**Extension of Hall-symbols of crystallographic space groups to magnetic space groups**

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**Abstract**

The Hall-symbols for describing unambiguously the generators of space groups have been extended to describe whatever setting of the 1651 types of magnetic space groups (Shubnikov groups).

**Introduction**

The international symbols for crystallographic space groups (CSG) are based in the Hermann-Mauguin (HM) symbols containing information about the Bravais lattice of the translation subgroup and the nature and orientation of the symmetry elements. However, as pointed out by S.R. Hall (Hall, 1981), there is no explicit information about the position of the origin determining partially the translational part of the symmetry operators. This has, consequently, an important drawback: the impossibility of generating the full set of symmetry operators from the symmetry elements appearing in the symbol, in the general case, coinciding with the lists of the *International Tables for Crystallography* (ITC). The computer programs, not using a database, doing that task have always some special code to handle special cases in order to generate strictly the same set of operators. The proposal by S.R. Hall (Hall, 1981) solves this problem in an elegant way allowing, at the same time, the possibility of describing reasonable arbitrary settings of CSG. The word “reasonable” is used here to emphasize the fact that the symmetry operators (with rotational part described always by an integer matrix) may be very complicated if a unit cell setting, for instance, is selected in which a three-fold axis in along an arbitrary [uvw] direction. However, even in these cases, the Hall symbol, as treated in (S. R. Hall and R. W. Grosse-Kunstleve, 2010, hereafter referred as HG-K) can be written by using the standard symbols described below followed by a setting change included at the end of the symbol.

In this document, we propose to modify slightly the original Hall-symbols in order to extend them for describing magnetic space groups (MSG). For that, we had to change the original symbol ' to ^  (see below). The reason is that the prime symbol ' is reserved for the time reversal operator to be consistent with the currently used symbols Belov-Neronova-Smirnova (BNS) and Opechowski-Guccione (OG) notations for MSG. Here we will not discuss the derivation of MSG or their relation with families of CSG that may be found in literature (see, for instance, D. Litvin, 2001); we shall be concentrated in the new nomenclature aspects.

The aim of the Hall symbols is to have an unambiguous descriptor of a CSG or MSG and not to substitute the HM, BNS or OG (or whatever new proposal that may be adopted) symbols used for a standard setting. Once we have a correspondence with the standard setting of CSG or MSG, an appropriate computer program can parse the symbols and generate the full set of symmetry operators. We have to keep in mind that two different Hall symbols may represent the same space group type in a particular setting. This is obvious by the fact that we can use different generators to obtain the same full representative operators of the group type.

**New proposed Hall symbols**

This section summarises the notation described in HG-K, and the modifications we propose.

The Hall-symbols try to describe a CSG or a MSG using a short set of generators. Each generator corresponds to a Seitz operator acting on atom positions or on magnetic moments. The action of the operator **S**={R, θ|**t**} on the atom position **r** having a magnetic moment **m** is given by:

**r**’= R **r** + **t**  (1)

The corresponding magnetic moment is obtained as:

**m**’= θ det(R) R **m** (2)

The integer θ is equal to -1 is the operator is “primed” (combined with time reversal operator 1’) and equal to 1 in the “non-primed” case. The Hall-symbol is the combination of a series of Seitz operators with a lattice symbol and the presence or not of a centre of symmetry at the origin. Each operator is constituted by a particular notation for the operator S that describes unambiguously the matrix R, the value of θ and the translation vector **t**. The operator S may by represented as a Seitz matrix or in the Jones faithful notation. For instance, a Seitz matrix representing a primed two-fold screw rotation along the b-axis is given by:



Notice that we have no way to explicitly writing the time reversal operator within the matrix; only the knowledge of the presence of time reversal is used in applying effectively the equations (1) and (2). The same operator in Jones faithful notation can be written as: , which is more compact. We will use this last representation hereafter.

