

Density Matrix and Reduced Density Matrix of a TLS-HO Composite System

- Previously, we calculated the evolved state written in the eigenbasis.
- Now, we will calculate the density and reduced density matrices for the unevolved state and evolved state.

1. Density Matrix for basis states

- Recall $|\psi\rangle = \alpha|g, n+1\rangle + \beta|e, n\rangle$.
- The density matrix is the sum of the states written as projectors. Since we have a pure state (only one state associated with system):

$$\begin{aligned}\rho &= |\psi\rangle\langle\psi| \\ &= \alpha\alpha^*|g, n+1\rangle\langle g, n+1| + \beta\beta^*|e, n\rangle\langle e, n| \\ &\quad + \alpha\beta^*|g, n+1\rangle\langle e, n| + \alpha^*\beta|e, n\rangle\langle g, n+1|\end{aligned}$$

- Written in matrix form:

$$\rho = \begin{pmatrix} \alpha\alpha^* & \alpha\beta^* \\ \alpha^*\beta & \beta\beta^* \end{pmatrix} //$$

- Now, for the reduced density matrices, we 'trace' out a subsystem to get a density matrix in terms of the other.

$$\begin{aligned}\rho_{\text{HO}} &= \text{Tr}_{\text{TLS}}(\rho) = \langle g|\rho|g\rangle + \langle e|\rho|e\rangle \\ &= \alpha\alpha^*\langle g|g, n+1\rangle\langle g, n+1|g\rangle \\ &\quad + \beta\beta^*\langle e|e, n\rangle\langle e, n|e\rangle = \alpha\alpha^*|n+1\rangle\langle n+1| + \beta\beta^*|n\rangle\langle n|\end{aligned}$$

- And for the subsystem for TLS:

$$\begin{aligned}\rho_{\text{TLS}} &= \text{Tr}_{\text{HO}}(\rho) = \langle n | \rho | n \rangle + \langle n+1 | \rho | n+1 \rangle \\ &= \alpha \alpha^* |g\rangle \langle g| + \beta \beta^* |e\rangle \langle e|\end{aligned}$$

- So, in summary:

$$\rho = \begin{pmatrix} \alpha \alpha^* & \alpha \beta^* \\ \alpha^* \beta & \beta \beta^* \end{pmatrix}$$

$$\rho_{\text{HO}} = \alpha \alpha^* |n+1\rangle \langle n+1| + \beta \beta^* |n\rangle \langle n|$$

$$\rho_{\text{TLS}} = \alpha \alpha^* |g\rangle \langle g| + \beta \beta^* |e\rangle \langle e|$$

2. Density Matrix for evolved state

- We derived $|\psi(t)\rangle$ previously to be:

$$|\psi(t)\rangle = \begin{bmatrix} e^{-i(\omega_2 + k')t} \left(1 + \left(\frac{a+k'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a+k'}{c'}\right) \\ e^{-i(\omega_2 - k')t} \left(1 + \left(\frac{a-k'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a-k'}{c'}\right) \end{bmatrix}$$

$$\text{where } k' = \hbar \sqrt{\left(\frac{\omega_1^2 + \omega_2^2}{4} - \frac{\omega_1 \omega_2}{2} + G^2\right)}$$

$$a = \hbar \left(\frac{\omega_1 - \omega_2}{2}\right)$$

$$c' = \hbar G$$

ω_1 is the frequency associated to the TLS

ω_2 is the frequency associated to the HO.

- The evolved state is written in terms of the eigenbasis $\{|V_+\rangle, |V_-\rangle\}$. Furthermore, it is associated with an initial state of $|\psi(0)\rangle = |e, 0\rangle$.

• We can thus write $|\psi(t)\rangle$ as a superposition in the eigenbasis.

$$|\psi(t)\rangle = \eta_+ |v_+\rangle + \eta_- |v_-\rangle$$

$$\text{where } \eta_{\pm} = e^{-i(\omega_2 \pm k')t} \left(1 + \left(\frac{a \pm k'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a \pm k'}{c'}\right)$$

• We now carry out the same steps as for the basis states,

$$\rho(t) = |\psi(t)\rangle \langle \psi(t)|$$

$$= \eta_+ \eta_+^* |v_+\rangle \langle v_+| + \eta_- \eta_-^* |v_-\rangle \langle v_-|$$

$$+ \eta_+ \eta_-^* |v_+\rangle \langle v_-| + \eta_-^* \eta_+ |v_-\rangle \langle v_+|$$

$$\text{where } \eta_{\pm}^* = e^{i(\omega_2 \pm k')t} \left(1 + \left(\frac{a \pm k'}{c'}\right)^2\right)^{-\frac{1}{2}} \cdot \left(\frac{a \pm k'}{c'}\right)$$

• In matrix form:

$$\rho(t) = \begin{pmatrix} \eta_+ \eta_+^* & \eta_+ \eta_-^* \\ \eta_-^* \eta_+ & \eta_- \eta_-^* \end{pmatrix} //$$

• And the subsystems are:

$$\rho_+(t) = \eta_+ \eta_+^* |v_+\rangle \langle v_+|$$

$$\rho_-(t) = \eta_- \eta_-^* |v_-\rangle \langle v_-| //$$