

## Correcting the Reduced Density Matrix Calculation.

- Previously I had written the reduced density matrices by tracing out over the respective eigenstates of the system. This, however, has no physical meaning: I should instead trace out over the subsystems (TLS, HO) of the composite system itself.

### Rewriting the density matrix

- Previously, I had the following solutions to the time-evolved state and the eigenstates.

$$|\psi(t)\rangle = e^{-i(\omega_2 + \frac{\kappa'}{\hbar})t} \left(1 + \left(\frac{a + \kappa'}{c'}\right)^2\right)^{-\frac{1}{2}} \left(\frac{a + \kappa'}{c'}\right) |V_+\rangle \\ + e^{-i(\omega_2 - \frac{\kappa'}{\hbar})t} \left(1 + \left(\frac{a - \kappa'}{c'}\right)^2\right)^{-\frac{1}{2}} \left(\frac{a - \kappa'}{c'}\right) |V_-\rangle$$

$$|V_+\rangle = \left(1 + \left(\frac{a + \kappa'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(|g, n+1\rangle + \frac{a + \kappa'}{c'} |e, n\rangle\right)$$

$$|V_-\rangle = \left(1 + \left(\frac{a - \kappa'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(|g, n+1\rangle + \frac{a - \kappa'}{c'} |e, n\rangle\right)$$

$$\text{where } a = \hbar \left(\frac{\omega_1 - \omega_2}{2}\right), \quad \kappa' =$$

$$\kappa' = \hbar \left[ \left(\frac{\omega_1^2 + \omega_2^2}{2} - \frac{\omega_1 \omega_2}{2} + G^2 \right)^{\frac{1}{2}} \right]$$

$$\kappa' = \hbar \text{ where } n=0$$

$$c' = \hbar G$$

- In order to transform  $|\psi(t)\rangle$  back into the basis original of joint state; we insert the identity operator:

$$\mathbb{1} = \sum_i |\phi_i\rangle \langle \phi_i| = |g, n+1\rangle \langle g, n+1| + |e, n\rangle \langle e, n|$$

which 'select' those bases.

$$|\psi(t)\rangle = \mathbb{1} \cdot \sum_i c_i |v_i\rangle = (|g, n+1\rangle \langle g, n+1| + |e, n\rangle \langle e, n|) \\ \times (c_+ |v_+\rangle + c_- |v_-\rangle)$$

$$|\psi(t)\rangle = \alpha |g, n+1\rangle + \beta |e, n\rangle$$

Recalling that our time evolved state has  $n=0$ , let's find the ~~same~~ factors  $\alpha$  and  $\beta$ .

$$\alpha(t) = e^{-i(\omega_2 + \frac{k'}{\hbar})t} \left(1 + \left(\frac{k'+a}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a+k'}{c'}\right) \cdot \left(1 + \left(\frac{a+k'}{c'}\right)^2\right)^{-\frac{1}{2}} \\ + e^{-i(\omega_2 - \frac{k'}{\hbar})t} \left(1 + \left(\frac{a-k'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a-k'}{c'}\right) \cdot \left(1 + \left(\frac{a-k'}{c'}\right)^2\right)^{-\frac{1}{2}}$$

$$= e^{-i(\omega_2 + \frac{k'}{\hbar})t} \left(1 + \left(\frac{a+k'}{c'}\right)^2\right)^{-1} \cdot \left(\frac{a+k'}{c'}\right) \\ + e^{-i(\omega_2 - \frac{k'}{\hbar})t} \left(1 + \left(\frac{a-k'}{c'}\right)^2\right)^{-1} \cdot \left(\frac{a-k'}{c'}\right)$$

$$\beta(t) = e^{-i(\omega_2 + \frac{k'}{\hbar})t} \left(1 + \left(\frac{a+k'}{c'}\right)^2\right)^{-1} \cdot \left(\frac{a+k'}{c'}\right)^2 \\ + e^{-i(\omega_2 - \frac{k'}{\hbar})t} \left(1 + \left(\frac{a-k'}{c'}\right)^2\right)^{-1} \cdot \left(\frac{a-k'}{c'}\right)^2$$

We previously derived that for  $|\psi\rangle = \alpha |g, n+1\rangle + \beta |e, n\rangle$ :

$$\rho_{\pi\pi} = \alpha\alpha^* |n+1\rangle \langle n+1| + \beta\beta^* |n\rangle \langle n|$$

$$\rho_{\pi\pi} = \alpha\alpha^* |n+1\rangle \langle g| + \beta\beta^* |e\rangle \langle e|$$