The newly proposed Hall-symbols have the general form (compare to equation A1.4.2.4 in HG-K):

 (3)

The symbol **L** is the lattice symbol, **L**= P, A, B, C, I, R, H, F, X. The symbol may include a preceding minus sign for indicating that a centre of symmetry exists at the origin. The X symbol is used for non-conventional lattice centrings, in such a case the explicit generators of the lattice must be provided within the *p* provided operators. The operators related with the **L** symbols are described in Table 1.

The symbol *N* is 1, 2, 3, 4 or 6 for proper rotations and  (or in text-only characters -1, -2, -3, -4 or -6) for improper rotations (the symbol -2 corresponds to mirror or glide planes depending on the associated translation). The symbol *θ*  is “ '” when the operator is associated with time reversal or it is absent otherwise. The symbol *A* indicates the direction of the rotation axis. The possible values are x, y, z, ^, ” and \*, for rotation axis along **a**, **b**, **c**, **a**-**b** (or alternatively **b** - **c** or **c** - **a**), **a**+**b** (or alternatively **b** + **c** or **c** + **a**) and **a**+**b**+**c**, respectively. The **t** translation symbols are 1, 2, 3, 4, 5, 6, *a*, *b*, *c*, *n*, *u*, *v*, *w*, *d* and they are described in Table A1.4.2.3 in HG-K. This table is reproduced below as Table 2. These translations apply additively [e.g. *ad* signifies a (3/4, 1/4, 1/4) translation]. The change of basis operator **V** may be used as described in HG-K; however, we propose an alternative change of notation that is more clearly related to its purpose: a change of reference frame, involving the basis vectors, to which refer the symmetry operators.

**Default axes from the positional order of operators**

The superscripts indicating the direction of rotation axes may be suppressed by adopting the following rules (see section A1.4.2.3.1. in HG-K) for the symbol given in (3):

(*i*) the first rotation (or roto-inversion) axis direction is always along **c**

(*ii*) the second rotation (if |*N*| is 2) axis direction is along **a** if preceded by an |*N*| of 2 or 4, it is along **a**-**b** if preceded by an |*N*| of 3 or 6

(*iii*) the third rotation (if |*N*| is 3) has an axis direction of **a** + **b** + **c**

This implies that in the major part of cases the symbol “^” is not needed. In all other cases, the superscript of the provided operators should be explicitly written. The rotational part of operators () are given explicitly in Tables 3, 4 and 5.

Another additional convention is that anti-translations () should appear after the rotational operators and non-conventional lattice translations ().

**New change of basis operator V**

In HG-K, the authors extend the original shift of origin symbol to a complete change of basis, however they use a vector form in which a symbol close to Jones faithful representation of symmetry operators is used to represent this time a change of basis. Symmetry operators are implicitly referred to a reference frame {O, **a**, **b**, **c**}, where O represents a point in 3D space and **a**, **b**, **c** are the basis vectors. In previous forms of the Hall symbol (see HG-K), the operator **V** is represented in vectorial form **V**= (*v*1, *v*2, *v*3) with *vi* = *ri*1X+ *ri*2Y+ *ri*3Z+*ti*, where *rij* and *ti* are rational numbers. The operator can be written as a Seitz matrix as given in HG-K:

 (4)

In case of a simple change of origin, *rij*=0 if *i*≠*j* and *rii*=1, we conserve the same notation as proposed in (Hall, 1981) and HG-K, where the Hall symbol finishes with an integer vector in parenthesis of the form (*tx*, *ty*, *tz*), meaning *t*1= *tx* /12, *t*2= *ty* /12 and *t*3= *tz* /12.

The meaning of the change of basis operator **V** is the following: the new basis vectors **a**', **b**', **c**' are obtained as:

 (5)

Notice that the 3×3 submatrix is transposed to that of **V**. The new origin O' have coordinates (*t*1, *t*2, *t*3) with respect to the original {O, **a**, **b**, **c**} basis.

