partial C_i}{\partial t} &= -\frac{P_m A}{V} C_i + \frac{P_m A}{V} C_e \end{split}$$

$$\tau = \frac{V}{P_m A} = \frac{L^3}{\frac{L}{T} L^2} = T$$
$$= \frac{\frac{4}{3} \pi R^3}{P_m 4 \pi R^2} = \frac{R}{3 P_m} = \frac{1}{6} sec$$

$$\frac{\partial C_i}{\partial t} = -\frac{1}{t}(C_i - C_e)$$

$$\rightarrow C_i(t) = C_e - C_e e^{-t/\tau}$$

3 Bias Random walk

This is the case where the probability of going left is not the same than to go right. $K_+ \neq K_-$

$$F=\gamma v \rightarrow v = \frac{F}{\gamma}$$
 avec γ le coefficient de friction

$$j_{New} = vC$$

$$j = -D\frac{\partial C}{\partial c} + vc$$

$$\Delta x = v\Delta t$$
 and $N = C\Delta xA = cA\Delta tv$

$$\begin{split} j &= \frac{\#_{part}}{time.area} = \frac{CvA\Delta t}{\Delta tA} \\ &= Cv = c\frac{F}{\gamma} \end{split}$$

$$j_{tot} = -D\frac{\partial C}{\partial x} + Cv$$

Equilibrium situation:

$$0 = -D\frac{\partial C}{\partial x} + C\frac{F}{\gamma}$$

With the force F = cst

$$\begin{split} \frac{\partial C}{\partial x} &= \frac{CF}{\gamma}; \frac{\partial C}{\partial x} = \frac{CF}{\gamma D} \\ \frac{\partial C}{\partial x} &= \frac{C}{\lambda} \text{ as } \lambda = \frac{\gamma D}{F} \\ \frac{\partial C}{\partial x} &= \frac{C}{\lambda}; \frac{\partial C}{C} = \frac{\partial x}{\lambda} \end{split}$$

$$\int \frac{\partial C}{C} = \ln C - \ln C_0 = \ln \frac{C}{C_0}$$
$$\int \frac{\partial x}{\lambda} = \frac{x}{\lambda}$$

So:
$$\ln \frac{C}{C_0} = \frac{x}{\lambda}$$

$$C = C_0 e^{x/\lambda}$$

Same case as above, but with F = -Kx

$$j_{Tot} = 0 = -D\frac{\partial C}{\partial x} + C\frac{F}{\gamma} = -D\frac{\partial C}{\partial x} - \frac{CKx}{\gamma} = C$$
Solution:
$$D\frac{\partial C}{\partial x} = -\frac{KCx}{\gamma}; \frac{\partial C}{\partial x} = -\frac{KCx}{\gamma D}$$

$$\frac{dC}{dx} = \frac{-KCx}{\gamma D}; \frac{dC}{C} = -\frac{KCx}{\gamma D} dx;$$

$$\int \frac{dC}{C} = \ln C - \ln C_0$$

$$\int \frac{-Kx}{\gamma D} dx = -\frac{K}{\gamma D} \frac{x^2}{2}$$

$$\ln \frac{C}{C_0} = -\frac{Kx^2}{2\gamma D} \Rightarrow$$

$$C = C_0 e^{\frac{-Kx}{2\gamma D}}$$

We expect that : $\sigma^2 = \frac{\gamma D}{K}$

In a drop:

At a surface, the particles evaporate at a rate $\frac{1}{\tau}$

$$N_{surface} = C_0 A \lambda$$

$$j = \frac{N_{surface}}{A \tau} = \frac{C_0 \lambda}{\tau}$$

$$N(t + \Delta t) = N(t) - jA\Delta t$$

$$\Rightarrow \frac{dN}{Dt} = -jA$$
And: $C_0 = \frac{N}{V} = cst$; $N = C_0V$

$$\Rightarrow \frac{dN}{dt} = C_0 \frac{dV}{dt}$$

$$\frac{dV}{dt} = \frac{-jA}{C_0}$$

$$= \frac{C_0 \lambda A}{C_0 \tau}$$

$$= -\frac{\lambda A}{\tau}$$

$$\frac{dV}{dt} = 4\pi R^2 \frac{dR}{dt} = A \frac{dR}{dt}$$

$$\frac{dR}{dt} = \frac{-\lambda}{\tau}$$

$$R(t) = R_0 \times -\frac{\lambda}{\tau}t$$