Thus, our time-evolved state  $|\psi(t)\rangle = \alpha(t)|g, n+1\rangle + \beta(t)|e, n\rangle$  has subspaces:

$$\rho_{\text{HO}}^{\text{res}}(t) = \alpha(t)\alpha^*(t)|1\rangle\langle 1| + \beta(t)\beta^*(t)|0\rangle\langle 0|$$

$$\rho_{\text{TLS}}^{\text{res}}(t) = \alpha(t)\alpha^*(t)|g\rangle\langle g| + \beta(t)\beta^*(t)|e\rangle\langle e|$$

Let's calculate  $\alpha(t)\alpha^*(t)$  and  $\beta(t)\beta^*(t)$ :

$$\begin{aligned}\alpha(t)\alpha^*(t) &= \left( e^{-i(\omega_2 + \frac{k'}{\hbar})t} \eta_+ + e^{-i(\omega_2 - \frac{k'}{\hbar})t} \eta_- \right) \\ &\times \left( e^{i(\omega_2 + \frac{k'}{\hbar})t} \eta_+ + e^{i(\omega_2 - \frac{k'}{\hbar})t} \eta_- \right)\end{aligned}$$

$$\text{where } \eta_{\pm} = \left( 1 + \left( \frac{a \pm k'}{c'} \right)^2 \right)^{-1} \cdot \left( \frac{a \pm k'}{c'} \right)$$

$$\begin{aligned}\alpha(t)\alpha^*(t) &= \eta_+^2 + \eta_-^2 + e^{-i\frac{k'}{\hbar}t} \eta_+ \eta_- + e^{i\frac{k'}{\hbar}t} \eta_+ \eta_- \\ &= \eta_+^2 + \eta_-^2 + 2\eta_+ \eta_- \cos\left(\frac{k'}{\hbar}t\right)\end{aligned}$$

$$\begin{aligned}\beta(t)\beta^*(t) &= \left( e^{-i(\omega_2 + \frac{k'}{\hbar})t} \eta_+ \left( \frac{a+k'}{c'} \right) + e^{-i(\omega_2 - \frac{k'}{\hbar})t} \eta_- \left( \frac{a-k'}{c'} \right) \right) \\ &\times \left( e^{i(\omega_2 + \frac{k'}{\hbar})t} \eta_+ \left( \frac{a+k'}{c'} \right) + e^{i(\omega_2 - \frac{k'}{\hbar})t} \eta_- \left( \frac{a-k'}{c'} \right) \right)\end{aligned}$$

$$\begin{aligned}&= \eta_+^2 \left( \frac{a+k'}{c'} \right)^2 + \eta_-^2 \left( \frac{a-k'}{c'} \right)^2 + e^{-i\frac{k'}{\hbar}t} \eta_+ \eta_- \left( \frac{a+k'}{c'} \right) \left( \frac{a-k'}{c'} \right) \\ &\quad + e^{i\frac{k'}{\hbar}t} \eta_+ \eta_- \left( \frac{a-k'}{c'} \right) \left( \frac{a+k'}{c'} \right)\end{aligned}$$

$$= \eta_+^2 \left( \frac{a+k'}{c'} \right)^2 + \eta_-^2 \left( \frac{a-k'}{c'} \right)^2 + 2\cos\left(\frac{k'}{\hbar}t\right) \eta_+ \eta_- \left( \frac{a+k'}{c'} \right) \left( \frac{a-k'}{c'} \right)$$

In summary:

$$\alpha(t)\alpha^*(t) = \eta_+^2 + \eta_-^2 + 2\eta_+\eta_- \cos\left(\frac{k'}{\hbar}t\right)$$

$$\beta(t)\beta^*(t) = \eta_+^2\left(\frac{a+k'}{c'}\right)^2 + \eta_-^2\left(\frac{a-k'}{c'}\right)^2 + 2\eta_+\eta_-\left(\frac{a+k'}{c'}\right)\left(\frac{a-k'}{c'}\right)\cos\left(\frac{k'}{\hbar}t\right)$$

For the reduced states:

$$\rho_{H0}(t) = \alpha(t)\alpha^*(t)|g\rangle\langle g| + \beta(t)\beta^*(t)|0\rangle\langle 0|$$

$$\rho_{\pi s}(t) = \alpha(t)\alpha^*(t)|g\rangle\langle g| + \beta(t)\beta^*(t)|e\rangle\langle e|$$

which were derived from an initial state  $|\psi(0)\rangle = |e, 0\rangle$  where  $n=0$ .

$$\text{where } \eta_{\pm} = \left(1 + \left(\frac{a \pm k'}{c'}\right)^2\right)^{-1} \left(\frac{a \pm k'}{c'}\right)$$

$$a = \hbar \left(\frac{\omega_1 - \omega_2}{2}\right)$$

$$k' = \hbar \int \left( \frac{\omega_1^2 + \omega_2^2}{4} - \frac{\omega_1 \omega_2}{2} + G^2 \right)$$

$$c' = \hbar G$$