The symmetry operators provided in the Hall symbol (generators) should be transformed, if the operator **V** is provided, according to the equation:

 (6)

Where the new operators are now implicitly referred to the basis {O', **a**', **b**', **c**'}. We have maintained the notation of the Hall symbol using

In order to be compatible with this form, we propose to use the colon “:” as separator to tell the parsing program that we use the new symbolic form a **V** operator. So the equation (3) is written as:

 (7)

**Detailed examples**

The header “ASCII” given in the examples below means that we use characters from the ASCII character set available in computers and coded with 1 byte (8 bits), so subscripts and superscripts are not used. Instead, the translation symbols are just juxtaposed to the rotational symbols. When using the text (ASCII) form of the Hall symbols we use one or more spaces between symbols as in:



An example of legal Hall symbol for a magnetic group is  that in ASCII form is represented as

P 6c -2 1'c

or

P 6c -2 1c'

Notice that the prime symbol commutes with the translations symbols. The above Hall symbol corresponds to the MSG (in BNS notation) in the standard setting.

One can create a Hall symbol without knowing to which MSG standard type corresponds. A computer program can identify the standard MSG type and calculate the change of basis to obtain the standard setting. For instance, a symbol like -P 4n' 1u' () corresponds to the group P\_C4\_2/m () in a particular setting. To obtain the standard setting one has to apply the following basis change: -**a**/4-**b**/4, **a**/4-**b**/4, **c**;-1/8,-1/8,0.

Let us consider conventional settings of MSG as examples. MSG are usually divided in four types:

**Type-1** MSG, also called colourless groups, are those having θ=1 for all operators. They are isomorphous to the 230 CSG. The BNS and OG symbols coincide in this case with the HM symbols of CSG. The Hall symbols are the same as those of the CSG. Examples of the symbols can be found in Table A1.4.2.7 of HG-K.

**Type-2** MSG, also called paramagnetic groups, are those having two set of operators. The first set is the same as those of Type-1 groups and in the second set the operators have θ = -1 combined with all operators of the first set. Their number is also 230 and the BNS and OG symbols are those of the corresponding CSG followed by the time reversal operator. The Hall symbols are obtained from the Hall symbols of Type-1 groups by putting the time reversal operator as the last generator.

Examples:

BNS/OG BNS/OG Compact Hall Hall

symbol ASCII ASCII symbol ASCII

 P 41 21 2 1' P4\_12\_121'  P 4abw 2nw 1'

 P 42/m m c 1' P4\_2/mmc1'  -P 4c 2 1'

The full set of operators for the first group () can be obtained by the generators: -*y+*1/2, *x+*1/2*, z+*1/4*,* 1 (4abw ); *x+*1/2, -*y+*1/2, -*z+*3/4(2nw); *x*, *y*, *z*, -1 (1')

The full set of operators for the second group () can be obtained by the generators: -*x*, -*y*, -*z*, 1 (-1); -*y*, *x, z+*1/2*,* 1 (4c ); *x*, -*y*, -*z* (2); *x*, *y*, *z*, -1 (1')

**Type-3** MSG, also called type-1 black and white (BW1) groups, are those having the same lattice as that of the paramagnetic group, but half of operators are associated with time reversal. Their number is 674 and the BNS and OG symbols are the same.

Examples:

BNS/OG BNS/OG Compact Hall Hall

symbol ASCII ASCII symbol ASCII

 P c' c n' Pc'cn'  -P 2ab' 2ac'

 P 62' 2 2' P6\_2'22'  P 62' 2c' (0 0 1)

The full set of operators for the first group () can be obtained by the generators: -*x*+1/2,-*y*+1/2, *z*,-1 (2ab'); *x*+1/2,-*y*,-*z*+1/2,-1 (2ac'); -*x*,-*y*,-*z*, 1 (-1)

The full set of operators for the second group () can be obtained by the generators: *x*-*y*, *x*, *z*+1/3,-1 (62'); -*y*,-*x*,-*z*+2/3,-1 (2c'). Notice that the second operator without taking into account the change of origin would have the form: -*y*,-*x*,-*z*+1/2,-1 according to the rule (*ii*) and Table 4.

**Type-4** MSG, also called type-2 black and white (BW2) groups, are those containing anti-translations: lattice translations associated with time reversal. Of course, these anti-translations are not lattice vectors. Their number is 517 and the BNS and OG symbols are different. Here we have conserved the columns concerning the ASCII representations of the BNS and OG in compact form.

Examples:

BNS OG BNS OG Hall Hall

symbol symbol ASCII ASCII symbol ASCII

  F\_Sdd2 P\_Inn2  F 2 -2d 1'S

  P\_Cbcn C\_Pm'cm  -P 2n 2ab 1'C

  P\_c6\_3cm P\_2c6'm'm  P 6c -2 1'c

  F\_Sd-3 P\_Fn-3  -F 2uv 2vw 3 1n'

The full set of operators for the first group () can be obtained by the generators: -*x*, -*y*, *z*, 1 (2); -*x*+1/4, *y*+1/4, *z*+1/4, 1 (); *x*+1/2, *y*+1/2, *z*+1/2, -1 () and the F-centring translations *x*, *y*+1/2, *z*+1/2, 1; *x*+1/2, *y*, *z*+1/2,1 and  *x*+1/2, *y*+1/2, *z*, 1. In this example we are using the anti-translation symbol S in the Hall symbol that corresponds to **t**=(1/2, 1/2, 1/2) (*n* symbol) when the lattice centring is F and to (0, 0, 1/2) (*c* symbol) in all other cases. Of course, we can use the alternative Hall symbol that is equivalent.

The full set of operators for the second group () can be obtained by the generators: -*x+*1/2, -*y+*1/2, *z+*1/2, 1 (); *x+*1/2*,-y+*1/2*,-z,*1 (); *-x*, -*y*, -*z,* 1 (); *x+*1/2, *y+*1/2, *z,*-1 ().

The full set of operators for the third group () can be obtained by the generators: -*x*-*y*, *x*, *z*+1/2, 1 (); *y*, *x*, *z*, 1 (); *x*, *y*, *z*+1/2, -1 ().

The full set of operators for the fourth group () can be obtained by the generators: -*x*-*y*, *x*, *z*+1/2, 1 (); *y*, *x*, *z*, 1 (); *x*, *y*, *z*+1/2, -1 ().

We provide a supplementary text document in which all the standard MSG are listed together with the Hall symbols and the explicit generators used in Jones faithful form. The the newly proposed International symbols (to be discussed in the Commission for Magnetic Structures of the IUCr) for the standard settings are also in the list.

**Table 1: Lattice symbols L**(similar to table A1.4.2.2. in HG-K)

The lattice symbol **L** implies the following generators in Jones’ faithful notation. The identity is always implied even if not provided, the time reversal operator is provided as the last item, the different generators are separated by “;”. In the general case X we have used the notation 1**t** ≡ *x*+*t*1, *y*+*t*2, *z*+*t*3, 1 where **t**=(*t*1, *t*2, *t*3), with *t*i rational numbers.

P: *x*, *y*, *z*, 1 -P: -*x*, -*y*, -*z*, 1

A: *x*, *y*+1/2, *z*+1/2, 1 -A: -*x*, -*y*, -*z*, 1; *x*, *y*+1/2, *z*+1/2, 1

B: *x*+1/2, *y*, *z*+1/2, 1 -B: -*x*, -*y*, -*z*, 1; *x*+1/2, *y*, *z*+1/2, 1

C: *x*+1/2, *y*+1/2, *z*, 1 -C: -*x*, -*y*, -*z*, 1; *x*+1/2, *y*+1/2, *z*, 1

I: *x*+1/2, *y*+1/2, *z*+1/2, 1 -I: -*x*, -*y*, -*z*, 1; *x*+1/2, *y*+1/2, *z*+1/2, 1

R: *x*+2/3, *y*+1/3, *z*+1/3, 1; *x*+1/3, *y*+2/3, *z*+2/3, 1 -R: -*x*, -*y*, -*z*, 1; *x*+2/3, *y*+1/3, *z*+1/3, 1; *x*+1/3, *y*+2/3, *z*+2/3, 1

H: *x*+2/3, *y*+1/3, *z*, 1; *x*+1/3, *y*+2/3, *z*, 1 -H: -*x*, -*y*, -*z*, 1; *x*+2/3, *y*+1/3, *z*, 1; *x*+1/3, *y*+2/3, *z*, 1

F: *x*+1/2, *y*+1/2, *z*, 1; *x*, *y*+1/2, *z*+1/2, 1; *x*+1/2, *y*, *z*+1/2, 1 -F: -*x*, -*y*, -*z*, 1; *x*+1/2, *y*+1/2, *z*, 1; *x*, *y*+1/2, *z*+1/2, 1; *x*+1/2, *y*, *z*+1/2, 1

X: 1**t**1 ; 1**t**2; …1**t**p -X: -*x*, -*y*, -*z*, 1; 1**t**1 ; 1**t**2; …1**t**p

**Table 2: Translation symbols t** (same as table A1.4.2.3. in HG-K)

Alphabetical symbols (given in the first column) specify translations along a fixed direction. Numerical symbols (given in the third column) specify translations as a fraction of the rotation order |N| and in the direction of the implied or explicitly defined axis. Putting several symbols together means the addition of the translation vectors corresponding to given symbols.

Translation Translation Subscript Fractional

Symbol Vector Symbol Translation

*a* 1/2, 0, 0 1 in 31 1/3

*b* 0, 1/2, 0 2 in 32 2/3

*c* 0, 0, 1/2 1 in 41 1/4

*n* 1/2, 1/2, 1/2 3 in 43 3/4

*u* 1/4, 0, 0 1 in 61 1/6

*v* 0, 1/4, 0 2 in 62 1/3

*w* 0, 0, 1/4 4 in 64 2/3

*d* 1/4, 1/4, 1/4 5 in 65 5/6

**Table 3: Operators** (similar to table A1.4.2.4. in HG-K)

Jones faithful operators for proper rotations along the three principal unit-cell directions are given below. For improper rotations (-2, -3, -4 and -6) the symbols are identical except that all signs are reversed. The identity 1 (*x*,*y*,*z*) and inversion centre -1 (-*x*,-*y*,-*z*) are omitted because they do not depend on the axis. The same holds for time reversal.

Rotation A-symbol Rotation order and Jones faithful symbol

Axis 2 3 4 6

**a** *x* *x*, -*y*, -*z* *x*, -*z*, *y*-*z x*, -*z*, *y x*, *y*-*z*, *y*

**b** *y* -*x*, *y*, -*z**z*-*x*, *y,* -*x z*, *y,* -*x z*, *y, z*-*x*

**c** *z* -*x*, -*y*, *z* -*y*, *x*-*y, z* -*y*, *x, z x*-*y*, *x, z*

**Table 4: Operators** (similar to table A1.4.2.5. in HG-K)

Jones faithful operators for proper two-fold rotations along the diagonal unit-cell directions.

Preceding Rotation Axis Jones faithful symbol

rotation symbol

*Nx* 2^ **b**-**c** -*x*, -*z*, -*y*

2” **b**+**c** -*x*, *z*, *y*

*Ny* 2^ **a**-**c** -*z*, -*y*, -*x*

2” **a**+**c** *z*, -*y*, *x*

*Nz* 2^ **a**-**b** -*y*, -*x*, -*z*

2” **a**+**b** *y*, *x*, -*z*

**Table 5: Operator** 3\*(similar to table A1.4.2.6. in S. R. Hall and R. W. Grosse-Kunstleve, 2010)

Jones faithful operator for proper three-fold rotation along the body diagonal unit-cell direction.

Rotation Axis Jones faithful symbol

symbol

3\* **a**+**b**+**c** *z*, *x*, *y*

**References**

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